



Title	Convergence of multigrid method for edge-based finite-element method
Author(s)	Watanabe, K.; Igarashi, H.; Honma, T.
Citation	IEEE TRANSACTIONS ON MAGNETICS, 39(3), 1674-1676 https://doi.org/10.1109/TMAG.2003.810356
Issue Date	2003-05
Doc URL	http://hdl.handle.net/2115/5898
Rights	© 2003 IEEE. Personal use of this material is permitted. However, permission to reprint/republish this material for advertising or promotional purposes or for creating new collective works for resale or redistribution to servers or lists, or to reuse any copyrighted component of this work in other works must be obtained from the IEEE. " IEEE, IEEE TRANSACTIONS ON MAGNETICS, Volume 39, Issue 3, 2003 Page(s):1674 - 1676
Type	article
File Information	ITM39-3.pdf



[Instructions for use](#)

Convergence of Multigrid Method for Edge-Based Finite-Element Method

K. Watanabe, H. Igarashi, and T. Honma, *Member, IEEE*

Abstract—This paper discusses robustness of the multigrid (MG) method against distortion of finite elements. The convergence of MG method becomes considerably worse as the finite elements become flat. It is shown that the smoother used in the MG method cannot effectively eliminate the high-frequency component of the residue for flat elements, and this gives rise to deterioration in the convergence. Moreover, the multigrid method with conjugate gradient (CG) smoother is shown to be more robust against mesh distortion than that with Gauss–Seidel smoother.

Index Terms—Convergence, eigenvalue, multigrid (MG).

I. INTRODUCTION

THE MULTIGRID (MG) method has been applied to electromagnetic field problems so far [1] to show that it can significantly reduce computational time in comparison with conventional linear solvers such as the incomplete Cholesky conjugate gradient (ICCG) solver. However, it has been pointed out [2] that the convergence of the MG method becomes considerably worse as finite elements become flat. It is important to develop the robust MG method for practical use. In this paper, we pay attention to the property of the smoother that plays a crucial role in the MG method. We investigate the robustness of the MG method with different smoothers. Moreover, the residue of a linear system is decomposed into the Fourier components, and each convergence is numerically investigated for different flatness of elements to clarify the reason for such an effect.

II. FORMULATION

A. Magnetostatic Problem

Let us consider magnetostatic field governed by

$$\nabla \times \nu \nabla \times \mathbf{A} = \mathbf{J}_0 \quad (1)$$

$$\nabla \cdot \mathbf{J}_0 = 0 \quad (2)$$

where ν is the magnetic reluctivity, \mathbf{A} is the vector potential, and \mathbf{J}_0 is the current density. The current vector potential

$$\nabla \times \mathbf{T} = \mathbf{J}_0 \quad (3)$$

is introduced for satisfaction of (2). Equation (1) now leads to

$$\nabla \times \nu \nabla \times \mathbf{A} = \nabla \times \mathbf{T}. \quad (4)$$

Finite-element discretization of (4) results in the system of linear equations

$$[K]\{x\} = \{b\} \quad (5)$$

where $[K]$ is a positively semidefinite matrix which is the discrete counterpart of the operator in the left-hand side of (4), $\{x\}$ and $\{b\}$ denote column vectors corresponding to \mathbf{A} and \mathbf{T} , respectively.

B. Multigrid

It is known that the linear solvers such as Gauss–Seidel and conjugate gradient (CG) methods tend to eliminate the high-frequency components of the residue in (5) more rapidly than the low-frequency components. The MG method is based on this property, that is, the high-frequency residual components are eliminated on a fine mesh by small numbers of iterations of the linear solver (smoother). The remaining residual components are then projected onto a more coarse mesh, in which they now have high frequency that can again be eliminated by small numbers of iterations. The MG method solves (5) successively performing these processes. This procedure is usually called the coarse grid correction. Although there are many variations in the MG method, all these variations are based on the coarse grid correction. The procedure of the two-grid V-cycle method that is the simplest MG method is described later.

