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Citation	Journal of Physics D: Applied Physics, 36(16), 1994-2000 https://doi.org/10.1088/0022-3727/36/16/311 Journal of Physics D: Applied Physics, 36(16), 2003, p.1994-2000
Issue Date	2003-08-21
Doc URL	http://hdl.handle.net/2115/766
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Type	article (author version)
File Information	sugawara-jpd-36.pdf



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Electron acceleration in gas by impulse electric field and its application to selective promotion of an electron–molecule reaction

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Abstract. We have simulated electron acceleration under short intense impulse electric fields of $0.1\text{--}1\text{ kV cm}^{-1}$ strength and $\sim 1\text{ ns}$ duration in N_2 at 65.5 Pa . The electron acceleration is completed in a practically collisionless condition when the impulse duration is comparable with or shorter than the electron mean free time. On the basis of this feature, we have proposed a way to produce a group of monoenergetic electrons. The electron energy after acceleration is controllable by the impulse strength and duration, and we have demonstrated selective promotion of the excitation producing $\text{N}_2(\text{A}^3\Sigma_u^+)$ metastable excited molecules. The selectivity has been quantified as the efficiency of the electron energy utilized for the target reaction, and efficiencies of 20–30% have been achieved in this simulation. These values are higher than those of electron swarms in DC equilibrium, in which the efficiency is at most 10%. The electron behaviour under impulse electric fields and the dependence of the energy efficiency on the impulse profile have been investigated.

PACS numbers: 52.80.-s Electric discharges

Published source: *Journal of Physics D: Applied Physics*, **36**(16), pp. 1994–2000 (2003)

1. Introduction

In applications of weakly ionised plasmas, high voltage pulses are expected as an advantageous means of dynamic plasma control utilizing non-equilibrium electron behaviour. Penetrante *et al* (1997) mentioned the use of fast-rising voltage pulses (< 10 ns risetime) as an effective way to increase the mean electron energy for high yields of electron–molecule reactions in decomposition of air pollutants. Macheret *et al* (2002) evaluated the input power required to sustain air plasmas by repetitive voltage pulses (0.1–5 ns duration), and reported its lower budgets compared with the DC mode. In investigations of dielectric barrier discharges for ultraviolet light sources, as well, the advantage of the use of short voltage pulses has been referred to as a key to obtain high radiation efficiencies (e.g., Vollkommer and Hitzschke (1997, 1998), Carman and Mildren (2003)).

Sugawara *et al* (1999) pointed out a theoretical possibility that electron acceleration under well-controlled impulse electric fields leads to selective promotion of a specific electron–molecule reaction. Electrons in gas can be accelerated by an impulse electric field in a practically collisionless condition when the impulse duration is comparable with or shorter than the electron mean free time. The electron energy after acceleration is controllable by the impulse field strength and duration, and monoenergetic electrons are obtained when the electron temperature before acceleration is low enough. Let us call this technique the impulse-field electron acceleration (IFEA) method. Impulse fields required for the IFEA method are, e.g., in the order of ~ 1 kV cm $^{-1}$ strength and ~ 1 ns duration at gas pressures of around 1 Torr. Recent progress in pulsed-power apparatus is achieving a satisfactory capability for this requirement. Impulse voltages of 10–100 kV amplitude and ~ 10 ns duration are available nowadays. Anikin *et al* (1998) observed ionisation waves in N $_2$ induced by impulse voltages of 15.5 kV amplitude and 25 ns duration in a pressure range of 0.1–16 Torr. Related to this work, Pancheshnyi *et al* (1999) measured the population of excited N $_2$ particles and Starikovskaia and Starikovskii (2001) modelled the electron energy distribution function. Katsuki *et al* (2001) reported on streamer discharges in atmospheric-pressure air induced by 10 ns impulses with voltages up to 140 kV. Although there may currently be technical difficulties in precise control of the impulse profiles, investigation of the electron behaviour under a short impulse electric field would be informative for applications of impulse-driven plasmas.

In this paper, we simulate electron behaviour during the IFEA and subsequent electron–molecule collisions. As the target reaction to be promoted, we take the excitation to produce N $_2(A^3\Sigma_u^+)$ metastable excited molecules. The efficiency of the electron energy utilized for the target reaction is compared with those of electron swarms in drift equilibrium under DC electric fields and monoenergetic electron beams. We investigate the dependence of the efficiency on the impulse profile.

2. Theory

In the beginning, we consider collision-free electron acceleration. After that, electron-molecule collisions during the acceleration is taken into account.

