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Role of ozone in the solar cycle modulation of the North Atlantic Oscillation

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Abstract

The effect of ozone on the 11-year solar cycle modulation of the winter-mean North Atlantic Oscillation (NAO) is examined through analyses of observed meteorological and ozone data from 1978 to 2000. It is found that a significant ozone anomaly associated with the winter NAO is created in winter in high solar (HS) years only and the anomaly persists from spring to summer, creating a large temperature anomaly in the lower stratosphere through radiative heating. Such a temperature anomaly in the stratosphere creates anomalous temperature of opposite sign at lower heights and anomalous zonal wind in the polar area of the troposphere in summer. The associated surface signal is very similar to the summer Arctic Oscillation (Summer-AO). The mechanism for the formation of the Summer-AO is also discussed.
1. Introduction

As solar radiation is almost the only energetic input on the climate system, it has been believed for a long time that the solar variation has a large impact on the climate [e.g., Haigh, 2003; Pap and Fox, 2004]. In fact, analysis of recent observations shows that the variation of solar flux associated with the 11-year cycle is as much as several tenths of a percent for shorter ultra-violet wavelengths [Rottman, 1988; Lean et al., 1997] and an associated significant temperature change was found in the upper stratosphere [Hood et al., 1993; McCormack and Hood, 1996].

Recent observations also indicate that tropospheric climate “modes” related to the stratosphere show significant structural modulation with the 11-year solar cycle. Kodera [2002, 2003] and Ogi et al. [2003] found that the winter mean North Atlantic Oscillation (NAO) shows significant modulation associated with the solar cycle. They found that variability of the sea level pressure (SLP) associated with the winter-mean NAO becomes hemispheric as for the “Arctic Oscillation” (AO) [Thompson and Wallace, 1998] and it extends to the upper stratosphere and persists until the following
summer in high solar (HS) years, whereas it is a tropospheric local variability and has
less persistence in the low solar (LS) years. Further, Kuroda and Kodera [2005] found
the late-winter mean Southern Annular Mode (SAM) shows a similar solar cycle
modulation as the winter NAO. They also suggest that the longer persistence of SAM in
HS years originates from the ozone anomaly in the lower stratosphere, which is created
by the modulation of Brewer-Dobson circulation in late winter. Such a modulation of
SAM was also well simulated by a chemistry-climate model [Kuroda and Shibata,
2006].

Concerning the persistence of the winter NAO signal into the summer in HS
years, Ogi et al. [2003] suggested that surface conditions such as snow, sea ice, and sea
surface temperature (SST) maintain a memory of the winter NAO. This surface-memory
scenario is in contrast to the ozone-memory scenario [Kuroda and Kodera, 2005]. The
purpose of the present study is to examine whether ozone in the stratosphere becomes a
source of the structural modulation of the winter NAO as it is for the SAM in the
southern hemisphere.
2. Data and analysis method

The meteorological data we used in this study are from the 40-year reanalysis data of the European Centre of Medium-range Weather Forecasts (ERA40) [Uppala et al, 2005]. We used 23 years of data from 1978 to 2000. In all analyses monthly-mean data have been used. Quantities containing second order effects such as the E-P flux or residual velocity are calculated from daily data and then monthly averaged. Residual velocities above 100-hPa are estimated by solving the transformed Eulerian mean equation with zero vertical residual velocity at 0.5 hPa, as in Kuroda and Kodera [2001], and they are combined with those calculated directly from analyzed winds and eddies at the 100-hPa level and below.

The ozone data we used is the monthly data that described in Randel and Wu [2007]. This has been combined from two satellite data of Stratospheric Aerosol and Gas Experiment (SAGE) I and II [e.g., Wang et al., 2002] and ozonesonde data at the polar region. As these satellite observations cover latitude only from about 55°S-55°N, ozonesonde data up to 30-hPa and extrapolation of the satellite data is made poleward of 60° latitude. Continuous data is then made by multi-regression method including the decadal trend, solar cycle, and Quasi-Biennial Oscillation (QBO) from gappy data. See
Randel and Wu [2007] for more detail. As the data covers from 1979 to 2005, and the data in December 1978 is absent, we made this from 5-year average of December data from 1979 to 1983.

