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# On-the-fly molecular dynamics study of the excited-state branching reaction of $\alpha$ -methyl-*cis*-stilbene

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## ABSTRACT

The branching reaction of  $\alpha$ -methyl-*cis*-stilbene (*cis*-mSB) into its *trans*-mSB and 4a,4b-dihydrophenanthrene (DHP) forms upon  $\pi\pi^*$  excitation was examined theoretically by exploring the excited-state potential energy surface and using on-the-fly molecular dynamics simulations at the spin-flip time-dependent density functional theory (SF-TDDFT) level of theory. The branching ratio of trajectories was calculated as DHP:*twist* = 11:29, where *twist* denotes a mid-region between the *cis*-form and *trans*-form, indicating that the *trans*-mSB is a dominant product. The branching mechanism was analyzed by comparison with the corresponding theoretical studies on stilbene (SB) and 1,1'-dimethyl-stilbene (dmSB). The present computations elucidate the origin of variations in the branching ratio in the photoreactions of *cis*-SB, *cis*-mSB, and *cis*-dmSB. We also found that, because of loss of the slow component of the decay to the ground state, *cis*-mSB shows a faster decay rate to the ground state than *cis*-SB and *cis*-dmSB.

## KEYWORDS

stilbene, excited-state dynamics, reaction pathways, on-the-fly MD

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**Notes:** The authors declare no competing financial interest.

**Graphical abstract:**

## I. Introduction

It is now recognised that the conical intersection (CI) of the potential energy surfaces (PESs) of the adiabatic electronic states plays a significant role in photoreactions [1–3]. Recent progress in theoretical chemistry has made it feasible to elucidate the mechanism and dynamics of photoreactions involving a nonradiative decay through CIs by the highly sophisticated quantum chemical approaches [4–6], ab initio surface hopping molecular dynamics with the nonadiabatic coupling terms [7,8], and systematic exploration of minimum energy CI (MECI) structures based on the global reaction route mapping (GRRM) strategy [9–11]. A systematic search for low-lying MECIs enables the determination of the energetically favourable internal conversion paths from the Franck–Condon region. Recently, the difference of fluorescence quantum yields of aromatic molecules has been discussed based on the barrier heights along the internal conversion paths [12,13].

Stilbene (SB) is a typical molecule that shows *cis–trans* photoisomerisation about the central C=C bond on the  $\pi\pi^*$  excitation, and the photoisomerisation process of SB has been investigated by many experimental [14–25] and theoretical [26–32] studies. *trans*-SB has a planar geometry of  $C_{2h}$  symmetry, and the isomerization pathway in the  $\pi\pi^*$  excited state for leaving the fluorescent region accompanies a small barrier, and thus, the lifetime was reported to be relatively long (10 ~ >100 ps) [18,22]. On the other hand,  $\pi\pi^*$ -excited *cis*-SB exhibits ultrafast decay to the ground state with a lifetime of around 1.0 ps [20,21] *via* the CI points of the singlet excited-state ( $S_1$ ) and ground-state ( $S_0$ ) potential energy surfaces (PESs) in the *twist* region where the central C=C bond is twisted; this is known as the phantom state ( $P^*$ ) [15]. Theoretical studies on the decay mechanism of the  $\pi\pi^*$ -excited *cis*-SB have confirmed the twisting motion of the central C=C bond in the  $S_1$  state, and this motion is known as the "hula-twist" [23]. Femtosecond Raman spectroscopy with quantum chemical calculations has clarified the initial nuclear dynamics of the  $\pi\pi^*$ -excited *cis*-SB based on the transit change of the vibrational spectra [24]. Femtosecond time-resolved fluorescence showed that the lifetime

of the excited *cis*-SB has two components at 0.23 and 1.2 ps [25].

