



Title	The distance function and defect energy
Author(s)	Aviles, P.; Giga, Y.
Citation	Hokkaido University Preprint Series in Mathematics, 293, 1-21
Issue Date	1995-6-1
DOI	10.14943/83440
Doc URL	<a href="http://hdl.handle.net/2115/69044">http://hdl.handle.net/2115/69044</a>
Type	bulletin (article)
File Information	pre293.pdf



[Instructions for use](#)

**THE DISTANCE FUNCTION AND  
DEFECT ENERGY**

**P. Aviles and Y. Giga**

**Series #293. June 1995**

**HOKKAIDO UNIVERSITY**  
**PREPRINT SERIES IN MATHEMATICS**

- # 270 A. Arai, Gauge theory on a non-simply-connected domain and representations of canonical commutation relations, 18 pages. 1994.
- # 271 S. Jimbo, Y. Morita and J. Zhai, Ginzburg Landau equation and stable steady state solutions in a non-trivial domain, 17 pages. 1994.
- # 272 S. Izumiya, A. Takiyama, A time-like surface in Minkowski 3-space which contains light-like lines, 7 pages. 1994.
- # 273 K. Tsutaya, Global existence of small amplitude solutions for the Klein-Gordon-Zakharov equations, 11 pages. 1994.
- # 274 H. Kubo, On the critical decay and power for semilinear wave equations in odd space dimensions, 22 pages. 1994.
- # 275 N. Terai, T. Hibi, Alexander duality theorem and second Betti numbers of Stanley-Reisner rings, 2 pages. 1995.
- # 276 N. Terai, T. Hibi, Stanley-Reisner rings whose Betti numbers are independent of the base field, 12 pages. 1995.
- # 277 N. Terai, T. Hibi, Computation of Betti numbers of monomial ideals associated with cyclic polytopes, 11 pages. 1995.
- # 278 N. Terai, T. Hibi, Computation of Betti numbers of monomial ideals associated with stacked polytopes, 8 pages. 1995.
- # 279 N. Terai, T. Hibi, Finite free resolutions and 1-skeletons of simplicial  $(d - 1)$ -spheres, 3 pages. 1995.
- # 280 N. Terai, T. Hibi, Monomial ideals and minimal non-faces of Cohen-Macaulay complexes, 6 pages. 1995.
- # 281 A. Arai, N. Tominaga, Analysis of a family of strongly commuting self-adjoint operators with applications to perturbed d'Alembertians and the external field problem in quantum field theory, 44 pages. 1995.
- # 282 T. Mikami, Asymptotic behavior of the first exit time of randomly perturbed dynamical systems with a repulsive equilibrium point, 29 pages. 1995.
- # 283 K. Iwata, J. Schäfer, Markov property and cokernels of local operators, 17 pages. 1995.
- # 284 T. Nakazi, M. Yamada, Riesz's Functions In Weighted Hardy And Bergman Spaces, 20 pages. 1995.
- # 285 K. Hidano, K. Tsutaya, Scattering theory for nonlinear wave equations in the invariant Sobolev space, 32 pages. 1995.
- # 286 A. Arai, Strong coupling limit of the zero-energy-state density of the Dirac-Weyl operator with a singular vector potential, 8 pages. 1995.
- # 287 T. Nakazi, Factorizations of outer functions and extremal problems, 15 pages. 1995.
- # 288 A. Kishimoto, The Rohlin property for automorphisms of UHF algebras, 15 pages. 1995.
- # 289 K. Goto, A. Yamaguchi and I. Tsuda, Nine-bit states cellular automata are capable of simulating the pattern dynamics of coupled map lattice, 24 pages. 1995.
- # 290 Y. Giga, Interior derivative blow-up for quasilinear parabolic equations, 16 pages. 1995.
- # 291 F. Hiroshima, Functional Integral Representation of a Model in QED, 48 pages. 1995.
- # 292 N. Kawazumi, A Generalization of the Morita-Mumford Classes to Extended Mapping Class Groups for Surfaces, 11 pages. 1995.

# THE DISTANCE FUNCTION AND DEFECT ENERGY

PATRICIO AVILES

DEPARTMENT OF MATHEMATICS AND PHYSICS

OXFORD UNIVERSITY, U.K.

OR

ETH. CH-8592, ZURICH

SWITZERLAND

YOSHIKAZU GIGA

DEPARTMENT OF MATHEMATICS

HOKKAIDO UNIVERSITY

SAPPORO 060, JAPAN

## 1. INTRODUCTION

It is important to measure the energy of jump discontinuities of a unit length gradient field  $\nabla\varphi$  in a bounded Lipschitz domain in  $\mathbb{R}^n$ .

Such problems arise in the modelling of smectic liquid crystals [SK], [AG1] or of the blistering of thin films [OG]. The quantity measuring the energy of the jump discontinuities, the defect of  $\nabla\varphi$ , is

$$J^\beta(\varphi) = \int_{\Sigma} |\nabla\varphi^+ - \nabla\varphi^-|^\beta d\mathcal{H}^{n-1}$$

where  $\beta > 0$ ; we call it a defect energy. Here  $\Sigma$  is the set of jump discontinuities of  $\nabla\varphi$  and  $\nabla\varphi^\pm$  is the trace of  $\nabla\varphi$  of each side of  $\Sigma$ ;  $\mathcal{H}^{n-1}$  is the  $n - 1$  dimensional Hausdorff measure which is the surface element when  $\Sigma$  is smooth.

There may be a lot of Lipschitz solutions of the eikonal equation

$$|\nabla\varphi| = 1 \text{ in } \Omega \text{ with } \varphi = 0 \text{ on } \partial\Omega,$$

but the distance function

$$d = d(x, \partial\Omega) = \inf\{|x - y|; y \in \partial\Omega\}$$

is the unique viscosity solution of the problem [CIL]. In other words the theory of viscosity solutions selects a solution of the eikonal equation. There is a fundamental question whether the distance function minimizes  $J^\beta$  among all *nonnegative* solutions of the eikonal equation.

If the space dimension  $n$  equals one,  $J^\beta$  just measures a constant multiple of the number of jumps of  $\nabla\varphi$ . There is no solution of the eikonal equation having no defect satisfying the zero boundary condition. Thus, the distance function is a (unique) minimizer of  $J^\beta$  since it has only one jump of the derivative of  $\varphi$ . However, for multidimensional case, the situation is different.

In this paper, we focus on the case  $\beta = 1$  because of independent interest related to the total variation of the Hessian

$$I(\varphi) = \int_{\Omega} |\nabla^2\varphi|.$$

This integral is closely related to  $J^1$ . Indeed, if  $\varphi$  is piecewise linear, more precisely,  $\nabla^2\varphi = 0$  (as a measure) outside  $\Sigma$ , then

$$I(\varphi) = \int_{\Sigma} |\nabla^2\varphi| = \int_{\Sigma} |\nabla^+\varphi \cdot \nu - \nabla^-\varphi \cdot \nu| d\mathcal{H}^{n-1}$$

where  $\nu$  is the approximate normal of  $\Sigma$  [G]. Since the tangential component of  $\nabla\varphi$  is approximately continuous,  $|\nabla^+\varphi \cdot \nu - \nabla^-\varphi \cdot \nu| = |\nabla^+\varphi - \nabla^-\varphi|$  if  $|\nabla\varphi| = 1$ . Thus,  $I(\varphi) = J^1(\varphi)$  for piecewise linear  $\varphi$ . Our principal results are

- (i) the distance function is the unique minimizer of  $I(\varphi)$  among all nonnegative (Lipschitz) solutions of the eikonal equation  $|\nabla\varphi| = 1$  in  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$  provided that  $\Omega$  is convex and  $n = 2$ . The values of minimum equals  $\mathcal{H}^{n-1}(\partial\Omega)$ .

- (ii) there is a simply connected nonconvex domain  $\Omega$  in  $\mathbb{R}^2$  such that the distance function is not a minimizer of  $J^1$  nor  $I$ .

This suggests that the selection mechanism of the ground state by  $I$  or  $J^1$  is different from that in the theory of viscosity solutions, in general.

