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Author(s)	小布施, 秀明; Bera, S.; Ludwig, A. W. W.; Gruzberg, I. A.; Evers, F.
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Statistics of Conductances and Subleading Corrections to Scaling near the Integer Quantum Hall Plateau Transition

H. OBUSE^{1,2}, S. BERA³, A. W. W. LUDWIG⁴, I. A. GRUZBERG⁵ and F. EVERS^{1,6,7}

¹ *Institut für Nanotechnologie, Karlsruhe Institute of Technology (KIT), 76021 Karlsruhe, Germany*

² *Department of Applied Physics, Hokkaido University, Sapporo 060-8628, Japan*

³ *Institut Néel, CNRS and Université Joseph Fourier, 38042 Grenoble, France*

⁴ *Department of Physics, University of California, Santa Barbara, CA 93106, USA*

⁵ *Department of Physics, The Ohio State University, 191 W. Woodruff Ave., Columbus, OH 43210, USA*

⁶ *Institut für Theorie der Kondensierten Materie, Karlsruhe Institute of Technology (KIT), 76128 Karlsruhe, Germany*

⁷ *DFG-Center for Functional Nanostructures, Karlsruhe Institute of Technology (KIT), 76131 Karlsruhe, Germany*

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Abstract – We study the critical behavior near the integer quantum Hall plateau transition by focusing on the multifractal (MF) exponents X_q describing the scaling of the disorder-average moments of the point contact conductance T between two points of the sample, within the Chalker-Coddington network model. Past analytical work has related the exponents X_q to the MF exponents Δ_q of the local density of states (LDOS). To verify this relation, we numerically determine the exponents X_q with high accuracy. We thereby provide, at the same time, independent numerical results for the MF exponents Δ_q for the LDOS. The presence of subleading corrections to scaling makes such determination directly from scaling of the moments of T virtually impossible. We overcome this difficulty by using two recent advances. First, we construct pure scaling operators for the moments of T which have precisely the same leading scaling behavior, but no subleading contributions. Secondly, we take into account corrections to scaling from irrelevant (in the renormalization group sense) scaling fields by employing a numerical technique (“stability map”) recently developed by us. We thereby numerically confirm the relation between the two sets of exponents, X_q (point contact conductances) and Δ_q (LDOS), and also determine the leading irrelevant (corrections to scaling) exponent y as well as other subleading exponents. Our results suggest a way to access multifractality in an experimental setting.

The integer quantum Hall (IQH) effect has been an exciting area of research in condensed matter physics for three decades [1, 2]. Recently, renewed interest put the IQH plateau transition, an Anderson (de-)localization transition driven by disorder [3], in the focus of intense experimental [4–9] and theoretical research [10–18]. Of particular interest is the scaling of transport and other properties at the transition, including the scaling of moments of the local density of states (LDOS) $\rho(\mathbf{r})$.

So far, there is no widely accepted analytical theory of the critical phenomena near the transition¹. Meanwhile, numerical simulations have revealed rich critical behavior of the LDOS that exhibit multifractality in the form of an

infinite set of exponents (see [3, 19] for a recent review) that describe the scaling of the moments $\langle \rho^q(\mathbf{r}) \rangle \sim L^{-\Delta_q}$ with the system size L [20–22]. Unfortunately, the probability distribution for the LDOS, and hence Δ_q , are difficult to access experimentally, but there are notable attempts in this direction [23, 24]. Transport measurements are typically easier to perform, and this makes understanding relations between multifractality of the critical LDOS and conductances quite important.

In Refs. [25] and [26] the authors have introduced the notion of the point contact conductance (PCC) $T(\mathbf{r}_1, \mathbf{r}_2)$ between two points \mathbf{r}_1 and \mathbf{r}_2 of the sample, and considered the scaling of its moments at criticality (we will use the short-hand notations $T \equiv T(\mathbf{r}_1, \mathbf{r}_2)$ and $r \equiv |\mathbf{r}_1 - \mathbf{r}_2|$

¹But see Ref. [17] for recent developments.

in the following), described by

$$\langle T^q \rangle \sim r^{-X_q}. \quad (1)$$

The following relation between the transport exponents X_q and the LDOS exponents Δ_q was derived [25–27]:

$$X_q = \begin{cases} 2\Delta_q, & \text{for } q < 1/2, \\ 2\Delta_{1/2}, & \text{for } q \geq 1/2. \end{cases} \quad (2)$$

While this relation was derived using the Chalker-Coddington (CC) network model [28] for the IQH transition, it connects two sets of universal critical exponents, and as such is expected to be generally valid, independent of the microscopic model. In fact, it is believed that similar relations hold for Anderson transitions in all symmetry classes [3, 29].

In this Letter, we subject eq. (2) to a sensitive test by determining numerically the transport exponents X_q and then comparing them with earlier numerics for Δ_q obtained from the scaling behavior of the LDOS (or wavefunction) moments [12, 13]. For the high-precision comparison that we are aiming at, a careful treatment of subleading power-laws is required. As recent work [14, 16, 18] demonstrates, the IQH transition poses a particular challenge in this respect. It is well known [3] that corrections to scaling decay near the IQH fixed point with an irrelevant exponent, $|y|$, that is certainly smaller than one, probably smaller than 0.5. This unusually small value requires us to keep several terms in power series expansions off scaling functions, see below. Therefore, the IQH fixed point is much more difficult to access than the three-dimensional Anderson transition where much larger values ($|y| \approx 1.5$ [30] and $|y| \approx 3.3$ [31]) were reported in the orthogonal class. Even though subleading terms do not influence the true asymptotic scaling behavior they are still important in practice, in particular because they determine the size of the critical region.

We remind the reader that two types of corrections to scaling generally exist. First, a particular physical observable may be a combination of several pure scaling operators. It is known, for example, that T^q is not a pure scaling operator [25], while $\rho^q(\mathbf{r})$ is [29, 32]. Thus, eq. (1) should be understood as characterizing the leading long-distance behavior of the observable $\langle T^q \rangle$, subject to subleading corrections from the admixture of other pure scaling operators, characterized by certain subleading scaling dimensions γ_q . Secondly, even correlation functions of pure scaling operators exhibit what is called *irrelevant* corrections to scaling due to the fact that they are calculated using a critical Hamiltonian different from the fixed point Hamiltonian. Deviations from the fixed point are controlled by irrelevant exponents ² $y_i < 0$ which are the same for all scaling operators. The second goal of the paper is to determine the leading irrelevant (i.e. least irrelevant) exponent

²In D spatial dimensions, $y_i = (D - x_i)$ is the scaling dimension of the coupling constant of an operator of scaling dimension x_i added to the Hamiltonian.

y (which should be independent of q) as well as the scaling dimensions γ_q of the above mentioned subleading operators.

Keeping these goals in mind, we write the scaling function for moments of the PCC as

$$\langle T^q \rangle = c_q r^{-X_q} \left(1 + \sum_{n=1}^{N_p} a_q^{(n)} r^{ny} \right) + d_q r^{-\gamma_q} \left(1 + \sum_{n=1}^{N_s} b_q^{(n)} r^{ny} \right). \quad (3)$$

Here in addition to the leading scaling operator characterized by the scaling dimension X_q we have retained only one subleading operator characterized by the scaling dimension $\gamma_q > X_q$, and in both contributions we have retained only one irrelevant, i.e. the leading corrections to scaling exponent y . We have also truncated the expansions of the scaling function in powers of r^y at orders N_p and N_s for the two scaling operators, for the purpose of the numerical analysis below ³.

Because of the presence of the subleading scaling dimension γ_q in eq. (3), it is difficult to get an accurate estimate of the leading scaling dimension X_q from fitting numerical data for the moments $\langle T^q \rangle$ to eq. (3). To circumvent this difficulty, we will consider $p_q(T) \equiv P_{-q}(2/T - 1)$, where $P_q(x)$ is the associated Legendre function of the first kind [33]. As explained below, this quantity has the following properties: (i) its leading scaling behavior is the same as that of $\langle T^q \rangle$ for $q \leq 1/2$, and (ii) it is a pure scaling operator.

Property (i) follows if we consider long distances r , where $T \ll 1$. Then $x \equiv 2/T - 1 \gg 1$, and we can use the standard asymptotics ⁴

$$P_{-q}(x) \sim \begin{cases} x^{-q} \sim T^q, & q < 1/2, \\ x^{q-1} \sim T^{1-q}, & q > 1/2. \end{cases} \quad (4)$$

Property (ii) can be derived as follows. In Ref. [26] the following formula was derived within the CC model:

$$2\pi\nu \left\langle \rho_2 f\left(\frac{\rho_2}{\rho_1}\right) \right\rangle = \left\langle \int_{-\pi}^{\pi} \frac{d\phi}{2\pi} f\left(\frac{|1 + e^{i\phi}\sqrt{1-T}|^2}{T}\right) \right\rangle. \quad (5)$$

Here $\rho_i \equiv \rho(\mathbf{r}_i)$ is the LDOS at point \mathbf{r}_i , ν is the mean level density, and $f(z)$ is an arbitrary function. If we choose $f(z) = z^{-q}$ [27], the integral on the right-hand side of eq. (5) becomes ⁵

$$\frac{1}{\pi} \int_0^\pi \frac{d\phi}{[x + (x^2 - 1)^{1/2} \cos \phi]^q} = P_{-q}(x) = P_{q-1}(x), \quad (6)$$

³More generally, $\langle T^q \rangle$ should include several irrelevant exponents, y_1, y_2, \dots , and contributions from more subleading scaling operators characterized by exponents $\gamma_q^{(1)}, \gamma_q^{(2)}, \dots$, and these contributions should not be truncated.

⁴See sect. 14.8 (iii) in Ref. [33]. For $q = 1/2$ there is an extra logarithmic correction: $P_{-1/2}(x) \sim x^{-1/2} \ln x \sim T^{1/2} \ln \frac{1}{T}$.

⁵See eq. (14.12.7) in Ref. [33].

where the last equality is a symmetry property of P_q . Thus, we obtain the following relation:

$$2\pi\nu\langle\rho^q(\mathbf{r}_1)\rho^{1-q}(\mathbf{r}_2)\rangle=\langle p_q(T)\rangle. \quad (7)$$

The left-hand side of this relation is a correlation function of pure scaling operators [32] with dimensions $\Delta_q=\Delta_{1-q}$ [34], which demonstrates the point (ii) above.

The arguments above allow us to write

$$\langle p_q(T)\rangle=\bar{c}_q r^{-\bar{X}_q}\left(1+\sum_{n=1}^{N_p}\bar{a}_q^{(n)}r^{n\bar{y}}\right), \quad (8)$$

$$\bar{X}_q=2\Delta_q. \quad (9)$$

In contrast to eq. (2), \bar{X}_q should be equal to $2\Delta_q$ for any q . Also, unlike in eq. (3), there are no admixtures of subleading scaling operators in eq. (8), which makes fitting numerical data to eq. (8) much better controlled. This allows us to extract reasonable numerical estimates of exponents \bar{X}_q . However, the numerical values of \bar{X}_q obtained in this way turn out to be not precise enough for a high accuracy test of eq. (9). The limitations are set by statistical noise in the raw data, which is of the order of 0.1% of the relative standard error. Also, the range of available distances ($r=3\text{--}59$ lattice constants) is not sufficient to separate the different power-law contributions in eq. (8) from each other. Going one step further, we solve this problem by employing conformal invariance.

By using a logarithmic function we conformally map the 2D plane to a cylinder with circumference M [35–38]. In this quasi-one dimensional (Q1D) geometry the distance between point contacts along the cylinder is denoted as L . For $L\gg M$, the PCC T in the Q1D geometry should have the same scaling properties as the two-terminal conductance g of the cylinder of length L . Therefore, we compute numerically moments of g and fit them to the following scaling function [35, 38], to be contrasted with eq. (3):

$$\langle g^q\rangle=c'_q\exp\left[-\pi\left(X'_q+\sum_{n=1}^{N_p}a'^{(n)}M^{ny'}\right)\frac{L}{M}\right]+d'_q\exp\left[-\pi\left(\gamma'_q+\sum_{n=1}^{N_s}b'^{(n)}M^{ny'}\right)\frac{L}{M}\right]. \quad (10)$$

(Primed exponents from the Q1D geometry and unprimed exponents from the 2D geometry are the same, in principle. We distinguish them nevertheless, in order to emphasize that the numerical estimates that are obtained in practice for primed exponents are significantly more reliable.) Notice, that in the Q1D geometry the irrelevant terms appear as corrections to the exponents X'_q and γ'_q [35, 39, 40]. (This is obtained by the conformal perturbation theory directly in quasi-1D geometry, see, e.g., Refs. [41, 42].) This fact, together with the relative exponential suppression of the subleading term (due to $\gamma'_q>X'_q$) leads us to a way to reliably extract numerical values of all exponents. This method is especially effective in cases where the subleading (γ_q) and irrelevant (y) exponents happen

to be numerically close, so that their contributions tend to mask each other.

We can combine the mapping to the Q1D geometry with the use of Legendre functions, and we will demonstrate below that this leads to the best accuracy. However, for two terminal conductances in Q1D geometry, eq. (5) is, strictly speaking, not applicable. Thus, we do not expect $p_q(g)$ to be a pure scaling operator. Consequently, we use the following scaling function for this quantity

$$\langle p_q(g)\rangle=c'_q\exp\left[-\pi\left(\bar{X}'_q+\sum_{n=1}^{N_p}\bar{a}'^{(n)}M^{n\bar{y}'}\right)\frac{L}{M}\right]+d'_q\exp\left[-\pi\left(\bar{\gamma}'_q+\sum_{n=1}^{N_s}\bar{b}'^{(n)}M^{n\bar{y}'}\right)\frac{L}{M}\right]. \quad (11)$$

which allows for an admixture of a subleading scaling operator.

We obtain the leading, subleading, and irrelevant exponents by numerically fitting raw data for conductances to the scaling functions (3), (8), (10), and (11). The quality of fitting for each data set is evaluated in a standard way by employing the chi-square and goodness of fitting. The fitting is difficult for two reasons. First, as already mentioned, corrections to scaling decay slowly near the IQH fixed point. In general, several subleading terms need to be included in order to obtain a consistent result. Consequently, up to ten fitting parameters need to be included. With so many parameters, the cost function has several local minima even though the statistical uncertainty in our raw data is below 0.02% for conductances in Q1D. In order to find the optimal fit we have proposed in an earlier work to employ a “stability map” that we use also here [18]. The statistical errors that we use result from the χ^2 -analysis at the optimal point. Second, results of our fitting depend on the choice of the window of the distance between contacts or the system size (e.g. insets fig. 2). This reflects a residual dependence on higher order terms in the power expansion eq. (8) that we disregard in our fitting to limit the number of free parameters, and leads to the second type of numerical uncertainty of our estimates. In order to account for this situation, we introduce “practical error bars”. They constitute the sum of both mentioned uncertainties and thus are expected to give a conservative upper bound of the true errors.

Only upon combining all four essential steps—the use of the Legendre functions, the mapping to the Q1D-geometry, the stability map, and the practical error bars—are we able to achieve the numerical accuracy necessary to finally confirm the relation (9).

Before presenting our numerical results, we make one more comment. Whenever moments of a random quantity exhibit multifractal scaling characterized by exponents X_q , the Legendre transform $F(a)=aq-X_q$, $a=dX_q/dq$, is related to the probability distribution of this quantity [25, 35]. Indeed, if we assume the following form for the distribution of a PCC T : $\text{Prob}(T=r^{-a})dT\sim r^{F(a)}da$ and ignore correction terms, then the moments of T will

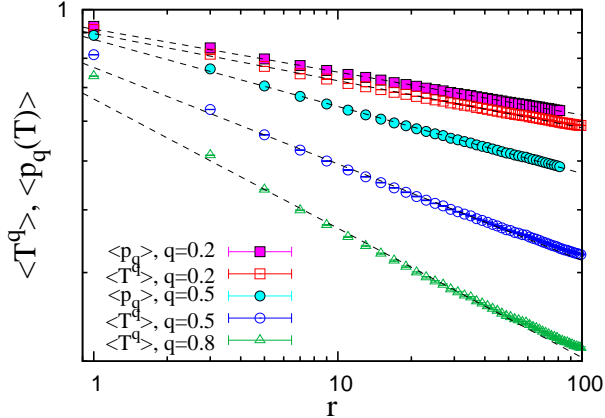


Fig. 1: The dependence of $\langle T^q \rangle$ and $\langle p_q(T) \rangle$ on r (measured in lattice spacings) along a horizontal line on a square of size $L = 480$ with periodic boundary conditions at $q = 0.2, 0.5$, and 0.8 . The dashed lines are guides to the eye. $\langle T^q \rangle$ is seen to deviate from the simple power law (1) suggesting the presence of subleading corrections that require us to use eqs. (3) and (8) to fit numerical data for relatively short distances r . At longer distances corrections in powers of r/L obscure the $L \rightarrow \infty$ asymptotics. As q increases, the difference in the slopes of $\langle T^q \rangle$ and $\langle p_q(T) \rangle$ also increases, leading to a difference in the exponents X_q and \bar{X}_q .

scale for large r as $\langle T^q \rangle \sim \int da r^{F(a)-aq} \sim r^{-X_q}$. Thus, for the probability distribution function of $\ln T$ we can write

$$\mathcal{P}(\ln T, r) = \mathcal{N} r^{F(a)}, \quad a = -\ln T / \ln r, \quad (12)$$

where \mathcal{N} is a normalization constant. Analogously, the LDOS exponents Δ_q lead (in 2D) to the so-called singularity spectrum [3] $f(\alpha) = (\alpha - 2)q - \Delta_q + 2$, where $\alpha - 2 = d\Delta_q/dq$. The relation (9) leads to

$$F(a) = 2[f(\alpha) - 2], \quad a = 2(\alpha - 2), \quad (13)$$

which can be used in eq. (12). In the same way, eq. (10) for the two-terminal conductance in Q1D leads to the following probability distribution for $\ln g$:

$$\mathcal{P}_{\text{Q1D}}(\ln g, L/M) = \mathcal{N}_{\text{Q1D}} e^{\pi F(a) \frac{L}{M}}, \quad a = -\frac{M \ln g}{\pi L}. \quad (14)$$

In practice, the normalization constant \mathcal{N}_{Q1D} is determined from the peak of $\mathcal{P}_{\text{Q1D}}(\ln g, L/M)$:

$$\mathcal{N}_{\text{Q1D}} = \mathcal{P}_{\text{Q1D}}(\ln g^{\text{typ}}, L/M), \quad (15)$$

where $\ln g^{\text{typ}}$ is the typical value of $\ln g$, since eq. (13) gives $F(a^{\text{typ}}) \propto f(\alpha_0) - 2 = 0$, where $a^{\text{typ}} = -M \ln g^{\text{typ}} / \pi L$ is related to α_0 from eq. (13), which gives a maximum of $f(\alpha)$, that is, $f(\alpha_0) = 2$.

Numerical analysis in 2D. – The PCC is computed numerically using the procedure described in Refs. [25,26]. For a given sample of linear size L , this requires a solution of a linear system of equations of dimension $\sim L^2$ for each position of the two point contacts. In this work we investigate square systems with periodic boundary conditions in

both directions, and place the two point contacts on the same horizontal row of links. Other boundary conditions will be treated elsewhere [43]. In our simulations $L = 480$ and we average over 10^5 samples.