Step 1 (Smoothing): The smoothing operation is applied to the system equation

$$[K_f]\{x\} = \{b\} \quad (6)$$

for the fine mesh to obtain approximate solution $\{\tilde{x}\}$, where $[K_f]$ denotes the system matrix defined on the fine mesh. In this step, the high-frequency components in the solution error are eliminated.

Step 2: The residual vector $\{r_f\}$ corresponding to the approximate solution $\{\tilde{x}\}$ is calculated

$$\{r_f\} = \{b\} - [K_f]\{\tilde{x}\}. \quad (7)$$

Step 3 (Restriction): The residual vector is projected onto a coarser mesh using the restriction matrix $[R]$

$$\{r_c\} = [R]\{r_f\} \quad (8)$$

where the component of the matrix $[R]$ is obtained by the following integration:

$$R_{ij} = \int_{E_j^f} \mathbf{w}_i^c \cdot d\mathbf{l} \quad (9)$$

Manuscript received June 18, 2002.

The authors are with the Division of Systems and Information Engineering, Graduate School of Engineering, Hokkaido University, Kita-ku Sapporo 060-8628, Japan (e-mail: wata@em-si.eng.hokudai.ac.jp).

Digital Object Identifier 10.1109/TMAG.2003.810356

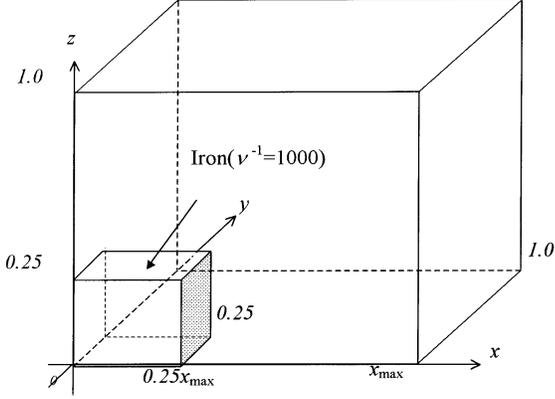


Fig. 1. Simple analysis model (1/8).

and E_j^f denotes j th edge in fine mesh, and w_i^c denotes the interpolation function corresponding to i th edge in the coarse mesh.

Step 4: The residual equation in coarse mesh is solved to obtain the error vector $\{e_c\}$ corresponding to the residual vector $\{r_c\}$

$$[K_c]\{e_c\} = \{r_c\} \quad (10)$$

where $[K_c]$ is the system matrix defined in the coarse mesh. It takes a short amount of time to solve (10) because there are a small number of unknowns in (10). Equation (10) cannot be solved by the direct solver such as Gauss-elimination method, because $[K_c]$ is singular. For this reason, the CG or ICCG method is used in this paper.

Step 5 (Prolongation): The error vector is projected onto the fine mesh using the prolongation matrix $[P]$

$$\{e_f\} = [P]\{e_c\} \quad (11)$$

where $[P]$ is usually chosen as the transpose of $[R]$.

Step 6: The solution $\{\tilde{x}\}$ obtained in Step 1 is corrected using error vector $\{e_f\}$

$$\{x^{\text{new}}\} = \{\tilde{x}\} + \{e_f\}. \quad (12)$$

Step 7 (Post-Smoothing): The smoothing operation is applied to the system equation again. This procedure is called post-smoothing. After post-smoothing, the convergence of the solution is tested. If the convergence condition is not satisfied, we go back to Step 2.

III. COMPARISON OF THE DIFFERENT SMOOTHERS

To investigate the robustness of the MG method against mesh distortion, we analyze a simple magnetostatic problem shown in Fig. 1. Only 1/8 of the model is considered due to the symmetry. The whole region ($0 \leq x \leq x_{\max}$, $0 \leq y, z \leq 1$) is divided into 768 tetrahedral elements for the coarse mesh. The fine mesh is automatically created from the coarse mesh as follows. First, the coarsest mesh is prepared by a mesh generator. The finer meshes are then obtained by dividing each coarse element into eight finer elements as shown in Fig. 2 [3].

