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An ideal emitter of entangled photon pairs combines the perfect symmetry of an atom with the convenient electrical trigger of light sources based on semiconductor quantum dots. Our source consists of strain-free GaAs dots self-assembled on a triangular symmetric (111)A surface. The emitted photons reveal a fidelity to the Bell state as high as 86(±2)% without postselection. We show a violation of Bell’s inequality by more than five times the standard deviation, a prerequisite to test a quantum cryptography channel for eavesdropping. Due to the strict nonlocal nature the source can be used for real quantum processing without any postprocessing. The remaining decoherence channel of the photon source is ascribed to random charge and nuclear spin fluctuations in and near the dot.

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The details of dot growth are reported elsewhere. We employed a standard molecular beam epitaxy machine. After growing an Al_{0.3}Ga_{0.7}As layer on the gallium-rich surface of a GaAs (111)A substrate, we supplied a 0.043 monolayer of gallium that formed Ga droplets at 400 °C. Then we supplied As$_3$ to crystallize the droplets into GaAs dots at 200 °C, followed by annealing at 500 °C. Several microscope observations which include in vacuo scanning tunneling microscopy and atomic force microscopy revealed the formation of dots with a truncated cone shape whose average radius and height were 16 and 1.4 nm, respectively. Figure 1(b) shows the atomic force microscope image of an investigated dot, which exhibits no lateral elongation. This is in stark contrast to dots grown on (100) surfaces, which exhibit significant elongation along [1-10].

As an excitation source we used a pulsed semiconductor laser with a high repetition frequency of 200 MHz. We simultaneously counted three photon channels, i.e., XX photons projected onto a given polarization state, X photons projected onto another polarization state (such as |L⟩), and its orthogonal complement (such as |R⟩). The use of three detectors enabled us to eliminate the influence of excitation fluctuations on coincidence visibility. The number of coincidence was analyzed with a time-to-digital converter. The typical integration time was 10 min for each polarization condition. All the experiments were performed at 9 K.

**Correlation measurements: Entanglement fidelity.** Figure 1(c) shows the photoluminescence (PL) spectrum of an isolated dot. It consists of four main lines, which are identified as being from the high-energy side, neutral excitons (X), positively charged excitons (X$^+$), neutral biexcitons (XX), and negatively charged excitons (X$^-$). For performing the correlation measurements, we select as-grown dots without a detectable FSS from the sample. Polarized PL was analyzed with a spectral resolution comparable to the radiative width, which is expected to be 1.2 μeV [560 ps in terms of lifetime; see Fig. 4(a)]. Small but nonzero FSS values are confirmed for most of the dots, and they are distributed around a mean value of 10 ± 5 μeV. This is noticeably smaller than both the typical values for Stranski-Krastanov grown dots, and those for droplet epitaxial GaAs dots grown on (100). In the investigated sample 5% of the dots show no detectable FSS. We have measured the photon correlations in more than ten selected dots and they all exhibit entanglement.

Figure 2 shows the results of photon correlation measurements in a typical dot. See Fig. 1 in the Supplemental Material for a spectral characterization in this dot. L, R, H, and V indicate projections along the left-handed circular, right-handed circular, linear laboratory horizontal, and vertical polarizations, respectively. D is linear diagonal with a polarization axis tilted by 45° from H, and A is antidiagonal where $A \perp D$. The top panel in Fig. 2(a) shows a coincidence histogram for L-polarized XX photons and R-polarized X photons (denoted by LR). The presence of a central peak confirms a radiative cascade. The XX and X photons are clearly correlated, resulting in a higher probability than that for detecting uncorrelated photons. The central peak disappears for a polarization combination of LL (second panel). Thus, the probability of observing both XX and X photons in L is close to zero. The same anticorrelation is confirmed for RR (third panel), but a positive correlation is recovered for RL (bottom panel). These results imply that the two-photon polarization state can be approximated by one of the Bell (maximally entangled) states,