Figure 1 shows the moments $\langle T^q \rangle$ and $\langle p_q(T) \rangle$ for the torus geometry. When one contact moves along a row, $T(r)$ is periodic in r with the period L . Hence, in this case the expansion (3) applies only at short distances where corrections of order r/L can be neglected. By comparing $T(r)$ traces for different L -values we found that a sufficient condition is $r/L \lesssim 0.1$ [44]. Due to this limitation the window of r values where we can hope to study the true asymptotic behavior is narrow, $r \lesssim 41 - 57$ for the system sizes available to us. Under these conditions a reliable fit to eqs. (3) and (8) is very difficult and can be achieved only using the stability map [44].

Results of this extensive analysis are displayed in fig. 2(a). It offers a comparison of the exponents X_q and \bar{X}_q (both divided by $2q(1-q)$) as obtained from fits of $\langle T^q \rangle$ and $\langle p_q(T) \rangle$ to eqs. (3) and (8), as well as the LDOS exponents Δ_q (divided by $q(1-q)$). All leading exponents are shown with practical error bars. The fitting $\langle p_q(T) \rangle$ to eq. (8) is straightforward, and we obtain results that are consistent with eq. (9), albeit with rather large error bars. By contrast, subleading terms in $\langle T^q \rangle$ interfere strongly in the vicinity of $q = 1/2$, and a controlled fit to eq. (3) is not possible with our data [44].

Numerical analysis in Q1D. – Practical error bars can be reduced by an order of magnitude for exponents obtained in the long cylinder (Q1D) geometry, where the scaling of the moments $\langle T^q \rangle$ and $\langle g^q \rangle$ should be identical in the limit $L \gg M$. The conductance g in Q1D is obtained by the transfer matrix method [45]. The width M is varied between 32 and 512. The range of L used for the fitting is $2M$ to $10M$. The number of samples for each M is 10^6 . The leading, subleading, and irrelevant exponents are all extracted from fits to eqs. (10) and (11) [44].

In fig. 2(b) we show the exponents X'_q and \bar{X}'_q describing the scaling of $\langle g^q \rangle$ and $\langle p_q(g) \rangle$. As expected, the error bars of \bar{X}'_q are dramatically smaller than those of \bar{X}_q in 2D, which makes a meaningful comparison with Δ_q possible. We thus obtain the first important numerical result of this work: exponents $\bar{X}'_q/2$ and Δ_q agree with accuracy better than 2% in the range $0 < q < 1$. This confirms the exponent relation (9). We interpret the small deviations visible outside the range $0 < q < 1$ as remnants of higher order corrections in eq. (8) not used in the fitting. As in 2D, errors in X'_q are much larger, the fitting remains uncontrolled near $q = 1/2$, and the validity of eq. (2) cannot be established from our data.

Next, we show two plots that highlight an essential difference between the two types of subleading corrections. fig. 3(a) shows results for $|y'|$ and $|\bar{y}'|$ from fitting $\langle g^q \rangle$ and $\langle p_q(g) \rangle$ to eqs. (10) and (11). We see that $|\bar{y}'|$ is essentially q -independent in the range $0 < q < 1$ where we trust our numerical method. This is what one expects from the

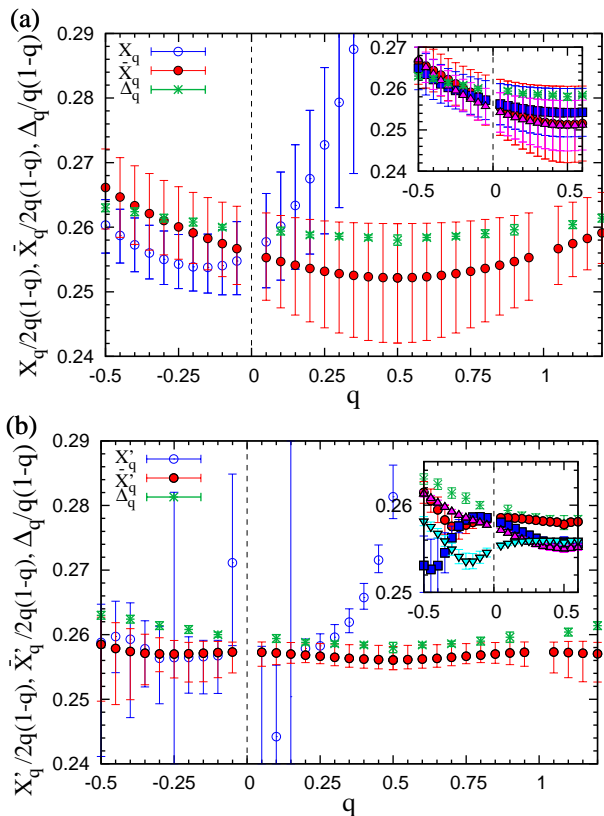


Fig. 2: (a) X_q and \bar{X}_q obtained from PCC in 2D eq. (3) with $N_p = 1$ and $d_q = 0$, see footnote⁶, and eq. (8) with $N_p = 1$. (b) X'_q and \bar{X}'_q obtained from the conductance g in Q1D and eqs. (10) and (11), both with $N_p = 2$, $N_s = 1$, see footnote⁷. We also show the LDOS exponents Δ_q from Ref. [13]. All exponents are divided by $2q(1-q)$, and, except for Δ_q , are shown with the *practical* error bars, see the main text and [44]. Note that the relative errors of the exponents \bar{X}'_q for $0 < q < 1$ are less than 2%. Inset in (a): \bar{X}_q calculated from data sets with different maximum distances r_{\max} ; $r_{\max}=41$ (\bullet), 49 (\blacksquare), and 57 (\blacktriangle). Minimum distance is fixed: $r_{\min}=3$. Inset in (b): \bar{X}'_q calculated from data sets with different maximum width M_{\max} ; $M_{\max} = 192$ (\bullet), 256 (\blacksquare), 384 (\blacktriangle), and 512 (\blacktriangledown). Minimum width $M_{\text{text}}=32$. Error bars represent standard (one-sigma) errors from χ^2 -fitting analysis.

definition of y as a property of the RG fixed point, and not of a particular observable. On the other hand, the subleading exponents γ_q are expected to depend on q in a way that is qualitatively similar to the leading ones X_q . Indeed, this is what is seen in fig. 3(b) where we show an entire spectrum of the subleading exponents $\bar{\gamma}'_q$ for the IQH.

Finally, we establish the validity of eqs. (13) and (14). Figure 4 shows the probability distribution functions for the random variable $\ln g$ on the Q1D cylinder with different aspect ratios $L/M = 5, 10$. The solid curves represent $\mathcal{P}_{\text{Q1D}}(\ln g)$ computed from eqs. (13) and (14) by using \mathcal{N}_{Q1D} specified in eq. (15) and the singularity spectrum $f(\alpha)$ as inputs. Here $f(\alpha)$ is calculated from the LDOS exponents Δ_q from Ref. [12]. We see a very reasonable

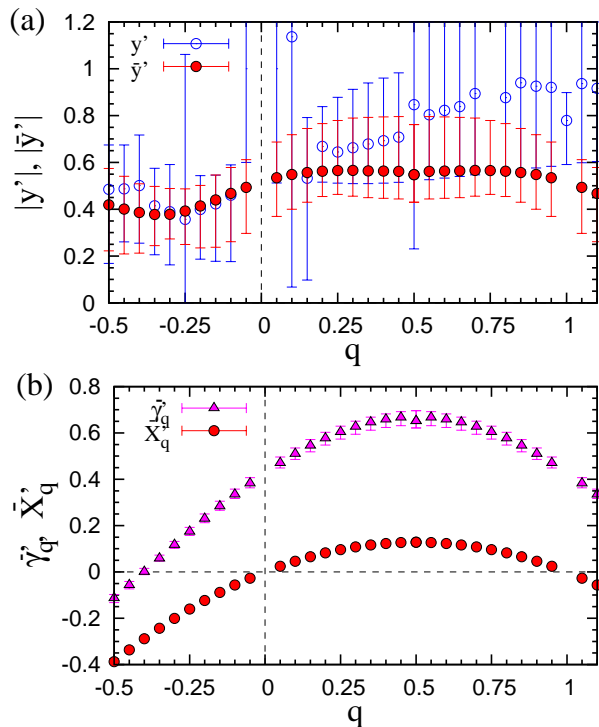


Fig. 3: (a) The irrelevant exponents $|y'|$ and $|\bar{y}'|$ obtained from fits of $\langle g^q \rangle$ and $\langle p_q(g) \rangle$ to eqs. (10) and (11). The large scatter in $|y'|$ reflects, once again, the difficulty to separate subleading exponents of the two kinds. The fact that $|\bar{y}'|$ is largely independent of q in the range $0 < q < 1$ can be viewed as evidence for the quality of the fit. Moreover, the value $|\bar{y}'| \approx 0.6 \pm 0.2$ is consistent with earlier estimates, $|y| \gtrsim 0.4$ [18]. (b) Subleading exponent $\bar{\gamma}'_q$ from the scaling analysis of $\langle p_q(g) \rangle$ and comparison with \bar{X}'_q (Data given with practical error bars.)

agreement between the curves and the symbols.

Remark on experiments. – We need not vary the system size but need only the fixed aspect ratio L/M to obtain the solid curves in fig. 4. In experiments it is not easy to change geometric quantities like L, M or r , used in equations above. At the same time, disorder realizations can presumably be changed, for example, by varying a gate voltage. By sweeping different disorder realizations, one can experimentally access the distribution function of, say, the two-terminal conductance in Q1D, $\mathcal{P}(\ln g, L/M)$, in a fixed geometry. By applying eqs. (13) and (14) to $\mathcal{P}(\ln g, L/M)$ inversely as we demonstrated above, this gives another way to access multifractality.

Conclusions. – In this Letter, we have numerically established relations between spectra of scaling exponents for disorder average moments of PCC and of the LDOS. These results were achieved by a careful finite-size scaling analysis accounting for subleading and irrelevant correc-

⁶We account for a single correction term to $\langle T^q \rangle$ only, i.e. $N_p=1$ and $d_q=0$ in eq. (3) because our data quality does not allow to distinguish the additional terms $\sim r^{-\gamma_q}$.

⁷We note that it is important to take $N_p \geq 2$ in order to obtain consistent results. With $N_p=1$, \bar{X}'_q and $2\Delta_q$ deviate and $|y|$ is too small: $|y| \sim 0.2$. See also Ref. [18].

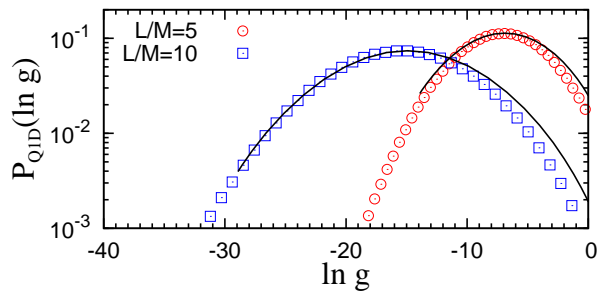


Fig. 4: Symbols show the probability distribution of the Q1D conductance, $\mathcal{P}_{\text{Q1D}}(\ln g)$, for $M = 128$ at $L/M = 5$ and 10 . The curves are computed by eq. (14) with the input of $f(\alpha)$ obtained from the LDOS exponents Δ_q from Ref. [12].

tions, and augmented by the use of the Legendre functions as well as of the stability map.

* * *

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Supplemental material for “Statistics of Conductances and Subleading Corrections to Scaling near the Integer Quantum Hall Plateau Transition”

In this supplemental material we present details of the numerical simulation for the two-terminal conductance in Q1D and the point contact conductance in 2D.

Two-terminal conductance in Q1D. –

Scaling analysis for two-terminal conductances in Q1D.

As we discuss in the main text of our paper, in the quasi-one dimensional (Q1D) system of length L and width M the scaling function for the q -th moment of the two-terminal conductance g , is approximated as

$$\begin{aligned} \langle g^q \rangle &= c'_q \exp \left[-\pi \left(X'_q + \sum_{n=1}^{N_p} a_q'^{(n)} M^{ny'} \right) \frac{L}{M} \right] \\ &+ d'_q \exp \left[-\pi \left(\gamma'_q + \sum_{n=1}^{N_s} b_q'^{(n)} M^{ny'} \right) \frac{L}{M} \right], \end{aligned} \quad (\text{S1})$$

where all the parameters can be used in the fitting procedure. Similarly, for the quantity $p_q(g) = P_{-q}(2/g - 1)$ which involves the Legendre function $P_{-q}(x)$, the corresponding scaling function is approximated by

$$\begin{aligned} \langle p_q(g) \rangle &\equiv \langle P_{-q}(2/g - 1) \rangle = \\ &\bar{c}'_q \exp \left[-\pi \left(\bar{X}'_q + \sum_{n=1}^{N_p} \bar{a}_q'^{(n)} M^{n\bar{y}'} \right) \frac{L}{M} \right] \\ &+ \bar{d}'_q \exp \left[-\pi \left(\bar{\gamma}'_q + \sum_{n=1}^{N_s} \bar{b}_q'^{(n)} M^{n\bar{y}'} \right) \frac{L}{M} \right] \end{aligned} \quad (\text{S2})$$

The coefficients c'_q, d'_q, \bar{c}'_q , and \bar{d}'_q exhibit a weak dependence on the system width M [S5], and we take this fact into account by keeping terms in the Taylor expansion of these coefficients in $1/M$ up to the first order:

$$\begin{aligned} c'_q &\rightarrow c'_q + \frac{c_q'^{(1)}}{M}, & d'_q &\rightarrow d'_q + \frac{d_q'^{(1)}}{M}, \\ \bar{c}'_q &\rightarrow \bar{c}'_q + \frac{\bar{c}_q'^{(1)}}{M}, & \bar{d}'_q &\rightarrow \bar{d}'_q + \frac{\bar{d}_q'^{(1)}}{M}. \end{aligned} \quad (\text{S3})$$

In our numerical simulations, we calculated two-terminal conductances in Q1D systems of widths $M = 32, 48, 64, 96, 128, 192, 256, 384, 512$ and maximum length $L_{\max} = 10M$. The number of samples for each M is 10^6 . We obtain the fitting parameters from the scaling function (S1)–(S3) by varying M and L . We prepare four data sets with the same minimum width $M_{\min} = 32$, but different maximum widths; $M_{\max} = 192, 256, 384, 512$. Since the nonlinear fitting to the functions (S1)–(S3) strongly depends on the initial values of the fitting parameters, we quantify the goodness of fitting by calculating the χ^2 value for each fitting trial. We had more than 1000 fitting trials for each input data set starting with different initial fitting parameters chosen at random, and found the most reliable fitting which gives the minimum χ^2 . This complicated analysis is clearly displayed by using the so-called “stability map” [S6].

The most reliable fitting was obtained when we chose $N_p = 2$ and $N_s = 1$ in Eqs. (S1) and (S2). The details of the most reliable fitting with different M_{\max} for $\langle g^q \rangle$ and $\langle p_q(g) \rangle$ are summarized in Tables S1–S4 and S5–S8,

respectively. The corresponding stability maps are shown in Figs. S1–S4 for $\langle g^q \rangle$ and Figs. S6–S9 for $\langle p_q(g) \rangle$.

From the stability maps, we find that the irrelevant exponents y' and \bar{y}' are broadly distributed in the interval $[0, 1]$. For larger values of M_{\max} , the χ^2 exhibits two minima as a function of y' and \bar{y}' . The global (deeper) minimum of the χ^2 as a function of \bar{y}' is located at $|\bar{y}'| \approx 1/2$ for all q in the range of $0 < q < 1$. We remark that at this minimum the coefficients $a_q'^{(1)}$ and $a_q'^{(2)}$ (also $\bar{a}_q'^{(1)}$ and $\bar{a}_q'^{(2)}$) take on opposite signs (except near $q \simeq 0$ for $\langle g^q \rangle$ where the scaling analysis becomes unstable). The other minimum, with a higher value of χ^2 , is located in the region of very small values of $|y'|$ or $|\bar{y}'|$, consistent with $|y| \sim 0$, as found previously by other groups. However, at this (higher) minimum with larger χ^2 , the values of $a_q'^{(2)}$ and $\bar{a}_q'^{(2)}$ become almost zero. Consistently, if we set $N_p = 1$ (in other words, $a_q'^{(2)}$ and $\bar{a}_q'^{(2)}$ are fixed to zero), we obtain $|y|$ close to zero, but a larger χ^2 . This observation is completely analogous to what we saw in the scaling analysis for the Lyapunov exponent in Ref. [S6].

We also remark that when fitting data for $\langle g^q \rangle$, the χ^2 as a function γ'_q possesses two local minima, with close values of χ^2 , when $q < 0.5$. One minimum is located at larger γ'_q , where most of the fitting results lie (a dense cloud of points). The other minimum at smaller γ'_q is comprised of a smaller number of points, see Figs. S1–S4. This makes the reliable determination of γ'_q difficult. In contrast, for $\langle p_q(g) \rangle$, the χ^2 clearly shows a well defined global minimum as a function of $\bar{\gamma}'_q$, as shown in Figs. S6–S9. Since the larger γ'_q is close to $\bar{\gamma}'_q$, we believe that the larger γ'_q might be closer to the correct results. However, a further careful analysis is needed to firmly establish this.

Practical error bars. The q dependence of the numerically obtained exponents X'_q, y' , and γ'_q is shown in Fig. S5. Fig. S10 shows the q dependence of the exponents \bar{X}'_q, \bar{y}' , and $\bar{\gamma}'_q$. Both figures show that the exponents obtained from different data sets do not agree with each other even when the error bars (estimated from the error-propagation theory, thin lines) are taken into account. We believe that the insufficient truncations for the irrelevant exponent y and the subleading scaling dimension γ_q of the scaling functions give rise to this inconsistency, while increasing the truncation orders is impossible with our numerical accuracy. Therefore, as a conservative upper bound for the errors, we introduce the “practical error bars” that represent the union of all “statistical error bars” obtained from the different data sets, as shown by the thick red lines in Figs. S5 and S10. The red cross symbol represents the average of the mean values obtained from all data sets.