To evaluate the flatness of the elements, two types of aspect ratio are defined: $a_1 = l_{\max}/l_{\min}$ and $a_2 = l_{\max}/d_{\min}$, where l_{\max} is the length of the largest edge, l_{\min} is the length of the smallest edge, and d_{\min} is the smaller distance between the

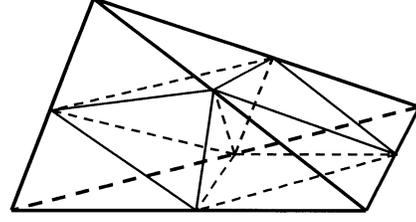


Fig. 2. Division to make the fine mesh.

TABLE I
ASPECT RATIO OF ELEMENTS

x_{\max}	Mesh	Aspect ratio (max./mean)	
		a_1	a_2
1	coarse	1.63/1.63	2.82/2.82
	fine	1.91/1.71	4.06/3.20
10	coarse	10.0/8.38	20.1/15.8
	fine	15.0/8.27	33.6/17.5
20	coarse	20.0/16.7	40.1/31.4
	fine	30.0/16.3	67.1/34.8
50	coarse	50.0/41.6	100/78.4
	fine	75.0/40.3	168/87.0

TABLE II
COMPARISON OF CALCULATION TIME

x_{\max}	Method	Time [s]
1	ICCG	0.8
	MG with GS smoother	0.3
	MG with CG smoother	3.3
10	ICCG	22.0
	MG with GS smoother	6.8
	MG with CG smoother	27.4
20	ICCG	172
	MG with GS smoother	59.4
	MG with CG smoother	49.5
50	ICCG	>300
	MG with GS smoother	182
	MG with CG smoother	102

vertex and that diagonal face in an element [4]. Table I shows the aspect ratio of these meshes. Finite elements become flatter as x_{\max} grows.

Table II shows the calculation time of the MG method with two different smoothers (Gauss-Seidel and CG smoother) as well as of the conventional ICCG method. The calculations are performed on a personal computer with Pentium III-1.26 GHz.

It is shown that the convergence of the MG method is strongly influenced by the mesh quality. Moreover, the MG method with CG smoother is shown to be more robust against mesh distortion than that with Gauss-Seidel smoother.

IV. DECOMPOSITION OF RESIDUAL VECTOR

Here, we consider the property of the eigenvalues in a system matrix. Finite elements become flatter as x_{\max} grows. It is known that convergence of the CG method becomes better (worse) when the condition number

$$k = \lambda_{\max}/\lambda_{\min} \quad (13)$$

TABLE III
CONDITION NUMBER AND EIGENVALUES IN SYSTEM MATRIX

x_{\max}	Condition number	Eigenvalue	
		λ_{\max}	λ_{\min}
1	1.85×10^4	100.7	5.438×10^{-3}
10	1.41×10^5	365.8	2.140×10^{-3}
20	6.68×10^5	723.4	1.083×10^{-3}
50	4.15×10^6	1803	4.345×10^{-4}

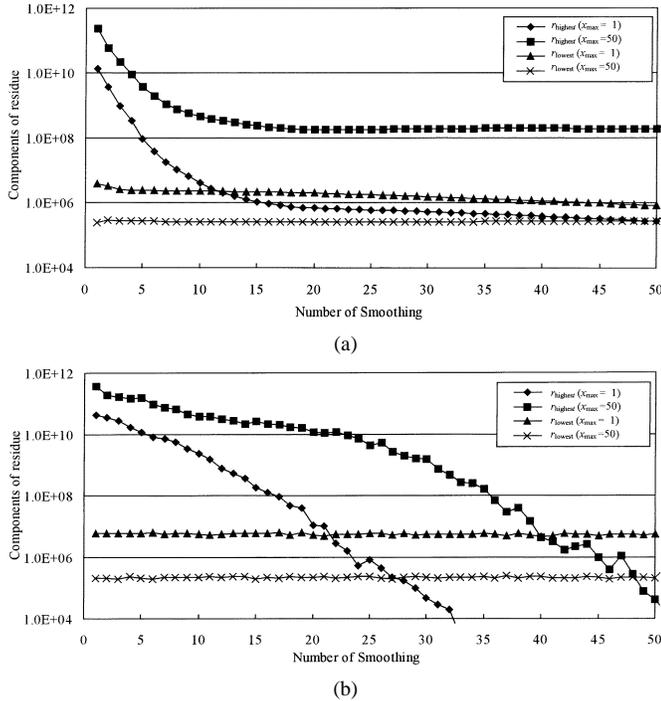


Fig. 3. Reduction of residual components. (a) Gauss-Seidel smoother. (b) CG smoother.