Letting $\mathbf{v} = (v_x, v_y, v_z)$ be the electron velocity and $E(t)$ be the electric field applied along the $-x$ -direction during a period $0 \leq t \leq T$, the velocity increase $\Delta v_x = v_x(T) - v_x(0)$ of an electron given by $E(t)$ is written as

$$\Delta v_x = \int_0^T \frac{e}{m} E(t) dt \equiv \frac{e}{m} E_{\max} T_{\text{eff}}. \quad (1)$$

Here, e and m are the electronic charge and mass. E_{\max} and T_{eff} are parameters introduced to characterize the strength and duration of $E(t)$; E_{\max} is the maximum value of $E(t)$, and T_{eff} is the effective impulse duration defined as $T_{\text{eff}} = \int_0^T E(t) dt / E_{\max}$. Δv_x is common for all of the electrons under $E(t)$. Thus the electron velocity distribution function (EVDF) $f(v_x, v_y, v_z, t)$ shifts parallelly in the $+v_x$ -direction in velocity space by the acceleration:

$$f(v_x, v_y, v_z, T) = f(v_x - \Delta v_x, v_y, v_z, 0). \quad (2)$$

When the electron temperature before acceleration is low, the electrons are concentrated around $\mathbf{v} = 0$ and the electrons after acceleration are also concentrated around $(\Delta v_x, 0, 0)$; monoenergetic electrons are obtained through this collective shift. With an appropriate choice of E_{\max} and T_{eff} , one can feed electrons into a specific velocity (energy) range in which target collisions occur efficiently. Especially in the case in which the target collision cross-section, q_{target} , lays alone in the range, the target process can be promoted selectively. Figure 1 illustrates an aspect of the IFEA aiming at q_{target} in such an energy range that is free from another cross-section q' which impairs the selectivity. Here, $f(v_x, t)$ in figure 1 is the projection of $f(v_x, v_y, v_z, t)$ to the v_x axis of velocity space:

$$f(v_x, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v_x, v_y, v_z, t) dv_y dv_z. \quad (3)$$

In the presence of gas molecules, some of the electrons drop out of the collective acceleration because of scattering. The ratio of the dropout electrons, i.e. the collision probability P_{coll} during $0 \leq t \leq T$, is given as

$$P_{\text{coll}} = 1 - \exp\left(-\int_0^T N q_{\text{total}}(v(t)) v(t) dt\right), \quad (4)$$

where q_{total} is the total collision cross-section as a function of $v = |\mathbf{v}|$ and N is the number density of gas molecules. This evaluation can be simplified as follows when the electron mean free time τ or the electron mean free path λ is known:

$$P_{\text{coll}} \approx 1 - \exp(-T/\tau) \quad \text{or} \quad P_{\text{coll}} \approx 1 - \exp(-L/\lambda), \quad (5)$$

where L is the flight distance of an electron during the acceleration. P_{coll} is satisfactorily small when IFEA is completed within $T < \tau$ or $L < \lambda$.

Here, a small N is also preferable for reducing P_{coll} . However, decrease in the collision frequency of the target reaction after acceleration and increase in the electron

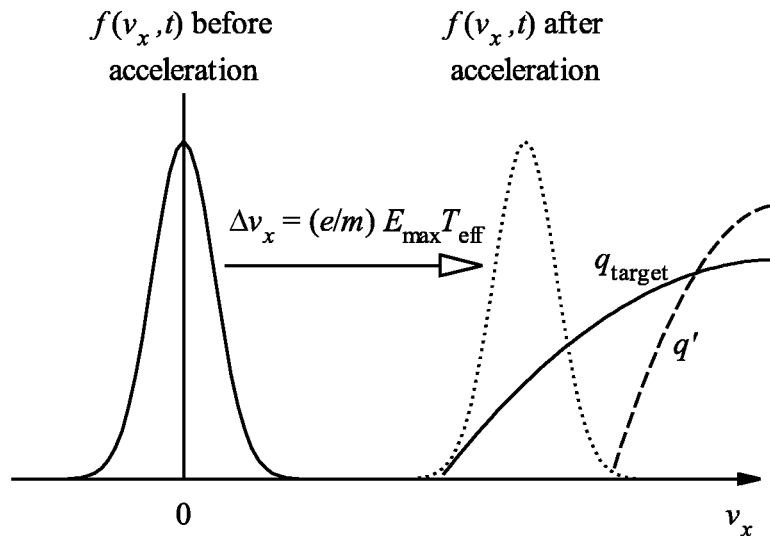


Figure 1. Schematic of electron acceleration by impulse electric field. $f(v_x, t)$, the distribution function of v_x ; q_{target} , the electron collision cross-section of the target process; and q' , another cross-section to be avoided so as not to impair the selectivity.

diffusivity, both of which lower the volume reaction rate of the target process, should be taken into account when one makes a practical choice of N .

