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$N = 4$ Supersymmetric Yang-Mills Theory on Orbifold- T^4/\mathbf{Z}_2

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Abstract

We derive the partition function of $N = 4$ supersymmetric Yang-Mills theory on orbifold- T^4/\mathbf{Z}_2 . In classical geometry, $K3$ surface is constructed from the orbifold- T^4/\mathbf{Z}_2 . Along the same way as the orbifold construction, we construct the partition function of $K3$ surface from orbifold- T^4/\mathbf{Z}_2 . The partition function is given by the product of the contribution of the untwisted sector of T^4/\mathbf{Z}_2 , and that of the twisted sector of T^4/\mathbf{Z}_2 i.e., $\mathcal{O}(-2)$ curve blow-up formula.

1 Introduction

Since the work of Vafa and Witten [14] have appeared, twisted $N = 4$ super Yang-Mills theory becomes a laboratory of testing S -duality conjecture, and it is now one of the concrete examples of the duality conjectures in various string theories. In [14] two major conjectures are made:

(i) The partition function of twisted $N = 4$ theory on 4-manifold is described as the summation of the Euler number of the ASD instanton moduli space with gauge group G .

(ii) Under S -duality transformation this partition function is transformed to the partition function with its dual gauge group \hat{G} . This duality is originated from Montonen-Olive duality [10].

For (i), it is confirmed that twisted $N = 4$ theory is described as a balanced topological field theory and that (i) is surely true if certain vanishing theorem holds [6], [3]. In the spirit of (i), matter coupled version is also considered and several analysis are done in [13].

Vafa and Witten derived the partition function satisfying both (i) and (ii) by physical method. (This partition function is called Vafa-Witten formula.) In their derivation, the most interesting part is the part which divides the evaluation of partition functions into two sectors: the bulk part and the cosmic string part. This observation arises from the mass perturbation of the theory. This perturbation is realized by a section of the canonical bundle of the base 4-manifold (complex surface) X , and there appears drastic difference between the zero-locus of K_X and the other part of X (they call these locus “cosmic string” and “bulk” respectively). Note that this observation was also applied to the topological Yang-Mills theory by Witten [15]. Next, they seem to identify the contribution from the cosmic string with the (generalized) blow-up formula speculated from the mathematical result of Yoshioka, and the one from the bulk with the Hilbert scheme of the points on the virtual manifold X_0 that has Euler number $3(\chi(X) + \sigma(X))$ and trivial canonical bundle. The principle of adding up these contributions is given by the celebrated S -duality conjecture.

With these results, it seems that the partition function of twisted $N = 4$ super Yang-Mills theory on any 4-manifold (at least algebraic) can be computed by combining Hilbert scheme of points of a certain surface with some blow-up formula. Then we try to interpret the orbifold construction of Kummer surface ($K3$ surface) from Abelian surface T^4 as the combination of these two contributions. In this construction, we first divide T^4 by \mathbf{Z}_2 . This process gives us the manifold S_0 with Euler number 8 with trivial canonical bundle and sixteen orbifold singularities. After that, we blow-up the sixteen orbifold singularities and obtain $K3$ surface. Then we speculate that we can realize these processes at the level of the partition function of twisted $N = 4$ SYM theory like the discussion given in [14]. But there are slight differences between our speculation and the original discussion in [14]. In the orbifold process, the blow-up operation should not change the canonical bundle mainly because T^4

and $K3$ have both trivial canonical bundles. So we must change the form of blow-up formula slightly in this case. Mathematically speaking, it is the difference between $\mathcal{O}(-1)$ curve blow-up and $\mathcal{O}(-2)$ curve blow-up. With these blow-up formulas, we can reconstruct the partition function of $K3$ surface from the partition function of S_0 , which is fundamentally equivalent to the Göttsche formula of S_0 . Coincidence with the well-known formula of $K3$ partition function in [14] is derived from the mathematical equalities among eta and theta functions. We think that our derivation is one example of orbifold construction of 4-manifold at the level of partition function of $N = 4$ SYM, and we expect many applications.

This paper is organized as follows.

In Section 2, we review the classical geometry of the orbifold construction of $K3$ from T^4/\mathbf{Z}_2 .

In section 3, we introduce Göttsche formula of the singular 4-fold $S_0 = T^4/\mathbf{Z}_2$ and construct the partition function of S_0 (we call this part untwisted sector).

In section 4, we propose the form of $\mathcal{O}(-2)$ curve blow-up formula (we call this part twisted sector) and derive the partition function of $K3$ by successive application of the $\mathcal{O}(-2)$ curve blow-up formula to the partition function of S_0 .

2 Classical Theory of T^4/\mathbf{Z}_2

In this section, we review the geometry of the manifold T^4/\mathbf{Z}_2 [4], [1].

First, we construct $K3$ surface from T^4 [1]. It is well known that compact complex Kähler 4-manifold with trivial canonical bundle $K_X = 0$ is only $K3$ or T^4 [4]. $T^4 = C^2/\mathbf{Z}^4$ is realized as a quotient space of the complex plane $(z_1, z_2) \in C^2$ divided by \mathbf{Z}^4 action, where \mathbf{Z}^4 action is $z_k \rightarrow z_k + 1, z_k \rightarrow z_k + i, k = 1, 2$. To obtain $K3$ surface, we first divide T^4 with \mathbf{Z}_2 action, where \mathbf{Z}_2 action is given by $(z_1, z_2) \rightarrow (-z_1, -z_2)$, (we denote $T^4/\mathbf{Z}_2 = S_0$ in the following). Then one gets sixteen fixed points of \mathbf{Z}_2 action

$$(0, 0), (0, \frac{1}{2}), (0, \frac{i}{2}), \dots, (\frac{1}{2} + \frac{i}{2}, \frac{1}{2} + \frac{i}{2}). \quad (2.1)$$

Let us consider the topology of T^4/\mathbf{Z}_2 . Since the \mathbf{Z}_2 action conserves the complex structure and leaves the Kähler form invariant, we expect S_0 to be a complex Kähler orbifold. Also, a moment's thought shows that any non-contractable loops of T^4 shrink to points after the \mathbf{Z}_2 identification. Thus $\pi_1(S_0) = 0$. Moreover, the holomorphic 2-form $dz_1 \wedge dz_2$ is invariant. Thus we obtain $K_{S_0} = 0$.