Table S1: The details of the most reliable fitting by the scaling analysis for the q th moment of Q1D with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 192$. The aspect ratio is varied from $L/M = 4$ to 10 and the number of total data point N is 376. The scaling function in Eq. (S1) with $N_p = 2$ and $N_s = 1$ is employed. χ^2_{\min} , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value.

q	X'_q	c'_q	γ'_q	d'_q	$ y' $	c''_q	a''_q	$a^{(2)}_q$	$d^{(1)}_q$	$b^{(1)}_q$	χ^2_{\min}/N	Q
-0.5	-0.396484 ± 0.000583	0.334988 ± 0.004119	-0.385740 ± 0.041406	-0.037842 ± 0.021902	0.615961 ± 0.003918	1.512426 ± 0.151471	-0.191301 ± 0.012190	0.864221 ± 0.923938	21.300731 ± 6.374798	4.329661 ± 0.376499	0.49	1.0
-0.4	-0.295630 ± 0.000281	0.456913 ± 0.002763	-0.245458 ± 0.031209	-0.023908 ± 0.013898	0.713652 ± 0.003596	1.300010 ± 0.102123	-0.186081 ± 0.004999	1.098436 ± 0.847415	16.920341 ± 4.070400	5.527381 ± 0.419233	0.44	1.0
-0.3	-0.203606 ± 0.000173	0.600194 ± 0.001332	-0.091364 ± 0.039858	0.003585 ± 0.015298	0.578024 ± 0.012599	0.835473 ± 0.057955	-0.091235 ± 0.022121	0.269477 ± 0.802268	14.502444 ± 4.063576	2.760727 ± 0.356532	0.40	1.0
-0.2	-0.123226 ± 0.001385	0.746903 ± 0.000531	0.276641 ± 0.038642	0.146892 ± 0.040293	0.384617 ± 0.078072	0.491493 ± 0.033404	-0.033404 ± 0.050944	0.035837 ± 0.168455	-2.953260 ± 1.094452	-0.406236 ± 0.139222	0.38	1.0
-0.1	-0.056740 ± 0.000159	0.885483 ± 0.000239	0.356603 ± 0.047604	0.079055 ± 0.029580	0.505266 ± 0.025619	0.250012 ± 0.015214	-0.020543 ± 0.007365	0.037323 ± 0.035740	-1.557074 ± 0.795048	-0.466589 ± 0.267361	0.45	1.0
0.1	0.045423 ± 0.000373	0.933196 ± 0.004217	0.059392 ± 0.001791	0.145973 ± 0.044286	1.819709 ± 0.042547	-12.443617 ± 1.166978	-2.059302 ± 0.843347	2962.0 ± 760.7	12.073200 ± 1.164921	-7.657070 ± 1.303236	0.41	1.0
0.2	0.082870 ± 0.000118	1.059712 ± 0.004121	0.189950 ± 0.010884	0.085366 ± 0.002654	0.809381 ± 0.029176	-0.283963 ± 0.170854	0.072965 ± 0.011911	-0.464585 ± 0.117745	0.203530 ± 0.207244	0.467173 ± 0.287221	0.35	1.0
0.3	0.109664 ± 0.000193	0.976775 ± 0.006798	0.215777 ± 0.006812	0.211366 ± 0.003975	0.885299 ± 0.026777	-0.164929 ± 0.307484	0.125748 ± 0.026669	-1.078541 ± 0.283988	0.373257 ± 0.294275	0.594414 ± 0.250716	0.28	1.0
0.4	0.128321 ± 0.000290	0.843614 ± 0.009104	0.234724 ± 0.005187	0.363946 ± 0.005022	0.935891 ± 0.024327	0.110747 ± 0.439973	0.180456 ± 0.048966	-1.871690 ± 16.035714	0.535342 ± 0.363941	0.719350 ± 1.004037	0.25	1.0
0.5	0.142707 ± 0.000423	0.740297 ± 0.007577	0.267093 ± 0.003670	0.499548 ± 0.005412	1.231197 ± 1.000671	-2.764682 ± 0.331021	0.259055 ± 2.237841	-16.035714 ± 88.658580	1.068127 ± 0.407653	-1.004037 ± 0.266620	0.30	1.0
0.6	0.151768 ± 0.000451	0.627423 ± 0.007929	0.284501 ± 0.003592	0.609876 ± 0.006450	1.120021 ± 0.119334	-3.141768 ± 0.433632	0.144743 ± 0.404229	-8.451723 ± 4.747448	0.700441 ± 0.441764	-0.872151 ± 0.236465	0.32	1.0
0.7	0.158097 ± 0.000488	0.530112 ± 0.007237	0.296865 ± 0.003263	0.683248 ± 0.007467	1.201198 ± 0.105530	-2.853750 ± 0.379234	0.160148 ± 0.548963	-14.984336 ± 7.321061	0.316265 ± 0.464925	-0.990592 ± 0.280980	0.39	1.0
0.8	0.162652 ± 0.000522	0.453488 ± 0.006546	0.307941 ± 0.003120	0.726259 ± 0.008520	1.264022 ± 0.102332	-2.528339 ± 0.334664	0.173028 ± 0.729360	-23.274193 ± 10.609417	-0.035229 ± 0.490302	-1.083277 ± 0.329744	0.46	1.0
0.9	0.166012 ± 0.000542	0.392766 ± 0.005839	0.317282 ± 0.003020	0.745728 ± 0.009373	1.347204 ± 0.117205	-2.177898 ± 0.283894	0.212830 ± 1.096200	-40.816932 ± 21.180793	-0.248847 ± 0.506087	-1.204268 ± 0.399400	0.52	1.0
1.1	0.170613 ± 0.000582	0.306800 ± 0.004818	0.332747 ± 0.003029	0.742630 ± 0.010748	1.466717 ± 0.144936	-1.683694 ± 0.222427	0.288416 ± 2.015521	-91.947556 ± 58.559267	-0.508610 ± 0.525351	-1.409160 ± 0.553678	0.62	1.0

Table S2: The details of the most reliable fitting by the scaling analysis for the q th moment of Q1D two-terminal conductances, (g^q) , with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 256$. The aspect ratio is varied from $L/M = 4$ to 10 and the number of total data point N is 486. The scaling function in Eq. (S1) with $N_p = 2$ and $N_s = 1$ is employed. χ^2_{\min} , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value.

q	χ'_q	c'_q	γ'_q	d'_q	$ y' $	$c_q^{(1)}$	$a_q^{(1)}$	$a_q^{(2)}$	$d_q^{(1)}$	$b_q^{(1)}$	χ^2_{\min}/N	Q
-0.5	-0.374025 ± 0.012324	0.346534 ± 0.002122	-0.084504 ± 0.074455	0.185819 ± 0.071099	0.251868 ± 0.082999	0.940797 ± 0.118917	-0.141092 ± 0.008548	0.144466 ± 0.058412	-1.251424 ± 2.040957	-0.025914 ± 0.209068	0.58	1.0
-0.4	-0.286756 ± 0.009950	0.473392 ± 0.001203	0.167188 ± 0.028322	0.102006 ± 0.008591	0.398169 ± 0.142818	-0.627689 ± 0.149515	-0.116813 ± 0.374925	0.160025 ± 1.175179	-1.503558 ± 0.288082	-1.679846 ± 0.124436	0.47	1.0
-0.3	-0.197819 ± 0.003366	0.603422 ± 0.000786	0.140468 ± 0.053621	0.132614 ± 0.034540	0.270860 ± 0.074159	0.608887 ± 0.057062	-0.46220 ± 0.021813	0.030099 ± 0.045291	-0.668488 ± 1.077597	-0.150138 ± 0.169112	0.35	1.0
-0.2	-0.124062 ± 0.000297	0.747853 ± 0.000364	0.303349 ± 0.036161	0.185041 ± 0.053405	0.480888 ± 0.016811	0.458064 ± 0.029080	-0.042083 ± 0.015642	0.069109 ± 0.069764	-3.587427 ± 1.403837	-0.563118 ± 0.199787	0.31	1.0
-0.1	-0.056940 ± 0.000663	0.885701 ± 0.000175	0.374825 ± 0.044380	0.089072 ± 0.033261	0.578410 ± 0.011101	0.241196 ± 0.013827	-0.025016 ± 0.004651	0.062067 ± 0.028487	-1.754863 ± 0.882509	-0.662210 ± 0.348013	0.33	1.0
0.1	0.042460 ± 0.001092	0.557804 ± 0.118796	0.052706 ± 0.001165	0.521044 ± 0.118722	2.320982 ± 0.137773	-11.656885 ± 0.528684	-13.155058 ± 25.287906	167457.1 ± 134345.3	11.299898 ± 0.531595	-17.740150 ± 5.152539	0.32	1.0
0.2	0.082486 ± 0.000112	1.060131 ± 0.003438	0.198163 ± 0.009352	0.088975 ± 0.002178	0.659897 ± 0.028616	-0.478111 ± 0.177592	0.046565 ± 0.004462	-0.171579 ± 0.039908	-0.027949 ± 0.180031	0.010017 ± 0.128064	0.32	1.0
0.3	0.109019 ± 0.000186	0.978669 ± 0.005436	0.221489 ± 0.005629	0.216708 ± 0.003258	0.668204 ± 0.031414	-0.697550 ± 0.271629	0.063620 ± 0.007652	-0.259034 ± 0.068076	0.064887 ± 0.266990	0.040643 ± 0.078697	0.30	1.0
0.4	0.127528 ± 0.000271	0.848067 ± 0.006943	0.239141 ± 0.004114	0.369736 ± 0.004101	0.676025 ± 0.033118	-0.839083 ± 0.332060	0.076953 ± 0.012272	-0.349509 ± 0.102073	0.237335 ± 0.330397	0.068236 ± 0.058631	0.31	1.0
0.5	0.140682 ± 0.000358	0.705505 ± 0.009290	0.253032 ± 0.004080	0.503682 ± 0.004469	0.937536 ± 0.020115	-0.332060 ± 0.599821	0.212552 ± 0.045467	-2.718920 ± 0.590415	0.710788 ± 0.358266	0.442181 ± 0.222669	0.41	1.0
0.6	0.148957 ± 0.000400	0.582303 ± 0.008183	0.265807 ± 0.003344	0.605345 ± 0.004775	0.895239 ± 0.018798	-0.001142 ± 0.468364	0.201420 ± 0.057284	-2.061302 ± 0.529787	0.980186 ± 0.404138	0.492228 ± 0.143559	0.42	1.0
0.7	0.154006 ± 0.000477	0.465265 ± 0.008247	0.269879 ± 0.003174	0.654282 ± 0.005845	1.044011 ± 0.008468	1.877978 ± 0.450299	0.475836 ± 0.116308	-6.498344 ± 1.420084	2.597988 ± 0.597372	1.893654 ± 0.271606	0.56	1.0
0.8	0.159207 ± 0.000507	0.417027 ± 0.006730	0.290285 ± 0.003201	0.701083 ± 0.006286	0.863284 ± 0.019653	-0.091367 ± 0.415482	0.191149 ± 0.081644	-1.834479 ± 0.574854	1.618341 ± 2.611038	0.459509 ± 0.829648	0.51	1.0
0.9	0.162161 ± 0.000567	0.355298 ± 0.006615	0.297012 ± 0.003434	0.704691 ± 0.007738	0.927948 ± 0.014348	0.402536 ± 0.415482	0.270416 ± 0.122227	-2.934130 ± 0.901206	2.611038 ± 0.682387	0.517534 ± 0.186203	0.57	1.0
1.0	0.164898 ± 0.000609	0.318221 ± 0.005372	0.308348 ± 0.003249	0.711567 ± 0.007860	0.813504 ± 0.021413	-0.210609 ± 0.319641	0.160795 ± 0.095576	-1.358406 ± 0.508608	1.988185 ± 0.601572	0.392484 ± 0.102036	0.59	1.0
1.1	0.166763 ± 0.000786	0.286850 ± 0.004490	0.317352 ± 0.003123	0.707403 ± 0.008028	0.685913 ± 0.030186	-0.555660 ± 0.240422	0.102332 ± 0.068681	-0.638286 ± 0.275390	1.478133 ± 0.523968	0.184979 ± 0.052628	0.62	1.0

Table S3: The details of the most reliable fitting by the scaling analysis for the q th moment of QID two-terminal conductances, $\langle g^q \rangle$, with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 384$. The aspect ratio is varied from $L/M = 4$ to 10 and the number of total data point N is 717. The scaling function in Eq. (S1) with $N_p = 2$ and $N_s = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value.

q	X'_q	c'_q	γ'_q	d'_q	$ y' $	$c_q^{(1)}$	$a_q^{(1)}$	$a_q^{(2)}$	$d_q^{(1)}$	$b_q^{(1)}$	χ_{\min}^2/N	Q
-0.5	± 0.395720	0.340006	-0.326910	-0.017778	0.671876	1.317255	-0.249383	1.356729	20.077285	4.696611	0.66	1.0
	± 0.000352	± 0.003285	± 0.038849	± 0.013783	± 0.002961	± 0.119577	± 0.005936	± 0.849714	± 6.070049	± 0.328406		
-0.4	-0.293761	0.461188	-0.236389	-0.021383	0.614671	1.138517	-0.156072	0.617612	19.336989	3.921922	0.52	1.0
	± 0.000139	± 0.001987	± 0.030214	± 0.010638	± 0.003517	± 0.074826	± 0.011188	± 0.569776	± 4.772048	± 0.217845		
-0.3	-0.203141	0.600304	-0.142294	-0.025818	0.541516	0.867346	-0.084613	0.220073	19.421414	3.105257	0.35	1.0
	± 0.000297	± 0.001009	± 0.027045	± 0.010032	± 0.007270	± 0.041900	± 0.030561	± 0.657579	± 4.667253	± 0.156559		
-0.2	-0.124533	0.745625	-0.057750	-0.020423	0.539831	0.595920	-0.050023	0.118521	13.909845	3.112354	0.27	1.0
	± 0.000111	± 0.000569	± 0.026918	± 0.007383	± 0.003667	± 0.024887	± 0.009911	± 0.197698	± 3.909845	± 0.162738		
-0.1	-0.056965	0.884568	-0.011977	-0.008031	0.575543	0.304396	-0.025089	0.066775	4.821933	3.660181	0.27	1.0
	± 0.000042	± 0.000329	± 0.030842	± 0.003290	± 0.008134	± 0.013567	± 0.006657	± 0.166176	± 1.362451	± 0.216194		
0.1	0.045894	1.067966	0.195670	0.014069	0.334925	-1.007886	0.008435	-0.007873	0.663214	-0.351494	0.28	1.0
	± 0.000053	± 0.001569	± 0.030526	± 0.000715	± 0.093046	± 0.152661	± 0.003236	-0.018802	± 0.180493	± 0.088396		
0.2	0.082196	1.063537	0.205673	0.089457	0.535820	-0.431712	0.031856	-0.065195	-0.051170	0.042202	0.29	1.0
	± 0.000091	± 0.002328	± 0.007443	± 0.002182	± 0.020155	± 0.105426	± 0.002575	± 0.013886	± 0.164434	± 0.053395		
0.3	0.108576	0.984334	0.228695	0.218769	0.533450	-0.717027	0.043372	-0.099664	-0.116867	0.004686	0.29	1.0
	± 0.000154	± 0.003887	± 0.004536	± 0.002963	± 0.023946	± 0.186956	± 0.003585	-0.022972	± 0.252757	-0.031933		
0.4	0.126913	0.854166	0.245322	0.374784	0.538326	-0.992159	0.054072	-0.146638	-0.136713	-0.002037	0.31	1.0
	± 0.000225	± 0.005171	± 0.003329	± 0.003617	± 0.027028	± 0.251016	± 0.004509	± 0.035294	± 0.323453	± 0.023722		
0.5	0.139277	0.714367	0.260212	0.517803	0.549803	-1.150657	0.064446	-0.206920	-0.213162	-0.001927	0.35	1.0
	± 0.000292	± 0.005661	± 0.002749	± 0.004112	± 0.029594	-0.278989	± 0.005549	± 0.051709	± 0.359718	± 0.020714		
0.6	0.147631	0.592085	0.273844	0.624854	0.564980	-1.178008	0.074186	-0.276204	-0.373175	-0.000696	0.41	1.0
	± 0.000347	± 0.005516	± 0.002456	± 0.004569	± 0.032071	-0.279778	± 0.006796	± 0.072206	± 0.370595	-0.020158		
0.7	0.153418	0.494535	0.286028	0.693488	0.581560	-1.122735	0.083005	-0.349747	-0.573348	0.000425	0.48	1.0
	± 0.000393	± 0.005098	± 0.002315	± 0.005111	± 0.034642	-0.268121	± 0.008255	± 0.096497	± 0.371057	± 0.020956		
0.8	0.157563	0.419028	0.296664	0.730811	0.598397	-1.030758	0.090876	-0.424499	-0.768317	0.001410	0.55	1.0
	± 0.000429	± 0.004631	± 0.002260	± 0.005699	± 0.037302	-0.252956	± 0.009943	-0.124317	± 0.369912	± 0.022619		
0.9	0.160643	0.360612	0.305810	0.745425	0.616105	-0.929237	0.098217	-0.502770	-0.931498	0.002489	0.62	1.0
	± 0.000459	± 0.004206	± 0.002257	± 0.006256	± 0.039903	-0.238435	± 0.011945	± 0.156345	± 0.370466	± 0.025042		
1.0	0.162995	0.314874	0.313608	0.744568	0.633286	-0.829949	0.104929	-0.580950	-1.054404	0.004000	0.68	1.0
	± 0.000485	± 0.003845	± 0.002287	± 0.006743	± 0.042303	-0.225571	± 0.014395	± 0.192205	± 0.373426	± 0.028053		
1.1	0.164827	0.278484	0.320239	0.733618	0.647788	-0.738353	0.110543	-0.649788	-1.138348	0.006048	0.74	1.0
	± 0.000509	± 0.003544	± 0.002338	± 0.007147	± 0.044373	± 0.214324	± 0.017403	-0.230573	± 0.378406	± 0.031454		

Table S4: The details of the most reliable fitting by the scaling analysis for the q th moment of Q1D two-terminal conductances, (g^q) , with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 512$. The aspect ratio is varied from $L/M = 4$ to 10 and the number of total data point N is 1025. The scaling function in Eq. (S1) with $N_p = 2$ and $N_s = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value.

q	X'_q	c'_q	γ'_q	d'_q	$ y' $	$c_q^{(1)}$	$a_q^{(1)}$	$a_q^{(2)}$	$d_q^{(1)}$	$b_q^{(1)}$	χ_{\min}^2/N	Q
-0.5	-0.386598	0.355199	1156121.4	-73655.7	0.400457	1.233223	-0.145448	0.284554	2863053.8	1768961.8	1.1	0.1
	± 0.000851	± 0.000426	± 0.000000	± 0.000000	± 0.000000	± 0.040532	± 0.143532	± 0.636583	± 0.000000	± 0.000000	0.52	1.0
-0.4	-0.285471	0.468068	0.202874	0.284177	0.278106	0.748748	-0.068749	0.051729	-5.903556	-0.554412		
	± 0.000753	± 0.000608	± 0.041075	± 0.087677	± 0.013802	± 0.073958	± 0.047575	± 0.111431	± 2.348020	± 0.115770	0.46	1.0
-0.3	-0.195259	0.603252	0.331848	0.120448	0.168540	0.555443	-0.029432	0.000781	-2.151591	-0.587345		
	± 0.000432	± 0.000461	± 0.043747	± 0.026305	± 0.006185	± 0.070884	± 0.012892	± 0.018963	± 0.723725	± 0.095318	0.43	1.0
-0.2	-0.120591	0.747082	0.372864	0.139630	0.190979	0.484278	-0.018086	0.001351	-2.783243	-0.405702		
	± 0.000140	± 0.000251	± 0.039187	± 0.033539	± 0.004331	± 0.028305	± 0.005386	± 0.008701	± 0.889503	± 0.090513	0.44	1.0
-0.1	-0.055282	0.885258	0.460095	0.054447	0.180447	0.251179	-0.006403	-0.001728	-1.159937	-0.452971		
	± 0.000086	± 0.000125	± 0.048905	± 0.015723	± 0.003900	± 0.014625	± 0.002073	± 0.003151	± 0.434422	± 0.107653	0.25	1.0
0.1	0.042051	1.070980	0.665873	0.012572	0.070900	-1.173374	0.008771	-0.001957	0.786909	-0.741842		
	± 0.002294	± 0.000516	± 0.078593	± 0.000620	± 0.003522	± 0.094736	± 0.007228	± 0.005814	± 0.120696	± 0.098851	0.24	1.0
0.2	-0.492833	0.090772	0.078921	1.073297	0.101720	0.216614	2.772094	-2.668405	-2.065717	0.008485		
	± 0.921458	± 0.002326	± 0.000135	± 0.001364	± 0.051380	± 0.175778	± 2.507368	± 1.483259	± 0.112880	± 0.000251	0.25	1.0
0.3	0.108860	0.987807	0.229516	0.217078	0.562664	-0.581562	0.046330	-0.113620	0.027827	0.066343		
	± 0.000109	± 0.003174	± 0.003909	± 0.002857	± 0.014958	± 0.147459	± 0.003595	± 0.020374	± 0.231225	± 0.028948	0.29	1.0
0.4	0.127492	0.855787	0.243509	0.371636	0.620295	-0.673724	0.066371	-0.227364	0.068935	0.078907		
	± 0.000158	± 0.004451	± 0.002851	± 0.003238	± 0.014752	± 0.210683	± 0.005992	± 0.039849	± 0.292004	± 0.027045	0.35	1.0
0.5	0.140026	0.712788	0.256927	0.513739	0.667955	-0.656912	0.086441	-0.383218	0.089494	0.101106		
	± 0.000207	± 0.005014	± 0.002348	± 0.003637	± 0.014518	± 0.248694	± 0.009667	± 0.068776	± 0.326587	± 0.028483	0.41	1.0
0.6	0.148437	0.587813	0.269673	0.619456	0.708522	-0.562059	0.106394	-0.575608	0.084932	0.128993		
	± 0.000249	± 0.004954	± 0.002097	± 0.004045	± 0.014092	± 0.263111	± 0.014764	± 0.107626	± 0.340412	± 0.032351	0.46	1.0
0.7	0.154193	0.488094	0.281016	0.685866	0.748709	-0.414572	0.128962	-0.826268	0.101280	0.168673		
	± 0.000283	± 0.004631	± 0.001982	± 0.004505	± 0.013204	± 0.266119	± 0.021752	± 0.161469	± 0.347094	± 0.039636	0.52	1.0
0.8	0.158237	0.411238	0.290897	0.720467	0.779904	-0.269262	0.149988	-1.075971	0.150156	0.212284		
	± 0.000313	± 0.004239	± 0.001943	± 0.005010	± 0.012416	± 0.119697	± 0.029543	± 0.219484	± 0.357585	± 0.047417	0.56	1.0
0.9	0.161171	0.351922	0.299270	0.732002	0.809492	-0.119697	0.173328	-1.357778	0.259230	0.269275		
	± 0.000339	± 0.003886	± 0.001955	± 0.005490	± 0.011384	± 0.255154	± 0.038760	± 0.290155	± 0.377828	± 0.058121	0.61	1.0
1.0	0.163260	0.303332	0.304956	0.723150	0.889849	0.287054	0.258561	-2.509821	0.905139	0.524088		
	± 0.000375	± 0.003755	± 0.002080	± 0.005950	± 0.008173	± 0.273767	± 0.059153	± 0.498663	± 0.460971	± 0.101022	0.64	1.0
1.1	0.164987	0.208512	0.311882	0.712968	0.864948	0.184440	0.231403	-2.034678	0.728268	0.446166		
	± 0.000391	± 0.003371	± 0.002084	± 0.006322	± 0.009014	± 0.245992	± 0.060894	± 0.470888	± 0.458256	± 0.089392		

