Evolution of Paramagnetic Quasiparticle Excitations Emerged in the High-Field Superconducting Phase of CeCoIn$_5$

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We present $^{115}$In NMR measurements in a novel thermodynamic phase of CeCoIn$_5$ in a high magnetic field, where exotic superconductivity exists with the incommensurate spin-density wave order. We show that the NMR spectra in this phase provide direct evidence for the emergence of the spatially distributed normal quasiparticle regions. The quantitative analysis for the field evolution of the paramagnetic magnetization and newly emerged low-energy quasiparticle density of states is consistent with the nodal plane formation, which is characterized by an order parameter in the Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) state. The NMR spectra also suggest that the spatially uniform spin-density wave is induced in the FFLO phase.

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The interplay between magnetism and unconventional superconductivity with a nontrivial Cooper pairing has been a topic of recent intense study. The quasi-two-dimensional heavy fermion superconductor CeCoIn$_5$ continues to excite great interest because it shows a number of fascinating superconducting properties. Its superconductivity at high fields is destroyed by the Pauli paramagnetic effect, as evidenced by the first-order phase transition at the upper critical field $H_c^2$ [3–5] and the anomalous flux line lattice form factor [6]. What is striking is that CeCoIn$_5$ exhibits a new thermodynamic phase transition at $(T^*, H^*)$ just below $H_{c2}$ [Fig. 1(a)] for both $H \parallel ab$ and $H \parallel c$ [7,8]. Closely related to the Pauli limited superconductivity, this high-field and low-temperature (HL) phase has been attributed to the Fulde-Ferrell-Larkin-Ovchinnikov (FFLO) state, in which the pair breaking arising from the Pauli effect is reduced by forming a new pairing state $(k \perp, -k + q \parallel)$ with nonzero $q$ between the Zeeman-split parts of the Fermi surface. One of the most fascinating aspects of the FFLO state is that Cooper pairs with finite center-of-mass momenta $hq$ develop an oscillating superconducting order parameter in real space such as $\Delta(r) \propto \sin(q \cdot r)$ and, as a result, nodal planes appear periodically perpendicular to the applied field [11].

The presence of the FFLO state in CeCoIn$_5$ has been supported by several experiments, including impurity [12], pressure [13], and ultrasound studies [14,15]. However, recent NMR [16–18] and neutron [19,20] data in parallel field demonstrated a long-range static magnetic order in the HL phase. The magnetic moment at Ce atoms is given by $\mu(r) = \mu_0 \cos(Q_S \cdot r)$ with $Q_S = 2\pi(\delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta \delta 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In sites. The axially symmetric In(1) is located in the CeIn layer, whereas In(2a) and In(2b) sites are located between Co and CeIn layers with the largest principal axis of electric field gradient parallel and perpendicular to the applied field, respectively [Fig. 1(c)]. The NMR frequency in the superconducting state is spatially distributed and is determined by the local magnetic field and hyperfine coupling to the conduction electron spins: \( \text{H} = \text{H}_{\text{eff}}(r) = \text{H} + \text{M}_0(r) + A_{\text{hf}}M_S(r) \), where \( \text{M}_0(r) \) is the local magnetization due to the orbital (diamagnetic screening current) effect, \( A_{\text{hf}} \) is the hyperfine coupling constant [26], and \( M_S(r) \) is the local spin magnetization. The Knight shift, given as \( K = A_{\text{hf}}M_S/H \), consists of the spin and orbital contributions, \( K = K_{\text{spin}} + K_{\text{orb}} \). The amplitude of \( K_{\text{orb}} \) is estimated at low field and low temperature in the superconducting state where \( K_{\text{spin}} \) vanishes. We note that in the vortex state of CeCoIn, line shapes at In(2a) and In(2b) sites with large amplitude of \( A_{\text{hf}} \) are predominantly determined by the hyperfine coupling \( (A_{\text{hf}}M_S > M_0) \).