To show (i) we first observe that

$$|\Delta\varphi| = |\nabla^2\varphi|$$

as measures if  $\varphi$  solves  $|\nabla\varphi| = 1$  and  $n = 2$ . This depends on the fact that  $\nabla^2\varphi$  has rank one which is easy to observe heuristically. Differentiating  $|\nabla\varphi| = 1$  implies that one of eigenvalues of  $\nabla^2\varphi$  always equal zero. To carry out this idea we appeal to the theory of functions of bounded variation [G]. Note that the singular part (w.r.t. the Lebesgue measure) of  $\nabla^2\varphi$  always has rank one [A1], [AG2]. Another key observation is

$$\int_{\Omega} |\Delta\varphi| \geq \int_{\Omega} -\Delta\varphi = \mathcal{H}^{n-1}(\partial\Omega)$$

if  $|\nabla\varphi| = 1$ ,  $\varphi \geq 0$  in  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$ . The last equality formally follows from integration by parts and the fact that  $|\nabla\varphi|$  agrees with inward normal derivative of  $\varphi$  on  $\partial\Omega$ . In section 2 we state these observations in a rigorous way allowing that  $\nabla^2\varphi$  is a measure. If  $\Omega$  is convex, the distance function  $d$  is concave in  $\Omega$  so that  $-\Delta d \geq 0$  in  $\Omega$  (in the distribution sense). Thus

$$\int_{\Omega} |\Delta d| = \int_{\Omega} -\Delta d = \mathcal{H}^{n-1}(\partial\Omega)$$

so that  $d$  minimizes  $I$  as well as  $\int_{\Omega} |\Delta\varphi|$ . It turns out that  $d$  is a unique minimizer among all  $\varphi$ ,  $|\nabla\varphi| = 1$ ,  $\varphi \geq 0$  in  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$ . The inequality

$$\int_{\Omega} |\Delta\varphi| \geq \int_{\Omega} -\Delta\varphi$$

is not sharp unless  $\Omega$  is convex. In other words the minimum of  $I$  is strictly greater than  $\mathcal{H}^{n-1}(\partial\Omega)$  (Theorem 2.5.) The proof of (ii) depends on an explicit construction of the domain  $\Omega$ .

As a corollary of (i) we get: if  $d$  is piecewise linear, more precisely,  $\nabla^2 d = 0$  outside the

defect as a measure, then  $d$  also minimizes  $J^1$  (among all nonnegative solutions of the eikonal equations) provided that the domain is convex. Note that such  $d$  exist if and only if the domain is a convex polygon as shown in Remark in §2.1.

Our counterexamples are interesting for the study of minimizers of singular perturbed variational problem

$$E_\varepsilon(\varphi) = \int_\Omega W(\nabla\varphi) + \varepsilon^2 \int_\Omega |\nabla^2\varphi|^2, \quad W(p) = (1 - |p|^2)^\sigma, \quad \sigma > 0$$

in a plane domain  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$ . Since the Euler-Lagrange equation is fourth order, we are entitled to impose another boundary condition. The natural choice seems to be  $\partial u/\partial\nu = -1$ , where  $\nu$  is the unit outward normal of  $\partial\Omega$ . We divide  $E_\varepsilon$  by  $\varepsilon$  so that we hope that the energy has a nonzero limit as  $\varepsilon \rightarrow 0$ .

Formal analysis for  $\sigma = 2$  done by [AG1], [OG] suggests that this problem has Gamma-limit

$$\tilde{J} = 2 \int_\Sigma |\nabla\varphi^+ \cdot \nu - \nabla\varphi^- \cdot \nu| \int_{-b}^b |1 - (a^2 + \tau^2)|^{\sigma/2} d\tau d\mathcal{H}^{n-1}, \quad a = |(\nabla\varphi)_{\text{tan}}|, \quad b = (1 - a^2)^{1/2},$$

where  $(\nabla\varphi)_{\text{tan}}$  denotes the tangential component of  $\nabla\varphi^+$  (or  $\nabla\varphi^-$ ). Since  $|\nabla\varphi| = 1$ , we see  $\tilde{J}$  is a positive constant times  $J^{\sigma+1}$ . This  $\tilde{J}$  (or  $J^{\sigma+1}$ ) is to be minimized subject to the same boundary condition as for  $E_\varepsilon$ , and the interior condition  $|\nabla\varphi| = 1$  a.e. in  $\Omega$ .

It is tempting to think that the minimizer  $\varphi_\varepsilon$  of  $E_\varepsilon$  tends to

$$\varphi_0(x) = d(x, \partial\Omega).$$

as  $\varepsilon \rightarrow 0$ . Similarly, it is tempting to think that this function might be a minimizer of  $\tilde{J}$  (or  $J^{\sigma+1}$ ). These conjectures are more or less explicit in [OG] (cf. [AG1] for  $\sigma = 2$ ).

An extended version of our examples (Theorem 3.5) says that the second conjecture is false for some nonconvex domain at least for  $\sigma < \beta_0 - 1$  with some  $\beta_0 > 1$  close to one. Unfortunately in our examples  $\beta_0$  is less than 3 so they do not solve the original conjecture for  $\sigma = 2$ . However, they are important because they show some possible pitfalls. In particular, they show that if these conjectures are true for  $\sigma = 2$ , then the reasons must be subtle since other equally

reasonable-sounding statements are false.

The limiting process of  $E_\varepsilon$  as  $\varepsilon \rightarrow 0$  is not at all clear compared with the case when  $\nabla\varphi$  is a scalar function. Such a convergence problem is studied in [KM1], [KM2] for  $\nabla\varphi$  when  $W$  has isolated equal minimums.

The authors are grateful to Professor Robert Kohn for his interest and valuable comments on this paper. Much of this work was done while the second author visited the Japan-American Mathematics Institute of the Johns Hopkins University. Its hospitality is gratefully acknowledged, as is support from Japan Society for the Promotion of Science. This work is partly supported by THE SUHARA MEMORIAL FOUNDATION and YAMADA SCIENCE FOUNDATION.

## 2. ESTIMATE OF TOTAL VARIATIONS OF GRADIENT FIELD OF LENGTH ONE

We are concerned with total variation of  $\nabla^2\psi$  in a bounded two dimensional domain  $\Omega$  when  $|\nabla\psi| = 1$ ,  $\psi \geq 0$  on  $\Omega$  and  $\psi = 0$  on the boundary  $\partial\Omega$ . Our principal result in this section is that the minimum of the total variation is attained (uniquely) at the distance function provided that  $\Omega$  is convex.

**2.1. Notation.** Let  $\Omega$  be a bounded domain in  $\mathbb{R}^n$  with Lipschitz boundary  $\partial\Omega$ . For a Lebesgue integrable function  $\varphi$ , i.e.  $\varphi \in L^1(\Omega)$ , let  $\nabla\varphi = (\partial_i\varphi)_{i=1}^n$  and  $\nabla^2\varphi = (\partial_i\partial_j\varphi)$  ( $1 \leq i, j \leq n$ ) a distributional gradient and Hessian of  $\varphi$ , respectively. Let  $X$  be the space of  $\varphi \in L^1(\Omega)$  such that  $\partial_i\varphi \in L^1(\Omega)$  ( $1 \leq i \leq n$ ) and  $\partial_i\partial_j\varphi$  is a finite Radon measure on  $\Omega$  ( $1 \leq i, j \leq n$ ). In other words,  $\partial_i\varphi$  is a function of (essentially) bounded variation, i.e.  $\partial_i\varphi \in BV(\Omega)$ . Let us recall fundamental decomposition of  $\nabla^2\varphi$  for  $\varphi \in X$ ; see e.g. [AG2]. Let  $\Omega_0$  be the largest subset in  $\Omega$  such that  $\nabla^2\varphi$  is absolutely continuous in  $\Omega_0$  and let  $\Sigma$  be the set of jump discontinuities of  $\nabla\varphi$ . Then

$$\nabla^2\varphi = \nabla^2\varphi|_{\Omega_0} + \nabla^2\varphi|_{(\Omega - \Omega_0 - \Sigma)} + \nu \otimes (\nabla\varphi^+ - \nabla\varphi^-)\mathcal{H}^{n-1}|_{\Sigma}.$$