Table S5: The details of the most reliable fitting by the scaling analysis for the Legendre function as a function of QID two-terminal conductances, $\langle p_q(g) \rangle$, with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 192$. The aspect ratio is varied from $L/M = 2$ to 10 and the number of total data points N is 420 except at $q = 0.5$. At $q = 0.5$, the maximum aspect ratio is limited to 8 due to our numerical accuracy and the number of total data points N is 324. The scaling function in Eq. (S2) with $N_p = 2$ and $N_s = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value. Note that the Legendre function satisfies $p_q(x) = p_{1-q}(x)$.

q	\bar{X}'_q	\bar{c}'_q	$\bar{\gamma}'_q$	\bar{a}'_q	$ \bar{y}' $	$\bar{c}''(1)$	$\bar{a}''(1)$	$\bar{c}''(2)$	$\bar{a}''(2)$	$\bar{d}''(1)$	$\bar{b}''(1)$	χ_{\min}^2	Q
0.5	0.128892 ± 0.000186	1.259646 ± 0.000539	0.663351 ± 0.032673	-0.070738 ± 0.004694	0.727002 ± 0.034635	-0.730887 ± 0.039244	0.086980 ± 0.019364	-0.400856 ± 0.162621	0.578356 ± 0.283376	0.578356 ± 0.283376	-1.153155 ± 0.743605	0.48	1.0
0.6	0.123889 ± 0.000137	1.248526 ± 0.000404	0.653068 ± 0.029383	-0.069492 ± 0.004502	0.762378 ± 0.032117	-0.727757 ± 0.029553	0.092810 ± 0.016732	-0.505492 ± 0.167836	0.552401 ± 0.291098	0.552401 ± 0.291098	-0.928184 ± 0.802835	0.47	1.0
0.7	0.108477 ± 0.000115	1.215405 ± 0.000354	0.621069 ± 0.026634	-0.060637 ± 0.003541	0.759376 ± 0.030471	-0.612117 ± 0.024617	0.080862 ± 0.013396	-0.433034 ± 0.130034	0.434811 ± 0.234375	0.434811 ± 0.234375	-0.858006 ± 0.712931	0.49	1.0
0.8	0.082704 ± 0.000085	1.161538 ± 0.000233	0.570880 ± 0.022989	-0.045972 ± 0.002286	0.736683 ± 0.028905	-0.438758 ± 0.017225	0.057818 ± 0.008676	-0.280092 ± 0.078295	0.270943 ± 0.156146	0.270943 ± 0.156146	-0.676493 ± 0.571802	0.51	1.0
0.9	0.046539 ± 0.000051	1.088872 ± 0.000118	0.504700 ± 0.019248	-0.025636 ± 0.001030	0.689558 ± 0.028965	-0.228973 ± 0.008590	0.028392 ± 0.003829	-0.112529 ± 0.029896	0.110458 ± 0.072334	0.110458 ± 0.072334	-0.420938 ± 0.397680	0.53	1.0
1.1	-0.056819 ± 0.000116	0.898015 ± 0.000119	0.330462 ± 0.013362	0.030151 ± 0.000698	0.533879 ± 0.038792	0.236937 ± 0.008180	-0.022552 ± 0.002781	0.046539 ± 0.014215	0.053905 ± 0.050159	0.053905 ± 0.050159	-0.054616 ± 0.143303	0.48	1.0
1.2	-0.123708 ± 0.000472	0.786433 ± 0.000255	0.228334 ± 0.011620	0.063849 ± 0.001093	0.444370 ± 0.052059	0.474406 ± 0.016664	-0.039654 ± 0.005330	0.057091 ± 0.021706	-0.103179 ± 0.077748	-0.103179 ± 0.077748	-0.010408 ± 0.084169	0.45	1.0
1.3	-0.200909 ± 0.001277	0.669308 ± 0.000451	0.118980 ± 0.010354	0.099579 ± 0.001318	0.395918 ± 0.066295	0.707446 ± 0.027767	-0.060041 ± 0.008268	0.076333 ± 0.033156	-0.179595 ± 0.094707	-0.179595 ± 0.094707	0.001151 ± 0.060417	0.45	1.0
1.4	-0.290603 ± 0.001644	0.552025 ± 0.000779	0.006154 ± 0.009057	0.134874 ± 0.001529	0.440857 ± 0.067659	0.920555 ± 0.045873	-0.099891 ± 0.013283	0.172678 ± 0.070931	-0.248832 ± 0.115358	-0.248832 ± 0.115358	0.008996 ± 0.063575	0.48	1.0
1.5	-0.392190 ± 0.001864	0.441646 ± 0.001255	-0.106291 ± 0.008823	0.165908 ± 0.001810	0.504153 ± 0.069851	1.090853 ± 0.072198	-0.160472 ± 0.021969	0.396645 ± 0.162391	-0.283190 ± 0.143105	-0.283190 ± 0.143105	0.017538 ± 0.080184	0.53	1.0

Table S6: The details of the most reliable fitting by the scaling analysis for the Legendre function as a function of Q1D two-terminal conductances, $\langle p_q(g) \rangle$, with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 256$. The aspect ratio is varied from $L/M = 2$ to 10 and the number of total data points N is 651 except at $q = 0.5$. At $q = 0.5$, the maximum aspect ratio is limited to 8 due to our numerical accuracy and the number of total data points N is 501. The scaling function in Eq. (S2) with $N_p = 2$ and $N_g = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value. Note that the Legendre function satisfies $p_q(x) = p_{1-q}(x)$.

q	\bar{X}'_q	\bar{c}'_q	$\bar{\gamma}'_q$	\bar{d}'_q	$ \bar{y}' $	\bar{c}''_q	\bar{a}''_q	\bar{a}'_q	\bar{d}''_q	\bar{b}'_q	χ_{\min}^2/N	Q
0.5	0.127866 ± 0.000315	1.261090 ± 0.000376	0.643272 ± 0.021914	-0.070508 ± 0.002875	0.507135 ± 0.038514	-0.799714 ± 0.026853	0.045831 ± 0.009781	-0.084169 ± 0.041462	0.662712 ± 0.165803	-0.416566 ± 0.198128	0.54	1.0
0.6	0.122687 ± 0.000281	1.249100 ± 0.000271	0.659716 ± 0.021146	-0.071088 ± 0.002907	0.500801 ± 0.034920	-0.758539 ± 0.020608	0.043466 ± 0.008335	-0.077862 ± 0.034498	0.642750 ± 0.172550	-0.384260 ± 0.193570	0.55	1.0
0.7	0.107596 ± 0.000195	1.215897 ± 0.000222	0.632244 ± 0.018603	-0.062683 ± 0.002293	0.528901 ± 0.029497	-0.637870 ± 0.017332	0.041063 ± 0.007099	-0.082668 ± 0.031778	0.561028 ± 0.139394	-0.456715 ± 0.194168	0.56	1.0
0.8	0.082213 ± 0.000112	1.161896 ± 0.000155	0.586949 ± 0.015510	-0.048142 ± 0.001488	0.561493 ± 0.023762	-0.456236 ± 0.012458	0.034334 ± 0.005101	-0.079043 ± 0.025003	0.417977 ± 0.094039	-0.540379 ± 0.189623	0.57	1.0
0.9	0.046383 ± 0.000049	1.089063 ± 0.000080	0.522848 ± 0.012617	-0.027068 ± 0.000682	0.588911 ± 0.019689	-0.237169 ± 0.006553	0.021005 ± 0.002605	-0.054337 ± 0.013829	0.215229 ± 0.045778	-0.561435 ± 0.178173	0.56	1.0
1.1	-0.056907 ± 0.000062	0.897894 ± 0.000086	0.348975 ± 0.008584	0.031561 ± 0.000495	0.558288 ± 0.020100	0.239629 ± 0.006973	-0.023878 ± 0.002508	0.054711 ± 0.012125	-0.162930 ± 0.037352	-0.309616 ± 0.110300	0.47	1.0
1.2	-0.123995 ± 0.000231	0.786445 ± 0.000185	0.242881 ± 0.007642	0.065509 ± 0.000794	0.474472 ± 0.027543	0.471068 ± 0.014267	-0.042281 ± 0.004695	0.068967 ± 0.018227	-0.215574 ± 0.062479	-0.127185 ± 0.068129	0.40	1.0
1.3	-0.199888 ± 0.001040	0.669903 ± 0.000323	0.123813 ± 0.007751	0.099677 ± 0.000949	0.351452 ± 0.046133	0.690151 ± 0.022413	-0.055131 ± 0.005466	0.056758 ± 0.018588	-0.139264 ± 0.078571	-0.000915 ± 0.039917	0.39	1.0
1.4	-0.283435 ± 0.003525	0.553779 ± 0.000545	-0.003166 ± 0.008384	0.132114 ± 0.001070	0.276741 ± 0.064327	0.878793 ± 0.033813	-0.081561 ± 0.002545	0.071765 ± 0.018261	0.015882 ± 0.096293	0.062143 ± 0.031161	0.45	1.0
1.5	-0.379616 ± 0.005078	0.445143 ± 0.000866	-0.123804 ± 0.008410	0.159553 ± 0.001250	0.289822 ± 0.067184	1.014620 ± 0.049782	-0.127824 ± 0.008976	0.136705 ± 0.025515	0.230040 ± 0.122484	0.108348 ± 0.032876	0.54	1.0

Table S7: The details of the most reliable fitting by the scaling analysis for the Legendre function as a function of QID two-terminal conductances, $\langle p_q(g) \rangle$, with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 384$. The aspect ratio is varied from $L/M = 2$ to 10 and the number of total data points N is 997 except at $q = 0.5$. At $q = 0.5$, the maximum aspect ratio is limited to 8 due to our numerical accuracy and the number of total data points N is 784. The scaling function in Eq. (S2) with $N_p = 2$ and $N_s = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value. Note that the Legendre function satisfies $p_q(x) = p_{1-q}(x)$.

q	\bar{X}'_q	\bar{c}'_q	$\bar{\sigma}'_q$	\bar{a}'_q	$ \bar{y}' $	$\bar{c}'_q(1)$	$\bar{a}'_q(1)$	$\bar{a}'_q(2)$	$\bar{d}'_q(1)$	$\bar{b}'_q(1)$	χ_{\min}^2/N	Q
0.5	0.127501 ± 0.000214	1.261022 ± 0.000273	0.652190 ± 0.016101	-0.069352 ± 0.002174	0.453816 ± 0.023808	-0.787101 ± 0.023206	0.039208 ± 0.006251	-0.054526 ± 0.021709	0.667292 ± 0.126287	-0.445857 ± 0.116736	0.54	1.0
0.6	0.122479 ± 0.000172	1.249190 ± 0.000203	0.664432 ± 0.014975	-0.069886 ± 0.002151	0.467256 ± 0.020249	-0.753561 ± 0.018122	0.039347 ± 0.005422	-0.059339 ± 0.019691	0.693239 ± 0.125898	-0.461968 ± 0.118097	0.53	1.00
0.7	0.107314 ± 0.000132	1.216132 ± 0.000167	0.631616 ± 0.013134	-0.061303 ± 0.001651	0.473782 ± 0.018269	-0.640731 ± 0.015055	0.034896 ± 0.004447	-0.053675 ± 0.016539	0.606057 ± 0.098988	-0.472567 ± 0.109018	0.52	1.0
0.8	0.081944 ± 0.000084	1.162250 ± 0.000117	0.579014 ± 0.010826	-0.046800 ± 0.001033	0.485959 ± 0.015765	-0.466603 ± 0.010624	0.027406 ± 0.003041	-0.044194 ± 0.011692	0.453100 ± 0.064637	-0.474084 ± 0.096816	0.51	1.0
0.9	0.046222 ± 0.000038	1.089350 ± 0.000060	0.511341 ± 0.008625	-0.026322 ± 0.000459	0.501638 ± 0.013319	-0.246559 ± 0.005500	0.016146 ± 0.001506	-0.027994 ± 0.006074	0.244495 ± 0.030382	-0.465251 ± 0.084861	0.51	1.0
1.1	-0.056757 ± 0.000043	0.897441 ± 0.000065	0.336311 ± 0.005644	0.031047 ± 0.000326	0.492275 ± 0.012582	0.254076 ± 0.005806	-0.019632 ± 0.001454	0.033602 ± 0.005727	-0.212346 ± 0.025268	-0.286249 ± 0.056218	0.47	1.0
1.2	-0.124094 ± 0.000109	0.785431 ± 0.000143	0.232415 ± 0.004752	0.065074 ± 0.000531	0.476356 ± 0.014107	0.497634 ± 0.012403	-0.042224 ± 0.003024	0.070737 ± 0.011741	-0.346189 ± 0.044895	-0.188538 ± 0.044628	0.42	1.0
1.3	-0.202217 ± 0.000212	0.668175 ± 0.000262	0.119088 ± 0.004180	0.100056 ± 0.000651	0.471735 ± 0.016643	0.721573 ± 0.021100	-0.071246 ± 0.004919	0.127772 ± 0.020713	-0.381763 ± 0.061094	-0.104912 ± 0.038431	0.44	1.0
1.4	-0.291488 ± 0.000349	0.551240 ± 0.000460	0.005018 ± 0.003985	0.134044 ± 0.000743	0.486075 ± 0.019264	0.897656 ± 0.034172	-0.112490 ± 0.007913	0.235457 ± 0.038894	-0.355429 ± 0.077932	-0.067673 ± 0.038444	0.51	1.0
1.5	-0.391936 ± 0.000524	0.441601 ± 0.000744	-0.106193 ± 0.004192	0.163659 ± 0.000870	0.508576 ± 0.021719	1.000598 ± 0.052708	-0.166680 ± 0.013288	0.411058 ± 0.073914	-0.303176 ± 0.098547	-0.072968 ± 0.043762	0.60	1.0

Table S8: The details of the most reliable fitting by the scaling analysis for the Legendre function as a function of Q1D two-terminal conductances, $\langle p_q(g) \rangle$, with the minimum width $M_{\min} = 32$ and the maximum width $M_{\max} = 512$. The aspect ratio is varied from $L/M = 2$ to 10 and the number of total data points N is 1458 except at $q = 0.5$. At $q = 0.5$, the maximum aspect ratio is limited to 8 due to our numerical accuracy and the number of total data points N is 1164. The scaling function in Eq. (S2) with $N_p = 2$ and $N_g = 1$ is employed. χ_{\min}^2 , and Q in the tables represent the minimum value of χ^2 , and the goodness of fit, respectively. In the table, the value with \pm means the error bar for its above value. Note that the Legendre function satisfies $p_q(x) = p_{1-q}(x)$.

q	\bar{X}'_q	\bar{c}'_q	\bar{d}'_q	$ \bar{y}' $	\bar{c}''_q	\bar{a}''_q	\bar{a}'_q	\bar{d}''_q	\bar{b}'_q	χ_{\min}^2/N	Q
0.5	0.127880	1.261637	-0.070035	0.502680	-0.805567	0.044965	-0.080693	0.887891	-0.695250	0.54	1.0
	± 0.000105	± 0.000216	± 0.001699	± 0.014023	± 0.022389	± 0.005253	± 0.022452	± 0.099765	± 0.101576		
0.6	0.122863	1.249876	-0.069912	0.521865	-0.778881	0.045937	-0.091743	0.948794	-0.749678	0.53	1.0
	± 0.000081	± 0.000164	± 0.001651	± 0.011684	± 0.017488	± 0.004616	± 0.021417	± 0.095329	± 0.103671		
0.7	0.107501	1.216620	-0.061013	0.502006	-0.658657	0.037769	-0.067574	0.754075	-0.629617	0.50	1.0
	± 0.000074	± 0.000134	± 0.001272	± 0.011748	± 0.014212	± 0.003632	± 0.015436	± 0.077517	± 0.088716		
0.8	0.081881	1.162523	-0.046199	0.466725	-0.477109	0.025794	-0.037911	0.488745	-0.459206	0.48	1.0
	± 0.000063	± 0.000093	± 0.000794	± 0.012180	± 0.009688	± 0.002352	± 0.008630	± 0.052529	± 0.068309		
0.9	0.046009	1.089469	-0.025704	0.413905	-0.252312	0.012375	-0.013410	0.218066	-0.281125	0.48	1.0
	± 0.000045	± 0.000048	± 0.000348	± 0.013681	± 0.004784	± 0.001107	± 0.003307	± 0.025304	± 0.047794		
1.1	-0.055885	0.897339	0.030164	0.285426	0.261068	-0.010860	0.005095	-0.141016	-0.065599	0.54	1.0
	± 0.000153	± 0.000051	± 0.000241	± 0.024099	± 0.004703	± 0.001203	± 0.002347	± 0.020916	± 0.021922		
1.2	-0.121720	0.785270	0.063496	0.261409	0.511558	-0.024145	0.011502	-0.205564	-0.018865	0.51	1.0
	± 0.000439	± 0.000113	± 0.000394	± 0.026830	± 0.009878	± 0.002189	± 0.004404	± 0.037777	± 0.017265		
1.3	-0.198862	0.668117	0.098168	0.295404	0.736462	-0.047657	0.036608	-0.201395	0.012571	0.48	1.0
	± 0.000542	± 0.000206	± 0.000486	± 0.022178	± 0.016758	± 0.002509	± 0.006996	± 0.053106	± 0.016293		
1.4	-0.287574	0.551334	0.131921	0.340645	0.914536	-0.082930	0.093218	-0.152016	0.027427	0.50	1.0
	± 0.000618	± 0.000357	± 0.000557	± 0.020483	± 0.026899	± 0.003590	± 0.013745	± 0.069479	± 0.017388		
1.5	-0.387254	0.441683	0.161359	0.373929	1.023250	-0.127047	0.181595	-0.098250	0.017967	0.55	1.0
	± 0.000791	± 0.000575	± 0.000651	± 0.021038	± 0.040883	± 0.006114	± 0.027293	± 0.088935	± 0.019809		