Since S_0 has sixteen \mathbf{Z}_2 orbifold singularities coming from the fixed points in (2.1), we have to resolve these singularities. As we explain below, we can resolve the orbifold singularities by $\mathcal{O}(-2)$ curve blow-up so that $K_X (= 0)$ is kept [1]. In general, blow-up is done by replacing a point with a curve. In this case, we replace an orbifold singularity with a $\mathcal{O}(-2)$ curve E . Let us first consider the blow-up of one \mathbf{Z}_2 singularity $\pi : \tilde{X} \rightarrow X$. Then we assume,

$$K_{\tilde{X}} = \pi^* K_X + n \cdot E. \quad (2.2)$$

Using the adjunction formula, we have

$$2g_E - 2 = (E + K_{\bar{X}}) \cdot E. \quad (2.3)$$

Since $g_E = 0$ and E is $\mathcal{O}(-2)$ curve, then we obtain,

$$K_{\bar{X}} \cdot E = (\pi^* K_X + nE) \cdot E = -2n = 0. \quad (2.4)$$

Thus we can conclude

$$K_{\bar{X}} = \pi^* K_X (= 0). \quad (2.5)$$

Hence $\mathcal{O}(-2)$ curve blow-up does not change the canonical bundle K_X . Using this $\mathcal{O}(-2)$ curve blow-up, we can resolve the sixteen orbifold singularities, and we get the smooth 4-manifold with trivial canonical bundle. This surface is nothing but the $K3$ surface.

These processes are equivalently replaced with the following processes [4]. First we remove 16 4-ball B_4 's from T^4 , which are neighborhood of 16 fixed points and divide $(T^4 - 16B_4)$ with \mathbf{Z}_2 action. Then the boundaries of $(T^4 - 16B_4)/\mathbf{Z}_2$ turn into sixteen RP^3 's. On the other hand, we have Ricci flat space T^*S^2 , that also has RP^3 as the boundary. Therefore, we can glue $(T^4 - 16B_4)/\mathbf{Z}_2$ with sixteen T^*S^2 's along sixteen RP^3 's smoothly. In this way, we obtain $K3$ surface, which is also a Ricci flat complex surface.

Next, we discuss the cohomology of T^4/\mathbf{Z}_2 . Especially we repeat the previous discussion with Hodge diamond.

First we write down the Hodge diamond of T^4 .

$$\begin{array}{ccccc}
& & h^{0,0} & & 1 \\
& h^{1,0} & & h^{0,1} & & 2 & 2 \\
h^{2,0} & & h^{1,1} & & h^{0,2} & = & 1 & 4 & 1 \\
& h^{2,1} & & h^{1,2} & & 2 & 2 \\
& & h^{2,2} & & & 1
\end{array} \quad (2.6)$$

The Hodge diamond of S_0 (modulo torsion) is obtained from the one of T^4 by picking up \mathbf{Z}_2 invariant forms of T^4 . We call this Hodge diamond untwisted sector.

$$\begin{array}{ccccc}
& & h^{0,0} & & 1 \\
& h^{1,0} & & h^{0,1} & & 0 & 0 \\
h^{2,0} & & h^{1,1} & & h^{0,2} & = & 1 & 4 & 1 \\
& h^{2,1} & & h^{1,2} & & 0 & 0 \\
& & h^{2,2} & & & 1
\end{array} \quad (2.7)$$

After sixteen blow-ups, we obtain additional sixteen $(1, 1)$ forms, and the Hodge diamond changes into,

$$\begin{array}{ccccc}
& & h^{0,0} & & 1 \\
& h^{1,0} & & h^{0,1} & & 0 & 0 \\
h^{2,0} & & h^{1,1} & & h^{0,2} & = & 1 & 20 & 1 \\
& h^{2,1} & & h^{1,2} & & 0 & 0 \\
& & h^{2,2} & & & 1 & .
\end{array} \quad (2.8)$$

Note that this Hodge diamond is the same as that of $K3$.

3 Untwisted Sector

In this section, we derive the partition function of the untwisted sector of T^4/\mathbf{Z}_2 [14], [11], [18], [8], [12].

3.1 Göttsche Formula

In this subsection, we introduce Göttsche formula of X [11],[8].

Following [8], we first introduce virtual Hodge number $e^{s,t}(X)$ and virtual Hodge polynomial $e(X; x, y)$ of X .

The virtual Hodge number is labeled by a pair of integers (s, t) and is given by

$$e^{s,t}(X) = \sum_k (-1)^k h^{s,t}(H_c^k(X, Q)), \quad (3.1)$$

and the virtual Hodge polynomial of X is given by

$$e(X; x, y) = \sum_{s,t=0}^2 e^{s,t}(X) x^s y^t. \quad (3.2)$$

Note that the virtual Hodge polynomial can be viewed as a convenient tool for computing the Hodge numbers of smooth projective varieties. Main reason of such characteristics comes from the following well-behavior:

- (i) When X is projective and smooth, $e(X; x, y)$ is the usual Hodge polynomial.
- (ii) If Y is a Zariski-closed subscheme of X , then

$$e(X; x, y) = e(Y; x, y) + e(X - Y; x, y). \quad (3.3)$$

- (iii) If $f : X \rightarrow Y$ is a Zariski-locally trivial bundle with fiber F , then

$$e(X; x, y) = e(Y; x, y) \cdot e(F; x, y). \quad (3.4)$$

For any algebraic surface X , Göttsche formula describes the generating function of the virtual Hodge polynomial of Hilbert scheme $X^{[n]}$:

$$\begin{aligned} Z^X(\tau; x, y) &\equiv \sum_{n=0}^{\infty} q^n e(X^{[n]}; x, y) \\ &= \prod_{m=1}^{\infty} \prod_{s,t=0}^2 (1 - x^{s+m-1} y^{t+m-1} q^m)^{(-1)^{s+t+1} h^{s,t}(X)}, \end{aligned} \quad (3.5)$$

where $h^{s,t}(X)$ stands for the Hodge numbers of X .