Here for a set  $Z$  and measure  $\mu$  we associate a new measure  $\mu \llcorner Z$  by

$$(\mu \llcorner Z)(B) = \mu(Z \cap B), B \subset \Omega.$$

The vector field  $\nu$  is the approximate unit normal of  $\Sigma$  and  $\nabla\varphi^\pm$  is the trace of  $\nabla\varphi$  on  $\Sigma$  in the direction of  $\pm\nu$ ;  $\mathcal{H}^k$  denotes the  $k$ -dimensional Hausdorff measure. The first term  $(\nabla^2\varphi)^{ab} = \nabla^2\varphi \llcorner \Omega_0$  is often called the *absolutely continuous* part of  $\nabla^2\varphi$ . We always identify  $(\nabla^2\varphi)^{ab}$  with corresponding locally Lebesgue integrable function in  $\Omega_0$ . The second term is often called the *mild* part and it lies on a non rectifiable set  $\Omega - \Omega_0 - \Sigma$  of Lebesgue measure zero. The sum of last two terms is called the *singular* part of  $\nabla^2\varphi$ . Let  $Y$  be the space of  $\varphi \in X$  such that  $\nabla\varphi$  has no absolute continuous part and no mild part. In other words

$$Y = \{\varphi \in X; \nabla^2\varphi = \nu \otimes (\nabla\varphi^+ - \nabla\varphi^-) \mathcal{H}^{n-1} \llcorner \Sigma\}.$$

Let  $A$  be the space of  $\varphi \in X$  such that  $|\nabla\varphi| = 1$  a.e. in  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$ . We need three subclasses of  $A$

$$A_+ = \{\varphi \in A; \varphi \geq 0 \text{ in } \Omega\}, \quad A^0 = A \cap Y, \quad A_+^0 = A_+ \cap A^0.$$

We consider two integrals for  $\varphi \in X$  which measure jumps of  $\nabla\varphi$ .

$$I(\varphi) = \int_{\Omega} |\nabla^2\varphi|,$$

$$J^\beta(\varphi) = \int_{\Sigma} |\nabla\varphi^+ - \nabla\varphi^-|^\beta d\mathcal{H}^{n-1} \quad \text{for } \beta > 0.$$

Since  $\nabla^2\varphi$  is a finite Radon measure, in the representation

$$\int_{\Omega} |\nabla^2\varphi| = \sup \left\{ \sum_{1 \leq i, j \leq n} \int_{\Omega} \theta^{ij} \partial_i \partial_j \varphi; \sum_{i, j} |\theta_{ij}|^2 \leq 1 \right\}$$

the test function  $\theta_{ij}$  is allowed to be  $\theta_{ij} \in C^1(\bar{\Omega})$  not necessarily compactly supported.

**Remark.** The set  $A_+^0$  and even  $A^0$  may be empty. In fact,  $A_+^0$  (and  $A^0$ ) is empty if  $\partial\Omega$  has a 'curved' part. Conversely,  $A_+^0$  is nonempty if  $\Omega$  is a polygon. The proof is by the induction of

numbers of vertices of  $\Omega$ . If  $\Omega$  is a triangle, the distance function  $d$  is certainly piecewise linear.

If  $\Omega$  is a polygon of  $m$  ( $> 3$ ) vertices, we set

$$\rho(x, \partial\Omega) = \min\{d(x, L(S)); L(S) \text{ is a straight line containing an edge } S \text{ of } \partial\Omega\}$$

$$d_* = \inf\{\rho(x, \partial\Omega); \text{there is at least three edges } S_1, S_2, S_3 \text{ of } \partial\Omega$$

$$\text{such that } \rho(x, \partial\Omega) = d(x, L(S_i)), i = 1, 2, 3 \text{ and } x \in \Omega\}.$$

It is easy to see that  $\rho$  is piecewise linear in

$$\Omega_* = \{x \in \Omega, d(x) < d_*\}.$$

By the choice of  $d_*$  the set  $K = \Omega - \Omega_*$  is a closed polygon with at most  $m - 1$  vertices; it may have no interior so that  $K$  is a set of points or segments. Let  $\Omega'$  be the interior of  $K$  so that it is a polygon with at most  $m - 1$  vertices. By the induction the function  $\rho(x, \partial\Omega')$  is piecewise linear in  $\Omega'$ . Since

$$\rho(x, \partial\Omega) = \rho(x, \partial\Omega') + d_* \quad \text{in } \Omega',$$

and  $\rho = \rho(x, \partial\Omega)$  is piecewise linear in  $\Omega_*$ , we see  $\rho$  is piecewise linear in  $\Omega$ . This shows that  $\rho \in A_+^0$  so that  $A_+^0$  is nonempty.

Note that  $\rho$  is the distance function  $d$  if and only if  $\Omega$  is a convex polygon. If  $\Omega$  is nonconvex,  $d$  is not piecewise linear, so  $d \in A_+^0$  if and only if  $\Omega$  is a convex polygon.

We conclude this remark by pointing out that there is a domain  $\Omega$  whose boundary is piecewise linear with infinite vertices such that  $\rho \in A_+^0$ . For example if we consider

$$\Omega = \{(x, y); |x| < 1, 1 > y > h(x)\}$$

$$h(x) = \begin{cases} \frac{1}{2^{\ell+1}} - |x - \frac{3}{2^{\ell+1}}| & \frac{1}{2^\ell} < x \leq \frac{1}{2^{\ell-1}}, \ell = 1, 2, \dots \\ 0 & x \leq 0 \end{cases}$$

then  $\rho \in A_+^0$ . See figure 1.

**2.2. Comparison Lemma of Hessian and Laplacian measure.** *Assume that  $n = 2$ . Then for  $\varphi \in A$ ,*

$$|\Delta\varphi| = |\nabla^2\varphi| \quad (\text{as measures}).$$

This is formally true since  $\nabla^2\varphi$  is rank one. Indeed, differentiating  $|\nabla\varphi|^2 = 1$  yields

$$\sum_{j=1}^n (\partial_i \partial_j \varphi) \partial_j \varphi = 0.$$

We shall justify this observation for general Hessian measure  $\nabla^2\varphi$  of  $\varphi \in A$ . We say that for  $\varphi \in X$  the rank of matrix of the Radon-Nikodym derivative

$$F(x) = \lim_{r \downarrow 0} \nabla^2\varphi(B_r(x)) / |\nabla^2\varphi|(B_r(x))$$

is the rank of  $\nabla^2\varphi$ , where  $B_r(x)$  denotes the closed ball of radius  $r$  centered at  $x \in \Omega$ . The rank of  $\nabla\varphi$  is defined for  $|\nabla^2\varphi|$ -almost every point  $x$  of  $\Omega$ . Since  $\nabla^2\varphi$  is absolutely continuous with respect to  $|\nabla^2\varphi|$ ,

$$|\Delta\varphi|(Z) = \int_Z |\text{trace } F| d\mu, \quad |\nabla^2\varphi|(Z) = \int_Z |F| d\mu$$

with  $\mu = |\nabla^2\varphi|$ , where  $|F|$  is the Hilbert-Schmidt norm of  $F$ , i.e.,  $|F|^2 = \sum_{ij} |F_{ij}|^2$ . If  $F$  is rank 1, then

$$|\text{trace } F| = |F|$$

so that  $|\Delta\varphi| = |\nabla^2\varphi|$ . Lemma 2.2 rigorously follows from the following two lemmas.

**Lemma [Al].** *If  $\varphi \in X$ , then the rank of the singular part of  $\nabla^2\varphi$  (i.e.  $\nabla^2\varphi - (\nabla^2\varphi)^{ab}$ ) is one.*

This is clear if  $\varphi \in Y$  because of representation of  $\nabla^2\varphi$ . Such property was proved for an important subset of the singular part by the authors [AG2] and conjectured there for all singular part. This difficult problem was solved by Alberti [Al]. We do not need to assume  $|\nabla\varphi| = 1$ .

