Identification of Magnetization Characteristics of Material from Measured Inductance Data

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This paper introduces a novel identification method of the magnetization properties of the core material for an inductor from measured inductance data. The proposed method allows us to obtain the magnetic hysteresis characteristics without a special measurement instrument for material BH properties. The proposed method determines the parameters included in the distribution function for the Preisach model from the measured L-I characteristics. It is shown that the identified initial magnetization curve and the minor loops are in good agreement with the original BH curves. The uniqueness for the identification problem is numerically verified.

Index Terms— Genetic Algorithm, Inductor, Inverse Problem, Magnetic Hysteresis, Preisach Model.

I. INTRODUCTION

FINITE element (FE) analysis of power inductors and reactors is based on the initial BH curves or hysteresis characteristics [e.g., 1]. It is, however, not always possible to obtain such data because a special measurement instrument for BH characteristics of magnetic material is necessary [2-4]. Moreover, it is sometimes hard to generate sufficiently strong magnetic field to measure the magnetic saturation properties. On the other hand, it is not difficult to measure the macroscopic magnetic field to measure the magnetic saturation properties. The edges in an inductor core even when a weak current is imposed. It can be measured using, e.g., a network analyzer. It is remarkable that magnetic saturation can occur near the coil edges in an inductor core even when a weak current is imposed. Moreover, the measured inductance is attributed to the distributed magnetic field that changes in time to draw major or minor loops in the inductor core. From these observations, a question arises; is it possible to identify the magnetization characteristics of the core material used in an inductor from the measured inductance data?

In this paper, we propose a method to identify the BH characteristics including magnetic hysteresis from the measured dependence of the inductance of an inductor on the amplitude and bias of the imposed current. The hysteresis property is assumed to be described by the Preisach model with a distribution function while the proposed method can be extended to other hysteresis models such as Play model. The parameters included in the distribution function are determined by the genetic algorithm (GA) so that the measured inductance properties are reproduced. We discuss the accuracy of the identified BH characteristics and uniqueness for the identification.

II. NUMERICAL METHOD

A. Preisach model

The Preisach model [5-7] is one of the hysteresis models which are widely used in FE analysis. In the Preisach model, the magnetic characteristic is represented by the superimposition of the basic hysteresis loops as shown in Fig. 1, where \( H_{u} \) and \( H_{v} \) are the upper and lower thresholds of the basic hysteresis loop. The magnetization \( M \) becomes positive when \( H \) increases to become greater than \( H_{u} \). Similarly, \( M \) becomes negative when \( H \) decreases to become smaller than \( H_{v} \). The Preisach model expresses \( M \) in the following form:

\[
M = \int_{H_{j}}^{H_{j+1}} \int_{H_{i}}^{H_{i+1}} K(H_{u}, H_{v})dH_{u}dH_{v} + M_{\text{min}} \tag{1}
\]

where \( K \) and \( M_{\text{min}} \) denote the distribution function and the minimum value of \( M \), respectively. Moreover, \( D \) is a triangular domain in the Preisach plane shown in Fig.2 which satisfies \(-H_{s} \leq H_{u} \leq H_{u} \leq H_{s}\). The interval \([\pm H_{u}, H_{s}]\) is divided into \( N_{p} \) subintervals \([H_{u}, H_{n-1}]\) \((0 < n < N_{p} - 1)\), the integral in (1) is partitioned into

\[
K'(i,j) = \int_{H_{j}}^{H_{j+1}} \int_{H_{i}}^{H_{i+1}} K(H_{u}, H_{v})dH_{u}dH_{v} \tag{2}
\]

Then, (1) is rewritten for numerical evaluation as

\[
M = \sum_{i=0}^{N_{p}-1} \sum_{j=0}^{i} K'(i,j)D(i,j) + M_{\text{min}} \tag{3}
\]

where \( D(i,j) \) is defined by

\[
D(i,j) = \begin{cases} 
1 : \text{positive magnetization} \\
0 : \text{negative magnetization} 
\end{cases} \tag{4}
\]

The initial magnetization curve as well as the major and minor BH loops can be expressed using the Preisach model. The AC magnetic property, shown in Fig. 3, of the material used for an inductor is computed from the slope of the minor loops governed by the Preisach model for the biased magnetic field \( H \) as shown in Fig. 4.
on the other hand, the DC inductance 

\[ L_{dc} = \frac{\int_{A} B \cdot J \, d\Omega}{I^2} \]  

where \( B \) is the magnetic flux density, \( A \) is the area, and \( J \) is the current density. For comparison, we compute the AC inductance \( L_{ac} \) using the peak value \( I_{peak} \) of an AC current from

\[ L = \int_{A} B \cdot J \, d\Omega \]  

On the other hand, the DC inductance \( L_{dc} \) is computed by the following procedure:

(i) For the magnitude of the magnetic flux density \( B \) at each finite element, the magnetic permeability is computed from Fig. 3.