In order to achieve the electron acceleration over the onset energy ϵ_{target} of the target reaction, the impulse field must be intense inversely proportional to the flight time or distance. For an electron with velocity $\mathbf{v} = (\Delta v_x, 0, 0)$, which is the representative (central) electron velocity after acceleration, the requirement on the kinetic energy $\epsilon = \frac{1}{2}m|\mathbf{v}|^2$ is given as follows from equation (1):

$$\epsilon = \frac{e^2}{2m}(E_{\text{max}}T_{\text{eff}})^2 > \epsilon_{\text{target}}. \quad (6)$$

3. Simulation conditions

3.1. Target reaction and energy efficiency

We simulate the IFEA in N_2 . Excitation to produce $\text{N}_2(\text{A}^3\Sigma_u^+)$ metastable excited molecules is chosen as the target process. This species is known for its long lifetime of over 10 s (Radzig and Smirnov 1985), and can be used as an energy source for moderate plasma–surface reactions; the excitation energy ϵ_{meta} is about 6.2 eV. Figure 2 shows the electron collision cross-sections of N_2 taken from Ohmori *et al* (1988). The target cross-section is indicated as q_{meta} . Its onset energy is the lowest among the electronic-state excitations of N_2 under consideration, and the second lowest onset energy is 7.4 eV. Let us call the energy range 6.2–7.4 eV ‘the target energy range’, in which q_{meta} lies alone, except for q_{el} . The presence of such a range contributes significantly to the selective promotion of the target reaction.

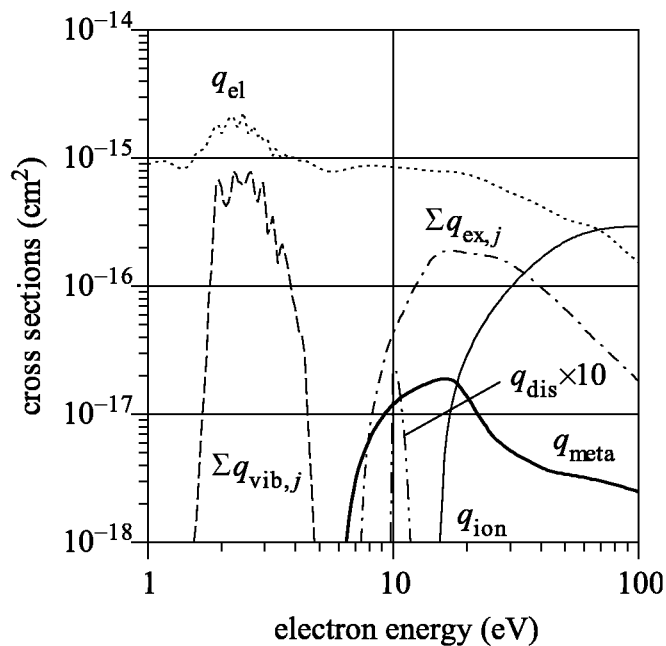


Figure 2. The electron collision cross-sections of N_2 (Ohmori *et al* 1988). q_{el} , elastic collision; $q_{vib,j}$, vibrational excitation; q_{meta} , excitation to metastable state $A^3\Sigma_u^+$; $q_{ex,j}$, the other electronic excitations; q_{dis} , dissociation; and q_{ion} , ionisation.

In order to evaluate the selectivity of the IFEA method, the energy efficiency η_{meta} is defined as being the ratio of the electron energy utilized for the target reaction to the total energy consumption by the collisions. The efficiency is compared between the following three modes: electrons under the IFEA ($\eta_{i,meta}$), electron swarms in DC drift equilibrium ($\eta_{e,meta}$), and monoenergetic electron beams ($\eta_{b,meta}$). Details of their conditions are described in following subsections.

3.2. Impulse mode

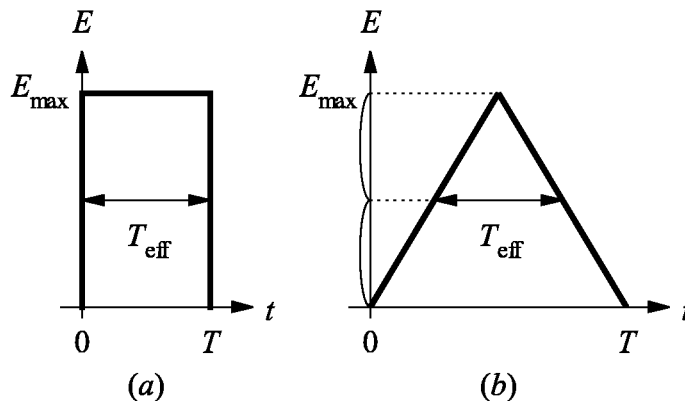
A Maxwellian distribution is assumed for the initial EVDF. The initial mean electron energy ϵ_{ini} is 1 eV in the central test conditions, and IFEA at $\epsilon_{ini} = 0.1$ eV is demonstrated additionally for comparison. N is set to be $1.77 \times 10^{16} \text{ cm}^{-3}$, which corresponds to a gas pressure of 66.5 Pa (0.5 Torr) at 0 °C. The electron mean free path $\lambda \approx (Nq_{total})^{-1}$ estimated at this N is about 0.06 cm, roughly approximating $q_{total} \approx 10^{-15} \text{ cm}^2$. The flight distance L along the x -direction which an electron needs to obtain ϵ_{meta} is set to be shorter than λ ; e.g. $L = \epsilon_{meta}/(eE)$ is about 0.031 cm at $E = 200 \text{ V cm}^{-1}$. Estimating from P_{coll} in equation (5), about half of the initial electrons survive collisionless up to this L .