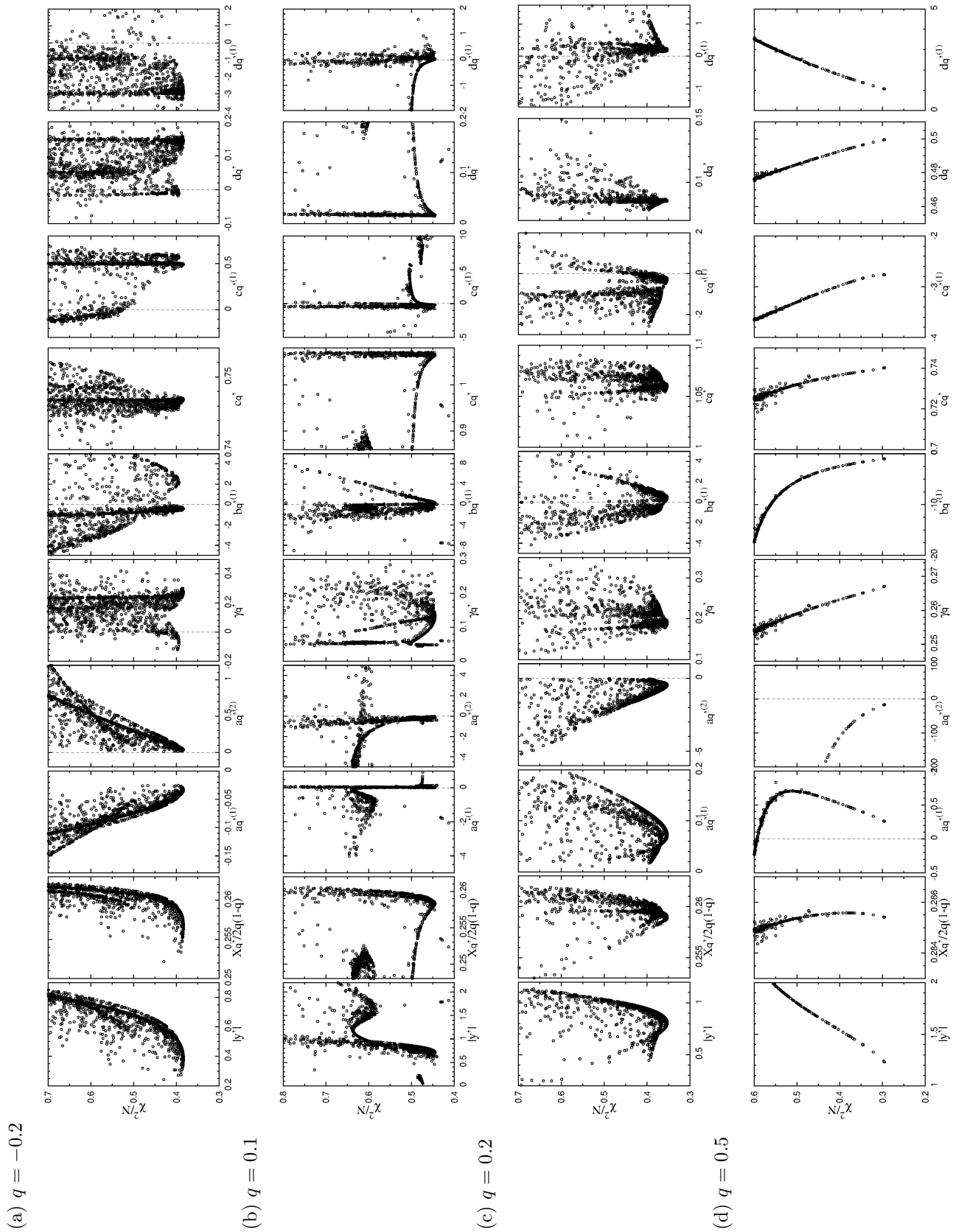


Fig. S1: Stability maps from the scaling analysis of $\langle g^q \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 192$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S1.

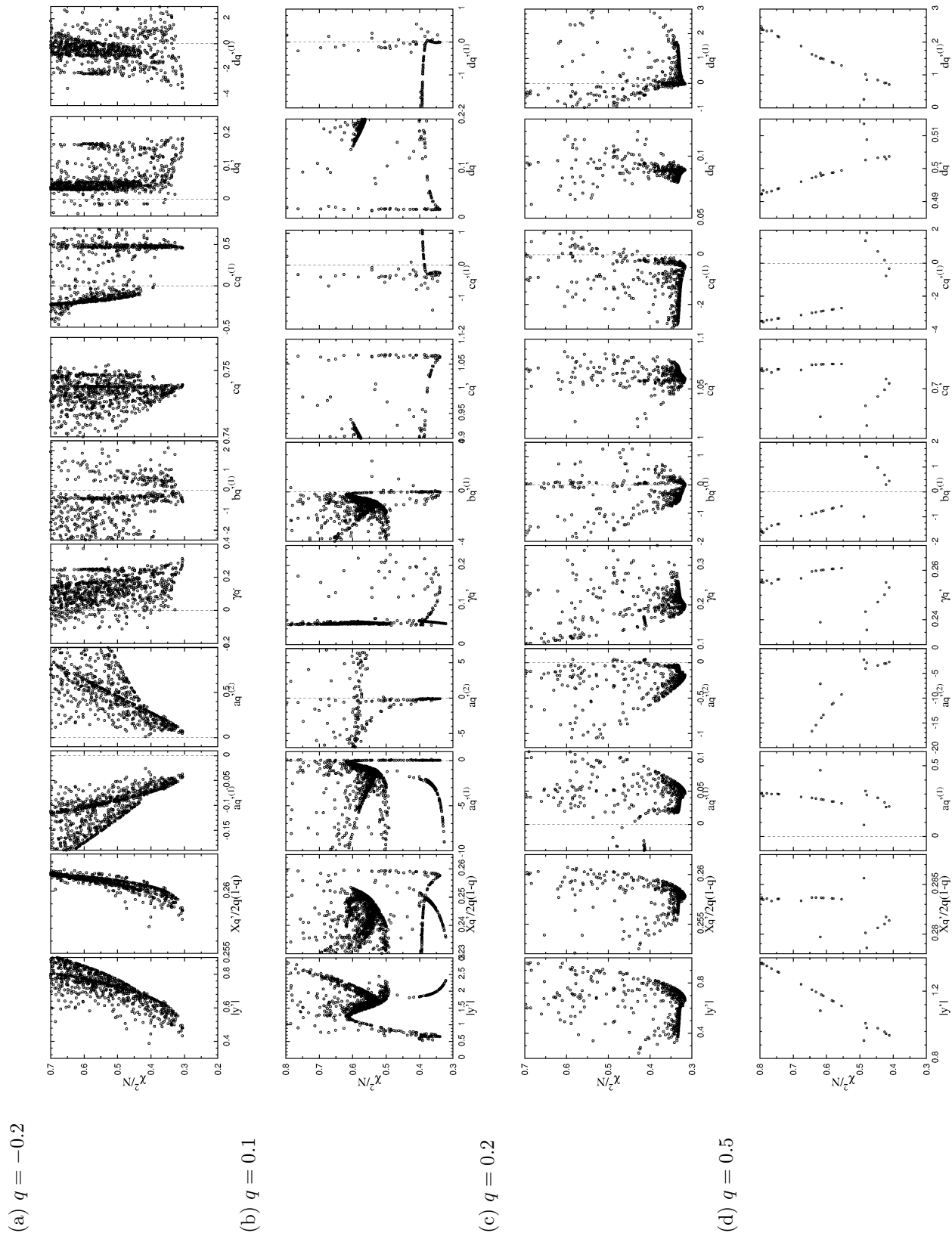


Fig. S2: Stability maps from the scaling analysis of $\langle g^q \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 256$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S2.

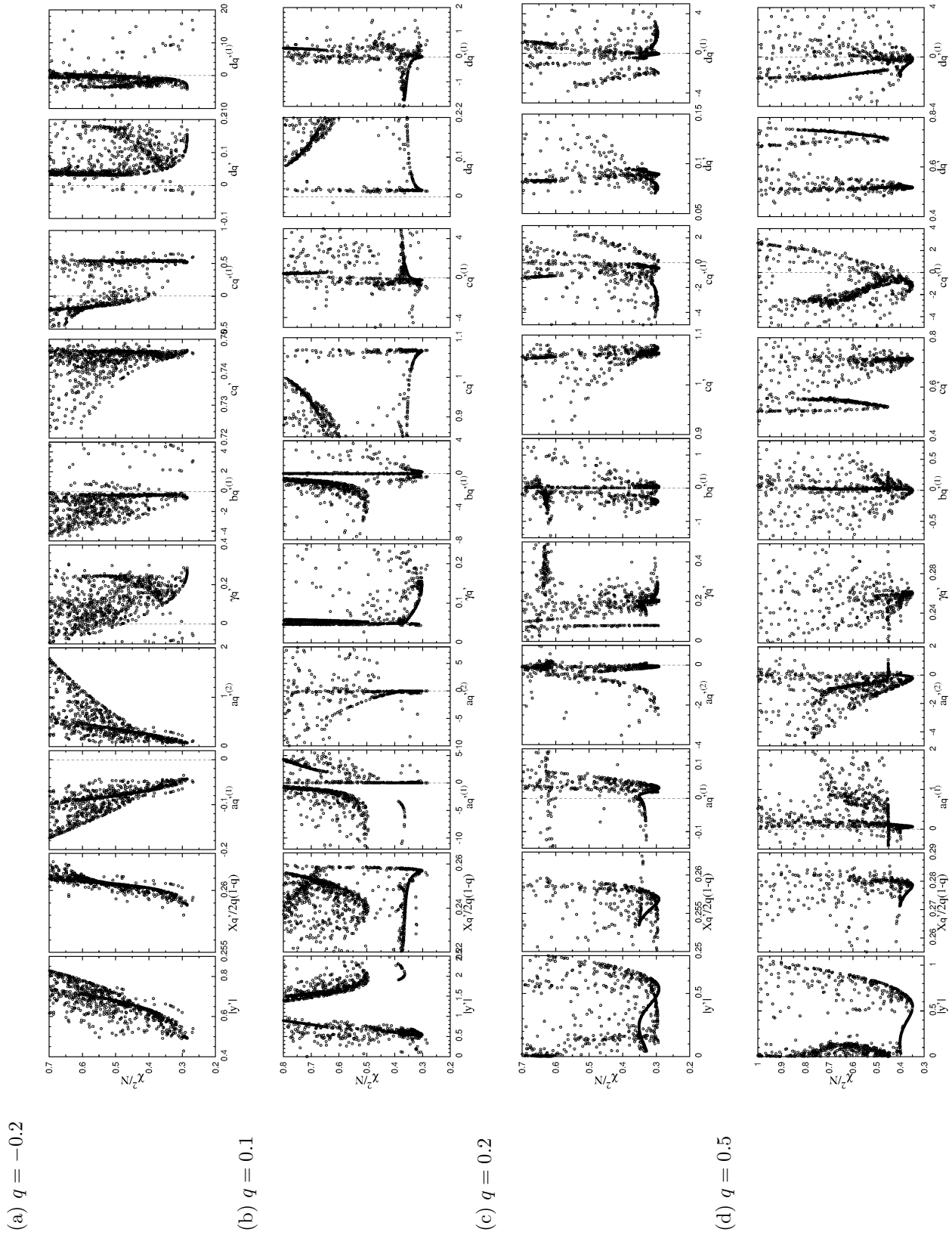


Fig. S3: Stability maps from the scaling analysis of (g^q) in case of $M_{\min} = 32$ and $M_{\max} = 384$ for (a) $q = -0.2$, (b) 0.1, (c) 0.2, and (d) 0.5. The values of fitting parameters at the global minimum are listed in Table. S3.

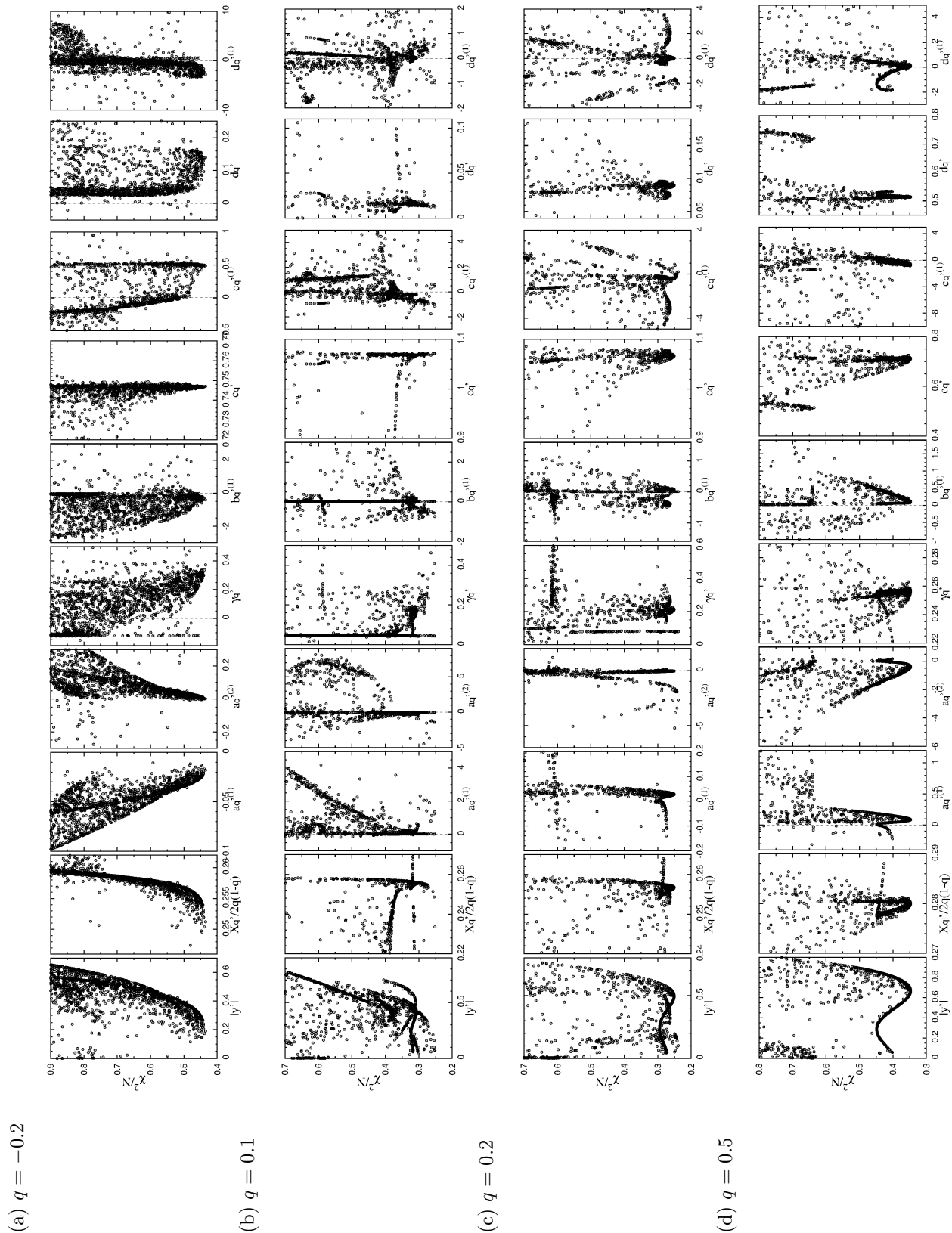
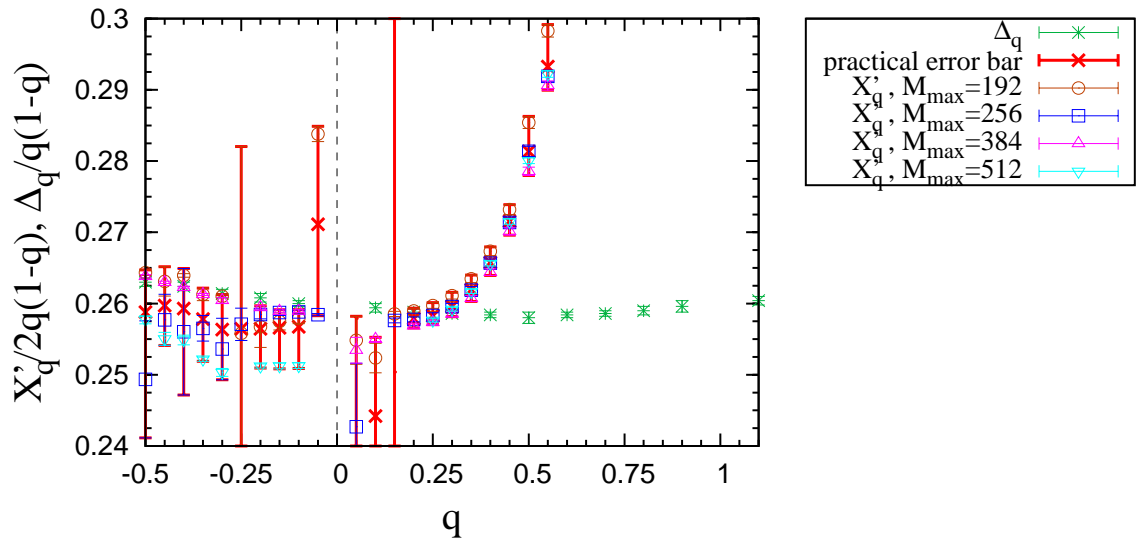
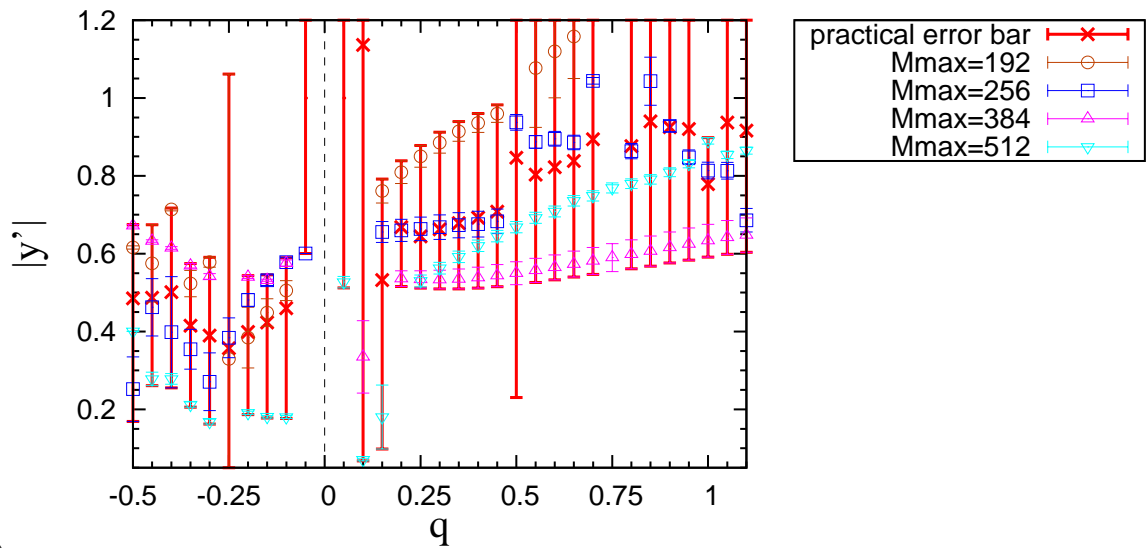


Fig. S4: Stability maps from the scaling analysis of $\langle g^q \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 512$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S4.