From meteorological and ozone data, we calculated also the solar heating rate with the radiation code used in the chemistry-climate model of the Meteorological Research Institute [Shibata et al, 2005]. Changes in the ultra-violet (UV) irradiance were adopted from data obtained by Lean et al [2001]. The NAO index used in the present study is the same as that used in Kodera [2002, 2003] and Ogi et al. [2003]. It is the difference in the normalized monthly mean SLP between Lisbon and Stykkisholmur as calculated in Hurrell [1996]. In this study, we used a winter-mean NAO index as the average from December to February (DJF). The time series of the winter-mean NAO index, together with the solar-cycle index, is presented in Figure 1.

First we classified each year as either HS or LS according to the December-to-March (DJFM) mean 10.7 cm solar radio flux: If the flux of a year is higher (lower) than the average, the year is categorized as a HS (LS) year. In this way, 10 (12) winters were selected as HS (LS) years (see lower panel of Figure 1). Correlation analysis was
then performed for the HS and LS years separately, based on DJF mean NAO indices in
a method similar to that used by Kodera [2002, 2003] and Ogi et al. [2003].

Though ERA40 data is available from the surface to 1 hPa, we have restricted
our analysis to the more reliable region from the surface to 10-hPa.

Most of the figures in this paper present correlation or regression with the DJF
mean NAO index. This means that the figures are relative to a positive change in the
DJF mean NAO index.

3. Results

Before performing the correlation analysis, it will be useful to examine
climatologies. To this end, we examined the climatologies of basic quantities for the HS
and LS years and their differences. Figure 2 presents the climatologies of two-month
mean zonal wind from December-January (DJ) to June-July (JJ) for HS and LS years
and their differences. It can be seen that climatological zonal winds are very similar for
HS and LS years. However, there are significant differences between them; the
meridional dipole structure of anomalous zonal wind in the stratosphere is prominent in DJ, as well as the polar tropospheric wind in AM. Winter signals indicate prominent poleward and downward movement of anomalous zonal wind with time if monthly data is used (not shown). The signal in DJ captures this. Such a signal should be closely related to the solar signal [Kuroda and Kodera, 2002; Kodera and Kuroda, 2002]. The variability of the zonal wind is a little larger in the LS than HS years in the stratosphere throughout the period. It is noted that the difference of zonal wind in JJ is quite small in the troposphere, but the variability is large in the high-latitude troposphere both in HS and LS years.

Figure 3 indicates the same plot except for ozone densities. It can be seen that ozone densities tend to become greater with increasing altitude. The densities also become greater with increasing latitude in the lower stratosphere but with decreasing latitude in the middle stratosphere. The variability is generally higher in winter and decreases from spring to summer. Though the climatological values of ozone densities are very similar between HS and LS years, there are significant differences between them, as depicted in the lower panels. In fact, ozone density in HS years tends to be about 2-4% greater in the stratosphere throughout the period [Soukharev and Hood,
2006]. It is noteworthy that the difference between the climatology of the HS and LS years is a little greater than the standard deviation of each year.

Figure 4 indicates the correlation of the zonal-mean zonal wind with the DJF-mean NAO index, calculated separately for “all”, HS, and LS years. Calculations are performed for two-month means from DJ to JJ. Here, correlation greater than 0.6 is contoured in steps of 0.1 and regions of correlation greater than 0.5 are shaded. As the 95% level of statistical significance corresponds to a correlation of 0.63 (0.57) for the sample size of HS (LS) years, the area of correlation greater than 0.6 corresponds to about the 95% significance level. Correlation of 0.5 corresponds to approximately the 90% level of significance. As the correlations of 0.5 and 0.6 correspond to much higher levels of significance for the case of “all years,” the same plot is compared with both HS and LS years.

In the correlation for “all years”, a prominent signal appears in DJ but not afterward. In DJ, a meridional dipole signal centered at 55°N and 35°N appears in the troposphere. Although the low-latitude signal is confined to the lower troposphere, high-latitude signal extends to about the 50-hPa level.
The situation was quite different when we performed the same analysis but stratified according to the solar cycle. In HS years, the larger area of significant correlation extends further into the upper stratosphere in DJ and after diminishing in the following months the signal appears again in JJ with the same polarity. Note that the positive correlations poleward of 50°N in JJ mean that a positive NAO in winter is associated with a stronger westerly vortex in the following summer. LS years, however, show neither the extension to the stratosphere in DJ nor reappearance in JJ.