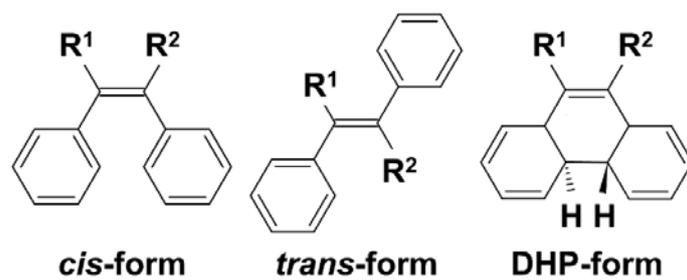
It is known that  $\pi\pi^*$ -excited *cis*-SB has a by-product, 4a,4b-dihydrophenanthrene (DHP). The branching ratio of the products from  $\pi\pi^*$ -excited *cis*-SB has been reported to be *cis:trans:DHP* = 55:35:10 [16,17,19]. Thus, the dominant product is *trans*-SB. Minezawa and Gordon [29] employed a spin-flip time-dependent density functional theory (SF-TDDFT) method [33–35] to locate the geometries of the minima in the  $S_1$  state, as well as minimum-energy conical intersection (MECI) points of the  $S_1$  and  $S_0$  states in the *twist* ( $(S_1)_{twist-min}$  and  $(S_1/S_0)_{twist}$ ) and DHP ( $(S_1)_{DHP-min}$  and  $(S_1/S_0)_{DHP}$ ) regions of the  $S_1$ -PES. It was also found that the photocyclisation is in competition with the photoisomerisation [29], which was later confirmed by a more sophisticated multireference method [30].

Harabuchi *et al.* [31] carried out reaction-path calculations on the  $S_1$ -PES and on-the-fly molecular dynamics (MD) simulations for the  $\pi\pi^*$ -excited *cis*-SB at the SF-TDDFT level of theory and discussed the reaction mechanism and dynamics. Along the steepest descent pathway starting from the Franck–Condon *cis*-structure, the two H atoms bound to the central C=C first move quickly so that each C-H-C<sub>6</sub>H<sub>5</sub> fragment in SB becomes planar and the  $S_1$ -PES becomes very flat; then, the pathway becomes sharply curved toward the DHP-form. It was found that the pathway shows a bifurcation in a very flat region of the  $S_1$ -PES: one direction is toward the DHP region as the pathway itself indicates, whereas the other is toward the *twist* region with a very small activation barrier. The excited-state reaction path bifurcation can be regarded as an extension of the ground-state reaction path bifurcation [36,37]. The branching ratio from the on-the-fly MD simulations on the  $\pi\pi^*$ -excited *cis*-SB indicated that the dominant product is *trans*-SB (*trans:DHP* = 35:13), which is consistent with the experimental data. This result suggests the importance of dynamics in the photoreaction.

In 2012, Berndt *et al.* [38] reported an experimental study of the photoreactions of 1,1'-dimethylstilbene (dmSB) in solution, using transit-absorption spectra. It was reported that the peak at 600 nm, which corresponds to the absorption band of the *cis*-form in the  $S_1$  state (CIS\*)

decreases with a lifetime of 0.2 ps, which is shorter than the corresponding lifetime of *cis*-SB (0.92 ps), while the peak at 330 nm corresponding to the absorption band of the P\* state shows a surprisingly long lifetime: 19 ps in hexane, which is much longer than the corresponding lifetime of *cis*-SB (1.2 ps) [38]. Harabuchi *et al.* [39] carried out theoretical computations of the excited-state PES and on-the-fly MD simulations for the reaction process of *cis*-dmSB in the S<sub>1</sub> state, and found that the long lifetime of the P\* state in dmSB can be ascribed to the large geometrical difference between (S<sub>1</sub>)<sub>twist-min</sub> and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist</sub> in the *twist* region and that, in the case of the ππ\*-excited *cis*-dmSB, the dominant product is not *trans*-dmSB but dmDHP. Very recently, Kokado *et al.* [40] examined the structural transformation of ππ\*-excited tetraphenylethene, in which the four H atoms in ethylene are substituted by four phenyl rings, both experimentally and theoretically and proposed that the rotation about the central C=C bond is essential for the target process. The effect of substitution on the photoreaction process is a significant factor that determines the decay lifetime and quantum yield of photoisomerisation for stilbene derivatives, and, thus, theoretical examination of the excited-state reaction process for new substituted species should provide valuable insights.