Figure 3 shows the waveforms of $E(t)$ assumed here. The rectangular models ideal impulses, and the triangular one simulates those with risetime. T_{eff} for these waveforms is equivalent to their full-width at half-maximum; $T_{eff} = T$ for the former, and $T_{eff} = \frac{1}{2}T$ for the latter. The values of E_{max} and T_{eff} we examine are listed in table 1 together with

Table 1. Waveform parameters and maximum energy efficiency.

Case	Waveform	E_{\max} (V cm ⁻¹)	T_{eff} (ns)	ϵ_{ini} (eV)	$\max(\eta_{i,\text{meta}})$ (%)
A	Rectangle	200	0.05–0.75	1.0	20.2
B	Rectangle	100	0.10–1.50	1.0	18.3
C	Triangle	100	0.10–1.50	1.0	17.7
D	Rectangle	1000	0.01–0.15	1.0	21.8
E	Rectangle	200	0.05–0.75	0.1	28.9

a part of the calculation results of $\eta_{i,\text{meta}}$. Cases A–C are to observe the dependence of $\eta_{i,\text{meta}}$ on the waveform. Cases D and E are examples to demonstrate further idealized IFEA with extremely short impulse and low electron temperature. The ranges of T_{eff} are chosen depending on given E_{\max} so that a parameter $E_{\max}T_{\text{eff}}$ (which gives the total acceleration) ranges 10–150 V cm⁻¹ ns. This range corresponds to the energy gain of 0.088–19.8 eV, which includes the target energy range of 6.2–7.4 eV.


Figure 3. Impulse waveforms modelled: (a), cases A, B, D and E; (b), case C.

The temporal variation of the EVDF is calculated by a propagator method (Sugawara *et al* 1998, 2001). It is a numerical technique to solve the Boltzmann equation. The electron–molecule collisions are counted throughout the acceleration and post-acceleration periods until the electron energy becomes less than the lowest onset energy among the inelastic collisions of N₂ under consideration. In this cross-section set, it is 0.29 eV for a vibrational excitation.

The efficiency $\eta_{i,\text{meta}}$ is calculated as

$$\eta_{i,\text{meta}} = \frac{k_{\text{meta}}\epsilon_{\text{meta}}}{\sum_j k_j\epsilon_j + \epsilon_{\text{el}}^*}. \quad (7)$$

Here, k_j is the collision count, and ϵ_j is the loss energy of an inelastic collision. Subscript j represents the j th kind of inelastic collision, and one of the j values corresponds to

‘meta’ for the target reaction. ϵ_{el}^* is the total loss energy ϵ_{el} of each elastic collision. The value of ϵ_{el} at an elastic collision depends on the incident electron speed $v = |\mathbf{v}|$ and the scattering angle. In this calculation, ϵ_{el} is approximated as follows, assuming isotropic electron scattering in the cold gas model:

$$\epsilon_{\text{el}}(v) = 2\frac{m}{M} \left(\frac{1}{2}mv^2 \right) \quad (8)$$

where the factor m/M represents the mass-to-mass ratio between an electron and an N_2 molecule. In addition, ϵ_{el}^* includes the electron energy remaining at the end of calculation; in other words, the residual energy is regarded as being lost by elastic collisions eventually.

3.3. DC equilibrium mode

The distribution function $f(v)$ of the electron speed $v = |\mathbf{v}| = \sqrt{v_x^2 + v_y^2 + v_z^2}$ in DC drift equilibrium is calculated at E/N of 1–1000 Td (1 Td = 10^{-17} V cm²) by a two-term Boltzmann equation code. $f(v)$ can be related to $f(v_x, v_y, v_z)$ through $f(v, \theta, \phi)$, which is the EVDF represented in a polar-coordinate form. Here, θ and ϕ are the polar and azimuthal angles, and $f(v)$, $f(v_x, v_y, v_z)$ and $f(v, \theta, \phi)$ in drift equilibrium are independent of time. Their relation is given as follows:

$$f(v) = \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} f(v, \theta, \phi) v^2 \sin \theta d\theta d\phi \quad (9)$$

where $f(v, \theta, \phi) = f(v_x, v_y, v_z)$ at $(v_x, v_y, v_z) = (v \cos \theta, v \sin \theta \cos \phi, v \sin \theta \sin \phi)$.