$$|\Psi⟩ = \frac{|LR⟩ + |RL⟩}{\sqrt{2}}.$$  

A key criterion for entanglement is the presence of a correlation independent of the chosen polarization basis.
violation of Bell's inequality. To verify this experimentally, are linked to nonlocality, which is only assured by the milestone, but the most powerful applications of entanglement photon pair emission from quantum dots (QDs) is an important plane of the Poincaré sphere (the polar angle of a polarization state that moves in the Supplemental Material). Sinusoidal oscillations in the angle dependence was also measured and shown in Fig. 2(b) It was experimentally controlled by the application of phase a maximum-likelihood technique. The presence of four real values at the corner of the matrix, with negligible values for the others, demonstrates the superior characteristics of our source. The matrix has a partial transpose with the minimum eigenvalue of $-0.36 < 0$, which clearly satisfies the Peres criterion of entanglement, which assures quantum inseparability. The density matrix allows us to evaluate the degree of coherence and the degree of mixedness of the measured state in terms of the tangle ($T$) and the linear entropy ($S_L$), respectively. From $T$ we derive one of the most basic measure of the entanglement of formation ($E_F$). Our results reveal that ($T,S_L,E_F$) = (0.53,0.32,0.63). Note that we achieve these values even without the postselection (or any local operation) of the photons.

Figure 2(b) shows coincidence histograms for rectilinear polarizations. A positive correlation appears for parallel polarizations ($HH, VV$), while it disappears for perpendicular polarizations ($HV, VH$). These results agree with the expression of the Bell state of Eq. (1) in a linear polarization basis, We define the correlation visibility $C = \langle n_\parallel - n_\perp \rangle / (n_\parallel + n_\perp)$, where $n_\parallel$ is the number of coincidences normalized with the two-photon flux for a co-polarized basis, and $n_\perp$ is that for a cross-polarized basis (see the Supplemental Material for the normalization procedure). An ideal source is expected to show $C = 1$ for any orthogonal basis set. Our results show that $C = 0.87 \pm 0.03$ for $R/L$ and $C = 0.78 \pm 0.03$ (0.77 \pm 0.03) for $H/V$ ($D/A$). The visibility for linear polarizations is found to be approximately independent of the polarization direction, which demonstrates the isotropic characteristic of our source [Fig. 2(a) in the Supplemental Material]. The higher $C$ value for the circular basis than for the linear bases originates from the hyperfine interaction of the exciton with nuclear spins (see the Supplemental Material for discussion). The entanglement fidelity is defined as the projection amplitude of a measured polarization state on a target Bell state, which is given by $f = (1 + C_R/L + C_H/V + C_D/A)/4$. Our results reveal that $f = 0.86(\pm0.02)$, which is much larger than the classical limit of 0.5, and rates among the best reported in previous studies on dot-based photon sources. Experimental violation of Bell’s inequality. Entangled photon pair emission from quantum dots (QDs) is an important milestone, but the most powerful applications of entanglement are linked to nonlocality, which is only assured by the violation of Bell’s inequality. To verify this experimentally, we have performed additional photon correlation experiments in nonorthogonal polarization bases. Figure 3(a) shows normalized coincidence counts as a function of the polarization angle of $X$ ($\theta_X$) at four different angle settings for $XX$ polarization ($\theta_{XX}$). Note that we define the angle of $\theta$ as the polar angle of a polarization state that moves in the $RLHV$ plane of the Poincaré sphere ($\theta = 0$ for $R$ and $\theta = 90^\circ$ for $H$). It was experimentally controlled by the application of phase retardance to each beam using liquid crystals. The azimuth-angle dependence was also measured and shown in Fig. 2(b) in the Supplemental Material. Sinusoidal oscillations in the coincidence counts provide evidence of quantum interference, distinct from classical correlation. The maximum violation of Bell’s inequality in the Clauser-Horne-Shimony-Holt form is expected to appear for polarization correlations with $\theta_{XX} = 0^\circ$, $90^\circ$, $180^\circ$, $270^\circ$ and $\theta_X = 45^\circ$, $135^\circ$, $225^\circ$, and $315^\circ$. We measure the coincidence counts for these settings, and estimate the $S$ parameter to be $2.33 \pm 0.06 > 2$. It clearly violates Bell’s inequality by more than five times the standard deviation, which is definite proof of the nonlocality of the measured photons. The importance of nonlocality of entanglement is seen in the Ekert protocol. If an eavesdropper measures either state of an entangled pair, it becomes local reality and nonlocality vanishes. This is why Bell’s inequality can serve as a secure test against eavesdropping. Other applications such as quantum teleportation and entanglement swapping also rely on nonlocality. Consequently, our device is a QD source which can be used for real applications without applying temporal or spectral filtering techniques.
FIG. 4. (Color online) (a) The decay of circularly polarized PL signals for the $X$ line after short-pulsed and quasiresonant excitation. Here we study the same dot as that used in the correlation measurement. (b) The degree of circular polarization, defined as $|I(\sigma) - I(\bar{\sigma})|/[I(\sigma) + I(\bar{\sigma})]$, where $I(\sigma)$ and $I(\bar{\sigma})$ are the co-circular and cross-circular intensities, respectively. The dashed line is an exponential fit to data, with an estimated decay time of $\Gamma_1^{-1} = 1.5$ ns. The inset shows the energy diagram of the exciton state. In the experiment, we used a short-pulsed parametric oscillator that emitted 4 ps pulses with a wavelength shifted by an optical phonon energy of 37 meV from the $X$ line. The excitation polarization was set as circular, and temporally modulated to maintain an equilibrium nuclear environment. Polarized PL was detected by a fast-response photomultiplier tube with a response time of 40 ps.