When we apply the formula to T^4 , we get

$$\begin{aligned}
& Z^{T^4}(\tau; \mathbf{x}, \mathbf{x}) \\
&= \sum_{n=0}^{\infty} q^n e((T^4)^{[n]}; \mathbf{x}, \mathbf{x}) \\
&= \prod_{m=1}^{\infty} \frac{(1 - x^{2m-1}q^m)^4 (1 - x^{2m+1}q^m)^4}{(1 - x^{2m-2}q^m)(1 - x^{2m}q^m)^6 (1 - x^{2m+2}q^m)}. \tag{3.6}
\end{aligned}$$

In the $K3$ case, we have,

$$\begin{aligned}
& Z^{K3}(\tau; \mathbf{x}, \mathbf{x}) \\
&= \sum_{n=0}^{\infty} q^n e((K3)^{[n]}; \mathbf{x}, \mathbf{x}) \\
&= \prod_{m=1}^{\infty} \frac{1}{(1 - x^{2m-2}q^m)(1 - x^{2m}q^m)^{22} (1 - x^{2m+2}q^m)}. \tag{3.7}
\end{aligned}$$

Now, we apply Göttsche formula to compute the contribution from the untwisted sector of T^4/\mathbf{Z}_2 .

$$\begin{aligned}
& Z^{S_0}(\tau; \mathbf{x}, \mathbf{x}) \\
&= Z^{T^4}(\tau; \mathbf{x}, \mathbf{x})|_{\mathbf{Z}_2 \text{ inv.}} \\
&= \prod_{m=1}^{\infty} \frac{(1 - x^{2m-1}q^m)^4 (1 - x^{2m+1}q^m)^4}{(1 - x^{2m-2}q^m)(1 - x^{2m}q^m)^6 (1 - x^{2m+2}q^m)} \Big|_{\mathbf{Z}_2 \text{ inv.}} \\
&= \prod_{m=1}^{\infty} \frac{1}{(1 - x^{2m-2}q^m)(1 - x^{2m}q^m)^6 (1 - x^{2m+2}q^m)} \tag{3.8}
\end{aligned}$$

Following [14], we define

$$\begin{aligned}
G(\tau) &\equiv q^{-\frac{1}{3}} Z^{S_0}(\tau; \mathbf{x} = \mathbf{y} = 1) \\
&= \frac{1}{\eta^8(\tau)}, \tag{3.9}
\end{aligned}$$

for later use. In the next subsection, we derive the partition function of the untwisted sector, which is fundamentally constructed from the formula (3.9).

3.2 Partition Function of the Untwisted Sector of S_0

In this subsection, we derive the partition function of S_0 of the untwisted sector. [8], [14], [18], [19].

General Structure of Vafa-Witten Conjecture

Following [8], we review the general structure of Vafa-Witten conjecture.

For $SO(3)$ theory with second Stiefel-Whitney class v on X , partition function is defined by,

$$Z_v^X(\tau) \equiv q^{-\frac{\chi(X)}{12}} \sum_k \chi(\mathcal{N}(v, k)) q^k, \quad (3.10)$$

where $\mathcal{N}(v, k)$ is the moduli space of anti-self-dual connections associated to $SO(3)$ -principal bundle with second Stiefel-Whitney class $v \in H^2(X, \mathbb{Z}_2)$ and fractional instanton number $k \in \mathbb{Z}/4$, and $\chi(X)$ is Euler number of X . For this partition function, Vafa and Witten conjectured

$$Z_v^X\left(-\frac{1}{\tau}\right) = 2^{-\frac{b_2(X)}{2}} \left(\frac{\tau}{i}\right)^{-\frac{\chi(X)}{2}} \cdot \sum_{u \in H^2(X, \mathbb{Z}_2)} (-1)^{u \cdot v} Z_u^X(\tau). \quad (3.11)$$

For later use, we introduce

$$\begin{aligned} Z_{SU(2)}^X(\tau) &\equiv \frac{1}{2} Z_0^X(\tau), \\ Z_{SO(3)}^X(\tau) &\equiv \sum_{u \in H^2(X, \mathbb{Z}_2)} Z_u^X(\tau). \end{aligned} \quad (3.12)$$

For these partition functions, the conjecture is reduced to the following formula:

$$Z_0^X\left(-\frac{1}{\tau}\right) = 2^{-\frac{b_2(X)}{2}} \left(\frac{\tau}{i}\right)^{-\frac{\chi(X)}{2}} Z_{SO(3)}^X(\tau). \quad (3.13)$$

$SO(3)$ Bundles on Spin Manifold X

In this part, we discuss the typical properties of $SO(3)$ bundles on the spin manifold X . Following [14], we first point out

$$k = n - \frac{v^2}{4}, \quad (3.14)$$

where n is second Chern class associated with rank two vector bundle, and k takes value on $k \in \mathbb{Z}/2$ according to the relation $v^2 = 0, 2 \pmod{4}$ on a spin manifold. As $H^2(X, \mathbb{Z})$ is $b_2(X)$ dimensional (and torsion free), v can take $2^{b_2(X)}$ values. In $SO(3)$ theory, we must sum over them.

There is no need to study separately $2^{b_2(X)}$ because X has a very large diffeomorphism group which permutes the possible values of v . One obvious diffeomorphism invariant of v is the value of v^2 modulo 4; if it is 0 we call even and if it is 2 we call odd. If v is odd, it is certainly non-zero, but for v even there is one more invariant: whether v is zero or not. It turns out that up to diffeomorphism, the invariants just stated are the only invariants of v . So on X there are really three partition functions to compute, namely the partition functions for $v = 0$ even but non-zero, and odd. We call these Z_0^X , Z_{even}^X , and Z_{odd}^X . Similarly, we write n_0^X , n_{even}^X , and n_{odd}^X for the number of values of v that are, respectively, trivial, even but non-trivial, and odd.

The number of each type is counted in Appendix A for $K3$ case. One can think of three types of partition functions:

$$\begin{aligned}
Z_0^X(\tau) &= q^{-\frac{\chi(X)}{12}} \sum_n \chi(\mathcal{N}(0, n)) q^n, \\
Z_{even}^X(\tau) &= q^{-\frac{\chi(X)}{12}} \sum_n \chi(\mathcal{N}(v_{even}, n)) q^n, \\
Z_{odd}^X(\tau) &= q^{-\frac{\chi(X)}{12}} \sum_n \chi(\mathcal{N}(v_{even}, n)) q^{n-\frac{1}{2}}.
\end{aligned} \tag{3.15}$$

Partition Function of the Untwisted Sector of S_0

In this part, we derive the partition function of the untwisted sector of S_0 as a concrete example. First we think of the moduli space of $X = T^4$ as an example of Corollary 1.7 in [18]. In [18], they treat the moduli space $\mathcal{M}_H(c_1, c_2)$ of rank two stable sheaves E with chern classes c_1, c_2 , but we want to treat the moduli space $\mathcal{N}(v, k)$ of $SO(3)$ or $SU(2)$ vector bundles with v, k . Thus following [8], we identify $\mathcal{M}_H(c_1, c_2) \equiv \mathcal{N}(v, k)$ with $c_1 = v \bmod 2$ and $k = c_2 - \frac{v^2}{4}$ in the following. According to [19], the moduli space of rank two stable sheaves E of V is given below. We introduce Mukai vector

$$V = ch(E) = 2 + c_1 + \frac{c_1^2 - 2c_2}{2}, \tag{3.16}$$

and its inner product

$$\begin{aligned}
\langle V^2 \rangle &= - \int_X \left(2 + c_1 + \frac{c_1^2 - 2c_2}{2} \right) \vee \left(2 + c_1 + \frac{c_1^2 - 2c_2}{2} \right) \\
&= 4c_2 - c_1^2.
\end{aligned} \tag{3.17}$$

Note that we use a symmetric bilinear form on $\oplus_i H^{2i}(X, Z)$:

$$\begin{aligned}
\langle x, y \rangle &= - \int_X (x \vee y) \\
&= \int_X (x_1 y_1 - x_0 y_2 - x_2 y_0),
\end{aligned} \tag{3.18}$$

where $x = x_0 + x_1 + x_2, x_i \in H^{2i}(X, Z)$ and $x \vee = x_0 - x_1 + x_2$.

