EFFECT OF SURFACE POTENTIAL ON THE RATE OF ELECTRON TRANSFER STEP OF HYDROGEN ELECTRODE REACTION ON METALS

By


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Abstract

The galvanostatic transient method has been applied to the hydrogen evolution reaction on gallium, gold, silver, nickel and platinum metals in aqueous sodium hydroxide and gold in sulfuric acid, in order to determine the overvoltage component $\eta_1$ which is caused by charging up of the electric double layer at the metal-solution interface.

It has been found that (i) the hydrogen overvoltage $\eta$ is composed of two independent components $\eta_1$ and the other component $\eta_0$ which is caused by the change of the surface potential of the electrode due to specifically adsorbed Na atoms in alkaline solutions or H atoms in acid which is produced by the electron transfer step, and (ii) the polarization curve of the overall reaction $\log i$ vs. $\eta$ is divided into two independent parts $\log i$ vs. $\eta_1$ and $\log i$ vs. $\eta_0$ which correspond respectively to the polarization curve of the electron transfer step $\text{Na}^+ + e^- \rightarrow \text{Na}(a)$ or $\text{H}^+ + e^- \rightarrow \text{H}(a)$ and to that of the recombination of adsorbed H atoms. In alkaline solutions $\eta_0$ controls the activity of H(a) through the decomposition of water molecule by Na(a).

On the basis of the FRUMKIN's electrode potential theory $\eta_0$ has been given as the difference of the chemical potential of Na(a) or H(a) $-\Delta \mu_{Na}$ or $-\Delta \mu_{H}$ between the reversible and polarization states which means a compensation effect between the changes of the surface potential and the free energy of adsorption of Na(a) or H(a). As the result of the application of the compensation effect to the activation free energy of the electron transfer step which is given by the HORIUTI-POLANYI's rule, it has been shown by the rate theory developed by HORIUTI in the frame of the transition state method that the rate of the electron transfer step can be expressed by the Tafel equation in terms of $\eta_1$ independent of the surface potential of the electrode in accordance with the experimental results.

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**) Research Institute for Catalysis, Hokkaido University, Sapporo, Japan.
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Introduction

The hydrogen electrode reaction (h.e.r.) on metals may occur as a sequence of several elementary steps and the hydrogen overvoltage \( \eta \) may be composed of a series of components:

\[
\eta = \sum_j \eta_j, \tag{1}
\]

where \( \eta_j \) means the overvoltage of the \( j \)-th elementary step which may be given by the chemical affinity \( \Delta G_j \) and the stoichiometric number \( \nu_j \) of the step:

\[
2F\eta_j = \nu_j \Delta G_j, \tag{2}
\]

in which \( F \) is the Faraday and the coefficient 2 comes from the number of electrons in the h.e.r., and the rate of each individual step should not be given in general as a function of the overvoltage of the overall reaction which is measurable in the steady state method, but should be expressed in terms of its own component of overvoltage which may be different in the physical conception among different steps.

As pointed out by FRUMKIN\(^3\), the establishment of the electrode potential is caused by a number of factors, i.e. the adsorption of ions, neutral atoms and molecules, the orientation of surface active substances or solvent dipoles and the Galvani potential of the electrode. All these factors may also result in the build up of the hydrogen overvoltage by charging up of the electric double layer at the metal-solution interface or by the change of the surface potential of the electrode.

The overvoltage caused by charging up of the double layer has been introduced by FRUMKIN in the rate expression of the h.e.r.\(^3\), and the polarization characteristics in many electrochemical reaction systems have been successfully explained on the basis of this theory.

On the other hand, the surface states of a metal hydrogen electrode usually show complicated change under polarization which are still at present far from the comprehensive understanding because of the experimental difficulties, and we have little informations on the effect of the surface potential on the rate of the electron transfer step of the h.e.r. However in such complicated systems the overvoltage component caused by charging up of the double layer, which will be denoted hereafter as \( \eta_1 \), may be determined separately by a galvanostatic transient method at steady states of polarization. In the present work the polarization characteristics of the electron transfer step determined in this way on high and low overvoltage
metals will be discussed in connection with the change of the surface states of the electrodes.

The principle of the determination of the values of $\eta_1$ from the galvanostatic transient curves has been described previously\(^4,5,22\), but will be briefly described here. The polarization resistance $r_1$ of the electron transfer through the electric double layer which is defined as the derivative of $\eta_1$ with respect to the net rate of the electron transfer can be determined at a steady state of polarization with the current density $i$ and the overvoltage $\eta$ from the perturbation of $\eta$ by a current impulse, the transient curves observed being analyzed by the following equations\(^4,5\):

$$\ln \left( -\frac{\Delta i}{\dot{\eta}} \right) = \frac{t}{\tau_1} + \ln C_D, \quad (3)$$

$$\tau_1 = C_D r_1, \quad (4)$$
in which $\Delta i$ is the height of the current pulse, $\dot{\eta}$ is the time derivative of the transient curve observed, $C_D$ is the differential capacity of the double layer and $\tau_1$ is the time constant of the electron transfer step. These equations have been derived assuming a simple equivalent circuit for the metal-solution interface as shown in Fig. 1. The assumption of this simple model for the electrode-solution interface will not affect the validity of deducing the kinetic parameters of the electron transfer step from the transient curves insofar as these curves are observed through a very early time window.

The value of $\eta_1$ can be obtained on the basis of the definition of $r_1$ from the following integral equation as a function of the current density $i$, taking into consideration that the net rate of the electron transfer step in a steady state equals $i$\(^5\):

$$\eta_1 = -\int_{0}^{i} r_1 di. \quad (5)$$

In this way we obtain the po-
polarization curve of the electron transfer step separately. From the comparison of the polarization curve of the electron transfer step \( \log i \) vs. \( \eta \), and that of the overall reaction \( \log i \) vs. \( \eta \), we can obtain one more polarization curve \( \log i \) vs. \( \eta_a \) where \( \eta_a = \eta - \eta_l \), which may offer informations on the overvoltage components due to the change of the surface potential.