**Lemma.** *If  $\varphi \in A$ , then rank of the absolutely continuous part  $(\nabla^2\varphi)^{ab}$  is less than or equal to  $n - 1$ .*

*Proof.* For  $\varphi \in X$  and  $j, 1 \leq j \leq n$  there is a representative of  $u = \nabla\varphi$  ( $\mathcal{L}^n$  - a.e.) so that the pointwise derivative  $\partial u/\partial x_j$  exists  $\mathcal{L}^n$ -a.e. and

$$\frac{\partial u}{\partial x_j}(x) = (\partial_j \nabla\varphi)^{ab}(x) \quad \Omega_0 \quad (\mathcal{L}^n - \text{a.e.}) \quad (1 \leq j \leq n),$$

where  $\mathcal{L}^n$  is the Lebesgues measure; see [AG3]. Note that the choice of  $u$  may depend on  $j$ .

Differentiating  $|\nabla\varphi|^2 = 1$  in the  $j$ -th direction yields

$$\begin{aligned} 0 &= 2 \sum_{i=1}^n \frac{\partial u_i(x)}{\partial x_j} \cdot u_i(x) = 2 \sum_{i=1}^n (\partial_j (\partial_i \varphi))^{ab}(x) u_i(x) \\ &= 2 \sum_{i=1}^n (\partial_j (\partial_i \varphi)(x))^{ab} (\partial_i \varphi)(x) \quad \text{for } \mathcal{L}^n\text{-a.e. } x \in \Omega_0, \end{aligned}$$

where  $u = (u_i)_{i=1}^n$ . Since  $|\nabla\varphi| \neq 0$  for a.e.  $x$ , this implies that  $(\nabla^2\varphi)^{ab}$  has a kernel for a.e.  $x$  so that rank of  $(\nabla^2\varphi)^{ab}$  is less than or equal to  $n - 1$ .  $\square$

**2.3. A key lemma.** For  $\varphi \in X$  assume that  $|\nabla\varphi| = 1$  in  $\Omega$  ( $\mathcal{L}^n$  - a.e.) and  $\varphi \geq 0$  in  $\Omega$  with  $\varphi = 0$  on  $\partial\Omega$ , i.e.,  $\varphi \in A_+$ . Then

$$(-\Delta\varphi)(\Omega) = \int_{\Omega} -\Delta\varphi \geq \mathcal{H}^{n-1}(\partial\Omega).$$

If  $\varphi(x) = \text{dist}(x, \partial\Omega)$  near  $\partial\Omega$ , then the equality holds.

This is easy if  $\varphi$  is regular so that  $\varphi$  is a distance function  $d(x, \partial\Omega)$  near  $\partial\Omega$ . Indeed, integrating by parts yields

$$\int_{\Omega} (-\Delta\varphi) d\mathcal{L}^n = - \int_{\partial\Omega} \frac{\partial\varphi}{\partial\nu} d\mathcal{H}^{n-1},$$

where  $\nu$  is the unit outward normal of  $\partial\Omega$ . Since  $\varphi = 0$  on  $\partial\Omega$  and  $|\Delta\varphi| = 1$  so that  $\nu = -\nabla\varphi/|\nabla\varphi|$ , the derivative  $-\partial\varphi/\partial\nu = 1$ . Thus, the equality

$$\int_{\Omega} (-\Delta\varphi) d\mathcal{L}^n = \mathcal{H}^{n-1}(\partial\Omega)$$

is proved.

*Proof.* 1. For  $\varphi \in X$  set

$$E_{\alpha} = \{x \in \Omega; \varphi(x) \geq \alpha\}, \alpha \geq 0.$$

Since  $\varphi \geq 0$  near  $\partial\Omega$ ,  $\bigcup_{\alpha>0} E_\alpha = \Omega$  and  $E_\alpha$  is decreasing in  $\alpha$ . Since  $\varphi$  is continuous,  $E_\alpha$  is closed for small  $\alpha \geq 0$ . Since  $\Delta\varphi$  is a finite Radon measure in  $\Omega$ , we obtain

$$\int_{\Omega} (-\Delta\varphi) = \lim_{\alpha \downarrow 0} \int_{E_\alpha} (-\Delta\varphi).$$

In particular

$$\int_{\Omega} (-\Delta\varphi) = \lim_{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \int_0^\varepsilon d\alpha \int_{E_\alpha} (-\Delta\varphi).$$

2. We shall prove the identity

$$\int_0^\varepsilon d\alpha \int_{E_\alpha} (-\Delta\varphi) = \mathcal{L}^n(\Omega - E_\varepsilon) \quad \text{for small } \varepsilon > 0. \quad (2.1)$$

Let  $\psi \in C^2(\Omega)$  with bounded gradient in  $\Omega$ . Since  $\varphi$  is Lipschitz near  $\partial\Omega$ , the co-area formula [G, S] yields

$$\int_{\Omega - E_\varepsilon} \nabla\psi \cdot \nabla\varphi d\mathcal{L}^n = \int_0^\varepsilon d\alpha \int_{L_\alpha} \left( \nabla\psi, \frac{\nabla\varphi}{|\nabla\varphi|} \right) d\mathcal{H}^{n-1} \quad \text{with } L_\alpha = \{x \in \Omega; \varphi(x) = \alpha\}. \quad (2.2)$$

Since  $\varphi$  is differentiable and  $|\nabla\varphi(x)| = 1$  for  $\mathcal{L}^n$ -a.e.  $x$  (near  $\partial\Omega$ ), by Fubini's theorem, for small ( $\mathcal{L}^1$ -)a.e.  $\alpha > 0$ ,

$$|\nabla\varphi(x_0)| = 1 \quad \text{for } \mathcal{H}^{n-1}\text{-a.e. } x_0 \text{ of } L_\alpha.$$

For such  $\alpha > 0$  we may assume that the level set  $L_\alpha$  is countably  $n-1$  rectifiable so that at  $x_0$  the approximate outer unit normal  $\nu_\alpha(x_0) = -\nabla\varphi(x_0)$ . We may also assume that  $E_\alpha$  is a set of finite perimeter so that the gradient  $\nabla\chi_{E_\alpha}$  of the characteristic function  $\chi_{E_\alpha}$  of  $E_\alpha$  is a finite Radon measure and that

$$\nabla\chi_{E_\alpha} = -\nu_\alpha \mathcal{H}^{n-1} \llcorner L_\alpha.$$

For these properties the reader is referred to the monographs [G, S]. For the above selected  $\alpha > 0$  we observe that

$$\int_{L_\alpha} \left( \nabla\psi \cdot \frac{\nabla\varphi}{|\nabla\varphi|} \right) d\mathcal{H}^{n-1} = \int_{\Omega} \nabla\psi \cdot \nabla\chi_{E_\alpha} = \int_{E_\alpha} (-\Delta\psi) d\mathcal{L}^n$$

by integration by parts. This together with (2.2) yields

$$\int_{\Omega - E_\varepsilon} \nabla\psi \cdot \nabla\varphi d\mathcal{L}^n = \int_0^\varepsilon d\alpha \int_{E_\alpha} (-\Delta\psi) d\mathcal{L}^n. \quad (2.3)$$

We would like to take  $\psi = \varphi$ . Since  $\varphi$  is not  $C^2$ , we need to approximate. Mollifying  $\varphi$  by a standard approximation as in [G, 1.17] we see that there is a sequence  $\psi_j \in C^2(\Omega)$  such that

$$\begin{aligned} \lim_{j \rightarrow \infty} \int_{\Omega} |\nabla \psi_j - \nabla \varphi| d\mathcal{L}^n &= 0, \quad \sup_{j \geq 1} \sup_{\Omega} |\nabla \psi_j| < \infty \\ \lim_{j \rightarrow \infty} \int_{\Omega} |\Delta \psi_j| d\mathcal{L}^n &= \int_{\Omega} |\Delta \varphi| \end{aligned}$$

since  $|\nabla \varphi|$  is bounded. In particular,  $\nabla \psi_j \rightarrow \nabla \varphi$  for  $\mathcal{L}^n$ -a.e.  $x$  by taking a subsequence if necessary. Moreover,

$$|\Delta \psi_j| \rightarrow |\Delta \varphi|, \quad \Delta \psi_j \rightarrow \Delta \varphi \quad \text{weakly as measures.}$$

We shall prove that  $\int_{E_\alpha} (-\Delta \varphi)$  is approximated by  $\int_{E_\alpha} (-\Delta \varphi_j) d\mathcal{L}^n$ . Since  $-\Delta \varphi$  is a nonnegative finite Radon measure and since  $L_{\alpha_1}$  and  $L_{\alpha_2}$  is disjoint for  $\alpha_1 \neq \alpha_2$ , we see

$$(-\Delta \varphi)(L_{\alpha'}) = 0$$

except at most countably many values of  $\alpha' > 0$ . For the above selected  $\alpha$

$$\mathcal{H}^{n-1}(\partial E_\alpha - L_\alpha) = 0$$

since  $E_\alpha$  is a set of finite perimeter [G].