We get the moduli space labeled by V

$$M_H^X(V) \cong \hat{X} \times (X)^{[\frac{\langle V^2 \rangle}{2}]} = \hat{X} \times (X)^{[2c_2 - \frac{c_1^2}{2}]}, \tag{3.19}$$

where H is some line bundle on X and \hat{X} is the dual of X [19].

Since S_0 has the trivial canonical bundle like $K3$ and T^4 , it doesn't have cosmic strings, which are given by zero locus of the section of the canonical bundle. Hence we assume the moduli space of V in S_0 case,

$$M_H(V) \cong (S_0)^{[\frac{\langle V^2 \rangle}{2}]} = (S_0)^{[2n - \frac{v^2}{2} - 1]}, \tag{3.20}$$

where we use the vanishing Picard group to drop the \hat{X} part and above identifications. Note that the shift in $[\cdot \cdot \cdot]$ by -1 comes from $\chi(\mathcal{O}_{S_0}) = \frac{1}{12} \int_{S_0} ((c_1(S_0))^2 + c_2(S_0)) = \frac{2}{3}$ instead of $\chi(\mathcal{O}_{T^4}) = 0$ (see also the definition of the moduli space in [18]). Using this we define the partition function of S_0 of the untwisted sector with even type,

$$\begin{aligned}
Z_{even}^{S_0}(\tau) &= q^{-\frac{2}{3}} \sum_{v^2 \equiv 0 \pmod{4}, n} e(M_H(V)) q^n \\
&= q^{-\frac{2}{3}} \sum_n e((S_0)^{[2n-1]}) q^n \\
&= q^{-\frac{1}{6}} \sum_m e((S_0)^{[m]}) \frac{(q^{\frac{1}{2}})^m - (-q^{\frac{1}{2}})^m}{2} \\
&= \frac{1}{2} G\left(\frac{\tau}{2}\right) - \frac{1}{2} e^{\frac{\pi i}{3}} G\left(\frac{\tau}{2} + \frac{1}{2}\right). \tag{3.21}
\end{aligned}$$

Similarly we define the partition function of S_0 of the untwisted sector for odd type,

$$\begin{aligned}
Z_{odd}^{S_0}(\tau) &= q^{-\frac{2}{3}} \sum_{v^2 \equiv 2 \pmod{4}, n} e(M_H(V)) q^{n-\frac{1}{2}} \\
&= q^{-\frac{2}{3}} \sum_n e((S_0)^{[2n-2]}) q^{n-\frac{1}{2}} \\
&= q^{-\frac{1}{6}} \sum_m e((S_0)^{[m]}) \frac{(q^{\frac{1}{2}})^m + (-q^{\frac{1}{2}})^m}{2} \\
&= \frac{1}{2} G\left(\frac{\tau}{2}\right) + \frac{1}{2} e^{\frac{\pi i}{3}} G\left(\frac{\tau}{2} + \frac{1}{2}\right). \tag{3.22}
\end{aligned}$$

For 0 type, we follow [14] and set,

$$Z_0^{S_0}(\tau) = NG(2\tau) + \frac{1}{2} G\left(\frac{\tau}{2}\right) - \frac{1}{2} e^{\frac{\pi i}{3}} G\left(\frac{\tau}{2} + \frac{1}{2}\right), \tag{3.23}$$

where N is not determined. Note that $G(2\tau)$ is obtained from $G\left(\frac{\tau}{2}\right)$ using the modular transformation $\tau \rightarrow -\frac{1}{\tau}$.

Strictly speaking, S_0 is not a spin manifold since the intersection matrix of S_0 is given by

$$\begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix} \oplus \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix} \oplus \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix}, \tag{3.24}$$

[20]. Hence it is possible for v^2 to take values 1 and 3 mod 4. This tells us that the general story in the previous part is not applicable directly. However, for our purpose to investigate the connection between $K3$ and S_0 , it is sufficient to forget the contribution from $v^2 = 1$ or 3 mod 4 (detailed explanation is done in 4.2).

In the rest of this part, we discuss the modular properties of these partition functions for later use.

For $\tau \rightarrow \tau + 1$,

$$\begin{aligned}
Z_{even}^{S_0}(\tau + 1) &= \frac{1}{2}G\left(\frac{\tau}{2} + \frac{1}{2}\right) - \frac{1}{2}e^{\frac{\pi i}{3}}e^{-\frac{\pi i}{12}8}G\left(\frac{\tau}{2}\right) \\
&= e^{-\frac{4\pi i}{3}}\left(\frac{1}{2}G\left(\frac{\tau}{2}\right) - \frac{1}{2}e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\right), \\
Z_{odd}^{S_0}(\tau + 1) &= \frac{1}{2}G\left(\frac{\tau}{2} + \frac{1}{2}\right) + \frac{1}{2}e^{\frac{\pi i}{3}}e^{-\frac{\pi i}{12}8}G\left(\frac{\tau}{2}\right) \\
&= -e^{-\frac{4\pi i}{3}}\left(\frac{1}{2}G\left(\frac{\tau}{2}\right) + \frac{1}{2}e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\right). \tag{3.25}
\end{aligned}$$

For τ to $-\frac{1}{\tau}$,

$$Z_{even}^{S_0}\left(-\frac{1}{\tau}\right) - Z_{odd}^{S_0}\left(-\frac{1}{\tau}\right) = -e^{\frac{\pi i}{3}}G\left(-\frac{1}{2\tau} + \frac{1}{2}\right) = \left(\frac{\tau}{i}\right)^{-4}\left(Z_{even}^{S_0}(\tau) - Z_{odd}^{S_0}(\tau)\right) \tag{3.26}$$

4 Twisted Sector

In this section, we introduce $\mathcal{O}(-2)$ blow-up formula and reconstruct $K3$ partition function [16], [8], [14].