The polarization characteristics of the electron transfer step will be discussed in the following sections comparing three kinds of polarization curves on gallium, gold, nickel, silver and platinum hydrogen electrodes in alkaline solutions and on gold in sulfuric acid.

**Experimental results**

(1) The gallium hydrogen electrode in sodium hydroxide solutions

The liquid gallium electrode of 99.9999\% purity (geometrical area 1 cm\(^2\)) shows a rest potential around \(-550\) mV of \( \eta \) in aqueous sodium hydroxide solutions under atmospheric hydrogen bubbling at 32°C, but in the hydrogen overvoltage region the polarization curves have a slope of 120 mV as shown in Fig. 2, and show a good coincidence with those obtained by BAGOZKAYA and POTAPOVA \(^6\).

The pH-dependence of \( \eta \) at \( i=10^{-3}\) amp/cm\(^2\) shows a gradient 125 mV as seen from Fig. 3, in accordance with the water molecule discharge me-

![Fig. 2. Polarization curves of the overall reaction on Ga in NaOH solutions at 32°C; thick line—BAGOZKAYA and POTAPOVA \(^6\).](image-url)

![Fig. 3. pH-dependence of overvoltage at 10\(^{-3}\) A/cm\(^2\); (△)—BAGOZKAYA and POTAPOVA \(^6\).](image-url)
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Fig. 4. \(\log(-\Delta i/\Delta t)\) vs. \(t\) relations obtained from potential-time curves on Ga in 0.01 N NaOH; the starting overvoltage in mV, (1) -1050; (2) -1008; (3) -965; (4) -932; (5) -895; (6) -720.

Fig. 5. \(\eta\)-dependence of the time constant of the electron transfer step \(\tau_1\) on Ga in NaOH solutions; (●) 0.01 N, (×) 0.075 N.

Fig. 6. The relation between the polarization resistance of the electron transfer step and the current density on Ga in 0.01 N NaOH solution.

Fig. 7. Polarization curves of the electron transfer step (Δ), and of the overall reaction (line) in NaOH solutions.
chanism proposed by Frumkin, Korshunov and Bagozkaya on gallium.

The experimental results of the galvanostatic transient study are shown in Figs. 4～8. It can be seen from Fig. 4 that the values of \( \log (\Delta i/\dot{\eta}) \) calculated from the transient curves change linearly with time in accordance with Eq. (3), and the double layer capacity has a constant value 16 \( \mu \text{F/cm}^2 \) independent of \( \eta \), but the values of \( \tau_i \) strongly depend on \( \eta \) as shown in Fig. 5. A typical \( \tau_i \) vs. \( i \) curve in 0.01 N NaOH is shown in Fig. 6.

The polarization curves of the electron transfer step \( \log i \) vs. \( \eta_i \) in NaOH solutions obtained from the integration of the \( \tau_i \) vs. \( i \) curves according to Eq. (5) are illustrated in Fig. 7 with a symbol of triangle, and the polarization curves of the overall reaction \( \log i \) vs. \( \eta \) in the same solutions are also shown in this Figure by straight lines. It can be seen that the two kinds of the polarization curves coincide with each other in every solution used. It follows from this fact that \( \eta \) equals \( \eta_i \), i.e., the hydrogen overvoltage is entirely caused by charging up of the double layer on gallium in NaOH solution.

This conclusion is also supported by the result of the comparison of the \( \log i \) vs. \( \eta \) curve of the overall reaction with the polarization curve of the electron transfer step \( \log i_t \) vs. \( \eta \) obtained from the decay of overvoltage after switching off the polarizing current, as will be seen in the following. In the course of the potential decay the rate of the electron transfer \( i_t \) may be given by the following equation if the overvoltage is caused by the change of the free charge density on the electrode surface:

\[
i_t = C_\theta \dot{\eta}.
\]

The values of \( \log i_t \) determined in this way are plotted against the values of overvoltage at which \( \dot{\eta} \) is calculated in the course of decay in Fig. 8. As seen from this Figure, \( \log i_t \) vs. \( \eta \) curve in the course of decay coincides with \( \log i \) vs. \( \eta \) curve in steady states.

In order to find the chemical species which accepts a metal electron in the double layer, the values

![Fig. 8. Polarization curves of the electron transfer step obtained from the potential decay after switching off the current (○) and (△), and of the overall reaction (line).](image-url)
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of \( \tau_1 \) at a constant potential drop between the metal surface and the Helmholtz plane are plotted against the pH or pNa of the solution in Fig. 9, the Frumkin correction being used for the calculation of the potential of the Helmholtz plane \( \psi_1 \). Taking into consideration the constancy of the value of \( C_0 \), the values of \( \tau_1 \) are used for convenience as a measure of the rate of the electron transfer step, since its value is measurable independent of the surface area of the electrode. As seen from this Figure, the values of \( \tau_1 \) at a constant potential drop across the compact double layer show no pH nor pNa dependence. It can be concluded from this fact that the metal electron transfers to a water molecule in the double layer. All the evidences obtained from the galvanostatic transient study confirm the mechanism of the water molecule discharge proposed by Frumkin, Korshunov and Bagozkaya and the hydrogen overvoltage is entirely caused by charging up of the double layer.