The antiferromagnetic staggered Ce moments due to IC SDW with \( \mu_0 || c \) induce the in-plane hyperfine field perpendicular and parallel to the applied field \( \text{H} \) at the In(2a) and In(2b) sites, respectively [Fig. 1(c)], through dipolar-type transferred hyperfine couplings [17]. Consequently, the In(2b) spectra should be shifted (into two peaks) by the antiferromagnetic staggered field, whereas In(2a) spectra should not. Figure 1(d) depicts the field evolution of In(2b) spectra at 50 mK as a function of Knight shift. In the normal state for \( \mu_0H \geq 11.8 \text{T} \), very sharp spectra with line width less than 50 kHz (11.8–12.0 T) are observed. For \( \mu_0H < 11.8 \text{T} \) the sharp In(2b) spectra broaden and split into two peaks with finite signal weight between them (10.1–11.75 T), which is characteristic of IC SDW long-range order along one spatial dimension. In the BCS phase below \( H^*( \approx 10 \text{T}) \), NMR spectra (8.2–9.7 T) becomes a single asymmetric line as expected in the usual vortex state. The inset of Fig. 1(d) shows the field dependence of the internal field \( H_{\text{int}} \) determined by the difference of the resonance frequency of the two peaks in the HL phase. Above \( H^* \), \( H_{\text{int}} \) increases with \( H \) and jumps to zero at the first-order \( H_{c2} \) transition after reaching maximum value of \( \approx 0.16 \text{T} \), which corresponds to \( \approx 0.15 \mu_B/\text{Ce} \). The vanishing of SDW order above \( H_{c2} \) and the magnitude of \( H_{\text{int}} \) are consistent with the previous NMR [16–18] and neutron results [19,20].

The field evolution of the In(2a) spectra is depicted in Fig. 2. The sharp resonance line in the normal state (12.0 T) becomes antisymmetric broad lines in the HL.
(10.0–11.5 T) and BCS (8.2–9.9 T) phases. The most salient feature of the HL-phase spectra is the emergence of the edge structure whose position coincides with the Knight shift in the normal state, as shown by the dotted line (also Fig. 2, inset). This provides direct evidence for the emergence of the normal quasiparticle region in the HL phase [27,28]. We note that as the edge position is independent of magnetic field, the effect of the hyperfine field due to the SDW ordering is negligibly small. This indicates that In(2a) spectra are predominantly affected by the local spin susceptibility arising from the normal quasiparticles.

Here we comment on the reproducibility of the data. The NMR spectra are reproducible for different rf powers and different cooling rates. We also measured the spectra on several different single crystals. Although the clear edge structure of In(2a) spectra is observed in all crystals, its sharpness slightly depends on the crystal. This appears to be related to the recent result that the HL phase is extremely sensitive to the nonmagnetic impurity [12]. The NMR spectra with less pronounced edge structure can also be found above $H^*$ in Ref. [18].

We stress that the clear-cut sharp edge structure is consistent with the recent calculation of the NMR spectra in the FFLO state based on the microscopic Bogoliubov-de Gennes equation [28]. For the quantitative analysis, we measured the temperature dependence of the In(2a) spectrum (Fig. 3). As shown by a thin black line, which is the spectrum at $T = 0.21$ K just above $T^*$, the shape of the main peak remains almost unchanged with temperature. This allows us to separate the spectrum in the edge region below 0.15 K displayed by the (green) shaded region in Fig. 3 by subtracting the BCS spectrum at 0.21 K shown by thin black lines from the whole spectra in the HL phase. The (green) shaded regions in Fig. 2 also indicate the spectra near the edge regions obtained by the similar method.

The total paramagnetic magnetization $M_p$ is evaluated from the whole NMR spectrum $P(K_{\text{spin}})$ by integrating the spin part of the Knight shift as $M_p(H) = \frac{H}{A_g} g(H)$, where $g(H) = \int g_{\text{spin}}(H) dK_{\text{spin}}/\int P(K_{\text{spin}}) dK_{\text{spin}}$. Figure 4(a) depicts the field dependence of $M_p$ obtained from the In(2a) spectra as well as $M_p(H)$ obtained from the same analysis for the In(1) spectra. Both $M_p(H)$ from the In(1) and In(2a) spectra increase linearly with $H$ in the HL phase. The difference of the paramagnetic magnetization between the HL and BCS phases, $\delta M_p(H) = M_p(H) - M_p^\text{BCS}(H)$, increases continuously from zero as $\delta M_p(H) \propto (H - H^*)$ in the HL phase, where $M_p^\text{BCS}(H)$ is obtained by a linear extrapolation from the BCS phase [dotted line in Fig. 4(a)]. We emphasize that $M_p(H)$ cannot be obtained from the bulk magnetization measurements [4] which contains the contribution of the antiferromagnetic Ce moments.

The DOS of the normal quasiparticles emerged in the HL phase is proportional to the area of (green) shaded regions in Figs. 2 and 3, as the intensity there should be proportional to the number of nuclei which detect the quasiparticle susceptibility. Figure 4(b) depicts the field dependence of the normalized DOS, $n_{\text{qp}}$, which is obtained by the area of the (green) shaded region divided by the area of the full spectrum at each field in Fig. 2. To estimate the DOS in an alternative way, the intensity $I_N$ of the spectrum at the normal state Knight shift (dotted line in Fig. 2) is also plotted in Fig. 4(b). Both $n_{\text{qp}}$ and $I_N$ increase in proportion to $\sqrt{H - H^*}$ in the HL phase.