4.1 $\mathcal{O}(-2)$ Curve Blow-up

In this subsection, we make an ansatz for the $\mathcal{O}(-2)$ blow-up formula for the partition function. We assume that the blow-up formula in this case is given by the following three functions:

$$\frac{\theta_2(\tau)}{\eta(\tau)^2}, \quad \frac{\theta_3(\tau)}{\eta(\tau)^2}, \quad \frac{\theta_4(\tau)}{\eta(\tau)^2}. \tag{4.1}$$

In [14], the blow-up formula for $\mathcal{O}(-1)$ curve is used to describe the effect of the existence of the cosmic string, that is the zero-locus of a section of the canonical divisor. This formula is proved by Li and Qin [8]. Of course, their results are inspired by the work of Yoshioka [16]. In the $\mathcal{O}(-1)$ curve case, the formula is given by,

$$\frac{\theta_2(2\tau)}{\eta(\tau)^2}, \quad \frac{\theta_3(2\tau)}{\eta(\tau)^2}. \tag{4.2}$$

Modification of $\theta(2\tau)$ into $\theta(\tau)$ can be estimated by following the proof of the $\mathcal{O}(-1)$ curve blow-up formula in [8], changing the condition of $E \cdot E = -1$ into $E \cdot E = -2$. But we don't have the explicit proof of this ansatz. We leave it to the future work. The reason for introducing $\theta_4(\tau)$ comes from the fact that $\theta_3(\tau)$ has both integral and half integral powers of q . We need $\theta_4(\tau)$ to project out half-integral (resp. integral) powers of q .

The factor $\frac{1}{\eta(\tau)^2}$ in (4.2) comes from the fact that we use Gieseker-Maruyama compactification of the moduli space. It appears as the effect of the additional

cohomology class E and of the fact that vector bundle we treat has two components. Hence it does not depend on the self intersection of the class E , and we do not change this factor.

With these ansatz, we define θ functions

$$\Theta_{\pm}(\tau) \equiv \frac{1}{2} \left(\theta_3^8(\tau) \pm \theta_4^8(\tau) \right), \quad (4.3)$$

$$\Theta_2(\tau) \equiv \frac{1}{2} \theta_2^8(\tau), \quad (4.4)$$

The reason for introducing 8th power of θ will be explained later, but as preparation of the discussion, we mention the structure of q -expansions of these functions. $\Theta_+(\tau)$, $\Theta_2(\tau)$ and $\Theta_-(\tau)$ have the following q -expansions

$$\begin{aligned} \Theta_+(\tau) &= 1 + q^n \text{ terms,} \\ \Theta_2(\tau) &= 2^7 q + q^n \text{ terms,} \\ \Theta_-(\tau) &= 2^4 q^{\frac{1}{2}} + q^{n+\frac{1}{2}} \text{ terms.} \end{aligned} \quad (4.5)$$

Note that $\Theta_-(\tau)$ is the only odd type. Using these θ functions, one can construct blow-up formulas by multiplying $\frac{1}{\eta^{16}(\tau)}$,

$$\begin{aligned} \tilde{Z}_+(\tau) &\equiv \frac{\Theta_+(\tau)}{\eta^{16}(\tau)}, \\ \tilde{Z}_-(\tau) &\equiv \frac{\Theta_-(\tau)}{\eta^{16}(\tau)}, \\ \tilde{Z}_2(\tau) &\equiv \frac{\Theta_2(\tau)}{\eta^{16}(\tau)}. \end{aligned} \quad (4.6)$$

These blow-up functions have the following modular properties

$$\begin{aligned} \begin{pmatrix} \tilde{Z}_+(\tau+1) \\ \tilde{Z}_-(\tau+1) \\ \tilde{Z}_2(\tau+1) \end{pmatrix} &= e^{-\frac{4\pi i}{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \tilde{Z}_+(\tau) \\ \tilde{Z}_-(\tau) \\ \tilde{Z}_2(\tau) \end{pmatrix}, \\ \begin{pmatrix} \tilde{Z}_+(-\frac{1}{\tau}) \\ \tilde{Z}_-(-\frac{1}{\tau}) \\ \tilde{Z}_2(-\frac{1}{\tau}) \end{pmatrix} &= \left(\sqrt{\frac{\tau}{i}} \right)^{-8} \begin{pmatrix} \frac{1}{2} & \frac{1}{2} & 1 \\ \frac{1}{2} & \frac{1}{2} & -1 \\ \frac{1}{2} & -\frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} \tilde{Z}_+(\tau) \\ \tilde{Z}_-(\tau) \\ \tilde{Z}_2(\tau) \end{pmatrix}. \end{aligned} \quad (4.7)$$

Thus, we can see that the above three functions form closed orbit under S -duality transformation. Note that if we require the length of the orbit to be three (in general, it becomes six), the only allowed power of θ is 8. Especially, we have to note that there is an equality:

$$\tilde{Z}_2\left(-\frac{1}{\tau}\right) \left(\tilde{Z}_+\left(-\frac{1}{\tau}\right) - \tilde{Z}_-\left(-\frac{1}{\tau}\right) \right) = \left(\sqrt{\frac{\tau}{i}} \right)^{-16} \tilde{Z}_2(\tau) \left(\tilde{Z}_+(\tau) - \tilde{Z}_-(\tau) \right). \quad (4.8)$$

4.2 $K3$ Partition Function

In this subsection, we reconstruct the $K3$ partition functions using the materials we have prepared [14].

From $Z_{even}^{S_0}(\tau)$ and $Z_{odd}^{S_0}(\tau)$ to $Z_{even}^{K3}(\tau)$ and $Z_{odd}^{K3}(\tau)$

In this part, we derive $Z_{even}^{K3}(\tau)$ and $Z_{odd}^{K3}(\tau)$ from $Z_{even}^{S_0}(\tau)$ and $Z_{odd}^{S_0}(\tau)$ using $\mathcal{O}(-2)$ curve blow-up formula (4.4). To determine $Z_{even}^{K3}(\tau)$ and $Z_{odd}^{K3}(\tau)$, we use the following two properties [14]. First property is that $Z_{even}^{K3}(\tau)$ and $Z_{odd}^{K3}(\tau)$ have q -expansions in the form of $(1 + q^n \text{ terms})$ and of $(q^{-\frac{1}{2}} + q^{\frac{1}{2}+n} \text{ terms})$ respectively. This follows from dimensional counting of the moduli space. The next property is that $Z_{even}^{K3}(\tau) - Z_{odd}^{K3}(\tau)$ is modular covariant under $\tau \rightarrow -\frac{1}{\tau}$ (see (b.10) in Appendix B). One can find

$$\begin{aligned} Z_{odd}^{S_0} \frac{1}{\eta^{32}(\tau)} &= q^{-\frac{1}{6}} q^{-\frac{32}{24}} + \dots \\ &= q^{-\frac{3}{2}} + \dots \end{aligned} \quad (4.9)$$

Thus to construct $Z_{even}^{K3}(\tau)$ from $Z_{odd}^{S_0}(\tau)$, one blow-ups only by $\tilde{Z}_2(\tau)\tilde{Z}_-(\tau)$. To construct $Z_{odd}^{K3}(\tau)$ from $Z_{odd}^{S_0}(\tau)$ one can blow-up by either $\tilde{Z}_2(\tau)\tilde{Z}_+(\tau)$ or $\tilde{Z}_2(\tau)\tilde{Z}_2(\tau)$ due to the first property. Due to the second property we choose the former. The same line of thought can be applied to the case of blowing up $Z_{even}^{S_0}(\tau)$ to $Z_{even}^{K3}(\tau)$ and to $Z_{odd}^{K3}(\tau)$.