(2) The gold hydrogen electrode in sulfuric acid and sodium hydroxide solutions

The mechanism of the hydrogen electrode reaction on gold has been discussed by several authors based mainly on the \( b \)-value of the Tafel slope \( 9,10,11 \), but the mechanism has not yet been generally accepted because of its wide variety among different authors. The mechanism of the establishment of the hydrogen overvoltage and the reaction mechanism will be discussed here from the comparison of the polarization characteristics between the individual step of the electron transfer and the overall reaction in sulfuric acid and sodium hydroxide solutions. Special attention was paid for the preparation and purification of the electrode metal and the solutions. A gold foil degassed or gold film prepared in vacuum was used as the working electrode and the electrolytic cell was made of silica glass.

(2 a) Polarization curves in sulfuric acid.
Sulfuric acid was prepared by dissolution of SO₂ gas into water in a vacuum system. The gold hydrogen electrode showed very high activity and good reproducibility after an anodic treatment at about 1.8 V (vs. R.H.E.) for a few seconds. The double layer capacity was found to be constant independent of the value of η and the concentration of the solution from the galvanostatic transient study, whereas the time constant τ₁ changed remarkably depending on η and the solution concentration.

Fig. 10 shows the polarization curves of the electron transfer step log i vs. η₁ determined by the galvanostatic transient method and log i vs. η curves of the overall reaction.[29] The polarization curve of the overall reaction in 1 N H₂SO₄ shows a straight line with the b-value of 29 mV, but in dilute solutions it shows a deviation from this line and the b-value tends to approach to 120 mV with the increase of the current density.

As seen from Fig. 10, the rate of the electron transfer step i₁ can be expressed by the following equation in terms of η₁:

$$\log i = -k \eta_1$$

Fig. 10. Polarization curves of the electron transfer step log i vs. η₁, and of the overall reaction log i vs. η on Au in sulfuric acid, (●) 1.3 N, (○) 0.67 N and (□) 0.30 N, and (---) theoretical curves calculated by Eq. (7).
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\[ i_1 = i_{10} \left\{ \exp \left( -\alpha F\eta_1/RT \right) - \exp \left( \beta F\eta_1/RT \right) \right\}, \quad (7) \]

with the transfer coefficient \( \alpha = \beta = 0.5 \). The exchange current density of the electron transfer step \( i_{10} \) is obtained from the polarization resistance at the reversible potential by the equation:

\[ i_{10} = RT/r_{10} F. \quad (8) \]

The polarization resistance \( r_{10} \) is obtained from the time constant \( \tau_{10} \) at the reversible potential according to Eq. (4). \( \log \tau_{10} \) is found to change linearly with the pH of the solution with a gradient 0.5 as shown in Fig. 11, and consequently \( i_{10} \) is proportional to the square root of the concentration of the hydronium ion. It may be concluded from these facts that the electron transfer step is the discharge of the hydronium ion, and its rate can be expressed by Eq. (7) in terms of the overvoltage component caused by charging up of the double layer.

As seen from the comparison of \( \eta_1 \) and \( \eta \) at a constant current density in Fig. 10, \( \eta_1 \) is negligibly small at low current densities, but with the increase of the current density it grows to such an extent that \( \eta_1 \) accounts for the considerable part of \( \eta \) particularly in dilute solutions.

![Fig. 11. Time constant \( \tau_{10} \) of the electron transfer step at the reversible potential on Au plotted against pH in acid, or pNa in alkaline solutions, (⊗) in NaOH+Na2SO4 solutions of pH 10.8 and 12.1.](image-url)
From the two kinds of polarization curves in Fig. 10, it is possible to deduce log \( i \) vs. \( \eta_a \) curves which may reveal the mechanism of the setting up of the surface potential. As seen from Fig. 12, the log \( i \) vs. \( \eta_a \) curves practically coincide with each other with the \( b \)-value of 29 mV and the value of \( \eta_a \) has a tendency to approach a saturation value of ca. −50 mV independent of the concentration of the solution. These facts may be explained by assuming that \( \eta_a \) is caused by the adsorption of hydrogen atoms on the electrode surface which result in the change of the surface potential due to dipoles with the positive end toward the solution side, and that its surface coverage is very low but it attains saturation at \( \eta_a \) about −50 mV.

From the comparison of the log \( i \) vs. \( \eta_1 \) and log \( i \) vs. \( \eta_a \) curves, it may be concluded that \( \eta_1 \) grows according to Eq. (7) but \( \eta_a \) attains a saturation value with the increase of the current density, or in other words, the rate-determining step switches over from the recombination of adsorbed hydrogen atoms to the discharge of hydronium ion. This conclusion contradicts to the mechanism of KUHN and BYRNE\(^{10}\) in which the transition of the rate-determining step occurs from the recombination of adsorbed hydrogen atoms.
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to the electrochemical desorption.

(2 b) *Polarization curves in sodium hydroxide solutions.*

The gold hydrogen electrode in aqueous NaOH shows a wide variety of the polarization curves of the overall reaction as shown by the shaded area in Fig. 13. However it gives good reproducible results on the log $i$ vs. $\eta_1$ curves, which enable us to deduce some definite conclusions on the kinetic behavior of the electron transfer step.

Fig. 13 shows examples of the polarization curves log $i$ vs. $\eta_1$. These curves can also be expressed by Eq. (7), and $i_{10}$ is now proportional to the square root of Na$^+$ ion independent of the pH of the solution, as seen from Fig. 11. It may be concluded that the electron transfer step is the discharge of Na$^+$ ion, in contrast to the case of the gallium hydrogen electrode.

The change of $\tau_1$ with $\eta$ is shown in Fig. 14. As seen from the Figure, $\tau_1$ remains at a constant value which equals $\tau_{10}$ in the wide range of overvoltage from 0 to $-300$ or $-400$ mV, i.e. the electron transfer step can be regarded as very close to its reversible state in spite of the remarkable

![Diagram](image-url)
growth of $\gamma$. This fact suggests that the electron transfer step occurs independent of the surface states of the electrode. This may be the reason why a definite polarization curve of good reproducibility can be obtained with respect to the electron transfer step in spite of the wide variety of the polarization curves of the overall reaction. The rate equation and the effect of the surface potential on the rate of the electron transfer step will be discussed in detail in a later section.