To examine the source of difference between the HS and LS years, we compared the winter NAO signals from December to April when the differences in the signals are more apparent. Figure 5 compares the E-P flux and its divergence (1st and 3rd rows), and ozone density and residual velocity (2nd and 4th rows) in HS (upper panels) and LS years (lower panels). It can be seen that a large negative correlation of ozone is developing in middle-latitude middle stratosphere to high-latitude lower stratosphere in January only in HS years. This means that a positive NAO index in winter is associated with a negative ozone anomaly in HS years. The formation of large correlation of ozone in HS years corresponds to the weakening of the Brewer-Dobson circulation, shown by
equatorward arrows around 30-40°N and 30-hPa in both December and January, and
upward arrows around 70°N and 200-30 hPa in January. Such anomalous residual
circulation should be driven by the positive anomaly of the E-P flux divergence above
these levels, which is more clearly formed in HS years than LS years. Such differences
in the E-P flux divergence should originate from differences in the E-P flux and
therefore their existence implies reduced upward wave propagation in the extratropical
stratosphere in HS years compared to LS years in January.

The extension of correlation of ozone in April in HS years also corresponds very
well with a weakening of the Brewer-Dobson circulation, shown by upward and
equatorward arrows around 70-80°N and 200-30 hPa, which will be driven by a positive
E-P flux divergence at the stratosphere.

Such a correlation of ozone persists through spring to summer in HS years,
whereas it is thoroughly absent in LS years. It should be noted as well that an apparent
temperature signal also persists from spring to summer in HS years, but it is not present
in LS years (not shown).
To see how a summer AO-like signal evolves with time in HS years, we compared regressed zonal-wind, temperature, ozone density, and solar heating from May to August (Figure 6). Corresponding to a large negative anomaly of ozone, the solar heating is significantly reduced in the stratosphere in May. Such solar heating corresponds well with anomalous temperatures in the stratosphere, although the area of significant correlation is small due to larger variability. This colder temperature signal in the stratosphere becomes more significant and persistent in June. In particular a vertical dipole structure of anomalous temperature, one in the stratosphere and the other in the troposphere, centered around 65°N is prominent. This dipole structure persists until August with slow equatorward propagation. The center of higher temperature in the troposphere is accompanied by an area of Eulerian downward flow. This is associated with anomalous Eulerian mechanical forcing centered around 300-hPa and 70°N (not shown). Anomalous zonal wind centered around 300-hPa and 75°N in June corresponds well with the anomalous higher temperature in the troposphere through the thermal wind relationship. Wave forcing around the tropopause is created through the modification of wave propagation by anomalous zonal wind that extends to the stratosphere. Thus wave forcing around the tropopause, the vertical dipole structure of
the temperature, and the zonal wind in the polar area are sustained by positive feedback in summer.

Figure 7 illustrates the horizontal structure of the regressed temperature in the lower troposphere and stratosphere from June to August. It can be seen that associated with the vertical dipole signal of the temperature seen in Figure 6, warmer areas in the lower troposphere exist just below the cooler areas in the lower stratosphere. It is also noteworthy that the horizontal pattern shows almost zonal wavenumber 3 structure. These clearly suggest existence of strong linkage between the stratosphere and troposphere in this period.

The surface signal in the polar region in the summer of HS years is very similar to the Summer Arctic Oscillation (Summer-AO) or Summer Northern Annular Mode proposed by Ogi et al. [2004]. Such AO-like structures will be created by the wave driven meridional circulation in a similar manner to the winter AO [Kuroda, 2005, 2007a]. Therefore we compared zonal-mean SLP (1st and 3rd rows) and potential surface pressure change (PPC) (2nd and 4th rows) caused by regressed wave forcings associated with the winter-mean NAO in HS years (upper panels) and these with Summer-AO
(lower panels) in Figure 8. PPC is defined as a surface pressure change if frictional and diabatic effects are absent, and it is evaluated by the zonal-mean quasi-geostrophic model on a sphere as used in Kuroda [2005]. See Appendix A for the detail of the model.

Red lines in Figure 8 indicate PPC calculated only by mechanical forcing (forcing due to eddy momentum flux) at 500-hPa and above. As can be seen from the comparison of the SLP and PPC, they are roughly proportional to each other for both the Summer-AO and the signal for the winter-mean NAO, although the signal with winter-mean NAO in June is slightly modified due to larger downward flow around 65°N as previously described. It takes about eight days to produce the SLP anomaly with PPS alone, which is closely related to the frictional damping effect on the meridional circulation at the surface. The similarity of the red line to the circle one means that the SLP signal associated with the winter-mean NAO in HS years is created mainly by the meridional circulation driven by the mechanical forcing in the upper troposphere as is the case for the Summer-AO. Note that the vertical axes of the upper panels are just half of the lower panels. Thus the amplitude of the SLP associated with the winter-NAO in HS years is about half size of the typical Summer-AO in summer.
This is because that the effect of the winter-NAO through ozone is only one source to create the Summer-AO.