Since  $-\Delta \varphi$  is absolutely continuous with respect to  $\mathcal{H}^{n-1}$  [G,S] we see

$$(-\Delta \varphi)(\partial E_\alpha - L_\alpha) = 0.$$

Since  $E_\alpha$  is closed so that  $\partial E_\alpha$  contains  $L_\alpha$ , we may assume

$$(-\Delta \varphi)(\partial E_\alpha) = \int_{\partial E_\alpha} (-\Delta \varphi) = 0$$

for the above selected  $\alpha$  by excluding the values of  $\alpha'$  with  $(-\Delta \varphi)(L_{\alpha'}) > 0$ . We now apply [G, Appendix A1] to get

$$\lim_{j \rightarrow \infty} \int_{E_\alpha} (-\Delta \psi_j) d\mathcal{L}^n = \int_{E_\alpha} (-\Delta \varphi).$$

Since  $\sup_{\Omega} |\nabla \psi_j|$  is bounded and  $\nabla \psi_j \rightarrow \nabla \psi$  a.e., we now obtain

$$\lim_{j \rightarrow \infty} \int_{\Omega - E_\varepsilon} (\nabla \psi_j \cdot \nabla \varphi) d\mathcal{L}^n = \int_{\Omega - E_\varepsilon} |\nabla \varphi|^2 d\mathcal{L}^n = \mathcal{L}^n(\Omega - E_\varepsilon).$$

Since

$$\sup_{\alpha, j} \int_{E_\alpha} |\Delta \psi_j|$$

is finite, the Lebesgue convergence theorem yields

$$\begin{aligned} \lim_{j \rightarrow \infty} \int_0^\varepsilon d\alpha \int_{E_\alpha} (-\Delta \psi_j) d\mathcal{L}^n &= \int_0^\varepsilon d\alpha \lim_{j \rightarrow \infty} \int_{E_\alpha} (-\Delta \psi_j) d\mathcal{L}^n \\ &= \int_0^\varepsilon d\alpha \int_{E_\alpha} (-\Delta \varphi). \end{aligned}$$

Plugging  $\psi = \psi_j$  in (2.3) and letting  $j \rightarrow \infty$  now yields (2.1).

3. Since  $|\nabla \varphi| = 1$  in  $\Omega$ , the set

$$\Omega - E_\varepsilon = \{x \in \Omega; 0 \leq \varphi(x) < \varepsilon\}$$

includes

$$F_\varepsilon = \{x \in \Omega; \text{dist}(x, \partial\Omega) < \varepsilon\}.$$

Since  $\partial\Omega$  is Lipschitz, we see

$$\lim_{\varepsilon \downarrow 0} \mathcal{L}^n(F_\varepsilon)/\varepsilon = \mathcal{H}^{n-1}(\partial\Omega).$$

By step 1 and 2,

$$\int_\Omega (-\Delta \varphi) = \lim_{\varepsilon \downarrow 0} \mathcal{L}^n(\Omega - E_\varepsilon)/\varepsilon \geq \lim_{\varepsilon \downarrow 0} \mathcal{L}^n(F_\varepsilon)/\varepsilon = \mathcal{H}^{n-1}(\partial\Omega).$$

The equality holds when  $\varphi(x) = d(x, \partial\Omega)$  near  $\partial\Omega$ .  $\square$

**2.4. Theorem on total variation of the Laplacian.** *Let  $\Omega$  be a bounded domain in  $R^n$  with Lipschitz boundary.*

(i)  $\int_\Omega |\Delta \varphi| \geq \mathcal{H}^{n-1}(\partial\Omega)$  for all  $\varphi \in A_+$

(ii) If  $\Omega$  is convex, the minimum of  $\int_\Omega |\Delta \varphi|$  over  $A_+$  is attained at  $\varphi_0(x) = \text{dist}(x, \partial\Omega)$  and the minimal value is  $\mathcal{H}^{n-1}(\partial\Omega)$ .

*Proof.* (i) This is a direct consequence of Lemma 2.3.

(ii) If  $\Omega$  is convex, then  $\varphi_0$  is a concave function in  $\Omega$ . This is easy; similar result is proved in

[C, p.53 Lemma]. In particular,  $-\Delta\varphi = |\Delta\varphi|$  as a measure. Thus, Lemma 2.3 yields

$$-\int_{\Omega} \Delta\varphi_0 = \mathcal{H}^{n-1}(\partial\Omega)$$

so  $\varphi_0$  is a minimizer of  $I$  over  $A_+$ .  $\square$

**2.5. Theorem on total variation of the Hessian.** *Let  $\Omega$  be a bounded domain in  $R^n$  with Lipschitz boundary.*

- (i)  $I(\varphi) = \int_{\Omega} |\nabla^2\varphi| dx \geq \mathcal{H}^{n-1}(\partial\Omega)$  for all  $\varphi \in A_+$  with  $n \leq 2$  or for all  $\varphi \in A_+^0$  for arbitrary  $n$ .
- (ii) If  $\Omega$  is convex, the minimum of  $I$  over  $A_+$  (with  $n \leq 2$ ) or over  $A_+^0$  is uniquely attained at  $\varphi_0(x) = \text{dist}(x, \partial\Omega)$  and

$$I(\varphi_0) = \mathcal{H}^{n-1}(\partial\Omega).$$

- (iii) If  $\Omega$  is not convex,  $I(\varphi) > \mathcal{H}^{n-1}(\partial\Omega)$  for all  $\varphi \in A_+$ , (with  $n \leq 2$ ) or for all  $\varphi \in A_+^0$ .

*Proof.* (i) By Lemma 2.2 we see

$$|\Delta\varphi| = |\nabla^2\varphi| \quad \text{for } \varphi \in A \quad \text{with } n = 2;$$

this equality is also true for  $\varphi \in A^0 = A \cap Y$  or for  $\varphi \in A$  with  $n = 1$ , since  $\nabla^2\varphi$  is rank one, Theorem 2.4 (i) now yields (i).

(ii) Since  $|\Delta\varphi| = |\nabla^2\varphi|$ , this follows from Theorem 2.4 (ii) except the uniqueness of the minimizer. Suppose that  $I(\varphi) = \mathcal{H}^{n-1}(\partial\Omega)$ , then

$$\int_{\Omega} |\nabla^2\varphi| = \int_{\Omega} |\Delta\varphi| = \int_{\Omega} -\Delta\varphi.$$

In particular  $-\Delta\varphi \geq 0$  as a measure. Since  $\nabla^2\varphi$  is rank 1, this means that  $\varphi$  is concave. A concave function  $\varphi$  in  $A_+$  is a viscosity solution of  $|\nabla\varphi| = 1$ ; see e.g. [L] so it is the distance function  $\varphi_0$ .

(iii) If there is  $\varphi$  such that  $I(\varphi) = \mathcal{H}^{n-1}(\partial\Omega)$ , we see  $\varphi$  is a concave distance function  $\varphi_0$  as in



(ii). However, such a distance function  $\varphi_0$  is concave in  $\Omega$  if and only if  $\Omega$  is convex. Thus, the strict inequality holds for a nonconvex domain.  $\square$

**Remark.** The uniqueness of minimizers of Theorem 2.4 (ii) would be true if  $\varphi \in A_+$  satisfying  $-\Delta\varphi \geq 0$  would be a viscosity solutions of  $|\nabla\varphi| = 1$  (without assuming that  $\varphi$  is concave.) However, we do not attempt to discuss this problem here.

### 3. COUNTEREXAMPLE

We shall construct a simply connected but nonconvex domain  $\Omega$  in  $\mathbb{R}^2$  such that the distance function does not minimize neither  $J^1$  nor  $I$  among the classes  $A, A_+, A^0, A_+^0$  defined in § 2.2.