(3) **The silver hydrogen electrode in sodium hydroxide solution**

A silver electrode of a bead shape of geometrical area $0.2 \text{ cm}^2$ was used after heat treatment in a stream of helium gas. The silver electrode prepared in this way showed good reproducibility in NaOH solutions.

Fig. 15 shows the log ($-\Delta i/\gamma$) vs. time curves obtained at different overvoltages in $5.6 \cdot 10^{-8} \text{ N NaOH}$. The values of $C_D$ given by $-\Delta i/\gamma$ at $t=0$ are plotted against $\gamma$ in Fig. 16. It can be seen from this Figure that the $\gamma$-dependence of $C_D$ on silver is different from that on gallium or gold. There exists a minimum in the $C_D$-curve in Fig. 16 at about $-0.9 \text{ V (vs. N.H.E.)}$ in agreement with the result of Leikis et al., which may corresponds to
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Fig. 15. \( \log(-\Delta v/\eta) \) vs. \( t \) curves on Ag in 5.63x10^{-3} N NaOH, the starting overvoltage of the transient in mV, (1) 0; (2) -100; (3) -200; (4) -240; (5) -300; (6) -330; (7) -400; (8) -490.

Fig. 16. Potential dependence of the double layer capacitance of Ag in 5.63x10^{-3} N NaOH obtained from Fig. 15.
Fig. 17. Time constant of the electron transfer step on Ag obtained from Fig. 15.

Fig. 18. log $i$ plotted against $\tau_1$ (○), $\eta$ (●), and $\eta_a$ (△) on Ag in 5.63 x 10^{-3} N NaOH.

The $\tau_1$ vs. $\eta$ curve shows a gradual decrease as in the case of gallium with no plateau, as seen from Fig. 17. Fig. 18 shows the log $i$ vs. $\eta_1$ curve obtained from Fig. 17 together with the log $i$ vs. $\eta$ of the overall reaction.
in a NaOH solution. As seen from this Figure, the log $i$ vs. $\eta_1$ curve can also be expressed by Eq. (7) and $\eta_1$ cannot be neglected as compared with $\eta$ except in the reversible region, and $\eta_1$ accounts for a greater part of $\eta$ at higher current densities. The log $i$ vs. $\eta_a$ curve obtained from the two kinds of polarization curves is also shown in Fig. 18 with the symbol of triangle, which has the $b$-value of 60 mV and the saturation value of $\eta$ around $-0.14$ V as in the case of gold in sulfuric acid.

It is interesting to compare the log $i$ vs. $\eta_a$ curve in Fig. 18 with the polarization curve on silver observed by KRISHTALIK and BYSTROV\(^{15}\) at low current densities in hydrochloric acid. It is found that the two polarization curves have the same $b$-value of 60 mV and are almost in the same region of the current density and overvoltage. It may be concluded from this fact that $\eta_a$ is attributed to the recombination of adsorbed hydrogen atoms, the existence of which has already been demonstrated by SHUMILOVA et al.\(^{16}\)

(4) **The nickel hydrogen electrode in alkaline solutions**

The nickel hydrogen electrode in alkaline solutions has been extensively studied and its polarization characteristics has been explained on the basis of the catalytic mechanism\(^{17}\) in which the recombination of adsorbed hydrogen atoms is rate-determining. However the kinetic study of the individual steps of the reaction by the galvanostatic transient method has required some modification as to the mechanism of the establishment of the hydrogen overvoltage and the reaction mechanism.

(4 a) *The polarization curves in sodium hydroxide solutions.*

The polarization curves of the electron transfer step log $i$ vs. $\eta_1$ and those of the overall reaction log $i$ vs. $\eta$ in aqueous NaOH are shown in Fig. 19.\(^{18}\) As seen in this Figure, the Tafel line of the overall reaction has a break in any concentration which suggests the change of the reaction mechanism. In the region lower than the break point the Tafel line has the $b$-value of ca. 80 mV and is independent of the solution concentration, but after the break point the $b$-value changes to 120 mV and the Tafel line is lowered with the increase of the solution concentration.

The polarization curve of the electron transfer step log $i$ vs. $\eta_1$ in Fig. 19 can also be expressed by Eq. (7) with $\alpha=\beta=0.5$, and the value of $\eta_1$ is negligibly small at current densities lower than that at the break point of the Tafel line of the overall reaction, at which it becomes detectable, and at higher current densities the Tafel line of the electron transfer step becomes parallel to that of the overall reaction, as seen from Fig. 19. On the other hand, log $\tau_{10}$ changes linearly with log $C_{Na^+}$ with a coefficient 0.5 in NaOH
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Fig. 19. Polarization curves of the electron transfer step \( \log i \) vs. \( \eta \) (filled symbols), and of the overall reaction (open symbols) on Ni in NaOH solutions.

or Na\(_2\)SO\(_4\) solutions as in the case of gold in alkaline solutions, as seen in Fig. 21.\(^{180}\) It may be concluded from these facts that the electron transfer step is composed of the discharge of Na\(^+\) ion.

The \( \log i \) vs. \( \eta_a \) curves calculated from \( \log i \) vs. \( \eta_l \) and \( \log i \) vs. \( \eta \) curves in Fig. 19 are shown in Fig. 20. As seen from this Figure, these polarization curves nearly coincide with each other and also with the polarization curve of the overall reaction in the region lower than the break point, but the value of \( \eta_a \) has a tendency to approach a saturation value with the increase of the current density as in the case of gold and silver. On the other hand, it has been found from the concentration dependence of the pseudo-capacitance of the nickel hydrogen electrode in alkaline solutions that the key intermediate which causes \( \eta_a \) has been determined to be Na atom adsorbed on the electrode surface, but not adsorbed hydrogen atoms.