4. Discussion and remarks

The analysis presented here suggests that the reappearance of the winter-NAO-index signal in summer in HS years originates from the formation of ozone anomaly in the lower stratosphere in winter. Since the winter ozone anomaly in the mid-latitude lower stratosphere persists until the following autumn [Fioletov and Shepherd, 2003], the ozone anomaly associated with the NAO in HS years becomes a persistent heat source and creates anomalous temperatures in the stratosphere in summer. This is the same as that occurs in the SAM [Kuroda and Kodera, 2005]. However, the SAM signal gradually disappears with the seasonal evolution in summer whereas the AO-like signal disappears quickly in spring but it reappears in summer for the NAO. So the key question we should resolve is why ozone anomaly does not work in spring but works well only in summer in HS years.

As had been seen in Figure 6, formation of Summer-AO signal appears firstly in June, but it is not very clear how the signal firstly appears in June. So we had analyzed
this period by 10-day mean data. Upper panels in Figure 9 show the regression of the Eulerian mechanical forcing with a Eulerian meridional circulation, whereas lower panels show the zonal-mean temperature with a residual circulation. The result shows that the higher temperature anomaly in the troposphere is created firstly in mid-June, associated with the downward Eulerian flow around 65 °N. The flow starts from the first 10-day of June, though it is weak, and corresponds well with anomalous mechanical forcing at the upper troposphere. The area of downward Eulerian flow exists just north of the downward residual flow in mid-June. It is interesting to note that the tropospheric downward flow seems to be connected with the stratospheric colder air if we look closely the residual velocity in the stratosphere for the second to third 10-day. This will explain why warmer areas exist just below the colder areas in Figure 7. In fact, if heavier cooler air is nearly vertically transported to the troposphere, it will create warmer zone in the troposphere. This can occur in June because climatologically the temperature gradient between the lower stratosphere and the upper troposphere becomes smaller and less stable as it evolves with time during this season (Figure 10). It is interesting to compare that though the stratification near the mid-latitude tropopause in the NH is less stable in summer, it is most stable in the SH. So formation of vertical
flow from the stratosphere in summer seems to be a key process, although this should be examined in more detail in a future study.

We had evaluated the PPC by using climatological stability of each month in Figure 8. We had examined the effect of these stabilities on the PPC by changing the stability to that from winter-mean value. The result indicates that the PPC at the pole becomes about 10% smaller if winter-mean values are used. It indicates that the low stability in summer enables wave forcing in the upper troposphere to easily change the lower tropospheric climate.

The ozone data we had analyzed in the present study was made by multi-regression method including the decadal trend, solar cycle, and Quasi-Biennial Oscillation (QBO) [Randel and Wu, 2007]. So using the data for the solar cycle problem as the present study may have problem because the data is made by using solar cycle itself. To verify whether large ozone signal in HS years is not special for the present data, we also performed the same analysis by another data. The other ozone data we used has been combined from two different datasets. One is the version 8 Solar Backscattered Ultraviolet (SBUV) data [Frith and Stolarski, 2005], which contains the
monthly zonal mean mixing ratio at stratospheric pressure levels of 50- to 0.5-hPa, as well as the column ozone between pressure levels, from November 1978 to 2003. The other is the original SAGE II data from October 1984 to 2000 [e.g., Wang et al., 2002]. By merging these two datasets, we produced a combined data set extending from November 1978 to 2000, in which SBUV(2) is used in the stratosphere at 50-hPa level and above, and SAGE II is used below. As the SAGE II data start from October 1984, lower level data before that date was obtained from seasonally varying climatological profiles of SAGE II data, scaled to be equal to the lowest column ozone from the surface to the 64-hPa level in SBUV(2) data. All missing data except for the polar area were replaced with values obtained by interpolating in time even just after the Pinatubo eruption. The same result as Figure 6 except for using this ozone data is shown in Figure 11. Signal for the LS years was found to be very weak (not shown). It is found that the signal is not so strong as that obtained in the present study but it is still present. It is interesting to note that the signal is very similar to the signal of temperature. Ozone signal is still present only in HS years in the alternative dataset.