(a)



(b)



(c)

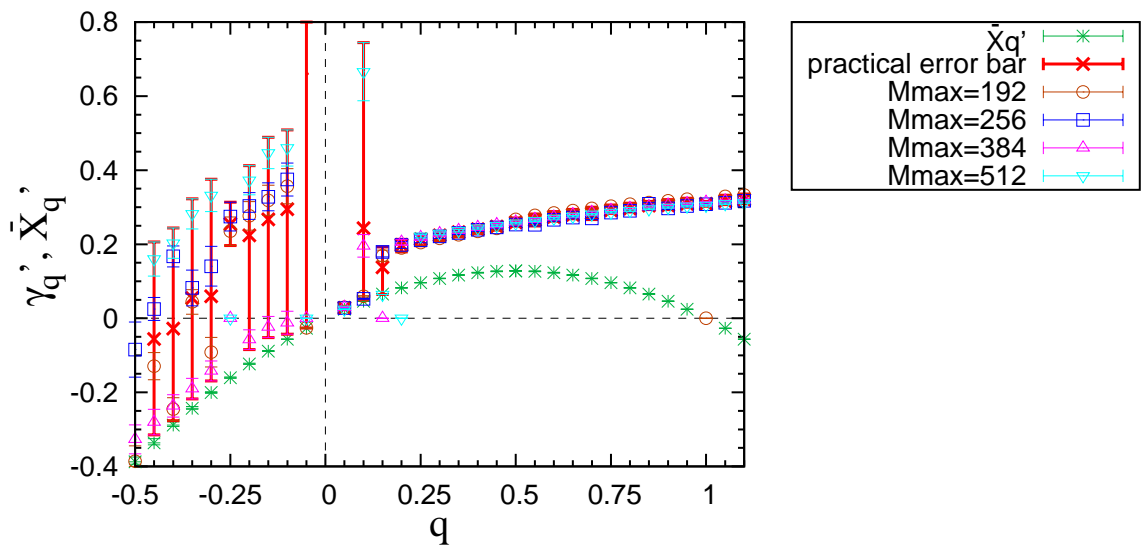


Fig. S5: q dependence of (a) X'_q , (b) y' , and (c) γ'_q from $\langle g^q \rangle$. The error bars by thin and thick lines represent the error bars estimated from the error-propagation theory and the practical error bars, respectively.

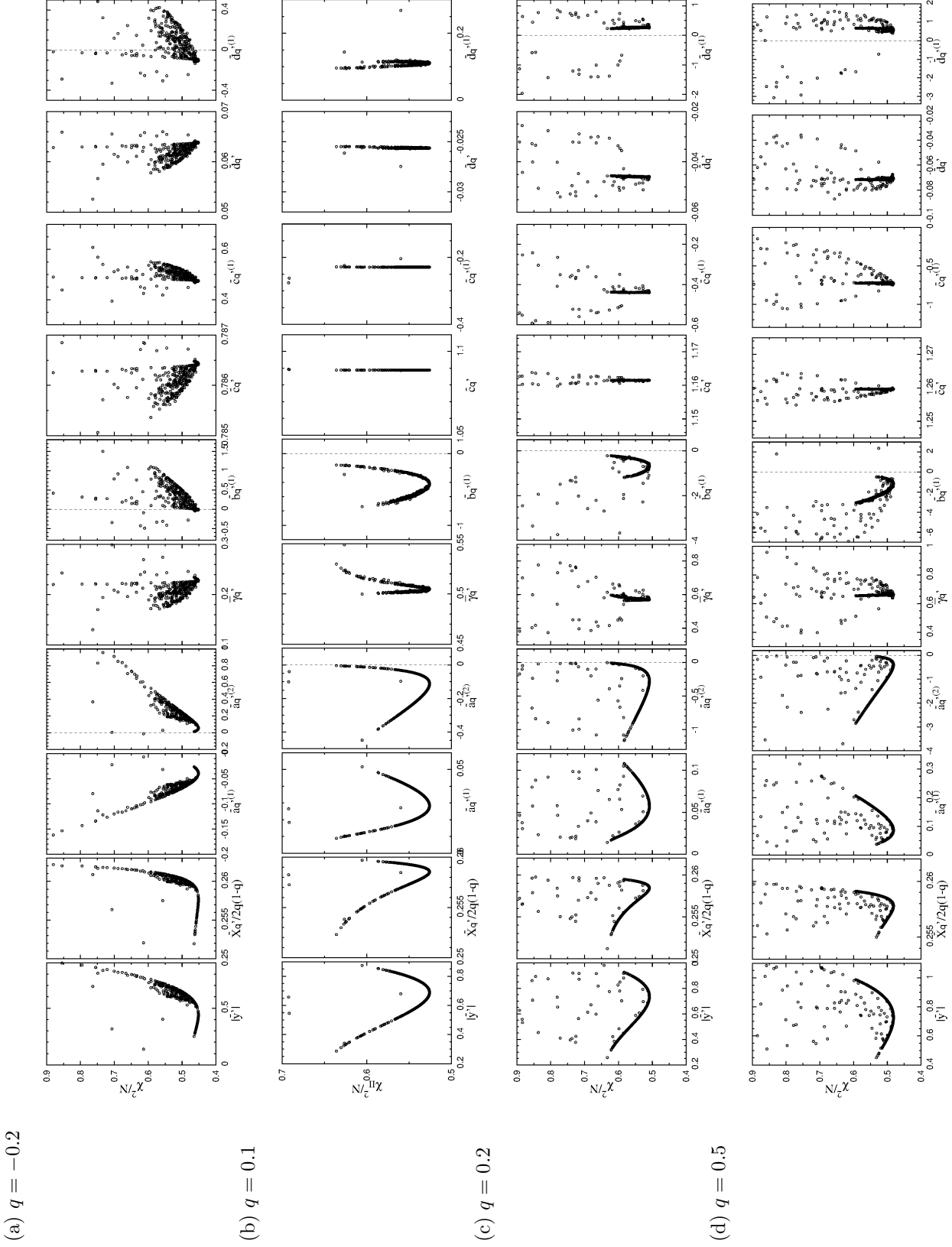


Fig. S6: Stability maps from the scaling analysis of $\langle \rho_c(g) \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 192$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S5.

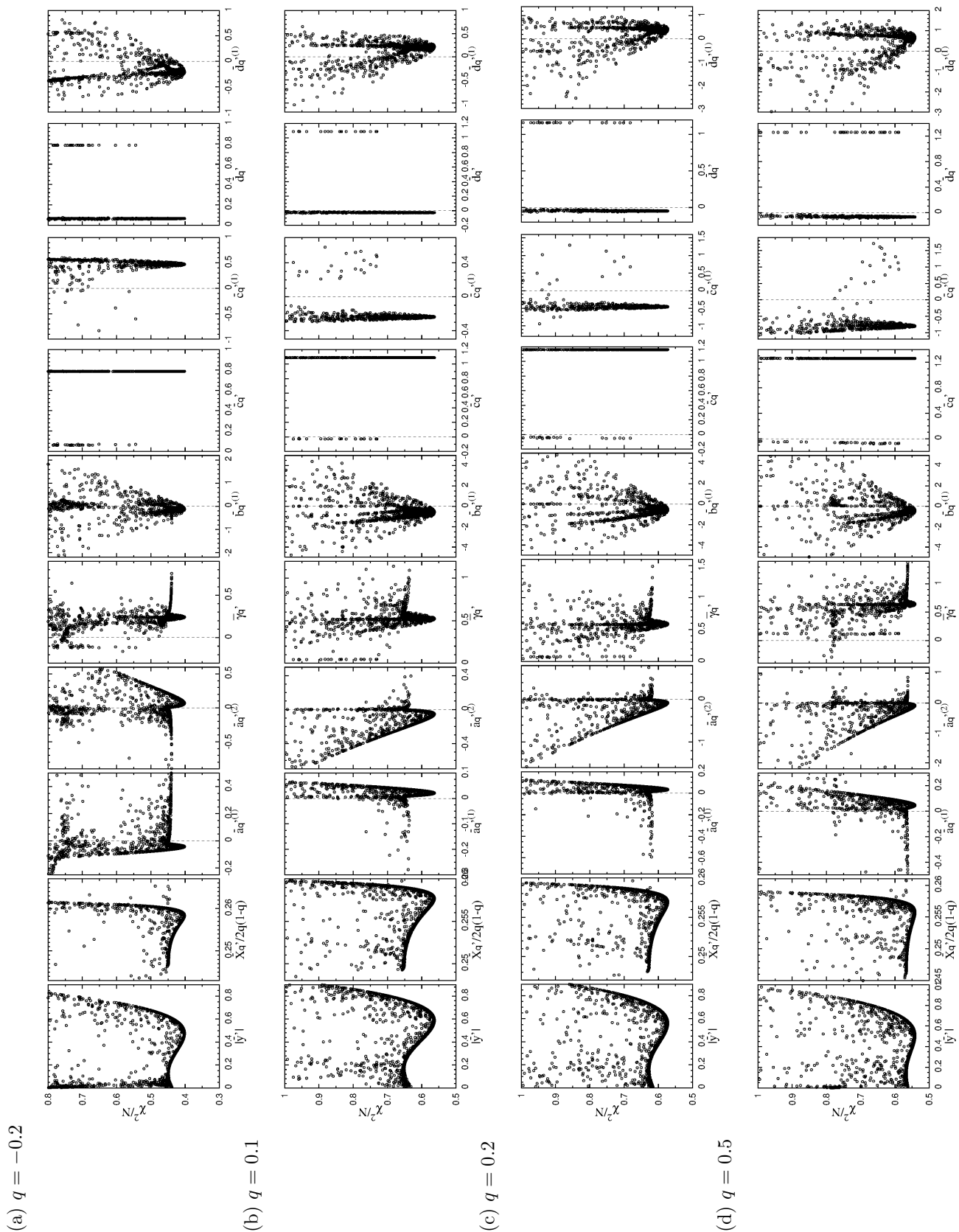


Fig. S7: Stability maps from the scaling analysis of $\langle p_q(g) \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 256$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S6.

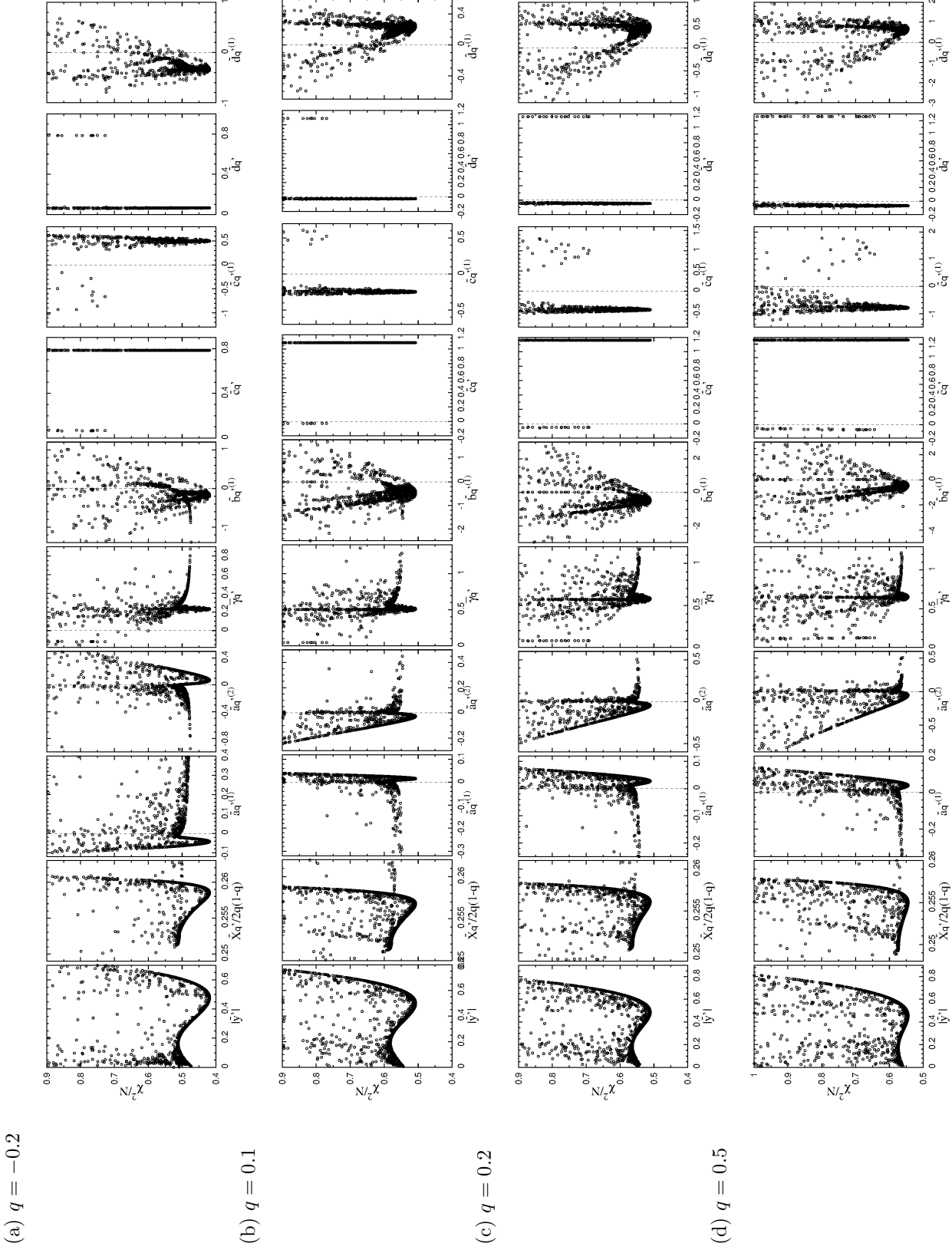


Fig. S8: Stability maps from the scaling analysis of $\langle \rho_x(g) \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 384$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S7.

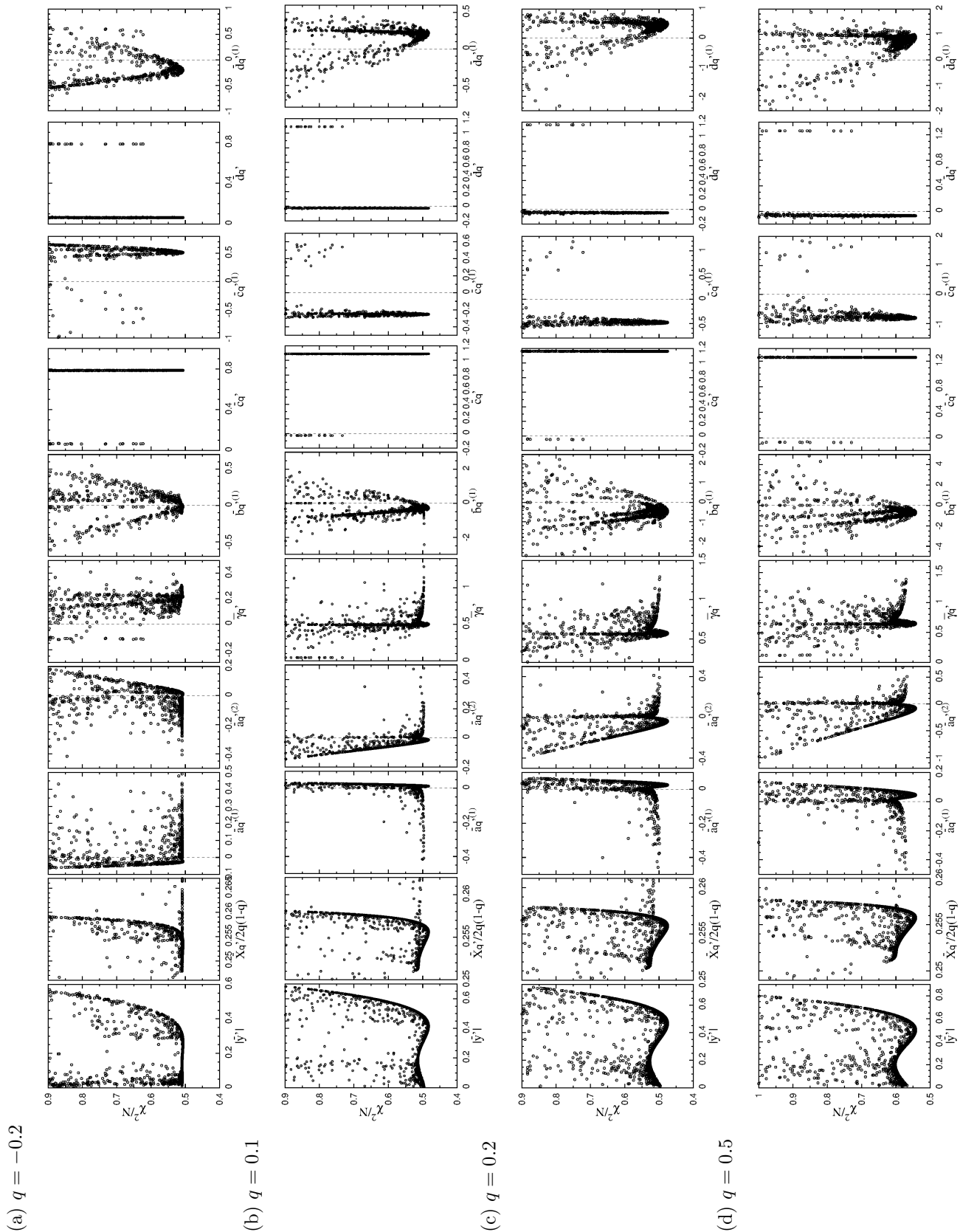
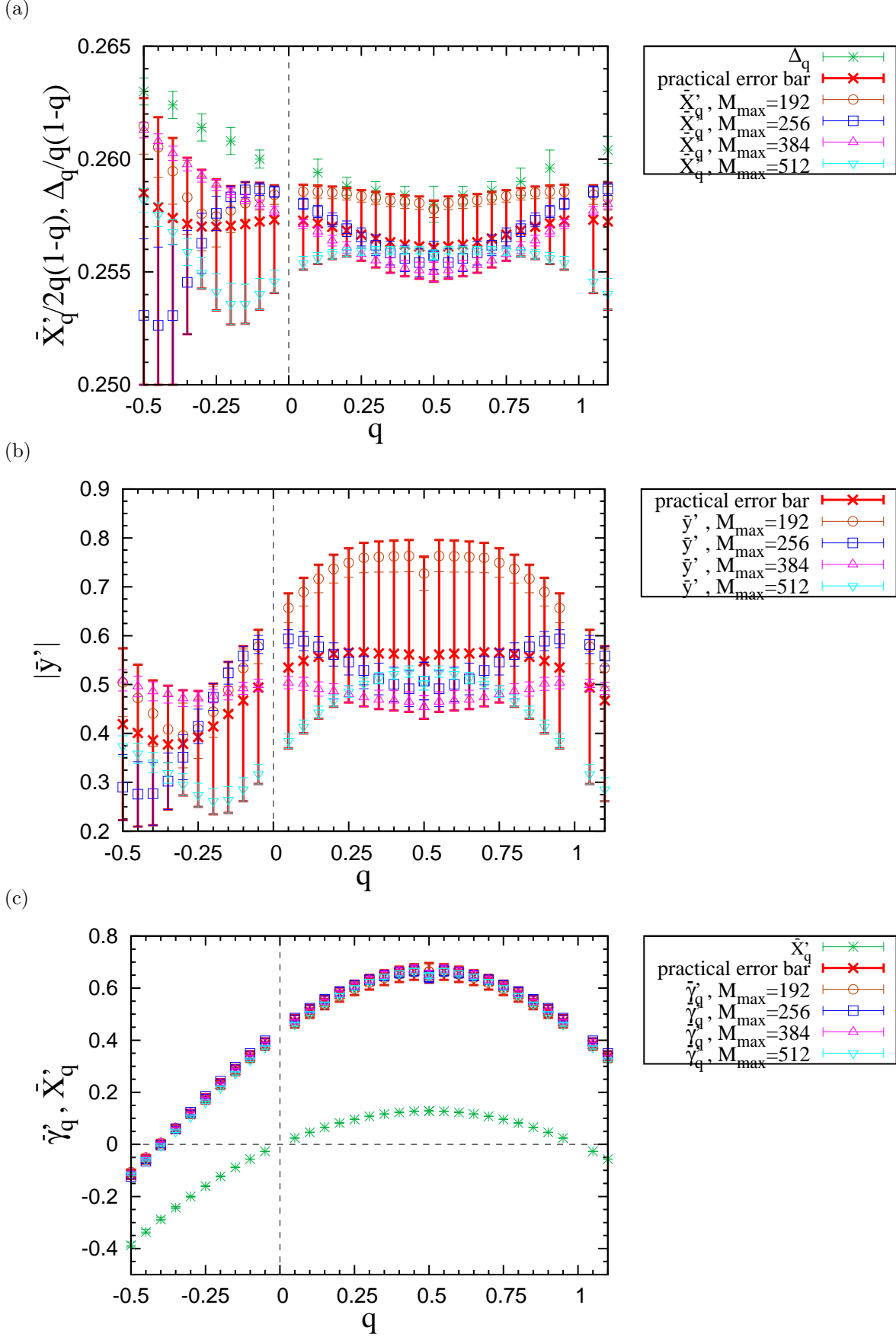