(ii) Assuming a small input current, the linear FE analysis is performed using the magnetic permeability evaluated in step (ii), and \( L_{dc} \) is computed from (6).

III. IDENTIFICATION METHOD

In this study, the BH characteristic of the material used by an inductor is identified from the inductance values, \( L_{ac} \) and \( L_{dc} \) defined in II, using GA, the detail of which is described in [7]. The distribution function \( K(H_u, H_v) \) in (1) is here expressed by the Gaussian function with an elliptic exponent that includes three parameters \( A, \sigma_1, \sigma_2 \) as follows:

\[ K = \frac{A}{2\pi \sigma_1 \sigma_2} \exp \left[ -\frac{(H_u - H_l)^2}{2 \sigma_1^2} - \frac{(H_u + H_l)^2}{2 \sigma_2^2} \right] \]  

The magnetic flux density \( B \) is expressed by

\[ B = \mu_0 \mu_{out} H + M \]  

where \( \mu_0 \mu_{out} \) is the relative permeability in the fully saturated region and \( M \) is obtained by the Preisach model based on (7). We determine the unknown parameters \( A, \sigma_1, \sigma_2, \mu_{out} \) in (7) and (8) by solving the optimization problem defined by

\[ \min F = \frac{1}{N_i} \sum_{i=1}^{N_i} \left( \frac{L_{dc}(i) - L_{dc0}(i)}{L_{dc0}(i)} \right)^2 + \frac{1}{N_j} \sum_{j=1}^{N_j} \left( \frac{L_{ac}(j) - L_{ac0}(j)}{L_{ac0}(j)} \right)^2 \]  

where \( i, j \) run over the number of sampling points \( N_i, N_j \), and \( L_{dc0} \) and \( L_{ac0} \) represent the original and computed DC inductances, respectively, and the same convention is used for the AC inductance. The cost function defined in (9) is minimized by GA for a set of optimization variables \( A, \sigma_1, \sigma_2, \mu_{out} \) in which the fitness is evaluated as follows:

(i) The initial magnetization curve and AC magnetic property shown in Fig. 3 for the current parameters \( \sigma_1, \sigma_2, \mu_{out} \) are computed using the Preisach model based on (7) and (8).

(ii) The \( L_{ac} - L_{dc} \) and \( L_{dc} - L_{bias} \) characteristics are computed using the method mentioned in II.B.

(iii) The cost function \( F \) that measures the difference between the computed and original inductance values, \( L_{dc} - L_{dc0} \) and \( L_{ac} - L_{ac0} \), is evaluated by (9). In GA, the individuals have a different set of \( A, \sigma_1, \sigma_2, \mu_{out} \) that correspond to the \( L_{ac} - L_{peak} \) and \( L_{dc} - L_{bias} \) characteristics. They evolve to reproduce the original inductance characteristics.
IV. NUMERICAL RESULT

A. Preparation of Inductance Characteristics

We consider here the ferrite inductor shown in Fig. 5 for a numerical example. Table I summarizes the specific parameters of the inductor. The measured major BH curve of ferrite used for the magnetic core is shown in Fig. 6. The initial magnetization curve obtained from Fig. 6 is used to compute \( L_{ac} - I_{peak} \) characteristic. Moreover, we determine the parameters \((\mu, B_0, \sigma_1, \sigma_2)\) in (7) and (8) from the measured curve in Fig. 6 by the curve fitting that solves

\[
\min E = \frac{1}{N} \sum_{i} \left( \frac{B(i) - B_0(i)}{B(i)} \right)^2
\]

(10)

using GA, where, \(B_0\) and \(B\) denotes the magnetic flux density shown in Fig. 6 and that computed by the Preisach model. The resultant parameters are summarized in the right column of Table III. Based on this result, \( L_{dc} - I_{bias} \) characteristic is computed using the method mentioned in II.B.