The large correlation of ozone generated by modulation of the Brewer-Dobson circulation in winter was the key process to convey the winter-NAO signal to summer
in HS years. This can be traced back to the fact that the stratospheric wave propagation was more strongly correlated to the winter-NAO index in HS years (Figure 5). One possible source of this stronger correlation in HS years is that it originates from larger amplification of the planetary wave from the surface due to higher mean ozone in the stratosphere (Figure 3), which creates larger contrast of solar heating in the longitudinal direction of planetary wave and produces stronger wave-mean flow interaction in the stratosphere in HS years. This should be examined in more detail in a future study. See Kuroda et al. [2007] for such analysis in the SH.

In this paper, we considered ozone as a memory of the winter-NAO to the following summer. In contrast, Ogi et al. [2003] considered surface condition such as snow, sea ice and SST as a memory. To investigate the role of ozone and surface condition on the memory of the winter-NAO to the following summer, controlled numerical experiment will be needed in a future study.

Two large volcanic eruptions (El Chichon and Pinatubo) took place in HS years during our period of analysis. To examine the effect of volcanic eruption on the result,
we performed the same calculation after removing the years following the volcanic eruptions. The results demonstrate that the effect was very small (not shown).

We did not consider effects of QBO, ENSO in the present analysis. However, previous analysis indicates that solar influence on climate is much affected by these phenomena [e.g., Labitzke and van Loon, 1999]. Although our present analysis is statistically too short to examine such an effect, analysis of 43-years of data from 1958 to 2000 shows that solar cycle modulation of the NAO is more prominent in the westerly phase of 50-hPa QBO wind. See Kuroda [2007b] for more detail.

Acknowledgments

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Appendix A

The model used in this study is the zonal-mean quasi-geostrophic model on a sphere that was used in Plumb [1982], Haynes and Shepherd [1989], and Kuroda [2005]:

\[
\frac{\partial u}{\partial t} - 2\Omega v \sin \phi = F + G
\]

\[2\Omega u \sin \phi = -\frac{1}{a} \frac{\partial \Phi}{\partial \phi}\]

\[
\frac{\partial \Phi}{\partial p} = -\frac{RT}{p}
\]

\[
\frac{\partial T}{\partial t} - \Gamma \omega = Q + S
\]

\[
\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} (v \cos \phi) + \frac{\partial \omega}{\partial p} = 0
\]

(A.1)

where

\[
F = -\frac{1}{a \cos^2 \phi} \frac{\partial}{\partial \phi} \left[ u'v' \cos^2 \phi \right]
\]

\[
Q = -\frac{1}{a \cos \phi} \frac{\partial}{\partial \phi} \left[ v'T \cos \phi \right]
\]

(A.2)

are eddy mechanical and thermal forcings, \( G \) and \( S \) are frictional forcing and diabatic heating, \( \phi \) is the latitude, \( \omega \) is the vertical pressure velocity, \( \Gamma = -\partial T_o/\partial p + \kappa T_o/p \) is the stability of the basic atmosphere, \( \Omega \) is the angular velocity of the Earth, \( a \) is the radius of the Earth, field variables with primes denote departure from zonal mean,
overbar denotes zonal averaging, and other symbols follow the usual convention [for example, see Andrews et al., 1987].

Boundary conditions of the equation should be

\[ \frac{D\Phi}{Dt} = 0, \]  

(A.3)

for the lower boundary and \( \omega = 0 \) for the upper boundary, and lateral conditions are

\[ \nu = 0 \] at \( \phi = \pm \pi / 2 \), where \( D / Dt \) is a material derivative.

As Equations (A.1) can be put into an elliptical differential equation of \( \omega \)

\[
\frac{1}{\cos \phi \partial \phi} \left( \cos \phi \partial \omega \right) + \frac{4\Omega^2 a^2 p \partial^2 \omega}{R \Gamma \partial p^2}
\]

\[ = \frac{2\Omega \alpha p}{R \Gamma \cos \phi \partial \phi} \left[ \cos \phi \partial \left( F + G \right) \right] - \frac{1}{\Gamma \cos \phi \partial \phi} \left[ \cos \phi \partial \left( Q + S \right) \right],
\]

(A.4)

the meridional circulation can be calculated from this and a continuum equation (the last equation of (A.1)) with boundary conditions. The surface pressure change is then obtained from \( \omega \) on the lower boundary.