The $H$-linear dependence of the paramagnetic magnetization and $\sqrt{H}$ dependence of the low-energy DOS both provide key information about the order parameter that characterizes the emergence of the normal quasiparticles in the HL phase. We stress that both field dependencies are

![FIG. 3](color online). Temperature evolution of the In(2a) spectrum in HL (0.05–0.18 K) and BCS (0.21 K) phases at $\mu_0 H = 11.1$ T. The peak intensity of each spectrum is normalized. Thin black lines indicate the spectrum at $T = 0.21$ K just above $T^*$. The (green) shaded regions in the low-temperature data indicate the quasiparticle spectrum formed in the HL phase.

![FIG. 4](color online). (a) Field dependence of the paramagnetic magnetization $M_p$ of the In(2a) (filled circles) and the In(1) (open circles) sites. Inset: Schematic quasiparticle structure in the FFLO state. The nodal planes with period of $2\pi/q$ appear perpendicular to the Abrikosov vortex lattice. (b) Field dependence of the DOS of the paramagnetic quasiparticles, $n_{\text{qp}}$ and $I_N$, extracted from the In(2a) spectra. The solid line is a fit to $\sqrt{H - H^*}$ dependence.
with the absence of satellite peaks in the neutron scattering region far away from the nodal planes. This is consistent IC SDW, in which the magnetic order is present even in the BCS spectrum. This suggests rather uniform In(2b) spectra in the HL phase shown in Fig. 1(d) (IC SDW spectrum with two peaks). In contrast, the state, the NMR spectra at the In(2b) site in the HL phase in the presence of such a spatially nonuniform IC SDW moment is induced around the FFLO nodal planes [24]. These results appear to indicate that coupling between FFLO and IC SDW becomes weaker with tilting [8,32]. These results are sufficient to discuss the FFLO state with SDW order. It has been theoretically proposed that the IC SDW moment is induced around the FFLO nodal planes [24]. In the presence of such a spatially nonuniform IC SDW state, the NMR spectra at the In(2b) site in the HL phase consist of both contributions of the regions far outside (BCS spectrum with one peak) and around the nodal planes (IC SDW spectrum with two peaks). In contrast, the present In(2b) spectra in the HL phase shown in Fig. 1(d) consist entirely of the split one with no discernible component of the BCS spectrum. This suggests rather uniform IC SDW, in which the magnetic order is present even in the region far away from the nodal planes. This is consistent with the absence of satellite peaks in the neutron scattering experiments [20].

We point out that the present results are incompatible with the pair-density wave scenario [21,22]. The emergence of the normal quasiparticles in the HL phase is not expected in this scenario. Moreover, the order parameter \( q \), which depends on \( H \) in accord with the FFLO second order transition, is at odds with this scenario that assumes the Cooper pair with the SDW modulation wave vector \( Q_s \), which is field independent [19].

It has been reported that the magnetic moment induced by the SDW disappears when the magnetic field is tilted away from the \( ab \) plane by \( 17^\circ \) [31]. Moreover, a possible FFLO phase with no magnetic order appears in \( H \parallel c \) [8,32]. These results appear to indicate that coupling between FFLO and IC SDW becomes weaker with tilting \( H \) from the \( ab \) plane. Recently, it has been suggested that in CeCoIn\(_5\) with \( d_{2x^2-y^2} \)-wave symmetry the strong Pauli paramagnetism plays an important role for the SDW formation as well as the FFLO state particularly in parallel field [33,34]. Further investigation is required to understand the perplexing relationship between the coexisting FFLO and IC SDW states.

In summary, NMR measurements demonstrate the emergence of a spatially distributed normal quasiparticle region in the HL phase of CeCoIn\(_5\) in parallel field. The field evolution of the paramagnetic magnetization and low-energy quasiparticle DOS can be described well by the order parameter associated with the nodal plane formation via the FFLO second order phase transition. The NMR spectra also reveal that the spatially uniform SDW coexists with the FFLO nodal planes.

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[26] The \( \lambda_{ab} \) values of 1.8, 5.5, and \(-3.07/\lambda_{ab} \) for In(1), In (2a), and In(2b) site, respectively, are obtained from the \( -\chi \) plot below \( T = 2 \) K by using \( K(T) \) in the normal state at \( \mu_B H = 12 \) T and \( \chi(T) \) [4] for \( H \parallel ||0|| \).
[27] K. Kakuyanagi et al., Phys. Rev. Lett. 94, 047602 (2005). There the spectra at the In(1) site were restricted in a narrow frequency range. Accordingly, those results were insufficient to discuss the FFLO state with SDW order.