Therefore, we are led to the following:

$$\begin{aligned} Z_{even}^{K3}(\tau) &= N \left(Z_{odd}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_-(\tau) + Z_{even}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_+(\tau) \right), \\ Z_{odd}^{K3}(\tau) &= N \left(Z_{odd}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_+(\tau) + Z_{even}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_-(\tau) \right). \end{aligned} \quad (4.10)$$

Note that $Z_{even}^{S_0}(\tau) - Z_{odd}^{S_0}(\tau)$ (3.26) and $\tilde{Z}_2(\tau)(\tilde{Z}_+(\tau) - \tilde{Z}_-(\tau))$ (4.8) are also modular covariant under $\tau \rightarrow -\frac{1}{\tau}$.

Here we remark two points concerning q -expansions. The first point is that all materials of Z_*^{K3} naturally have the form either q^n or $q^{n+\frac{1}{2}}$, since Z_*^{K3} itself has the form either q^n or $q^{n+\frac{1}{2}}$. This is why we only introduce $Z_{even}^{S_0}, Z_{odd}^{S_0}$ and $Z_0^{S_0}$ in 3.2 and \tilde{Z}_\pm and \tilde{Z}_2 in 4.1. The second point is that there should be one which has the form beginning with q , since the gap between $q^{-\frac{3}{2}}$ in (4.9) and 1 in Z_{even}^{K3} can only be filled with the combination q and $q^{\frac{1}{2}}$. Therefore, we introduce 8-th power of θ_2 in (4.4).

$Z_0^{K3}(\tau)$ and $Z_{SO(3)}^{K3}(\tau)$

Following the discussion in [14] for $Z_0^{K3}(\tau)$, we set the form of $Z_{even}^{K3}(\tau)$ as follows:

$$Z_0^{K3}(\tau) = \frac{N_1}{4} G(2\tau)\tilde{Z}_a(\tau) + N \left(Z_{even}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_-(\tau) + Z_{odd}^{S_0}(\tau)\tilde{Z}_2(\tau)\tilde{Z}_+(\tau) \right), \quad (4.11)$$

where $\tilde{Z}_a(\tau)$ is made of $\tilde{Z}_\pm(\tau), \tilde{Z}_2(\tau)$ so that $Z_0^{K3}(\tau)$ is even type.

As the last step, we have to determine the unknown factors N, N_1 and the form of $\tilde{Z}_a(\tau)$. Here it is convenient to introduce

$$\begin{aligned}\tilde{Z}_b(\tau) &\equiv 2\tilde{Z}_2(\tau) (\tilde{Z}_+(\tau) + \tilde{Z}_-(\tau)) = \frac{\theta_2^8(\tau)\theta_3^8(\tau)}{\eta^{32}(\tau)}, \\ \tilde{Z}_c(\tau) &\equiv 2\tilde{Z}_2(\tau) (\tilde{Z}_+(\tau) - \tilde{Z}_-(\tau)) = \frac{\theta_2^8(\tau)\theta_4^8(\tau)}{\eta^{32}(\tau)}.\end{aligned}\quad (4.12)$$

Using these functions, we have the following formulas:

$$\begin{aligned}Z_{even}^{K3}(\tau) &= \frac{N}{4} \left(G\left(\frac{\tau}{2}\right)\tilde{Z}_b(\tau) - e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\tilde{Z}_c(\tau) \right), \\ Z_{odd}^{K3}(\tau) &= \frac{N}{4} \left(G\left(\frac{\tau}{2}\right)\tilde{Z}_b(\tau) + e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\tilde{Z}_c(\tau) \right), \\ Z_0^{K3}(\tau) &= \frac{N_1}{4}G(2\tau)\tilde{Z}_a(\tau) + \frac{N}{4} \left(G\left(\frac{\tau}{2}\right)\tilde{Z}_b(\tau) - e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\tilde{Z}_c(\tau) \right).\end{aligned}\quad (4.13)$$

We also define $Z_{SO(3)}^K(\tau)$

$$\begin{aligned}Z_{SO(3)}^{K3}(\tau) &= Z_0^{K3}(\tau) + n_{even}Z_{even}^{K3}(\tau) + n_{odd}Z_{odd}^{K3}(\tau) \\ &= \frac{N_1}{4}G(2\tau)\tilde{Z}_a(\tau) + 2^{20}NG\left(\frac{\tau}{2}\right)\tilde{Z}_b(\tau) - 2^9Ne^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\tilde{Z}_c(\tau),\end{aligned}$$

where we use the numbers $n_0 = 1$, $n_{even} = \frac{2^{22}+2^{11}}{2} - 1$, $n_{odd} = \frac{2^{22}-2^{11}}{2}$ (see Appendix A).