On the basis of the arguments mentioned above it may be concluded that the hydrogen electrode reaction occurs as the sequence of the following elementary steps:
Electron Transfer Step of Hydrogen Electrode Reaction

(1) \[ \text{Na}^+ + e^- \rightarrow \text{Na}(\text{a}) \]
(II) \[ \text{Na}(\text{a}) + \text{H}_2\text{O} \rightarrow \text{H}(\text{a}) + \text{OH}^- + \text{Na}^+ \]
(III) \[ 2\text{H}(\text{a}) \rightarrow \text{H}_2 \]

in which the step (II) can be regarded as practically in equilibrium, but the

Fig. 20. \( \log i \) vs. \( \eta_a \) curves obtained from Fig. 19.

Fig. 21. Concentration dependence of the time constant \( \tau_{10} \) on Ni in alkaline solutions.
Fig. 22. \( \eta \)-dependence of the time constant \( \tau_1 \) on Ni in NaOH solutions.

Fig. 23. Polarization curves of the electron transfer step \( \log i \) vs. \( \eta_1 \) and of the overall reaction on Ni in Ba(OH)\(_2\) solutions.

Fig. 24. \( \log i \) vs. \( \eta_1 \) curves on Ni in alkaline solutions containing Ba\(^{2+}\) ion.
transition of the rate-determining step occurs from step (III) to (I) with the increase of the current density, on which basis the break of the Tafel line of the overall reaction may be explained.

The $\eta$-dependence of $\tau_1$ is shown in Fig. 22. As seen from this Figure, a plateau appears in the $\tau_1$ vs. $\eta$ curve as in the case of gold, which also

![Fig. 25. Polarization curves of the electron transfer step in the reverse direction $\log i$ vs. $\eta_1$ on Ni in Ba(OH)$_2$ solutions.](image)

![Fig. 26. Surface concentration of Ba or Ca atom under polarization on Ni determined by atomic absorption spectroscopy.](image)
means that the electron transfer step is close to the reversible state in spite of the remarkable growth of $\eta$.

(4 b) The polarization curves in barium hydroxide solutions

In aqueous Ba(OH)$_2$ the polarization curves of the overall reaction are quite similar to those in aqueous NaOH as shown in Fig. 23$^{19}$, and the rate of the electron transfer step can also be expressed by Eq. (7) with $\alpha=0.68$, as seen from Fig. 24.$^{19}$

The rate of the electron transfer step in the reverse direction $\bar{i}$ can be determined as a function of $\eta_1$ from the following equations:

$$\bar{i} = \bar{i}_0 \exp \left( -\frac{aF\eta_1}{RT} \right), \quad \alpha = 0.68. \quad (9)$$

The polarization curves log $\bar{i}$ vs. $\eta_1$ determined in this way are shown in Fig. 25. The transfer coefficient $\beta$ in the reverse direction is found to be 1.37 from Fig. 25 and the sum $\alpha + \beta$ is quite close to 2 which may be obtained as the result of the discharge of Ba$^{n}$ ion. In fact, the intermediate species of the electrolytic hydrogen evolution in aqueous Ba(OH)$_2$ has been found to be Ba atom by atomic absorption spectroscopy$^{20}$, and its surface concentration varies with $\eta$ as shown in Fig. 26. It may be concluded that the hydrogen electrode reaction on nickel in Ba(OH)$_2$ solutions occurs in the same frame of the reaction mechanism as in the aqueous NaOH with the intermediate species of Ba atom which causes the change of the surface potential of the electrode.

(5) The platinum hydrogen electrode in alkaline solutions

The polarization characteristics of the electron transfer step has been found quite similar to those on nickel in alkaline solutions$^{4,5}$. The value of $\tau_0$ in NaOH, LiOH$^{20}$ and CsOH$^{20}$ is also proportional to the square root of the cation concentration as shown in Fig. 27, and it does not show a remarkable distinction among alkali ions, but slightly increases in the order Li$^+ <$ Na$^+ <$ Cs$^+$, for instance, 75, 90 and 120 $\mu$sec in 0.1 N concentration respectively.

The polarization curves of the electron transfer step log $i$ vs. $\eta_1$ can also be expressed by Eq. (7) and its exchange rate is proportional to the square root of the cation concentration, but the log $i$ vs. $\eta_n$ curves do not depend on the solution concentration and have the $b$-value of 60 mV, as exemplified in Fig. 28 in the case of CsOH solutions.

The adsorption capacitance determined from the potential decay after switching off the current has shown a strong dependence on the cation concentration with a sharp maximum at the potential ca. $-800$ mV (vs. N.H. E.$^{21,22,23}$).
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Fig. 27. Concentration dependence of $\tau_{10}$ on Pt in LiOH (■), NaOH (○), CsOH (△), and LiOH+Li₂SO₄ (□), NaOH+Na₂SO₄ (●), CsOH+Cs₂SO₄ (▲).

Fig. 28. (a)—Polarization curves of the electron transfer step log $i$ vs. $\eta_i$, (b)—Polarization curves log $i$ vs. $\eta_a$ on Pt in CsOH, (○) pH 12.8, (△) pH 11.86 and (□) 10.86.
Fig. 29. \( \eta \)-dependence of the time constant \( \tau \) on Pt in CsOH of the same concentration as in Fig. 28.

All these evidences for the platinum hydrogen electrode in alkaline solutions are in favor of the mechanism on nickel described in the previous section. The \( \tau \) vs. \( \eta \) curve also has a plateau as shown in Fig. 29.