**3.1. Choice of Domain.** For a positive constant  $\ell$  let  $D_\ell$  be a square of the form

$$D_\ell = \{(x, y); |y| < \ell, 0 < x < 2\ell\}.$$

Let  $D$  be a unit square of the form

$$D = \{(x, y); |y| < 1/2, -1 < x < 0\}.$$

Let  $\Omega_\ell$  be the interior of the union of  $\bar{D}$  and  $\bar{D}_\ell$ , where the bar denotes the closure. We shall always assume that  $\ell > 1/2$ . Clearly,  $\Omega$  is a bounded simply connected but nonconvex domain.

**3.2. Defects.** We consider three solutions of the eikonal equation

$$|\nabla\varphi| = 1 \quad \text{in } \Omega_\ell \quad \text{with } \varphi = 0 \quad \text{on } \partial\Omega_\ell$$

of the form:

$$\varphi_0(x) = \text{dist}(x, \partial\Omega_\ell)$$

$$\varphi_1(x) = \begin{cases} -\text{dist}(x, \partial D) & \text{for } x \in \bar{D} \\ \text{dist}(x, \partial D_\ell) & \text{for } x \in \bar{D}_\ell \end{cases}$$

$$\varphi_2(x) = \begin{cases} \text{dist}(x, \partial D) & \text{for } x \in \bar{D} \\ \text{dist}(x, \partial D_\ell) & \text{for } x \in \bar{D}_\ell. \end{cases}$$

The set of jump discontinuities of  $\nabla\varphi_1$  consists of

$$\begin{aligned} L_1 : y = \ell - x, 0 < x < 2\ell, & \quad L_2 : y = x - \ell, 0 < x < 2\ell \\ L_3 : y = -\frac{1}{2} - x, -1 < x < 0, & \quad L_4 : y = \frac{1}{2} + x, -1 < x < 0. \end{aligned}$$

The magnitude of jumps

$$j = |\nabla\varphi^+ - \nabla\varphi^-|$$

on each defects are  $\sqrt{2}$ . For later convenience we decompose defects of  $\varphi_1$ :

$$\begin{aligned} \Sigma_1^+ &= L_3 \cap R_1, & \Sigma_1^- &= L_4 \cap R_1 \\ S_1^+ &= L_4 \cap R_2, & S_1^- &= L_3 \cap R_2 \\ \Sigma_2^+ &= L_1 \cap R_3, & \Sigma_2^- &= L_2 \cap R_3 \\ S_2^+ &= L_1 \cap R_4, & S_2^- &= L_2 \cap R_4 \\ \Sigma_3^+ &= L_2 \cap R_5, & \Sigma_3^- &= L_1 \cap R_5 \end{aligned}$$

with

$$\begin{aligned} R_1 &= \{-1 < x < -1/2\}, & R_2 &= \{-1/2 < x < 0\} \\ R_3 &= \{0 < x < \ell - 1/2\}, & R_4 &= \{\ell - 1/2 < x < \ell\} \\ R_5 &= \{\ell < x < 2\ell\}. \end{aligned}$$

The defect of  $\varphi_0$  consists of

$$\begin{aligned} \Sigma_i^\pm \quad i = 1, 2, 3 \quad \text{and} \\ \Gamma_1 &= \{y = 0\} \cap R_2, & \Gamma_2 &= \{y = 0\} \cap (R_3 \cup R_4), \\ \Gamma^+ &= \{x^2 + (y - \frac{1}{2})^2 = (\ell - y)^2, y > 0\} \cap R_4 \end{aligned}$$

$$\Gamma^- = \{x^2 + (y + \frac{1}{2})^2 = (\ell + y)^2, y < 0\} \cap R_4.$$

see figure 2. The defect of  $\varphi_2$  consists of  $\Sigma_i^\pm$  ( $i = 1, 2, 3$ )  $S_i^\pm$  ( $i = 1, 2$ ) and

$$C = \{x = 0, |y| < 1/2\}$$

with jump  $j = 2$  on  $C$ .

**3.3. Computation of defect energy.** We shall estimate the difference  $J^\beta(\varphi_0) - J^\beta(\varphi_1)$ . By symmetry with respect to  $y = 0$  we observe that

$$\begin{aligned} J^\beta(\varphi_0) &= 2 \sum_{i=1}^3 J^\beta(\varphi_0, \Sigma_i^+) + \sum_{i=1}^2 J^\beta(\varphi_0, \Gamma_i) + 2J^\beta(\varphi_0, \Gamma^+) \\ J^\beta(\varphi_1) &= 2 \sum_{i=1}^3 J^\beta(\varphi_1, \Sigma_i^+) + 2 \sum_{i=1}^2 J^\beta(\varphi_1, S_i^+) \end{aligned}$$

where

$$J^\beta(\varphi, B) = \int_{\Sigma \cap B} j^\beta d\mathcal{H}^{n-1} \quad \text{with } j = |\nabla \varphi^+ - \nabla \varphi^-|.$$

Since jump  $j$  is the same both for  $\varphi_0$  and  $\varphi_1$  on  $\Sigma_i^+$ ,

$$J^\beta(\varphi_0, \Sigma_i^+) = J^\beta(\varphi_1, \Sigma_i^+) \quad i = 1, 2, 3.$$

Thus

$$J^\beta(\varphi_0) - J^\beta(\varphi_1) = \sum_{i=1}^2 J^\beta(\varphi_0, \Gamma_i) + 2J^\beta(\varphi_0, \Gamma^+) - 2 \sum_{i=1}^2 J^\beta(\varphi_1, S_i^+).$$

**Proposition.** (i)  $J^\beta(\varphi_0, \Gamma_1) = 2^{\beta-1}$ .

(ii)  $\lim_{\ell \rightarrow \infty} J^\beta(\varphi_0, \Gamma_2) = \int_0^\infty \frac{dx}{(x^2+1/4)^{\beta/2}} \geq \frac{1}{\beta-1} 2^{\beta-1}$  for  $\beta > 1$ .

For  $\beta \leq 1$ ,  $\lim_{\ell \rightarrow \infty} J^\beta(\varphi_0, \Gamma_2) = \infty$ .

(iii)  $J^\beta(\varphi_0, \Gamma^+) \rightarrow 2^{(\beta-1)/2}$  as  $\ell \rightarrow \infty$ .

(iv)  $J^\beta(\varphi_1, S_i^+) = 2^{(\beta-1)/2}$ ,  $i = 1, 2$ .

*Proof.* (i) Clearly,  $j = 2$  on  $\Gamma_1$ . Since the length of  $\Gamma_1$  is  $1/2$ ,  $J^\beta = 2^{\beta-1}$ .

(ii) Since the level curve of  $\varphi_0$  intersecting  $(x, 0)$  ( $0 < x < \ell$ ) is the circle centered  $(0, 1/2)$  for

$y > 0$ , the normal component of  $\nabla\varphi_0^+$  equals

$$\frac{1}{2} \frac{1}{(x^2 + 1/4)^{1/2}}.$$

Thus

$$J^\beta(\varphi_0, \Gamma_2) = \int_0^\ell \frac{dx}{(x^2 + 1/4)^{\beta/2}} \geq \int_0^\ell \frac{dx}{(x + 1/2)^\beta} = \frac{1}{\beta - 1} (x + \frac{1}{2})^{1-\beta} \Big|_0^\ell$$

and letting  $\ell \rightarrow \infty$  completes the proof.