Fig. S9: Stability maps from the scaling analysis of $\langle p_q(g) \rangle$ in case of $M_{\min} = 32$ and $M_{\max} = 512$ for (a) $q = -0.2$, (b) 0.1 , (c) 0.2 , and (d) 0.5 . The values of fitting parameters at the global minimum are listed in Table. S8.



Point-contact conductance in 2D. –

Boundary condition effects. At an Anderson transition in two dimensions (2D), the q -th moment of the two point-contact conductance (PCC) between two points separated by distance r in the plane is expected to behave as $\langle T(r)^q \rangle \propto r^{-X_q}$, plus power-law corrections due to irrelevant exponents and subleading scaling dimensions. However, since we calculate the PCC in a system with periodic boundary conditions (the geometry of a torus), and the expected simple power law cannot appear for large r . To avoid this effect caused by the periodic boundary conditions, we have focused on relatively short distances. To determine the appropriate range of r that we use in our scaling analysis, we compared the PCC for two different system sizes.

Figure S11 shows the r dependence of the moments $\langle T^q \rangle$ in 2D with the different system size; $L = 480$ (filled dots) and 1024 (open dots). Since $\langle T^q \rangle$ for both systems [practically coincide for $r < 57$], we can safely use the numerical data for $r < 57$ in the fitting.

Scaling analysis for PCC in 2D. Since the T^q is not a pure scaling operator, as we explained in the main paper, the appropriate scaling function in 2D is a combination of the leading power laws with exponent X_q , and correction terms with the irrelevant exponent y and the subleading dimension γ_q :

$$\begin{aligned} \langle T^q \rangle = & c_q r^{-X_q} \left(1 + \sum_{n=1}^{N_p} a_q^{(n)} r^{ny} \right) \\ & + d_q r^{-\gamma_q} \left(1 + \sum_{n=1}^{N_s} b_q^{(n)} r^{ny} \right). \end{aligned} \quad (\text{S4})$$

The use of the Legendre functions allows to omit the corrections with subleading scaling dimensions for the quantity $p_q(T) \equiv P_{-q}(2/T - 1)$:

$$\langle p_q(T) \rangle = \bar{c}_q r^{-\bar{X}_q} \left(1 + \sum_{n=1}^{N_p} \bar{a}_q^{(n)} r^{-ny} \right). \quad (\text{S5})$$

We fit the numerical data in the range $r < 57$ (chosen as explained in the previous section) to these power laws to extract exponents and coefficients. Because of the difficulty of fitting the data to multiple power laws with limited data points, we can introduce only a single correction term to obtain the reliable fitting. Therefore, we truncate Eq. (S5) at $N_p = 1$, and Eq. (S4) at $N_p = 1, d_q = 0$. This simplification makes the simultaneous determination of y and γ_q in Eq. (S4) impossible.

The results of the fitting are given in Tables S9 and S10. We also show the corresponding stability maps [S6] in Figs. S12 and S13.

Finally, we estimate the practical error bars for X_q and \bar{X}_q in the same way as in the Q1D case. These are shown in Fig. S14.

Table S9: The details of the most reliable fitting by the scaling analysis for q th moment of 2D PCC, $\langle T^q \rangle$, with the minimum r , $r_{\min} = 3$ and the maximum r , $r_{\max} = 41$, (b) $r_{\max} = 49$, and (c) $r_{\max} = 57$. The scaling function in Eq. (S4) with $N_p = 1$ and $d_q = 0$ is employed. N , χ^2_{\min} , and Q in the tables represent the number of data point used in the fitting, the minimum χ^2 , and the goodness of fit, respectively.

r_{\min}	r_{\max}	N	q	χ^2_{\min}/N	Q	$X_q/2$	$X_q/2q(1-q)$	c_q	$ y $	$a_q^{(1)}$
3	41	20	-0.5	0.491	0.876	-0.39062 ± 0.00579	0.26041 ± 0.00386	1.4544 ± 0.0280	2.6689 ± 1.5472	-1.0132 ± 1.5362
3	41	20	-0.4	0.380	0.960	-0.28871 ± 0.00410	0.25777 ± 0.00366	1.3234 ± 0.0187	2.1724 ± 0.9425	-0.4464 ± 0.3766
3	41	20	-0.3	0.324	0.982	-0.19943 ± 0.00292	0.25568 ± 0.00374	1.2240 ± 0.0128	1.8622 ± 0.6894	-0.2359 ± 0.1312
3	41	20	-0.2	0.307	0.986	-0.12216 ± 0.00201	0.25450 ± 0.00419	1.1417 ± 0.0086	1.6139 ± 0.5522	-0.1241 ± 0.0481
3	41	20	-0.1	0.295	0.989	-0.05601 ± 0.00111	0.25458 ± 0.00505	1.0686 ± 0.0046	1.3995 ± 0.4645	-0.0528 ± 0.0140
3	41	20	0.1	0.258	0.995	0.04683 ± 0.00151	0.26017 ± 0.00837	0.9334 ± 0.0061	1.0612 ± 0.3575	0.0457 ± 0.0041
3	41	20	0.2	0.243	0.996	0.08543 ± 0.00356	0.26697 ± 0.01113	0.8675 ± 0.0142	0.9359 ± 0.3223	0.0913 ± 0.0066
3	41	20	0.3	0.234	0.997	0.11675 ± 0.00627	0.27797 ± 0.01493	0.8020 ± 0.0243	0.8364 ± 0.2939	0.1408 ± 0.0179
3	41	20	0.4	0.230	0.997	0.14171 ± 0.00966	0.29523 ± 0.02013	0.7372 ± 0.0361	0.7589 ± 0.2699	0.1971 ± 0.0383
3	41	20	0.5	0.232	0.997	0.16124 ± 0.01370	0.32248 ± 0.02740	0.6740 ± 0.0489	0.6996 ± 0.2489	0.2615 ± 0.0680
3	41	20	0.6	0.238	0.997	0.17622 ± 0.01827	0.36713 ± 0.03806	0.6136 ± 0.0615	0.6549 ± 0.2301	0.3344 ± 0.1073
3	41	20	0.7	0.248	0.996	0.18747 ± 0.02321	0.44636 ± 0.05527	0.5568 ± 0.0731	0.6217 ± 0.2129	0.4157 ± 0.1562
3	41	20	0.8	0.261	0.995	0.19572 ± 0.02838	0.61161 ± 0.08868	0.5045 ± 0.0828	0.5973 ± 0.1973	0.5045 ± 0.2142
3	41	20	0.9	0.275	0.993	0.20157 ± 0.03363	1.11984 ± 0.18681	0.4569 ± 0.0905	0.5797 ± 0.1829	0.5997 ± 0.2805
3	41	20	1.1	0.308	0.986	0.20806 ± 0.04399	-0.94573 ± 0.19995	0.3760 ± 0.0997	0.5582 ± 0.1580	0.8047 ± 0.4348
r_{\min}	r_{\max}	N	q	χ^2_{\min}/N	Q	$X_q/2$	$X_q/2q(1-q)$	c_q	$ y $	$a_q^{(1)}$
3	49	24	-0.5	0.427	0.964	-0.38930 ± 0.00531	0.25954 ± 0.00354	1.4605 ± 0.0270	2.4458 ± 1.2788	-0.8301 ± 1.0249
3	49	24	-0.4	0.335	0.992	-0.28727 ± 0.00379	0.25649 ± 0.00338	1.3297 ± 0.0184	1.9559 ± 0.7597	-0.3775 ± 0.2484
3	49	24	-0.3	0.290	0.997	-0.19849 ± 0.00261	0.25448 ± 0.00335	1.2279 ± 0.0121	1.7097 ± 0.5636	-0.2132 ± 0.0939
3	49	24	-0.2	0.272	0.998	-0.12177 ± 0.00171	0.25368 ± 0.00356	1.1433 ± 0.0076	1.5349 ± 0.4624	-0.1188 ± 0.0383
3	49	24	-0.1	0.256	0.999	-0.05594 ± 0.00089	0.25428 ± 0.00405	1.0688 ± 0.0038	1.3792 ± 0.3958	-0.0524 ± 0.0124
3	49	24	0.1	0.229	0.999	0.04708 ± 0.00108	0.26154 ± 0.00597	0.9344 ± 0.0044	1.1158 ± 0.3096	0.0460 ± 0.0050
3	49	24	0.2	0.223	1.000	0.08632 ± 0.00240	0.26976 ± 0.00749	0.8709 ± 0.0094	1.0139 ± 0.2804	0.0903 ± 0.0062
3	49	24	0.3	0.223	1.000	0.11876 ± 0.00399	0.28277 ± 0.00950	0.8094 ± 0.0152	0.9315 ± 0.2572	0.1359 ± 0.0082
3	49	24	0.4	0.230	0.999	0.14538 ± 0.00583	0.30288 ± 0.01215	0.7502 ± 0.0212	0.8664 ± 0.2383	0.1842 ± 0.0152
3	49	24	0.5	0.241	0.999	0.16711 ± 0.00788	0.33422 ± 0.01576	0.6940 ± 0.0273	0.8155 ± 0.2227	0.2356 ± 0.0266
3	49	24	0.6	0.256	0.999	0.18478 ± 0.01008	0.38496 ± 0.02101	0.6411 ± 0.0331	0.762 ± 0.2095	0.2900 ± 0.0413
3	49	24	0.7	0.274	0.998	0.19912 ± 0.01238	0.47410 ± 0.02947	0.5920 ± 0.0383	0.7458 ± 0.1983	0.3469 ± 0.0589
3	49	24	0.8	0.294	0.997	0.21075 ± 0.01472	0.65859 ± 0.04600	0.5469 ± 0.0427	0.7224 ± 0.1887	0.4056 ± 0.0788
3	49	24	0.9	0.314	0.994	0.22017 ± 0.01707	1.22319 ± 0.09481	0.5058 ± 0.0464	0.7042 ± 0.1804	0.4655 ± 0.1008
3	49	24	1.1	0.357	0.987	0.23400 ± 0.02170	-1.06363 ± 0.09864	0.4346 ± 0.0518	0.6788 ± 0.1669	0.5864 ± 0.1496
r_{\min}	r_{\max}	N	q	χ^2_{\min}/N	Q	$X_q/2$	$X_q/2q(1-q)$	c_q	$ y $	$a_q^{(1)}$
3	57	28	-0.5	0.428	0.980	-0.39155 ± 0.00387	0.26103 ± 0.00258	1.4501 ± 0.0195	2.8735 ± 1.5219	-1.2255 ± 1.8972
3	57	28	-0.4	0.307	0.998	-0.28842 ± 0.00284	0.25752 ± 0.00254	1.3246 ± 0.0138	2.1404 ± 0.7722	-0.4364 ± 0.3115
3	57	28	-0.3	0.250	1.000	-0.19881 ± 0.00208	0.25488 ± 0.00267	1.2265 ± 0.0098	1.7633 ± 0.5260	-0.2210 ± 0.0959
3	57	28	-0.2	0.240	1.000	-0.12164 ± 0.00145	0.25341 ± 0.00302	1.1439 ± 0.0066	1.5080 ± 0.4096	-0.1170 ± 0.0342
3	57	28	-0.1	0.239	1.000	-0.05573 ± 0.00080	0.25331 ± 0.00364	1.0697 ± 0.0036	1.3068 ± 0.3400	-0.0508 ± 0.0100
3	57	28	0.1	0.236	1.000	0.04658 ± 0.00105	0.25880 ± 0.00584	0.9324 ± 0.0045	1.0083 ± 0.2580	0.0456 ± 0.0033
3	57	28	0.2	0.237	1.000	0.08506 ± 0.00242	0.26580 ± 0.00757	0.8660 ± 0.0102	0.9022 ± 0.2316	0.0917 ± 0.0051
3	57	28	0.3	0.241	1.000	0.11646 ± 0.00413	0.27728 ± 0.00984	0.8007 ± 0.0168	0.8199 ± 0.2107	0.1414 ± 0.0116
3	57	28	0.4	0.249	1.000	0.14182 ± 0.00616	0.29546 ± 0.01282	0.7374 ± 0.0239	0.7573 ± 0.1936	0.1963 ± 0.0232
3	57	28	0.5	0.260	1.000	0.16214 ± 0.00842	0.32428 ± 0.01684	0.6769 ± 0.0310	0.7106 ± 0.1793	0.2567 ± 0.0391
3	57	28	0.6	0.273	0.999	0.17833 ± 0.01083	0.37152 ± 0.02257	0.6202 ± 0.0375	0.6763 ± 0.1672	0.3221 ± 0.0589
3	57	28	0.7	0.289	0.999	0.19120 ± 0.01331	0.45524 ± 0.03168	0.5679 ± 0.0430	0.6515 ± 0.1567	0.3914 ± 0.0818
3	57	28	0.8	0.305	0.998	0.20144 ± 0.01577	0.62949 ± 0.04929	0.5204 ± 0.0474	0.6338 ± 0.1477	0.4634 ± 0.1072
3	57	28	0.9	0.322	0.998	0.20959 ± 0.01818	1.16441 ± 0.10100	0.4776 ± 0.0508	0.6213 ± 0.1400	0.5369 ± 0.1345
3	57	28	1.1	0.356	0.995	0.22136 ± 0.02271	-1.00620 ± 0.10324	0.4054 ± 0.0546	0.6067 ± 0.1276	0.6841 ± 0.1925

Table S10: The details of the most reliable fitting by the scaling analysis for the Legendre function as a function of 2D PCC, $\langle p_q(T) \rangle$, with the minimum r , $r_{\min} = 3$ and the maximum r , $r_{\max} = 41$, (b) $r_{\max} = 49$, and (c) $r_{\max} = 57$. The scaling function in Eq. (S5) with $N_p = 1$ and $N_s = 0$ is employed. N , χ_{\min}^2 , and Q in the tables represent the number of data point used in the fitting, the minimum χ^2 , and the goodness of fit, respectively. Note that the Legendre function satisfies $p_q(x) = p_{1-q}(x)$.