The role of the surface potential in the rate equation of the electron transfer step

It has already been established by Frumkin\textsuperscript{24} and Krishtalik\textsuperscript{25} from the theoretical point of view and experimentally by Fedorovich et al.\textsuperscript{26} that for the comparison of the rate of the electron transfer between metals of different work function on which no change of the surface potential is observed under polarization, the rate of the reaction should be compared at a constant potential referred to a definite reference electrode, but not at a potential referred to the zero charge potential of the electrode, or in other words, the work function of the metal electrode into vacuum does not appear explicitly in the rate equation of the electron transfer reaction, but the work function into solution plays an important role in the promotion of the rate of the electron transfer reaction.
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However, the situation seems to be quite different in the case of the hydrogen evolution reaction on low overvoltage metals. As already described in the previous sections, the overvoltage component appears on low overvoltage metals which is caused by the change of the surface potential due to the specific adsorption of the intermediate species on the electrode surface, for instance, H (a) on gold in sulfuric acid, or Na (a) on gold, nickel and platinum in sodium hydroxide solutions. The rate of the electron transfer step in these systems, however, can simply be expressed experimentally by the Tafel equation in terms of the overvoltage component due to charging up of the double layer, but not in terms of the overvoltage of the overall reaction referred to the reversible hydrogen electrode. These experimental results seem to be contradictory to the conclusion attained by Frumkin et al.

In the present section the two kinds of the overvoltage components which are caused respectively by charging up of the double layer and the change of the surface potential will be given quantitatively on the basis of the electrode potential theory developed by Frumkin and his school in the reversible systems, and the effect of the change of the surface potential due to the specific adsorption of the intermediate species on the rate of the electron transfer will be discussed on the basis of the rate theory developed by Horiuti in the frame of the transition state method for the heterogeneous systems. As a result it will be shown that the change of the free energy of adsorption of the intermediate species with the increase of the surface
coverage will be compensated in the expression of the activation free energy by the change of the surface potential caused by adsorbed dipoles as seen from Fig. 30, and consequently the surface potential does not appear in the rate equation of the electron transfer step.

(1) The two kinds of overvoltage components $\eta_1$ and $\eta_a$

As described in the previous sections, there appears a plateau in the $\tau_1$ vs. $\eta$ curves, for instance, on gold, nickel and platinum in sodium hydroxide solutions, as shown respectively in Figs. 14, 22 and 29, and the values of $\eta_1$ is negligibly small and the growth of overvoltage in this region can be attributed substantially to $\eta_a$, as seen from the comparison of the polarization curves of the electron transfer step and the overall reaction in Figs. 13 and 14. Therefore the electron transfer step can be regarded as practically in equilibrium in this polarization region:

$$Na^+ + e^- \rightarrow Na(a),$$

and the charge density on the electrode surface may be kept nearly constant although the electrode potential caused by $Na(a)$ shows a marked growth in this region.

In such a quasi-reversible system the Frumkin’s theory may be applicable for the estimation of the amount of the growth of the electrode potential caused by charging up of the double layer which will be denoted as $\delta \varphi_0$, and that caused by the change of the surface potential due to the adsorbed $Na(a)$ which will be denoted as $\delta \varphi$.

These two quantities have been found on the basis of the Frumkin’s theory to be mutually independent and given by the following equations respectively in terms of the chemical potentials of $Na^+$ ion in the solution and the adsorbed $Na$ atom:

$$F \delta \varphi_0 = \delta \mu_{Na^+}, \quad (11)$$
$$F \delta \varphi = -\delta \mu_{Na}, \quad (12)$$

These equations mean that the Nernst equation which represents the equilibrium relation of the electron transfer step in the differential from

$$\delta \mu_{Na^+} + \delta \mu_{Na} = \delta \mu_{Na(a)} \quad (13)$$
can be divided into two independent parts. As seen from Eq. (11), $\varphi_0$ is constant in a given solution of definite concentration of $Na^+$ ion.

$\eta_1$ is given by the difference of the electrostatic potential caused by charging up of the double layer between the polarized and reversible states,

$$\eta_1 = \varphi - \varphi_0, \quad (14)$$
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and \( \eta_a \) is given by the difference of \( \omega \) between the reversible and the polarized states according to Eq. (12) as,

\[
\eta_a = \omega - \omega_0,
\]
where suffix 0 means the reversible state.

In the frame of the reaction mechanism in sodium hydroxide solutions described in the previous section, \( \eta_1 \) turns to the affinity of the electron transfer step (I) \( \Delta G_1 \),

\[
\Delta G_1 = \mu_{Na} - \mu_{Na^+} - \mu_e = (\mu_{Na} - \mu_{Na_{a,0}}) - (\mu_e - \mu_{e,0})
\]
\[
= -F(\eta_a - \eta) = F\eta_1,
\]
and \( \eta_a \) turns to the affinity of the recombination of the adsorbed hydrogen atoms \( \Delta G_s \) through the equilibration of the step (II),

\[
\eta_a = \mu_{Na_{a,0}} - \mu_{Na} = \mu_{H_{a,0}} - \mu_{H} = \frac{1}{2} \Delta G_s .
\]

In the case of the discharge of hydronium ion which produces adsorbed hydrogen atom as an intermediate, the components of the electrode potential have been found to be expressed as \( \delta \)

\[
F\delta \varphi_0 = \delta \beta_{H^+},
\]
\[
F\delta w = -\delta \mu_H.
\]

(2) The rate equation of the electron transfer step

The effect of the change of the surface potential on the rate of the electron transfer step will be discussed here on the basis of Eq. (12) in the frame of the transition state method developed by HORIUTI for the heterogeneous reaction systems.