to a finite probability of observing depolarized photons. We can estimate the correlation visibility of photon pairs to be $\Gamma_1/(\Gamma_1 + \Gamma_2) \approx 0.7$, which is in fairly good agreement with the observed $C$ value. These findings indicate that our source is neither affected by incoherent noise associated with carrier recapturing\cite{17,20} nor light emission from other luminescent centers than the dot. The degree of entanglement is thus purely limited by the scattering of excitons. We ascribe the exciton depolarization to random charge and nuclear spin fluctuations in and near the dot. The slowly varying environment represents a remaining source of asymmetry that limits the degree of quantum interference even in a solid-state photon source based on symmetric dots, as discussed in the Supplemental Material.\cite{39}

In summary, we have demonstrated the generation of entangled photon pairs using a strain-free GaAs dot as a symmetric artificial-atom cascade on (111)A surfaces. A clear violation of Bell’s inequality is observed in correlation measurements that do not rely on postselection through filtering or tuning. Using our source in quantum cryptography applications would allow safe key distribution not based on a mathematical difficulty but on a fundamental physical law that protects the system, namely, the completeness of quantum mechanics.\cite{3} The purity of our source also paves the way for using droplet dot-based emitters to investigate the exact connection between the security of quantum cryptography and tests of quantum nonlocality\cite{24} in the solid state. We clarified the impact of the remaining exciton depolarization on the degree of entanglement in the emitted pairs. The influence of depolarization could be efficiently suppressed in the future by using the Purcell enhancement of the radiative rate\cite{17} or time-domain filtering with fast-response detectors.\cite{20} As droplet epitaxy is not strain driven, a large variety of dot-barrier material combinations can be grown. This opens up the possibility of extending the current emission wavelengths around 700 nm (that yields the highest photon counting efficiency in Si-based detectors) towards the wavelength regions suitable for long-distance fiber-based telecommunication.

Note added. Recently, entangled photon emission from dots grown on patterned (111)B substrates has been reported.\cite{36}

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