The efficiency $\eta_{\text{e,meta}}$ is defined in the same way as $\eta_{\text{i,meta}}$ in equation (7):

$$\eta_{\text{e,meta}} = \frac{\int_0^{\infty} \epsilon_{\text{meta}} N q_{\text{meta}}(v) v f(v) dv}{\sum_j \int_0^{\infty} \epsilon_j N q_j(v) v f(v) dv + \int_0^{\infty} \epsilon_{\text{el}}(v) N q_{\text{el}}(v) v f(v) dv}. \quad (10)$$

3.4. Monoenergetic electron beam mode

An initial energy ϵ_{ini} is assumed in common for all of the initial electrons, which are released straight towards the $+x$ -direction in N_2 of $N = 3.54 \times 10^{16}$ cm⁻³ as a monoenergetic electron beam. The electron flight and collision are simulated by a Monte Carlo method under $E = 0$ in boundary-free space. Electrons may deviate from the course of the beam after their first collisions; however, their motion and energy loss processes are traced until the electrons satisfy the same termination condition as in the impulse mode. The energy loss by elastic collision is taken into account as well as that by the inelastic collisions, and the electron scattering due to the collisions is assumed to be isotropic. We examine ϵ_{ini} of 1–20 eV, and the number of initial electrons n_{ini} is 10 000 at each ϵ_{ini} value. The efficiency $\eta_{\text{b,meta}}$ is defined with the total initial electron energy $n_{\text{ini}}\epsilon_{\text{ini}}$ as

$$\eta_{\text{b,meta}} = \frac{k_{\text{meta}}\epsilon_{\text{meta}}}{n_{\text{ini}}\epsilon_{\text{ini}}}. \quad (11)$$

Here, although we have specified the N value and the direction of the initial electron velocity for practical calculation, $\eta_{b,meta}$ is independent of them when $E = 0$ and the space simulated is boundary-free.

4. Results and discussion

4.1. Electron behaviour in velocity space

Figure 4 shows the EVDF before and after acceleration in case A at $T_{eff} = 0.5$ ns. Here, the EVDF is depicted as the projection to the v_x - v_y plane as $f(v_x, v_y, t) = \int_{-\infty}^{\infty} f(v_x, v_y, v_z, t) dv_z$. v_x and v_y are plotted with a scale normalized by v_1 (the electron speed associated with 1 eV), and $(v/v_1)^2$ gives the value of the electron energy in eV. The EVDF peak decays by collisions during acceleration. The dropout electrons are scattered isotropically around $\mathbf{v} = 0$, and form the distribution tail. They restart acceleration following the leading peak for the residual impulse period, but they are out of control by the IFEA method. Electrons forming the leading peak continue the collective motion in velocity space, and those without collisions reach the target energy range.

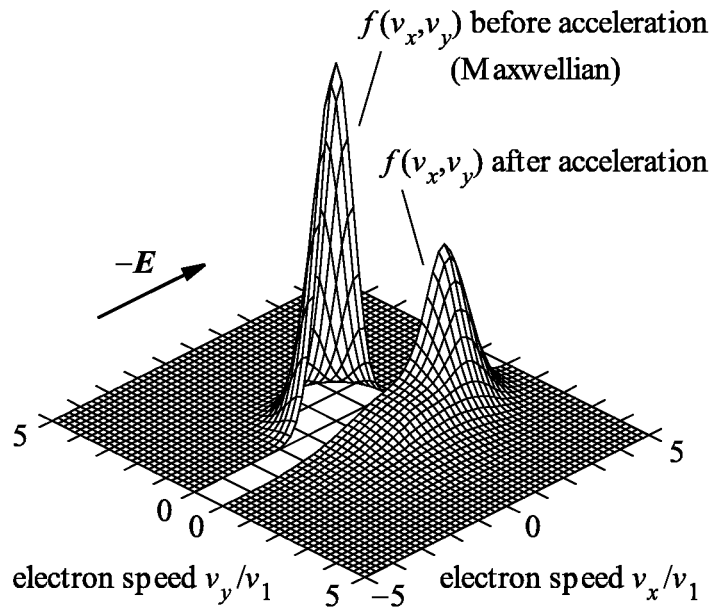


Figure 4. The EVDF $f(v_x, v_y, t)$ before and after acceleration (case A, $T_{eff} = 0.5$ ns). Each is symmetric in $v_y = 0$, then is depicted only for $v_y \geq 0$. The scale unit v_1 is the electron speed associated with 1 eV.

Figures 5 and 6 show the cross-sections of the EVDF at $v_y = 0$ in cases A and D. $E_{max}T_{eff} = 100$ V cm $^{-1}$ ns in both cases ($T_{eff} = 0.5$ ns in case A, and 0.1 ns in case D). The first six curves in each figure are those during acceleration, and the others are those after acceleration. The peak in case D decays less than in case A because of

fewer collision opportunities due to the shorter acceleration period. The decay is not simply exponential because q_{total} is energy dependent. Especially in figure 5, a large decay appears between $t = 0.2$ and 0.3 ns. This is due to scattering by the q_{el} and q_{vib} peaks lying around $\epsilon = 2$ eV ($v/v_1 = 1.4$). This influence particular to N_2 is called the barrier effect (Engelhardt *et al* 1964, Sugawara *et al* 2001), which functions to prevent slow electrons from accelerating. After the impulse field is turned off, the peak moves back towards $v_x = 0$ for energy loss by collision. The electrons concentrate around $v = 0$ again after a cooling period over $10 \mu\text{s}$.