becomes smaller (larger) [5], [6] where λ_{\min} is the nonzero smallest eigenvalue and λ_{\max} is the largest eigenvalue of the system matrix. Table III shows k for different values of x_{\max} . We can see that the condition number becomes larger as x_{\max} increases. The condition number is expected to characterize the convergence of not only the CG method but also the MG method. However, the condition number is not always available because it takes a long time to calculate the eigenvalues.

Next, to consider the cause of the poor convergence in more detail, the convergence of each Fourier component of the residue in the smoothing process of the MG method is plotted in Fig. 3. The residual components r_{lowest} and r_{highest} with the lowest and highest spatial frequency, respectively, are defined by

$$r_{\text{lowest}} = \{v_{\min}\}^T \{r\} \quad (14)$$

$$r_{\text{highest}} = \{v_{\max}\}^T \{r\} \quad (15)$$

where $\{v_{\min}\}$ and $\{v_{\max}\}$ are the eigenvectors corresponding to λ_{\min} and λ_{\max} , and $\{r\}$ is the residual vector after smoothing process.

First, we consider the results of the Gauss-Seidel smoother. When the whole region is nearly a cube ($x_{\max} = 1$), the r_{highest} rapidly reduces within small numbers of iterations as expected. When the element becomes flat ($x_{\max} = 50$), r_{highest} rapidly reduces within small numbers of iterations again. However, there remains the relatively large residue and it hardly decreases any longer. This means that the smoother cannot effectively eliminate the residue for flat elements so that the MG method with Gauss-Seidel smoother requires a number of iterations.

Next, we consider the CG smoother. Although the convergence of r_{highest} is affected by the flatness of element, r_{highest} decreases almost linearly with the iteration. This means that the CG smoother with enough iteration can eliminate the residue even for flat elements. They are consistent with the results shown in Table II.

The residue r_{lowest} in both Gauss-Seidel and CG smoother seemingly unchange because they reduce very slowly. Therefore, r_{lowest} should be reduced by the coarse grid correction.

V. CONCLUSION

This paper discusses dependence of the MG method on the shape of finite elements. The convergence of the MG method becomes considerably worse as the finite elements become flat. It is shown that the smoother used in the MG method cannot effectively eliminate the high-frequency component of the residue for flat elements, and this gives rise to deterioration in the convergence. Moreover, the MG method with CG smoother is shown to be more robust against mesh distortion than that with Gauss-Seidel smoother.

ACKNOWLEDGMENT

The authors would like to thank A. Kameari for helpful discussions.

REFERENCES

- [1] R. Hiptmair, "Multigrid method for Maxwell's equations," *SIAM J. Numer. Anal.*, vol. 36, pp. 204–225, 1998.
- [2] A. Kameari, "Application of geometrical multigrid method to electromagnetic computation by finite element method," (in Japanese), Tech. Rep. IEEJ, SA-01-11, RM-01-79, 2000.
- [3] M. Schinnerl, J. Schoberl, and M. Kaltenbacher, "Nested multigrid methods for the fast numerical computation of 3D magnetic fields," *IEEE Trans. Magn.*, vol. 36, pp. 1557–1560, July 2000.
- [4] F. X. Zgainski, Y. Marechal, and J. L. Coulomb, "An *a priori* indicator of finite element quality based on the condition number of the stiffness matrix," *IEEE Trans. Magn.*, vol. 33, pp. 1748–1751, Mar. 1997.
- [5] H. Igarashi, "On the property of the curl-curl matrix in finite element analysis with edge elements," *IEEE Trans. Magn.*, vol. 37, pp. 3129–3132, Sept. 2001.
- [6] S. Kaniel, "Estimates for some computational techniques in linear algebra," *Math. Comput.*, pp. 369–378, 1966.