r_{\min}	r_{\max}	N	q	χ_{\min}^2/N	Q	$\bar{X}_q/2$	$\bar{X}_q/2q(1-q)$	\bar{c}_q	$ \bar{y} $	$\bar{a}_q^{(1)}$
3	41	20	0.5	0.244	0.996	0.06283 ± 0.00231	0.25131 ± 0.00924	0.8454 ± 0.0174	0.9932 ± 0.4077	0.1050 ± 0.0086
3	41	20	0.6	0.245	0.996	0.06037 ± 0.00218	0.25154 ± 0.00907	0.8512 ± 0.0164	1.0024 ± 0.4070	0.1012 ± 0.0083
3	41	20	0.7	0.248	0.996	0.05296 ± 0.00181	0.25221 ± 0.00860	0.8686 ± 0.0138	1.0298 ± 0.4051	0.0897 ± 0.0078
3	41	20	0.8	0.254	0.995	0.04052 ± 0.00127	0.25327 ± 0.00792	0.8984 ± 0.0098	1.0749 ± 0.4028	0.0699 ± 0.0072
3	41	20	0.9	0.264	0.994	0.02292 ± 0.00064	0.25466 ± 0.00715	0.9415 ± 0.0051	1.1370 ± 0.4014	0.0408 ± 0.0053
3	41	20	1.1	0.289	0.990	-0.02839 ± 0.00063	0.25812 ± 0.00576	1.0775 ± 0.0055	1.3091 ± 0.4111	-0.0566 ± 0.0117
3	41	20	1.2	0.299	0.988	-0.06242 ± 0.00127	0.26009 ± 0.00529	1.1797 ± 0.0116	1.4193 ± 0.4325	-0.1360 ± 0.0351
3	41	20	1.3	0.314	0.985	-0.10225 ± 0.00198	0.26217 ± 0.00507	1.3160 ± 0.0197	1.5523 ± 0.4816	-0.2506 ± 0.0827
3	41	20	1.4	0.373	0.963	-0.14803 ± 0.00289	0.26434 ± 0.00516	1.5023 ± 0.0318	1.7327 ± 0.5962	-0.4309 ± 0.1998
3	41	20	1.5	0.489	0.878	-0.19993 ± 0.00415	0.26657 ± 0.00554	1.7653 ± 0.0517	2.0378 ± 0.8839	-0.7879 ± 0.6150
(a)										
r_{\min}	r_{\max}	N	q	χ_{\min}^2/N	Q	$\bar{X}_q/2$	$\bar{X}_q/2q(1-q)$	\bar{c}_q	$ \bar{y} $	$\bar{a}_q^{(1)}$
3	49	24	0.5	0.230	0.999	0.06354 ± 0.00146	0.25415 ± 0.00585	0.8504 ± 0.0108	1.1179 ± 0.3654	0.1050 ± 0.0128
3	49	24	0.6	0.229	0.999	0.06102 ± 0.00139	0.25425 ± 0.00579	0.8558 ± 0.0103	1.1229 ± 0.3641	0.1013 ± 0.0126
3	49	24	0.7	0.229	0.999	0.05346 ± 0.00118	0.25457 ± 0.00564	0.8722 ± 0.0089	1.1377 ± 0.3605	0.0901 ± 0.0117
3	49	24	0.8	0.230	0.999	0.04081 ± 0.00086	0.25509 ± 0.00540	0.9005 ± 0.0067	1.1623 ± 0.3555	0.0706 ± 0.0098
3	49	24	0.9	0.233	0.999	0.02302 ± 0.00046	0.25580 ± 0.00512	0.9422 ± 0.0037	1.1960 ± 0.3503	0.0413 ± 0.0062
3	49	24	1.1	0.252	0.999	-0.02835 ± 0.00051	0.25773 ± 0.00462	1.0778 ± 0.0045	1.2881 ± 0.3484	-0.0563 ± 0.0104
3	49	24	1.2	0.265	0.998	-0.06214 ± 0.00108	0.25891 ± 0.00451	1.1822 ± 0.0104	1.3463 ± 0.3598	-0.1318 ± 0.0278
3	49	24	1.3	0.284	0.997	-0.10151 ± 0.00178	0.26029 ± 0.00457	1.3231 ± 0.0188	1.4225 ± 0.3938	-0.2342 ± 0.0599
3	49	24	1.4	0.333	0.992	-0.14682 ± 0.00271	0.26218 ± 0.00484	1.5152 ± 0.0319	1.5573 ± 0.4849	-0.3856 ± 0.1372
3	49	24	1.5	0.428	0.963	-0.19869 ± 0.00390	0.26492 ± 0.00520	1.7802 ± 0.0514	1.8545 ± 0.7350	-0.6836 ± 0.4295
(b)										
r_{\min}	r_{\max}	N	q	χ_{\min}^2/N	Q	$\bar{X}_q/2$	$\bar{X}_q/2q(1-q)$	\bar{c}_q	$ \bar{y} $	$\bar{a}_q^{(1)}$
3	57	28	0.5	0.244	1.000	0.06275 ± 0.00153	0.25099 ± 0.00611	0.8447 ± 0.0121	0.9748 ± 0.2986	0.1048 ± 0.0076
3	57	28	0.6	0.243	1.000	0.06028 ± 0.00145	0.25116 ± 0.00603	0.8504 ± 0.0115	0.9816 ± 0.2979	0.1010 ± 0.0074
3	57	28	0.7	0.241	1.000	0.05285 ± 0.00121	0.25167 ± 0.00578	0.8677 ± 0.0097	1.0018 ± 0.2961	0.0894 ± 0.0070
3	57	28	0.8	0.239	1.000	0.04040 ± 0.00087	0.25248 ± 0.00542	0.8974 ± 0.0071	1.0353 ± 0.2938	0.0696 ± 0.0061
3	57	28	0.9	0.238	1.000	0.02282 ± 0.00045	0.25357 ± 0.00499	0.9407 ± 0.0038	1.0820 ± 0.2921	0.0404 ± 0.0042
3	57	28	1.1	0.237	1.000	-0.02822 ± 0.00046	0.25651 ± 0.00417	1.0790 ± 0.0043	1.2171 ± 0.2983	-0.0549 ± 0.0083
3	57	28	1.2	0.235	1.000	-0.06201 ± 0.00093	0.25838 ± 0.00386	1.1833 ± 0.0092	1.3128 ± 0.3149	-0.1300 ± 0.0244
3	57	28	1.3	0.245	1.000	-0.10164 ± 0.00144	0.26062 ± 0.00368	1.3218 ± 0.0154	1.4456 ± 0.3563	-0.2370 ± 0.0585
3	57	28	1.4	0.301	0.999	-0.14753 ± 0.00205	0.26345 ± 0.00365	1.5074 ± 0.0240	1.6639 ± 0.4617	-0.4128 ± 0.1530
3	57	28	1.5	0.423	0.982	-0.20023 ± 0.00278	0.26697 ± 0.00371	1.7615 ± 0.0362	2.1017 ± 0.7676	-0.8317 ± 0.5931
(c)										

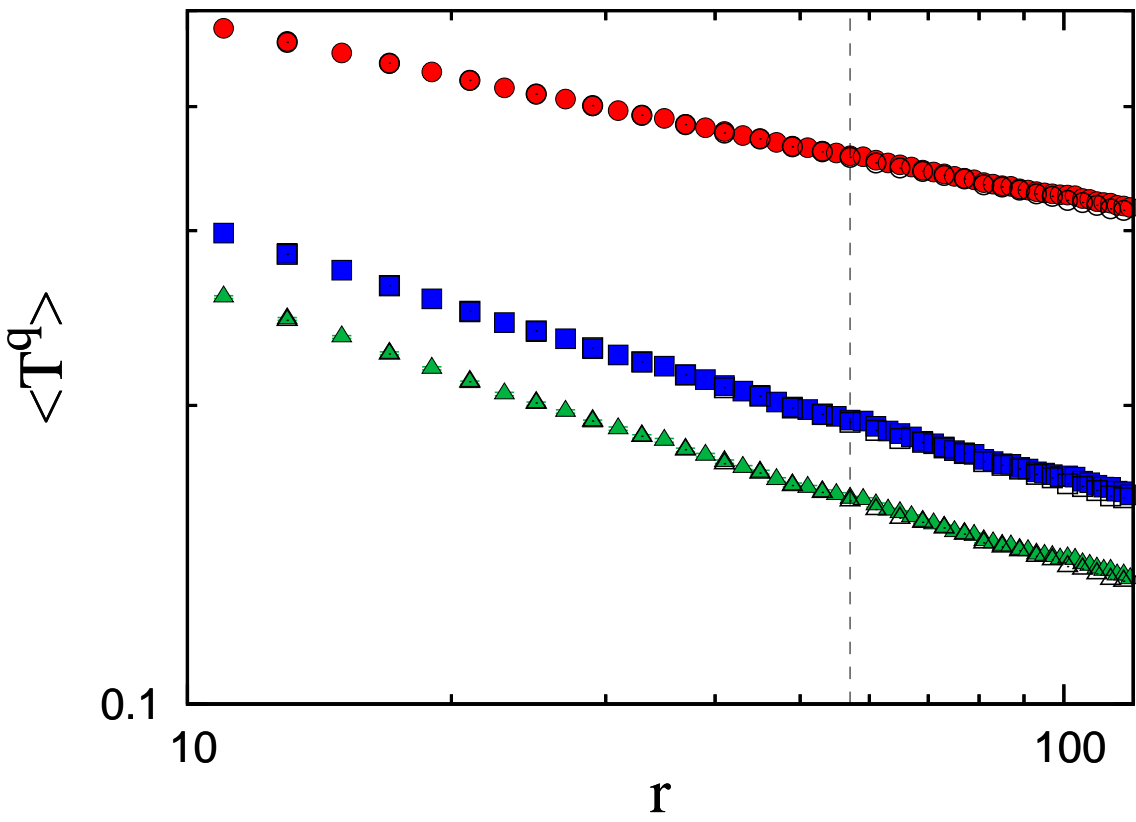


Fig. S11: q th moment of PCC in 2D, $\langle T^q \rangle$, for the system $L = 480$ (filled dots) and $L = 1024$ (open dots) for $q = 0.5$ (circles), 1 (rectangles), and 1.2 (triangles). The vertical dashed lines represent $r = 57$.

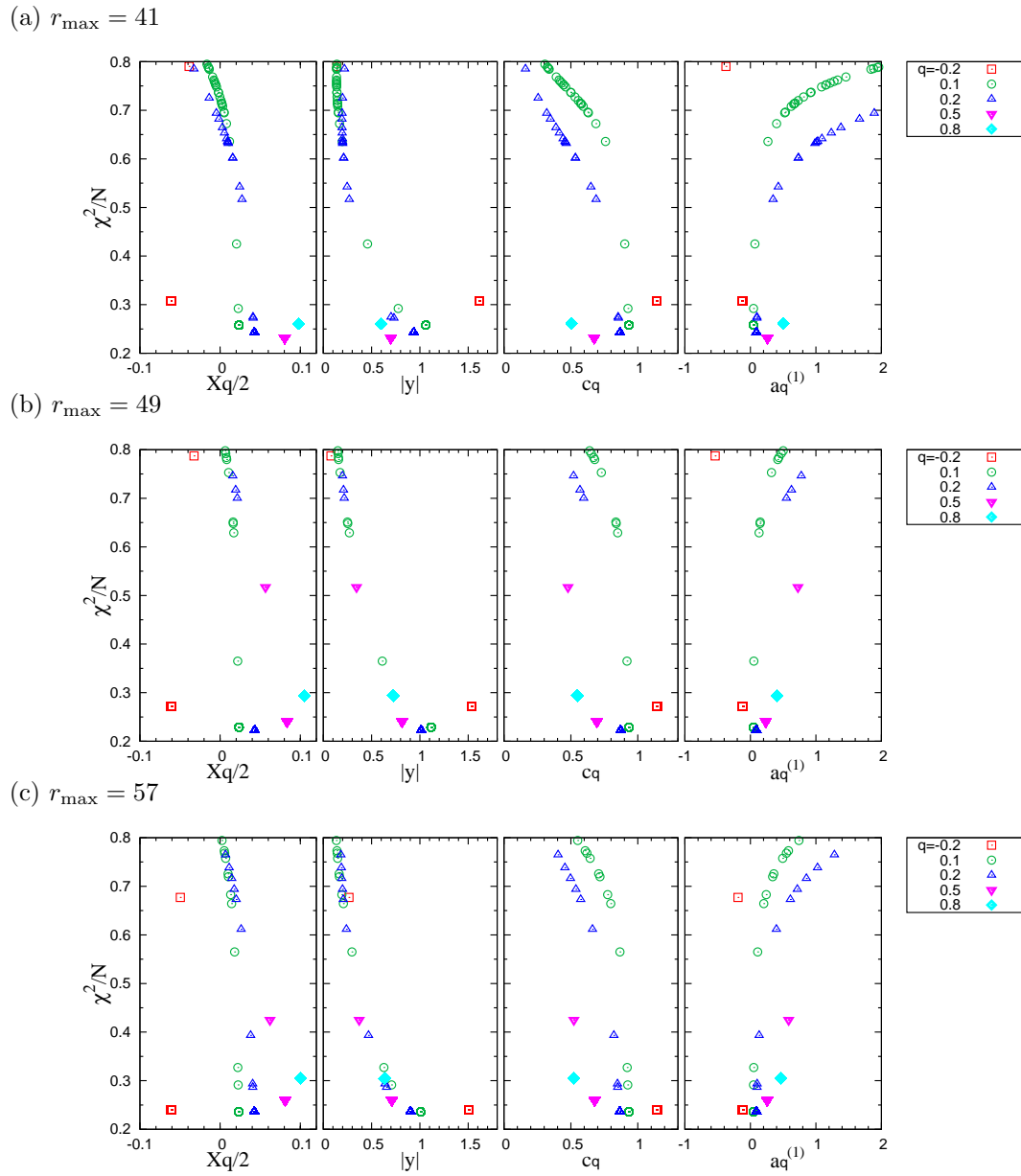


Fig. S12: Stability maps from the scaling analysis of $\langle T^q \rangle$ for various q in case of (a) the maximum distance $r_{\max} = 41$, (b) $r_{\max} = 49$, and (c) $r_{\max} = 57$. The value of q is listed in the legend at the right. The values of fitting parameters at the global minimum are listed in Table. S9.

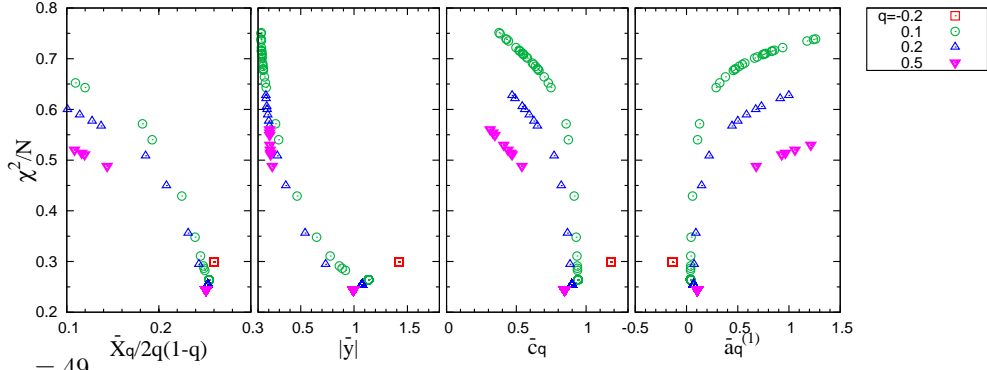
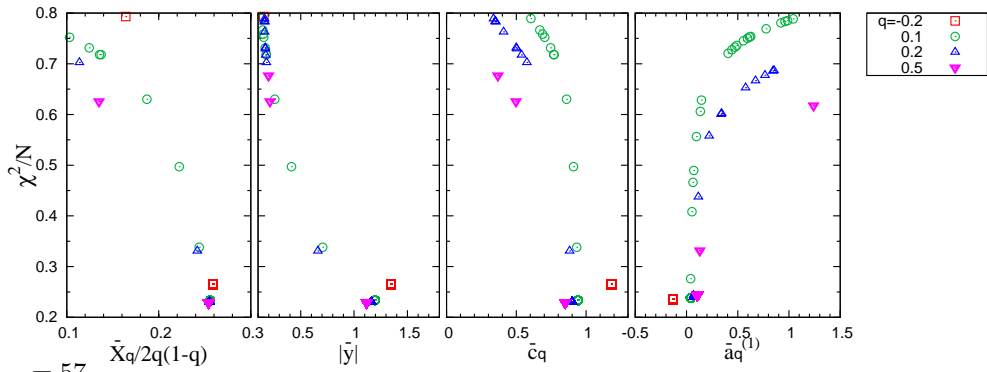
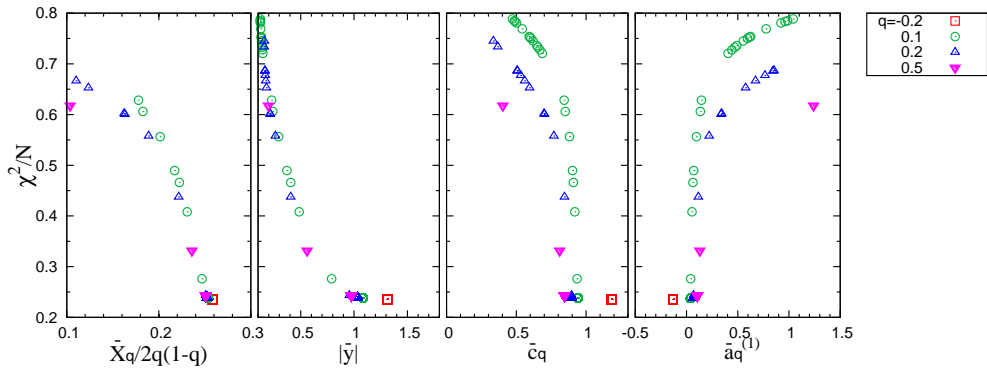
(a) $r_{\max} = 41$

 (b) $r_{\max} = 49$

 (c) $r_{\max} = 57$


Fig. S13: Stability maps from the scaling analysis of $\langle p_q(T) \rangle$ for various q in; case of (a) the maximum distance $r_{\max} = 41$, (b) $r_{\max} = 49$, and (c) $r_{\max} = 57$. The value of q is listed in the legend at the right. The values of fitting parameters at the global minimum are listed in Table. S10.

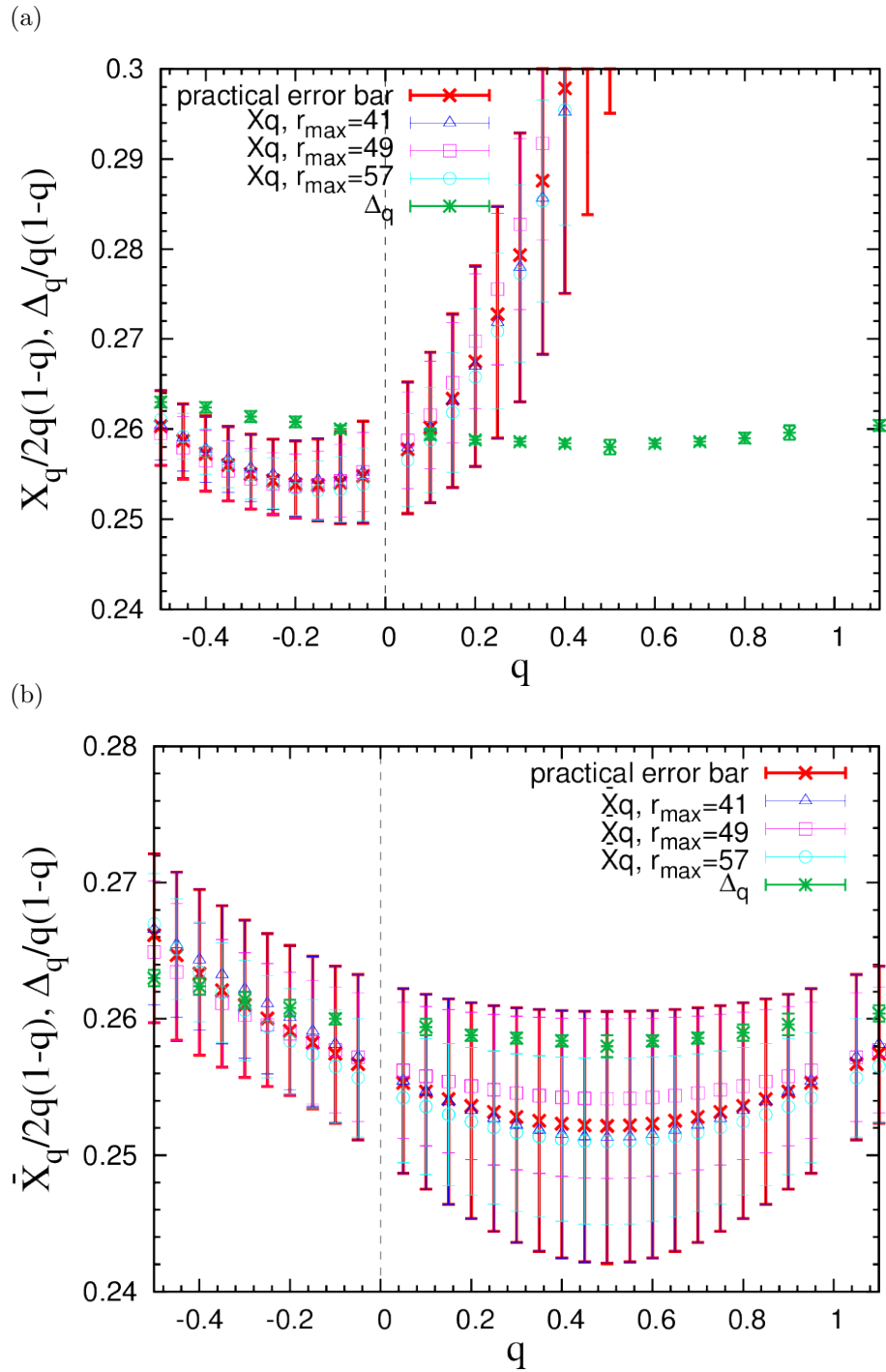


Fig. S14: The q dependence of (a) X_q from $\langle T^q \rangle$ and (b) \bar{X}_q from $\langle p_q(T) \rangle$. The thin and thick lines represent the error bars estimated from the error-propagation theory and the practical error bars, respectively.

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