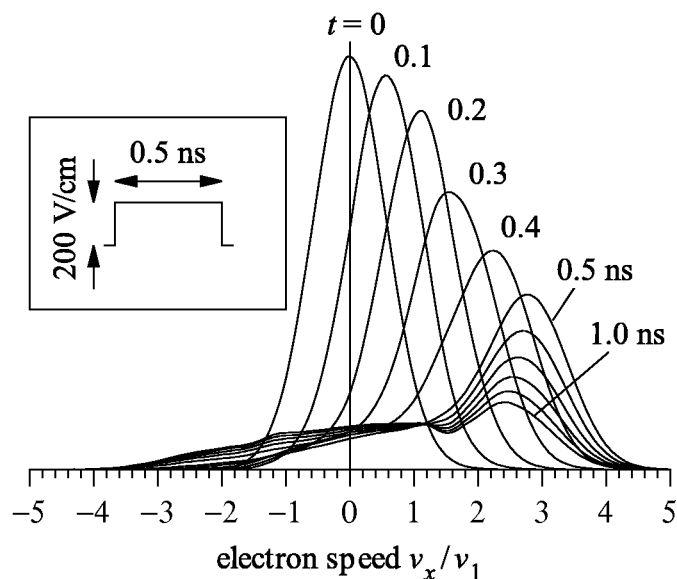


Figure 5. The EVDF $f(v_x, v_y, t)|_{v_y=0}$ at $t = 0-1.0$ ns (0.1 ns step); case A, $T_{\text{eff}} = 0.5$ ns.

Figure 7 shows loci of the EVDF peak in cases A–D at $E_{\text{max}}T_{\text{eff}} = 100 \text{ V cm}^{-1} \text{ ns}$. The relative peak height represents the ratio of the electrons remaining free from collisions. The decay becomes larger as T increases. Compared with cases A–C, case D at a shorter T is much less influenced by collisions. Electrons can evade the barrier effect of the q_{el} and q_{vib} peaks when the impulse is short and intense.

4.2. Comparison of energy efficiencies in different modes

Figures 8–10 show $\eta_{i,\text{meta}}$, $\eta_{e,\text{meta}}$ and $\eta_{b,\text{meta}}$ together with the energy consumption ratios of other collision processes. In figure 8, $\eta_{i,\text{meta}}$ reaches 20%. This value is higher than the maximum value of $\eta_{e,\text{meta}}$ of about 10% in figure 9. The selectivity in the equilibrium mode is low for dispersed EVDF because different kinds of inelastic collisions occur simultaneously in a wide energy range. In contrast, electrons in the impulse mode are more concentrated around the target energy range, which leads to higher energy efficiencies. This is indicated by the result that the maximum values of $\eta_{i,\text{meta}}$ appear around a specific value of $E_{\text{max}}T_{\text{eff}} = 84 \text{ V cm}^{-1} \text{ ns}$ which corresponds to the energy gain

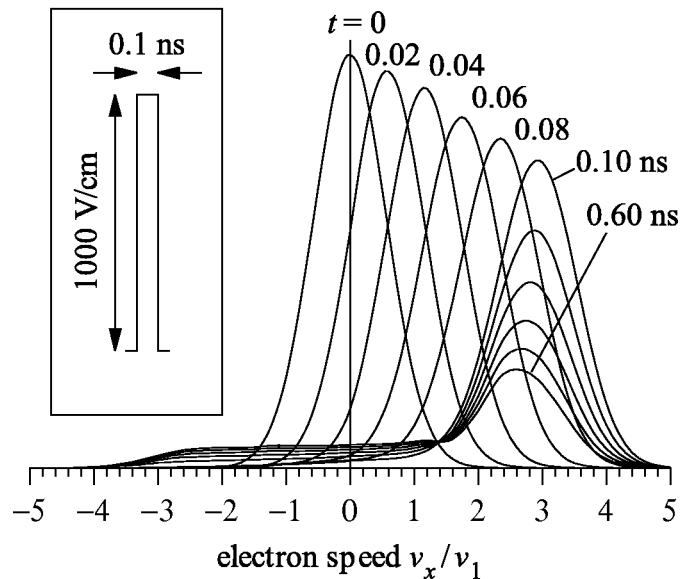


Figure 6. The EVDF $f(v_x, v_y, t)|_{v_y=0}$ at $t = 0$ – 0.1 ns (0.02 ns step) and $t = 0.1$ – 0.6 ns (0.1 ns step); case D, $T_{\text{eff}} = 0.1$ ns.