The electron transfer step in the sodium hydroxide solutions will be taken as an example for the rate expression:

\[
Na^+ + e^- \rightarrow Na(a), \quad (1)
\]
in which \( Na^+ \) is the sodium ion on the Helmholtz-plane in the electric double layer at the metal-solution interface. The absolute rate of the step (I) in the right or reverse direction will be given by the following equation in the frame of the transition state method:

\[
\tilde{\gamma} = \gamma_0 \frac{kT}{\hbar} \exp\left( - \frac{\rho^* - \rho_{H}}{RT} \right)
\]
or

\[
\tilde{\gamma} = \gamma_0 \frac{kT}{\hbar} \exp\left( - \frac{\rho^* - \rho_{H}}{RT} \right)
\]
where \( e_0 \) is the elementary charge of an electron, \( \kappa \) is the transmission coefficient, \( k \), \( T \) and \( h \) are respectively the Boltzmann constant, absolute temperature and the Planck constant, \( \beta^* \), \( \beta_i \) and \( \beta_f \) are respectively the electrochemical potential of the activated complex and the initial and final states,

\[
\beta_i = \beta_{Na^+} + \beta_e \tag{20}
\]

\[
\beta_f = \mu_{Na} . \tag{21}
\]

\( \beta^* \), \( \beta_{Na^+} \) and \( \mu_{Na} \) are given by the statistical mechanical theory of Horiuti \( ^2, ^3, ^4 \) as:

\[
\beta^* = e^* - RT \ln G \theta^*_0 , \tag{22}
\]

\[
\mu_{Na} = \varepsilon_{Na^+} + RT \ln \frac{\theta}{1 - \theta} , \tag{23}
\]

\[
\beta_{Na^+} = \varepsilon_{Na^+} + RT \ln C_{Na^+D} \tag{24}
\]

where \( e^* \), \( \varepsilon_{Na^+} \) and \( \varepsilon_{Na} \) denote respectively the reversible work per mole required to bring an activated complex, \( Na(a) \) and \( Na^+ \) ion to the definite adsorption site from their standard states, \( \theta^*_0 \) is the surface coverage of the vacant site for the activated complex which equals unity practically, \( G \) is the number of adsorption site for the activated complex per unit area, \( \theta \) is the surface coverage of \( Na(a) \), and \( C_{Na^+D} \) is the concentration of \( Na^+ \) ion on the Helmholtz-plane which is given by the concentration of \( Na^+ \) ion in the solutions as:

\[
C_{Na^+D} = C_{Na^+} e^{-\frac{\varepsilon_{Na^+}}{RT}} , \tag{25}
\]

where \( \phi_1 \) is the electrostatic potential at the Helmholtz-plane referred to that in the bulk of the solution.

On the other hand, the second term of the right hand side of Eq. (23) can be neglected as compared with the first term when \( \theta \) is not close to zero nor unity, i.e.,

\[
|\varepsilon_{Na} - RT \ln \frac{\theta}{1 - \theta} | . \tag{26}
\]

In the frame of this approximation the rate equation is written by Eqs. (18) \( \sim \) (24) as:

\[
\dot{\gamma}_1 = e_0 k T \frac{k}{h} G \cdot C_{Na^+} e^{-\frac{\varepsilon_{Na^+}}{RT}} \cdot \exp \left[ - \left( e^* - \varepsilon_{Na^+} - \mu_e \right) / RT \right] , \tag{27}
\]

in which \( e^* - \varepsilon_{Na^+} - \mu_e \equiv \Delta \varepsilon \) is the activation free energy of the reaction. According to the theory of Horiuti and Polanyi \( ^{30} \) for the change of the activation free energy \( \delta \left( \Delta \varepsilon \right) \), we have
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\[ \delta (\Delta \theta) = \alpha \delta (\varepsilon_{Na} - \varepsilon_{Na^+} - \mu_e) \]

and in the integral form

\[ \Delta \theta = \Delta \theta_0 + \alpha (\varepsilon_{Na} - \varepsilon_{Na^+} - \mu_e), \tag{28} \]

where \( \Delta \theta_0 \) is a constant which means the activation free energy at the electrode potential \( \mu_e = \varepsilon_{Na} - \varepsilon_{Na^+} \), and it does not depend on the surface potential of the electrode and equals in the right and reverse directions.

On the other hand, we can obtain the following relation between \( \varepsilon_{Na} \) and \( \omega \) from Eqs. (12) and (23) within the accuracy of Eq. (26):

\[ F \Delta \omega + \Delta \varepsilon_{Na} = 0. \tag{29} \]

This equation means a compensation effect between the change of the surface potential and that of the free energy of the adsorbed atom at a definite adsorption site, being equally reflected respectively in the free energy curves of the initial and final states, as illustrated in Fig. 30.

Furthermore \( \mu_e, \varepsilon_{Na} \) and \( \varepsilon_{Na^+} \) can be expressed as,

\[ \mu_e = -F(\varphi + \omega_0 + \Delta \omega), \tag{30} \]

\[ \varepsilon_{Na} = \varepsilon_{Na,0} + \Delta \varepsilon_{Na}, \tag{31} \]

\[ \varepsilon_{Na^+} = \varepsilon_{Na^+,0} + F\psi_1, \tag{32} \]

where \( \varepsilon_{Na,0} \) is \( \varepsilon_{Na} \) on the bare surface, \( \omega_0 \) is the work function of the base metal into vacuum and \( \varepsilon_{Na^+,0} \) is \( \varepsilon_{Na^+} \) at \( \psi_1 = 0 \). Putting Eqs. (29)~(32) into Eq. (28), we obtain the activation free energy as:

\[ \Delta \theta = \Delta \theta_0 + \alpha (\varepsilon_{Na,0} - \varepsilon_{Na^+,0}) + F\omega_0 + aF\varphi, \tag{33} \]

and the rate equation from Eqs. (27) and (33) as:

\[ \dot{I}_t = \dot{k} C_{Na^+} e^{-\frac{1 - a \nu_1}{RT_1}} e^{-\frac{\omega_0}{RT}}, \tag{34} \]

where

\[ \dot{k} = e_0 kT \frac{kT}{h} G \exp \left[ -\alpha (\varepsilon_{Na,0} - \varepsilon_{Na^+,0})/RT \right] e^{-\frac{\Delta \theta_0}{RT}} e^{-\frac{\omega_0}{RT}}, \tag{35} \]

is a constant which does not depend on the surface potential of the electrode. It can be seen from Eq. (34) that the rate of the electron transfer step does not depend on the surface potential of the electrode caused by the specific adsorption of the intermediate species, but depend only on the electrostatic potential caused by charging up of the double layer as the result of the compensation between the adsorption free energy and the surface potential. This is the reason why the polarization curve of the electron transfer step can be
determined separately independent of the surface potential of the electrode.