**Figure 1** shows the *cis*-, *trans*-, and DHP-forms of stilbene derivatives; by changing R<sup>1</sup> and R<sup>2</sup> to different groups, various stilbene derivatives can be obtained. In this study, we have chosen *α*-methyl-*cis*-stilbene (*cis*-mSB) as a target species, which should have a character between those of *cis*-SB and *cis*-dmSB, although only *cis*-mSB has asymmetric groups for R<sup>1</sup> and R<sup>2</sup>. In the same way as previous studies on *cis*-SB and *cis*-dmSB, we examined the S<sub>1</sub>-PES and on-the-fly MD simulations for the ππ\* excited *cis*-mSB at the SF-TDDFT level of theory and examined the reaction pathways in the S<sub>1</sub>-state, the branching ratio for products from MD simulations, and the decay lifetime to the ground state. Because there are no experimental studies on this species, our computations can be used as predictions for this new target molecule.



**Figure 1.** Three isomers of stilbene derivatives: stilbene for  ${}^1R = {}^2R = H$ ;  $\alpha$ -methyl-stilbene for  ${}^1R = \text{methyl group}$  and  ${}^2R = H$ ; 1,1'-dimethylstilbene for  ${}^1R = {}^2R = \text{methyl group}$ .

## II. Computational details

Geometry optimisations were performed on mSB at the SF-TDDFT level of theory to locate the minima in the  $S_0$  and  $S_1$  states (corresponding to *cis*-, *trans*-, *twist*-, and DHP-forms), the transition state (TS) structures that connect the DHP and *twist* regions in the  $S_1$  state, and the  $S_1/S_0$ -MECIs in DHP and *twist* regions [33–35]. The obtained geometries were verified as minima or transition states by normal mode analysis. The intrinsic reaction coordinate (IRC) [36] was calculated from the  $S_1$ -TS to confirm the connectivity of the two minima in the DHP and *twist* regions. The steepest descent path in the  $S_1$  state was calculated from the Franck–Condon (FC) structure in mass-weighted coordinates (referred to as the meta-IRC path). The SF-TDDFT calculations were performed with the BHHLYP functional and 6-31G(d) basis set using the GAMESS program [41], whereas all the minima, TSs, MECIs, IRCs, and meta-IRC paths were calculated by employing the GRRM14 program [42] with GAMESS. The branching plane update method [43] was used in geometry optimisation for the  $S_1/S_0$ -MECIs.

On-the-fly MD simulations were performed for the  $\pi\pi^*$ -excited *cis*-mSB at the SF-TDDFT level using the SPPR program [44]. To avoid mixed singlet and triplet states in SF-TDDFT calculations around the CI region along the trajectory, the  $T_{SF}$ -index method [10] was employed. The initial conditions for the on-the-fly MD simulations were determined by normal mode sampling of the ground state equilibrium geometry of *cis*-mSB. The atomic coordinates and velocities were generated randomly by adding an energy of  $kT$  to each normal mode under the Boltzmann distribution at 300 K. The time step was set to 0.2 fs, and 40 trajectories were calculated until the energy difference between the  $S_0$  and  $S_1$  states became less than 0.2 eV or the simulation time reached 1 ps.