B. Identified Results

The \( L_{ac} - I_{peak} \) and \( L_{dc} - I_{bias} \) characteristics computed from the measured data shown in Fig. 6 and fitted by the proposed method mentioned in III are plotted in Fig. 7 and Fig. 8. The parameters identified from the inductance characteristics are summarized in the right column of Table III. Although there are small discrepancies for small input currents in \( L_{ac} - I_{peak} \) characteristic, satisfactory coincidence is obtained. The convergence history of GA for solving problem (9) is shown in Fig. 9. We can find in Fig. 9 that the cost function \( F \) becomes sufficiently small after 1000 generations. The initial magnetization curve and the BH loops identified from the inductance characteristics shown in Figs. 7 and 8 are plotted with the original curves obtained from the measured data in Fig. 6 and the Preisach model in Fig. 10 and 11. It can be seen that both are in good agreement. Figure 12 shows the hysteresis loss with respect to the amplitude of magnetic flux density. The hysteresis loss is overestimated by the proposed method in comparison with the measured value. This discrepancy comes from the errors in the BH loops shown in Fig. 11; the ratio \( S_1/S_2 \), where \( S_1 \) and \( S_2 \) denote the area of the measured and identified minor loops, respectively, is almost the same as the ratio of the measured and computed hysteresis losses shown in Fig. 12.

C. Discussions

The above-mentioned method is based on the assumption that the BH characteristics can be described by the Preisach model with the distribution function (7). If the core material does not obey this model, the present identification would be unsuccessful. If the distribution function of the material can be described by an explicit function of parameters, the proposed method can be extended to identify the parameters. The extension of the proposed method to identification of implicit distribution functions is remained as an open question.

\begin{table}[h]
\centering
\caption{Inductor Parameters}
\begin{tabular}{|c|c|}
\hline
Parameter & Value  \\
\hline
\hline
\end{tabular}
\end{table}

\begin{table}[h]
\centering
\caption{Analysis Conditions}
\begin{tabular}{|c|c|}
\hline
Parameter & Value  \\
\hline
\hline
\end{tabular}
\end{table}

\begin{table}[h]
\centering
\caption{Preisach Parameters}
\begin{tabular}{|c|c|c|}
\hline
Parameter & Value  \\
\hline
\hline
\end{tabular}
\end{table}
D. Uniqueness of Solution
We consider here if the BH characteristics can be uniquely determined from the inductance characteristics. When assuming the distribution function of (7), the BH characteristics are described by the parameters \((A, \sigma_1, \sigma_2, \mu_0)\). To verify if these parameters are uniquely determined from the inductance characteristics, we performed additional five identifications starting from different initial populations generated randomly. The identified parameters and the final value of \(F\) are summarized in Table IV. It can be concluded from these results that there exists uniqueness to the solution at least for the data considered here. The theoretical proof of the uniqueness is remained for a future work.

V. Conclusion
We have proposed a new method to identify the magnetization characteristics of ferrite used for the magnetic core of an inductor from the inductance characteristics. This approach is useful because special measurement instruments for BH loops are not needed. The BH characteristics are computed from a measured major loop. The inductance characteristics are obtained from those BH data. Then, the BH characteristics have been identified from the inductance characteristics. The identification result is satisfactory. The uniqueness for the identification result is numerically verified. The extension of the proposed method to implicit distribution functions and theoretical proof of the uniqueness have been remained as open questions.

REFERENCES

| TABLE IV | IDENTIFIED PARAMETERS AND RESIDUAL |
|---|---|---|---|---|---|---|
| \(A\) | \#1 | \#2 | \#3 | \#4 | \#5 | \#6 |
| 0.82 | 0.82 | 0.81 | 0.83 | 0.82 | 0.81 |
| \(\sigma_1\) | 29.52 | 29.66 | 28.98 | 30.06 | 29.43 | 29.42 |
| \(\sigma_2\) | 21.60 | 21.61 | 21.56 | 21.65 | 21.45 | 21.53 |
| \(\mu_0\) | 173.96 | 173.94 | 173.27 | 173.83 | 173.95 | 174.30 |
| Final \(F\) | 3.99×10^3 | 3.99×10^3 | 3.72×10^3 | 4.39×10^3 | 3.89×10^3 | 3.86×10^3 |

Fig. 9. Convergence history of GA

Fig. 10. Identified initial magnetization curve

(a) \(H_m = 40\)
(b) \(H_m = 80\)
(c) \(H_m = 160\)
(d) \(H_m = 240\)

Fig. 11. Identified BH loops, where \(H_m\) denotes the amplitude of \(H\)

Fig. 12. Hysteresis loss with respect to amplitude of magnetic flux density