(iii) Recall that  $\Gamma^+$  is a set of points whose distance to  $(0, 1/2)$  equals the distance to the line  $y = \ell$ . The curve  $\Gamma^+$  is a parabola given by

$$(2\ell - 1)y = -x^2 + \ell^2 - 1/4.$$

The jump

$$j(x) = 2 \frac{1}{(y'(x)^2 + 1)^{1/2}}$$

and the length element equals  $(y'(x)^2 + 1)^{1/2} dx$ . If  $p$  is the largest zero of  $y(x)$ , i.e.

$$p = (\ell^2 - 1/4)^{1/2},$$

then

$$J^\beta(\varphi_0, \Gamma^+) = \int_{\ell-1/2}^p \left( 2 \frac{1}{(y'(x)^2 + 1)^{1/2}} \right)^\beta ((y')^2 + 1)^{1/2} dx.$$

Notice that

$$y'(x)^2 + 1 = \left( \frac{-2x}{2\ell - 1} \right)^2 + 1, \quad \ell - 1/2 < x < \ell$$

$$\rightarrow 2 \quad \text{as } \ell \rightarrow \infty$$

$$p - (\ell - 1/2) = (\ell^2 - 1/4)^{1/2} - (\ell - 1/2) = \frac{1}{2} + \frac{-2/4}{(\ell^2 - 1/4)^{1/2} + \ell}$$

$$\rightarrow \frac{1}{2} \quad \text{as } \ell \rightarrow \infty.$$

We thus conclude that

$$J^\beta(\varphi_0, \Gamma^+) = \int_{\ell-1/2}^p 2^\beta ((y')^2 + 1)^{\frac{1-\beta}{2}} dx \rightarrow 2^\beta 2^{(1-\beta)/2} \cdot \frac{1}{2} = 2^{(\beta-1)/2} \quad \text{as } \ell \rightarrow \infty.$$

(iv) Since  $j = \sqrt{2}$  and the length of  $S_i^+$  equal  $1/\sqrt{2}$ , the result follows immediately.  $\square$

**3.4. Proposition.** *Let  $\Omega$  be the domain  $\Omega_t$  defined in §3.1. Let  $\varphi_0$  be the distance function of  $\partial\Omega$  and let  $\varphi_1$  and  $\varphi_2$  be solutions of the eikonal equation defined in §3.2.*

- (i)  $\lim_{t \rightarrow \infty} (J^\beta(\varphi_0) - J^\beta(\varphi_1)) = \infty$  for  $0 < \beta \leq 1$ .
- (ii)  $\lim_{t \rightarrow \infty} (J^\beta(\varphi_0) - J^\beta(\varphi_2)) = \infty$  for  $0 < \beta \leq 1$ .
- (iii)  $\lim_{t \rightarrow \infty} (J^\beta(\varphi_0) - J^\beta(\varphi_1)) > 0$  for all  $\beta > 0$ .
- (iv)  $\lim_{t \rightarrow \infty} (J^\beta(\varphi_0) - J^\beta(\varphi_2)) > 0$  for all  $0 < \beta < \beta_0$  with some  $\beta_0 > 4/3$ .

*Proof.* Applying the previous Proposition to the formula  $J^\beta(\varphi_0) - J^\beta(\varphi_1)$  yields

$$\begin{aligned} T &= \lim_{t \rightarrow \infty} (J^\beta(\varphi_0) - J^\beta(\varphi_1)) = 2^{\beta-1} + \lim_{t \rightarrow \infty} J^\beta(\varphi_0, \Gamma_2) + 2 \cdot 2^{(\beta-1)/2} \\ &\quad - 2(2^{(\beta-1)/2} + 2^{(\beta-1)/2}) \\ &= 2^{\beta-1} - 2 \cdot 2^{(\beta-1)/2} + \lim_{t \rightarrow \infty} J^\beta(\varphi_0, \Gamma_2). \end{aligned}$$

If  $\beta \leq 1$ , then this formula yield (i). Note that

$$J^\beta(\varphi_2) = J^\beta(\varphi_1) + J^\beta(\varphi_2, C) \quad \text{with} \quad C = \bar{D} \cap \bar{D}_t.$$

Since  $J^\beta(\varphi_2, C) = 1 \cdot 2^\beta$ , the proof of (ii) is now complete.

To show (iii) we may assume  $\beta > 1$  and use the estimate

$$\lim_{t \rightarrow \infty} J^\beta(\varphi_0, \Gamma_2) \geq \frac{1}{\beta-1} 2^{\beta-1}.$$

to get

$$T \geq 2 \cdot 2^{(\beta-1)/2} (f(\beta) - 1)$$

with  $f(\beta) = 2^{(\beta-3)/2} \frac{\beta}{\beta-1}$ . An elementary calculation shows that  $f$  takes the only minimum at  $\beta = \beta_1$  over all  $\beta > 0$ . The number  $\beta_1$  is the solution of

$$-\frac{1}{\beta_1-1} + \frac{1}{2}(\log 2)\beta_1 = 0 \quad \text{so that} \quad \beta_1 \geq 1.$$

Since  $\log 2 > 1/2$ ,  $(\beta_1 - 1)\beta_1 \leq 1/4$  so that  $\beta_1 - 1 \leq 1/2$ . Since  $\beta_1 \geq 1$  we now obtain

$$\begin{aligned} f(\beta) &\geq f(\beta_1) = 2^{(\beta_1-3)/2} \left(1 + \frac{1}{\beta_1 - 1}\right) \\ &> 2^{-1}(1 + 2). \end{aligned}$$

We thus conclude

$$J^\beta(\varphi_0) > J^\beta(\varphi_1) + 2 \cdot 2^{(\beta-1)/2} \cdot 1/2.$$

It remains to prove (iv). We may assume  $\beta > 1$ . Notice that  $J^\beta(\varphi_2, C) = 2^\beta$  to get

$$\begin{aligned} T &\geq \frac{\beta}{\beta-1} 2^{\beta-1} - 2 \cdot 2^{(\beta-1)/2} - 2^\beta \\ &= 2 \cdot 2^{(\beta-1)/2} \left( \left( \frac{1}{\beta-1} - 1 \right) 2^{(\beta-3)/2} - 1 \right). \end{aligned}$$

The right hand side is positive if  $\beta \leq 4/3$ .  $\square$

**3.5. Theorem.** *Assume that  $\Omega = \Omega_\ell$ .*

(i) ( $\beta \leq 1$ ) *For each  $M > 0$ , there is a constant  $\ell_0 = \ell_0(\beta)$  such that if  $\ell > \ell_0(\beta)$  then*

$$J^\beta(\varphi_0) \geq J^\beta(\varphi_2) + M \geq J^\beta(\varphi_1) + M.$$

*Moreover,  $\varphi_0$  does not minimize neither  $I$  nor  $J^\beta$  in  $A_+$ ,  $A$ .*

(ii) ( $\beta > 1$ ) *There is a constant  $\ell_1 = \ell_1(\beta)$  such that  $\varphi_0$  does not minimize  $J^\beta$  in  $A$  for  $\ell > \ell_1(\beta)$ . Moreover, if  $\beta \leq 4/3$  (or  $\beta < \beta_0$ ), then  $\varphi_0$  does not minimize  $J^\beta$  in  $A_+$ .*

*Proof.* (i) The first inequality follows from Proposition 3.4 (ii) and

$$J^\beta(\varphi_2) = J^\beta(\varphi_1) + 2^\beta$$

for sufficiently large  $\ell$ . Since  $\varphi_0 \in A_+ \subset A$ ,  $\varphi_1 \in A^0$  and  $\varphi_2 \in A_+^0$ , we now observe that

$$I(\varphi_0) \geq J^1(\varphi_0) > I(\varphi_2) = J^1(\varphi_2) > I(\varphi_1) = J^1(\varphi_1).$$

and

$$J^\beta(\varphi_0) > J^\beta(\varphi_2) > J^\beta(\varphi_1) \quad \text{for } \beta \leq 1.$$

(ii) Since  $\varphi_0 \in A_+$ ,  $\varphi_1 \in A^0$  and  $\varphi_2 \in A_+^0$ , Proposition 3.4 yields the desired conclusion.  $\square$