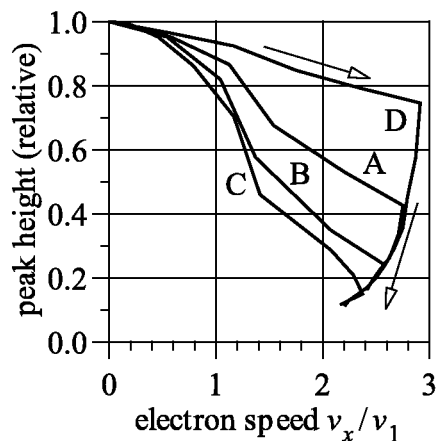


Figure 7. Comparison of loci of the EVDF peaks in cases A–D at $E_{\text{max}}T_{\text{eff}} = 100 \text{ V cm}^{-1} \text{ ns}$.

of $\epsilon_{\text{meta}} = 6.2 \text{ eV}$. At higher $E_{\text{max}}T_{\text{eff}}$ values, $\eta_{i,\text{meta}}$ decreases because of over-acceleration that induces inelastic processes other than the target reaction.

The merit of concentrated EVDF emerges dramatically in the beam mode. In figure 10, $\eta_{b,\text{meta}}$ is over 80% when ϵ_{ini} is in the target range 6.2–7.4 eV. Although elastic collision is more than 100 times as frequent as the target process in this energy range, the energy loss by elastic collision is not significant because of a small m/M of only about 2×10^{-5} . The electrons hold most of their initial energy until the target reaction occurs even after hundreds of elastic collisions. This result suggests improvement of $\eta_{i,\text{meta}}$ by making the EVDF in the impulse mode more monoenergetic. A possible way

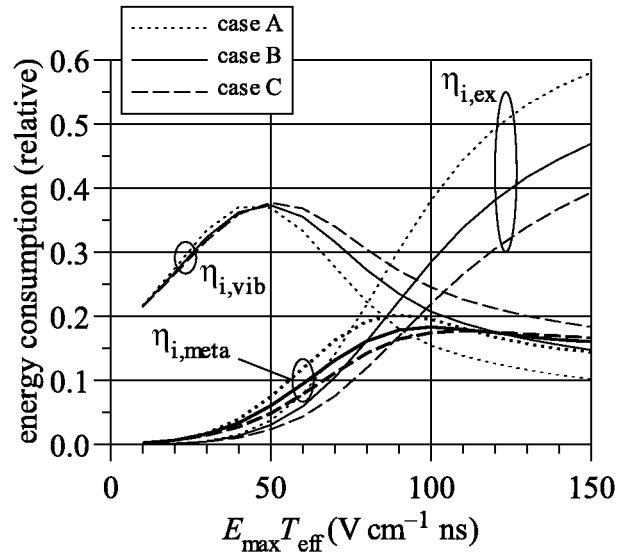


Figure 8. Energy consumption ratios in the impulse field mode in cases A, B and C (see table 1 for detailed conditions).

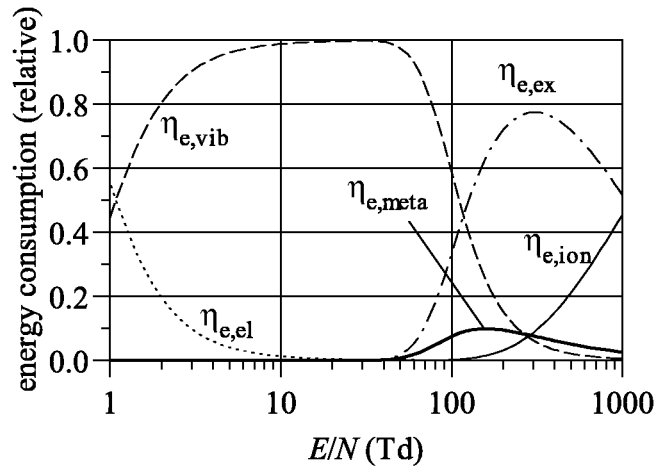


Figure 9. Energy consumption ratios in the DC drift equilibrium mode.

would be to reduce the scattering during the IFEA with shorter and intenser impulse. Another would be to make the electron temperature before the IFEA lower. Results of such attempts are presented in the next subsection.

4.3. Dependence of energy efficiency on impulse waveform and electron temperature

In figure 8, curves of $\eta_{i,meta}$ for cases A–C show that $\eta_{i,meta}$ increases as T becomes shorter in a parameter range of $E_{max}T_{eff} < 100 \text{ V cm}^{-1} \text{ ns}$. This tendency is understood as decrease of the electron dropout, as explained before. Because the dropout leaves electrons unaccelerated, this understanding is supported also by the decrease in $\eta_{i,vib}$; vibrational excitation occurs mainly at low energies. Keeping this tendency, a higher

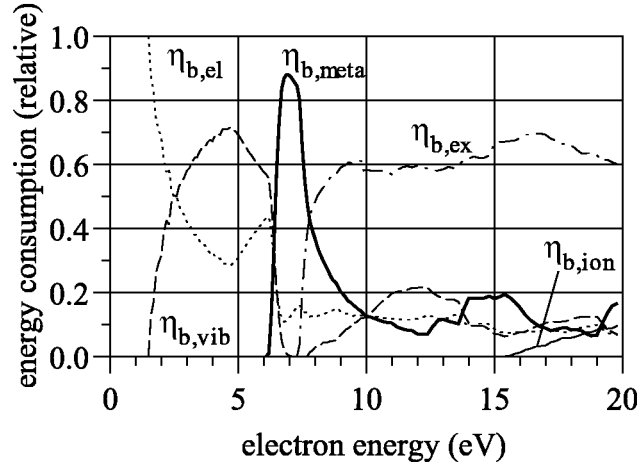


Figure 10. Energy consumption ratios in the monoenergetic electron beam mode. The horizontal axis represents the initial energy of electrons in the beam.