Following [14], the $\tilde{Z}_a(\tau)$ should be transformed to $\tilde{Z}_b(\tau)$ under transformation $\tau \rightarrow -\frac{1}{\tau}$. Therefore, we obtain the form,

$$\tilde{Z}_a(\tau) = \frac{\theta_3^8(\tau)\theta_4^8(\tau)}{\eta^{32}(\tau)}.\quad (4.14)$$

Then we apply the S -duality property (3.13),

$$Z_0^{K3}\left(-\frac{1}{\tau}\right) = 2^{-11}\tau^{-12}Z_{SO(3)}^K(\tau),\quad (4.15)$$

that reduces to the equality:

$$Z_0^{K3}\left(-\frac{1}{\tau}\right) = \tau^{-12} \left(\frac{N_1}{4}2^4G\left(\frac{\tau}{2}\right)\tilde{Z}_b(\tau) + \frac{N}{4}2^{-4}G(2\tau)\tilde{Z}_a(\tau) - \frac{N}{4}e^{\frac{\pi i}{3}}G\left(\frac{\tau}{2} + \frac{1}{2}\right)\tilde{Z}_c(\tau) \right).$$

Thus we get

$$N_1 = 2^7N.$$

Now we compare the results with the ones in [14]. We use the identities of theta and eta functions,

$$\tilde{Z}_a(\tau) = \frac{1}{\eta^{16}(2\tau)},$$

$$\begin{aligned}\tilde{Z}_b(\tau) &= \frac{2^8}{\eta^{16}(\frac{\tau}{2})}, \\ \tilde{Z}_c(\tau) &= e^{\frac{8\pi i}{3}} \frac{2^8}{\eta^{16}(\frac{\tau}{2} + \frac{1}{2})},\end{aligned}$$

and rewrite Z_0^{K3} , Z_{even}^{K3} , Z_{odd}^K and $Z_{SO(3)}^{K3}$. After all, our results turn into,

$$\begin{aligned}Z_0^{K3}(\tau) &= 2^7 N \left(\frac{1}{4} \frac{1}{\eta^{24}(2\tau)} + \frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2})} + \frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2} + \frac{1}{2})} \right), \\ Z_{even}^{K3}(\tau) &= 2^7 N \left(\frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2})} + \frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2} + \frac{1}{2})} \right), \\ Z_{odd}^{K3}(\tau) &= 2^7 N \left(\frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2})} - \frac{1}{2} \frac{1}{\eta^{24}(\frac{\tau}{2} + \frac{1}{2})} \right), \\ Z_{SO(3)}^{K3}(\tau) &= 2^7 N \left(\frac{1}{4} \frac{1}{\eta^{24}(2\tau)} + 2^{21} \frac{1}{\eta^{24}(\frac{\tau}{2})} + 2^{10} \frac{1}{\eta^{24}(\frac{\tau}{2} + \frac{1}{2})} \right).\end{aligned}\quad (4.16)$$

If we set $2^7 N = 1$, the results are the same as Vafa-Witten's [14]. Thus we have connected T^4 with $K3$ by orbifold construction at the partition function level.

5 Conclusion

In this paper, we reconstruct the partition function of twisted $N = 4$ SYM theory on $K3$ surface using the orbifold construction T^4/\mathbf{Z}_2 at the partition function level. But, we used two major assumptions:

- (i) The contribution from the untwisted sector is fundamentally described by Göttsche formula applied naively to even part of $H^*(T^4, \mathbf{C})$.
- (ii) $\mathcal{O}(-2)$ curve blow-up formula is given by $\theta_i(\tau)/(\eta(\tau))^2$.

Justification of these two things remains to be pursued. Roughly speaking, we have neglected the behavior of vector bundles on sixteen singularity points of S_0 . Since singularity points have complex codimension 2 in base manifold, we expected that they do not affect the theory severely unlike cosmic strings.

Now, we discuss the future direction of our computations. In this paper, we were only considering the partition function for $SU(2)$ or $SO(3)$, but there are methods to obtain the partition functions for the other gauge groups, such as $SU(N)$ or $SU(N)/\mathbf{Z}_N$ [14], [7], [2]. Our computations can be also generalized to the other gauge groups. One more possible generalization is changing \mathbf{Z}_2 into other discrete groups: for example, A-D-E type discrete subgroup of $SU(2)$. One aim of this generalization is to construct various blow-up formulas for A-D-E singularities. We think that this direction is interesting with respect to the connection with the work of Nakajima on ALE spaces. We also hope our approach will contribute to the production of many examples of S -duality conjecture of twisted $N = 4$ SYM theory.

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A Counting n_0 , n_{even} and n_{odd}

In this section, we derive the numbers n_0, n_{even} and n_{odd} of $K3$, following the discussion in [14]. We denote $v \in H^2(K3, Z_2)$ by

$$v = \sum_{i=1}^{22} \epsilon_i v_i, \quad (\text{a.1})$$

where $v_i \in H^2(X, Z)$ $i = 1, \dots, 22$ and $\epsilon_i = 0$ or 1 . Thus v can take 2^{22} values.

For $u, v \in H^2(K3, Z_2)$, intersection form $u \cdot v$ on $K3$ is $H^{\oplus 3} \oplus (-E_8)^{\oplus 2}$, but for our purpose to count the number of 0, even and odd types, it is sufficient to think that intersection form is $H^{\oplus 11}$, due to the discussion in [14]. Thus we separate $v_i \in H^2(K3, Z)$ as

$$(v_{2j-1}, v_{2j}), \quad j = 1, \dots, 11, \quad (\text{a.2})$$

and we set that intersection form of each piece is H , that is

$$\begin{pmatrix} v_{2j-1} \cdot v_{2j-1} & v_{2j-1} \cdot v_{2j} \\ v_{2j} \cdot v_{2j-1} & v_{2j} \cdot v_{2j} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}. \quad (\text{a.3})$$

For each piece v_{2j-1}, v_{2j} , we define

$$V_j = \epsilon_{2j-1} v_{2j-1} + \epsilon_{2j} v_{2j}. \quad (\text{a.4})$$

If $V_j^2 = 0 \pmod{4}$, then $V_j = 0, v_{2j-1}$ or v_{2j} . If $V_j^2 = 2 \pmod{4}$, then $V_j = v_{2j-1} + v_{2j}$. Using this, we count the numbers n_0, n_{even} and n_{odd} of $K3$.

$$\begin{aligned} n_0 &= \# \text{ of } v = 0 \\ &= 1, \end{aligned} \quad (\text{a.5})$$

$$\begin{aligned} n_{even} &= \# \text{ of } v^2 = 0 \pmod{4} \text{ but nonzero} \\ &= (\# \text{ of even choices of } v_{2j-1} + v_{2j} \text{ in } H^{\oplus 11}) - 1 \\ &= {}_{11}C_0 3^{11} + \dots + {}_{11}C_{10} 3^1 - 1 \\ &= \frac{1}{2}(2^{22} + 2^{11}) - 1, \end{aligned} \quad (\text{a.6})$$

$$\begin{aligned} n_{odd} &= \# \text{ of } v^2 = 2 \pmod{4} \\ &= \# \text{ of odd choices of } v_{2j-1} + v_{2j} \text{ in } H^{\oplus 11} \\ &= {}_{11}C_1 3^{10} + \dots + {}_{11}C_{11} 3^0 \\ &= \frac{1}{2}(2^{22} - 2^{11}). \end{aligned} \quad (\text{a.7})$$

B Modular Properties of $Z_v^{K3}(\tau)$

In this section, we derive the useful modular properties of $Z_v^{K3}(\tau)$ on $K3$ from (3.11).