## REFERENCES

- [Al] G. Alberti, *Rank one property for derivatives of functions with bounded variation*, Proc. Roy. Soc. Edinburgh Sect. A **123** (1993), 239-274.
- [AG1] P. Aviles and Y. Giga, *A mathematical problem related to the physical theory of liquid crystal configurations*, Proc. of the Centre for Mathematical Analysis, Australian National University (eds. J. Hutchinson & L. Simon) **12** (1987), 1-16.
- [AG2] P. Aviles and Y. Giga, *Singularities and rank one properties of Hessian measures*, Duke Math. J. **58** (1989), 441-467.
- [AG3] P. Aviles and Y. Giga, *Variational integrals on mappings of bounded variation and their lower semicontinuity*, Arch. Rational Mech. Anal. **115** (1991), 201-255.
- [C] F.H. Clarke, *Optimization and Nonsmooth Analysis*, John Wiley and Sons, New York, Chichester, Birsbane, Toronto, Singapore (1983).
- [CIL] M.G. Crandall, H. Ishii and P.-L. Lions, *User's guide to viscosity solutions of second order partial differential equations*, Bull. AMS, **27** (1992), 1-67.
- [G] E. Giusti, *Minimal surfaces and functions of bounded variations*, Birkhauser, Boston-Basel-Stuttgart (1984).
- [KM1] R. Kohn and S. Muller, *Surface energy and microstructure in coherent phase transition*, Comm. Pure Appl. Math., **47** (1994), 405-435.
- [KM2] R. Kohn and S. Muller, *Relaxation and regularization of nonconvex variational problems*, Rend. Sem. Mat. Fis. Univ. Milano, **62** (1992), 89-113.
- [L] P.L. Lions, *Generalized solutions of Hamilton-Jacobi equations*, Pitmann Advanced Publ. Program (1982), Boston.
- [OG] M. Ortiz and G. Gioia, *The morphology and folding patterns of buckling driven thin-film filters*, J. Mech. Phy. of Solids, **42** (1994), 531-559.
- [SK] J. Sethna and M. Kleman, *Spheric domains in smectic liquid crystals*, Phys. Rev. A. vol. **26** (1982), 3037-3040.
- [S] L. Simon, *Lectures on geometric measure theory*, Proc. of the Centre for Mathematical Analysis,

Australian National University 3 (1983).



Figure 1

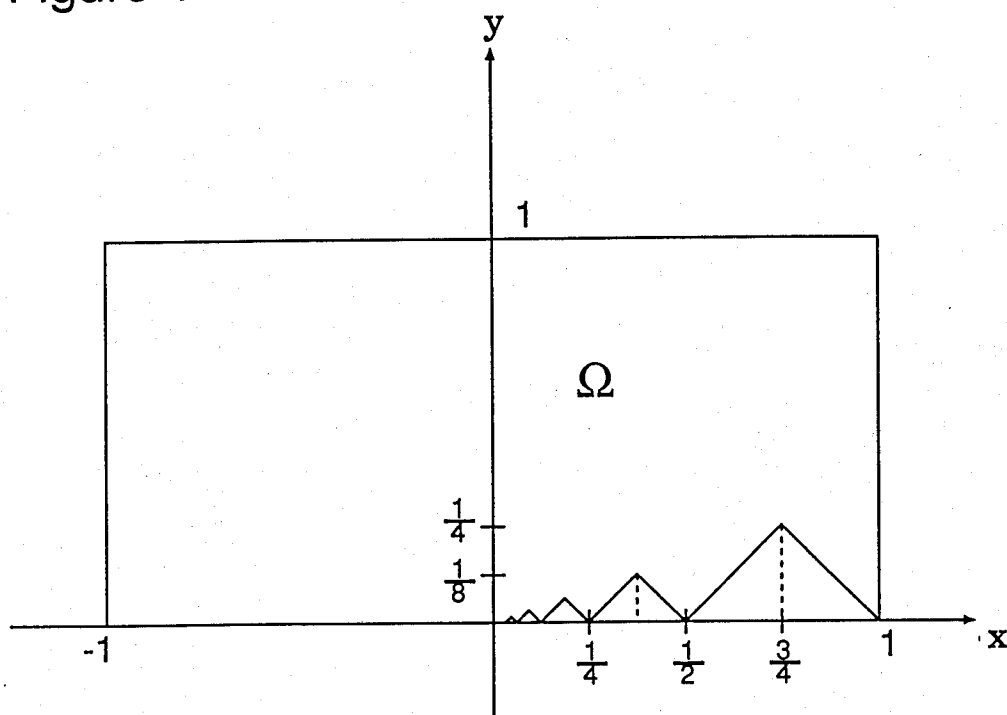
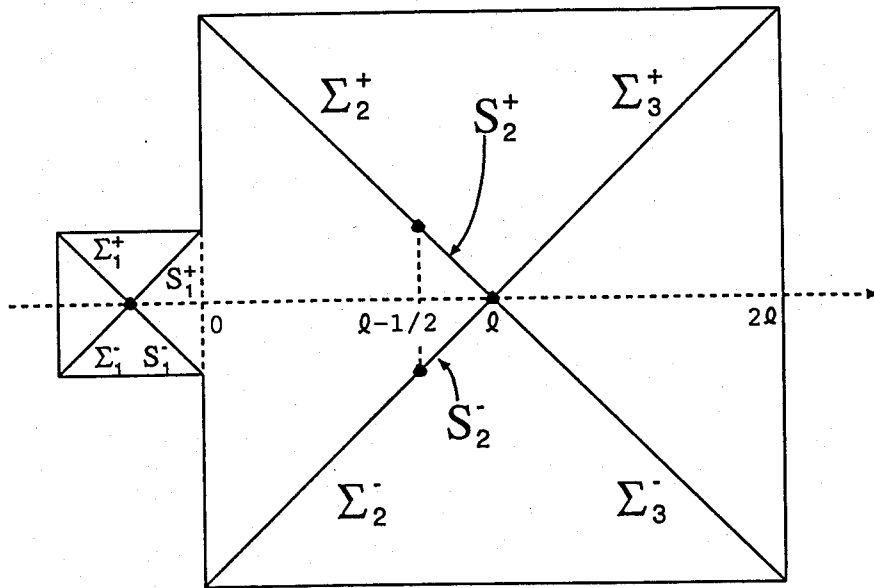
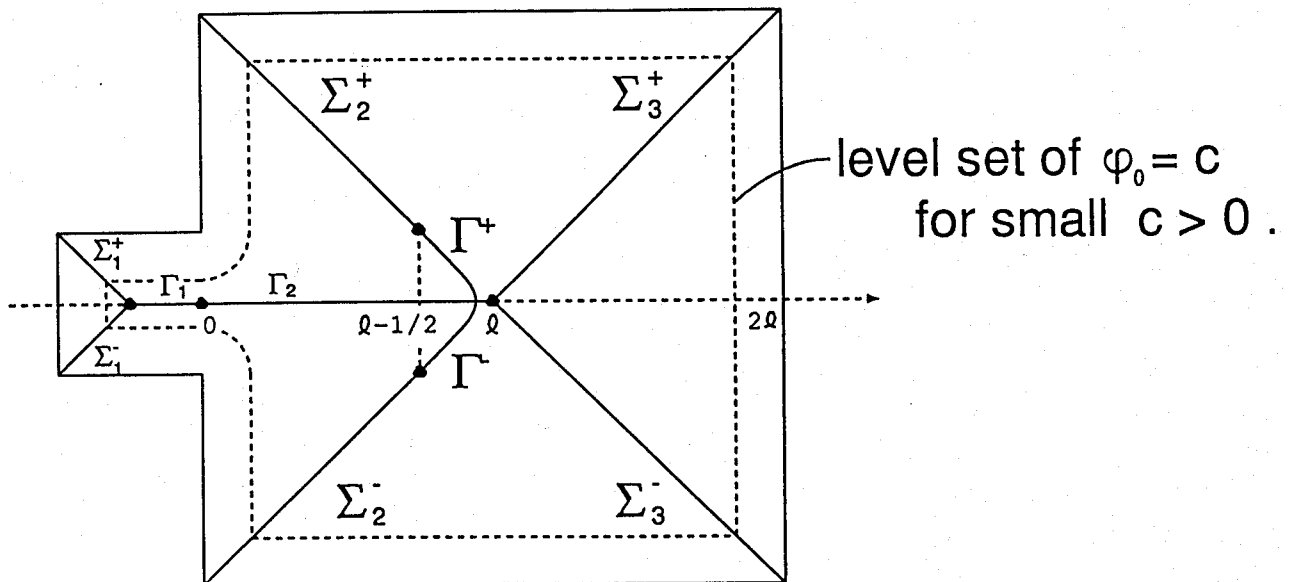


Figure 2



defects of  $\varphi_1$



defects of  $\varphi_0$