For practical method to solve these equations, see Appendix of Kuroda and Kodera (2004).
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References


Frith, S. M., R. S. Stolarski (2005), Merged profile ozone data from the SBUV/SBUV2 series of instruments, *Eos Trans. AGU*, 86(52), Fall Meet. Suppl., Abstract A23B-0943


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Figure 1, Time coefficients of winter-mean NAO index (upper panels) and standardized F10.7 index (lower panels) used in the present study. Closed (open) circle in the lower panels indicate High (Low) Solar year. Year corresponds to the months of January.

Figure 2, Climatological zonal-mean zonal wind for the HS years (upper panels), LS years (middle panels), and their differences (lower panels). The contour interval is 10 m/s in the upper and middle panels and 2 m/s in the lower panels. Shading in the upper and middle panels indicates the standard deviation (light shading 1 m/s, middle shading 2 m/s, and heavy shading 5 m/s). Shading in the lower panels indicates the area of 95% significance levels.

Figure 3, Same as Figure 2, except showing ozone density. The contour interval is 1 ppmv in the upper and middle panels and 0.05 ppmv in the lower panels. In the upper and middle panels, light shading indicates 0.05 ppmv, middle shading 0.1 ppmv, and heavy shading 0.2 ppmv.
Figure 4, Lagged correlation coefficients between December-February mean NAO index and the two-month mean zonal-mean zonal wind at each grid point for the period of 1978-2000 for all years (upper panels), HS years (middle panels), and LS years (lower panels). The contour interval is 0.1 and contours are drawn for absolute values greater than or equal to 0.6 and for zero. Shading is applied to regions where the absolute value of the correlation is greater than 0.5. Dashed lines indicate negative values.

Figure 5, Same as Figure 4, except showing the E-P flux and its divergence (1st and 3rd rows), and residual velocity and ozone density (2nd and 4th rows) for December to April for HS years (upper panels) and LS years (lower panels). Correlations of the E-P flux and residual velocity are shown by arrows. Only arrows whose absolute value of correlations is greater than 0.6 are plotted.

Figure 6, Same as Figure 4, except showing the lagged regression of the E-P flux (arrow) and zonal-mean zonal wind (1st row), temperature (2nd row), ozone density (3rd row), and solar heating (4th row) from May to August for HS years. Contour
intervals of the zonal wind, temperature, ozone density, and solar heating are 0.5 m/s, 0.2 K, 0.02 ppmv, and 0.005 K/day, respectively. Shading is applied to regions where the absolute value of correlation is greater than 0.5. Only E-P fluxes for which the absolute value of correlations is greater than 0.5 are plotted. They are scaled by the reciprocal square root of the pressure.

Figure 7, Lagged regression coefficients between December-February mean NAO index and the monthly mean temperature at 50-hPa (upper panels) and 1000-hPa (lower panels) from June to August for HS years. Contour intervals of 50- and 1000-hPa temperature are 0.2 K and 0.5 K, respectively.

Figure 8, Zonal-mean of the regressed SLP (upper panels) and the potential surface pressure change (PPC, dPs/dt) due to Eulerian wave forcings (lower panels). Quantities for June to August associated with winter-mean NAO index in HS years (upper panels) and the Summer-AO (lower panels) are shown. Red lines in 2nd and 4th rows indicate values due solely to mechanical forcing at 500-hPa and above. The unit of vertical axis is 1 hPa (hPa/day) for the SLP (PPC).
Figure 9, Same as Figure 4, except showing the lagged regression of the Eulerian meridional circulation (arrow) and Eulerian mechanical forcing (contour) (upper panels), and residual velocity (arrow) and zonal-mean temperature (contour) (lower panels) from first to last 10-day of June. Contour intervals of the Eulerian mechanical forcing (zonal-mean temperature) is 0.2 m/s/day (0.2 K). Only arrows with a larger correlation (greater than 0.5) are shown.

Figure 10, Line plots showing climatological temperature difference between 200-hPa and 300-hPa at 60°N (with open circles) and 60°S (with closed circles).

Figure 11, Same as Figure 6, except showing the lagged regression of ozone density (1st row), and solar heating (2nd row) from May to August for HS years of combined SBUV(/2) and SAGE II data.
Figure 1

Figure 2
Figure 3

Figure 4
Figure 5
Figure 6
Figure 9

Figure 10
Figure 11