$$Z_v^{K3}\left(-\frac{1}{\tau}\right) = 2^{-11} \left(\frac{\tau}{i}\right)^{-12} \cdot \sum_{u \in H^2(K3, Z_2)} (-1)^{u \cdot v} Z_u^{K3}(\tau). \quad (\text{b.1})$$

From the definition of $Z_{SO(3)}^{K3}(\tau) \equiv \sum_{u \in H^2(K3, Z_2)} Z_u^{K3}(\tau)$, it is apparent

$$Z_0^{K3}\left(-\frac{1}{\tau}\right) = 2^{-11} \left(\frac{\tau}{i}\right)^{-12} Z_{SO(3)}^{K3}(\tau). \quad (\text{b.2})$$

For the modular properties of even or odd type, it is necessary to count the numbers $n_{even}^{u \cdot v=1}$, $n_{even}^{u \cdot v=0}$, $n_{odd}^{u \cdot v=1}$ and $n_{odd}^{u \cdot v=0}$ defined below. For even v , $Z_{odd}^{K3}(\tau)$ transform

$$\begin{aligned} & Z_{odd}^{K3}\left(-\frac{1}{\tau}\right) \\ = & 2^{-11} \left(\frac{\tau}{i}\right)^{-12} \cdot \left(Z_0^{K3}(\tau) + \sum_{u: \text{even}} (-1)^{u \cdot v} Z_{even}^{K3}(\tau) + \sum_{u: \text{odd}} (-1)^{u \cdot v} Z_{odd}^{K3}(\tau) \right) \\ = & 2^{-11} \left(\frac{\tau}{i}\right)^{-12} \cdot \left(Z_0^{K3}(\tau) + (-n_{even}^{u \cdot v=1} + n_{even}^{u \cdot v=0}) Z_{even}^{K3}(\tau) + (-n_{odd}^{u \cdot v=1} + n_{odd}^{u \cdot v=0}) Z_{odd}^{K3}(\tau) \right), \end{aligned} \quad (\text{b.3})$$

where we separate summation $\sum_{u: \text{even}}$ to $u \cdot v = 1 \pmod 2$ type and $u \cdot v = 0 \pmod 2$ type and introduce the numbers $n_{even}^{u \cdot v=1}$ and $n_{even}^{u \cdot v=0}$. The numbers $n_{odd}^{u \cdot v=1}$ and $n_{odd}^{u \cdot v=0}$ are defined from summation $\sum_{u: \text{odd}}$ similarly. We count the numbers $n_{even}^{u \cdot v=1}$, $n_{even}^{u \cdot v=0}$, $n_{odd}^{u \cdot v=1}$ and $n_{odd}^{u \cdot v=0}$ on $K3$. Without loss of generality, we take odd $v (= v_1 + v_2)$.

$$\begin{aligned} n_{even}^{u \cdot v=1} &= \# \text{ of } u \cdot v = 1 \pmod 2 \text{ in even } u\text{'s} \\ &= \# \text{ of } u = v_1 \text{ or } v_2 + \text{even } u \text{ combinations} \\ &= 2({}_{10}C_0 3^{10} + \cdots + {}_{10}C_{10} 3^0) \\ &= 2^{20} + 2^{10}, \end{aligned} \quad (\text{b.4})$$

$$\begin{aligned} n_{even}^{u \cdot v=0} &= \# \text{ of } u \cdot v = 0 \pmod 2 \text{ in even } u\text{'s} \\ &= (\# \text{ of } u = 0 + \text{even } u \text{ combinations nonzero}) \\ &\quad + (\# \text{ of } u = v_1 + v_2 + \text{even } u \text{ combinations}) \\ &= ({}_{10}C_0 3^{10} + \cdots + {}_{10}C_{10} 3^0) - 1 + ({}_{10}C_1 3^9 + \cdots + {}_{10}C_9 3^1) \\ &= 2^{20} - 1, \end{aligned} \quad (\text{b.5})$$

$$n_{odd}^{u \cdot v=1} = \# \text{ of } u \cdot v = 1 \pmod 2 \text{ in odd } u\text{'s}$$

$$\begin{aligned}
&= \# \text{ of } u = v_1 \text{ or } v_2 + \text{odd } u \text{ combinations} \\
&= 2({}_{10}C_1 3^9 + \cdots + {}_{10}C_9 3^1) \\
&= 2^{20} - 2^{10}, \tag{b.6}
\end{aligned}$$

$$\begin{aligned}
n_{odd}^{u \cdot v=0} &= \# \text{ of } u \cdot v = 0 \pmod{2} \text{ in odd } u\text{'s} \\
&= (\# \text{ of } u = 0 + \text{odd } u \text{ combinations}) \\
&\quad + (\# \text{ of } u = v_1 + v_2 + \text{odd } u \text{ combinations}) \\
&= ({}_{10}C_0 3^{10} + \cdots + {}_{10}C_{10} 3^0) + ({}_{10}C_1 3^9 + \cdots + {}_{10}C_9 3^1) \\
&= 2^{20}. \tag{b.7}
\end{aligned}$$

Finally we obtain the modular property of $Z_{odd}^{K3}(\tau)$

$$\begin{aligned}
&Z_{odd}^{K3} \left(-\frac{1}{\tau} \right) \\
&= 2^{-11} \left(\frac{\tau}{i} \right)^{-12} \cdot \left(Z_0^{K3}(\tau) + (-2^{10} - 1) Z_{even}^{K3}(\tau) + 2^{10} Z_{odd}^{K3}(\tau) \right). \tag{b.8}
\end{aligned}$$

In the same way as $Z_{odd}^{K3}(\tau)$, we obtain the modular property of $Z_{even}^{K3}(\tau)$

$$\begin{aligned}
&Z_{even}^{K3} \left(-\frac{1}{\tau} \right) \\
&= 2^{-11} \left(\frac{\tau}{i} \right)^{-12} \cdot \left(Z_0^{K3}(\tau) + (2^{10} - 1) Z_{even}^{K3}(\tau) - 2^{10} Z_{odd}^{K3}(\tau) \right). \tag{b.9}
\end{aligned}$$

Combining $Z_{odd}^{K3}(\tau)$ and $Z_{even}^{K3}(\tau)$, we also obtain

$$Z_{even}^{K3} \left(-\frac{1}{\tau} \right) - Z_{odd}^{K3} \left(-\frac{1}{\tau} \right) = \left(\frac{\tau}{i} \right)^{-12} \cdot \left(Z_{even}^{K3}(\tau) - Z_{odd}^{K3}(\tau) \right). \tag{b.10}$$

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