$\eta_{i,meta}$ of 22% is obtained in case D in figure 11 under an extremely short impulse.

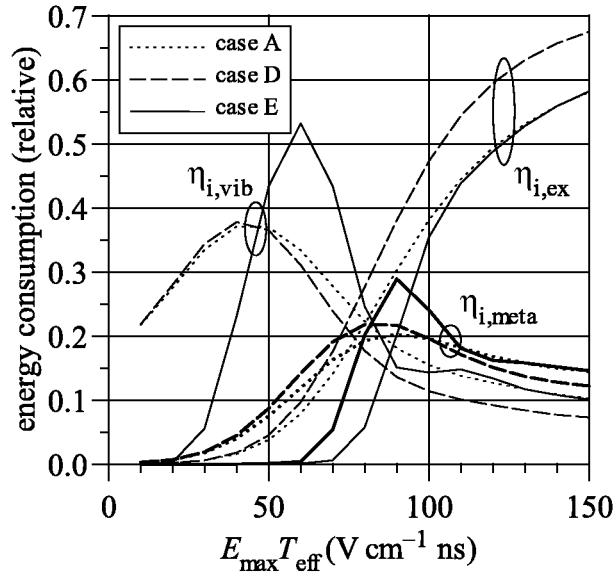


Figure 11. Energy consumption ratios in the impulse field mode in cases A, D and E (see table 1 for detailed conditions).

However, it is at a lower electron temperature that a significant improvement of the energy efficiency is achieved. In case E at $\epsilon_{ini} = 0.1$ eV, $\eta_{i,meta}$ reaches 29% (figure 11). As indicated by a sharp rise of $\eta_{i,meta}$, the selectivity in the impulse mode is high and sensitively dependent on $E_{max}T_{eff}$ at low electron temperature.

When the IFEA is driven in repetitive pulse mode, some efficient electron cooling mechanism would be necessary to keep a high selectivity by recovering a low electron temperature before applying succeeding impulses. For this purpose, admixture of

coolant species to the gas could be supposed. However, because the IFEA method is a technique to avoid the EVDF dispersion and has no active function to make the EVDF concentrated, we leave the cooling issue for future investigation elsewhere.

4.4. Possibility of large volume reaction by the IFEA method

As a potential application of the IFEA method, large volume reaction in a plasma bulk is considered here. Electrons in plasma bulks are usually at a low temperature and a high concentration, which is suitable for the IFEA method, and mass production of $N_2(A^3\Sigma_u^+)$ is expected. In order to perform IFEA in plasma bulk, we need satisfactory penetration of the impulse electric field into the bulk. The impulse risetime must be shorter than the plasma response time, otherwise the impulse field would be shielded by the plasma sheath. The period of plasma oscillation (i.e. the inverse of the plasma frequency f_p) can be a measure to evaluate the plasma response time:

$$\frac{1}{f_p} = 2\pi\sqrt{\frac{\epsilon_0 m}{e^2 n_e}}. \quad (12)$$

Here, ϵ_0 is the permittivity of free space, and n_e is the electron concentration. At a typical value of $n_e \sim 10^9 \text{ cm}^{-3}$ for N_2 radio frequency (13.56 MHz) plasmas, we obtain $1/f_p = 3.5 \text{ ns}$. The impulses investigated in this paper are shorter than this value, which suggests a fair prospect of the IFEA in plasma bulks.

5. Conclusion

We have presented a theoretical basis of the IFEA in gas, and have demonstrated selective production of $N_2(A^3\Sigma_u^+)$ by electron-impact excitation of N_2 as an example of the IFEA application.

Under a short intense electric field, electrons in gas can be accelerated in a practically collisionless condition, and the concentration of the electrons in the velocity distribution is kept during the IFEA. Controlling the impulse field strength and duration, one can feed a bunch of electrons to an energy range in which this target process (i.e. excitation of $N_2(A^3\Sigma_u^+)$) occurs efficiently. The electron energy efficiency has been defined to quantify the reaction selectivity, and it has been shown that the efficiency of the IFEA method reaches to 20–30%. These values are higher than that in electron swarms in DC drift equilibrium (i.e. less than 10%).

The high selectivity of the IFEA method is based on the concentration of the electron velocity distribution. Therefore, the selectivity is high and its dependence on the impulse property is sharp when the electron temperature before IFEA is low.

Acknowledgments

The authors wish to thank Professor S Katsuki of Kumamoto University for providing references and comments on ultra-short high voltage pulses.

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