In the special case where there is no change of the surface potential under polarization, Eq. (34) can be reduced simply to the equation derived by FRUMKIN\textsuperscript{24},

\[
\tilde{i}_1 = \tilde{k} e^{-\frac{F(1-\phi)_1}{R T}} \cdot e^{-\frac{q_e}{k T}} ,
\]

since in this case \( \beta_e \) equals simply \(-F(w_0 + \phi)_1\), and the rate of the electron transfer step on gallium can be expressed by this equation, as shown in the previous section.

The rate of the electron transfer step in the reverse direction can be expressed in the same procedure:

\[
\tilde{i}_1 = \tilde{k} e^{-\frac{F\phi_0}{R T}} \cdot e^{-\frac{\mu}{k T}}
\]

and

\[
\tilde{k} = e_k \frac{k T}{h} G \frac{\theta}{1-\theta} \cdot e^{-\frac{J_{Na}^*}{R T}} \cdot \exp\left[-\beta (\varepsilon_{Na^+} - \varepsilon_{Na})/R T\right] \cdot e^{-\frac{\mu_{Na^+}}{k T}}.
\]

It can readily be seen from Eqs. (34) and (37) that

\[
RT \ln \frac{\tilde{i}_1}{\tilde{k}_1} = -F\eta_1 ,
\]

and also from Eqs. (18) and (19):

\[
RT \ln \frac{\tilde{i}_1}{\tilde{k}_1} = \beta_{Na^+} + \beta_e - \mu_{Na} \equiv -\Delta G_1 ,
\]

and consequently we obtain

\[
F\eta_1 = \Delta G_1 ,
\]

which has already been derived in Eq. (16) from the thermodynamic consideration of the overvoltage components.

The exchange rate of the electron transfer step \( i_{10} \) can be expressed by Eq. (34) as:

\[
i_{10} = kC_{Na^+} \cdot e^{-\frac{aF\phi_1}{k T}} ,
\]

assuming that \( \phi_1 \)-potential may reasonably be regarded as constant in the case of low overvoltage metals, since the greater part of the overvoltage is caused by the surface potential. Taking into consideration Eq. (11), we obtain the concentration dependence from Eq. (42),

\[
i_{10} = \text{const} \cdot (C_{Na^+})^\delta ,
\]
which shows that in the case of the monovalent cation $i_0$ is proportional to the square root of the Na$^+$ ion concentration in accordance with the experimental results, since $\beta = 0.5$.

Flinn and Schuldiner\textsuperscript{30} has recently reported that the exchange rate of the electron transfer step of the h.e.r. on platinum in sulfuric acid remains constant with the increase of the reversible potential which is attained by diminishing the partial pressure of hydrogen and that the exchange rate of the electron transfer step is proportional to the square root of the concentration of the hydronium ion. These phenomena are quite similar to those on gold in sulfuric acid and can be explained on the basis of Eqs. (11') and (12') and the rate theory described above.

Thus the polarization characteristics of the electron transfer step which have been found experimentally in the case of low overvoltage metals can well be explained on the basis of the compensation effect between the free energy of adsorption and the surface potential in the frame of the transition state method.

Recently, many theoretical and experimental reports have been accumulated in the vacuum systems, which concern with the relation between the binding energy of the adsorption bond and the work function of the metal substrate or the adsorption states.\textsuperscript{31} If the adsorption bond is completely ionic, the increase of the surface potential with the increase of the surface coverage may result in the decrease in the binding energy of the adsorption bond as shown in the classical theory of Langmuir and Taylor\textsuperscript{32}, and consequently the compensation effect can be seen between these two quantities.

However, recent studies on monocrystal surfaces in the ultra high vacuum systems show that the nature of the adsorption bond cannot be explained by the classical theory and the covalent nature cannot be neglected even in the case of alkali adsorption on metals\textsuperscript{33,36}.

Although at present we cannot find a conclusive physical conception which can explain the relation between the binding energy of the adsorption bond and the work function of the substrate or the surface potential of the adsorption states in the vacuum systems, the change of the surface potential in the hydrogen electrode system which takes part in the setting up of the hydrogen overvoltage can be expressed by Eq. (12) in terms of the chemical potential of the intermediate adsorbate. The existence of the two independent components of the electrode potential as expressed in Eqs. (11) and (12) or in Eqs. (11') and (12') can explain the reason why the polarization curve of the hydrogen electrode reaction can be divided into two independent
parts.

Dogonadze, Kuznetsov and Levich\(^{37}\) have recently developed a quantum mechanical rate theory of the electron transfer step in the case of high overvoltage metal on which no specific adsorption is observed under polarization. In order to extend their theory to the case of low overvoltage metals, the compensation effect between the surface potential and the free energy of adsorption of an intermediate should probably be taken into consideration in the quantum mechanical treatment of the rate of the electron transfer step, since the potential energy curves of the initial and final states of the electron transfer step may equally be affected in the same extent by the change of the surface coverage of the intermediate.

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