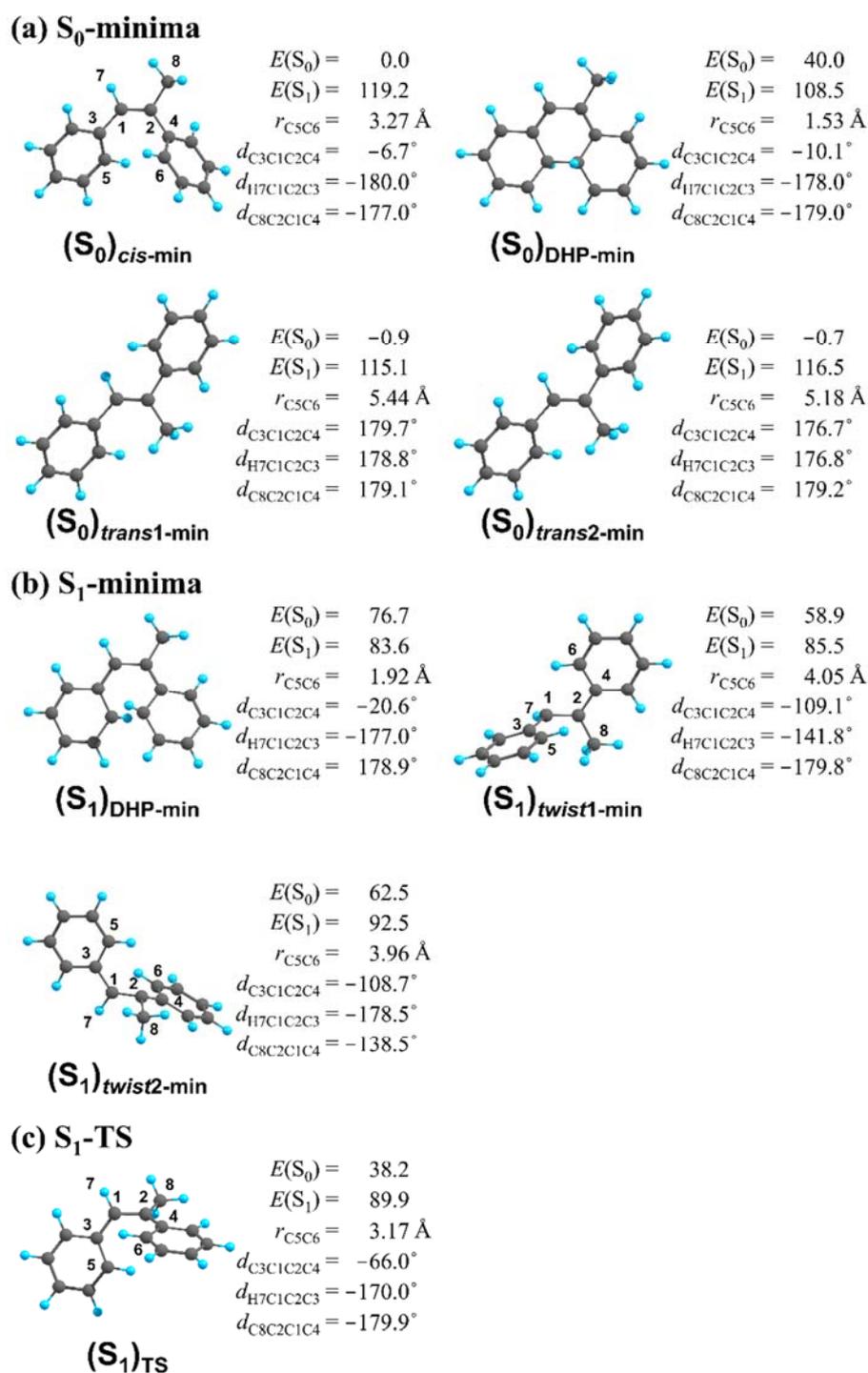
### III. Results and discussion

#### A. Geometries of minima, TSs, and S<sub>1</sub>/S<sub>0</sub>-MECIs

For mSB, we located four minima in the S<sub>0</sub> state ((S<sub>0</sub>)*cis*-min, (S<sub>0</sub>)*DHP*-min, (S<sub>0</sub>)*trans*1-min, and (S<sub>0</sub>)*trans*2-min), three minima in the S<sub>1</sub> state ((S<sub>1</sub>)*DHP*-min, (S<sub>1</sub>)*twist*1-min, and (S<sub>1</sub>)*twist*2-min), and one TS between (S<sub>1</sub>)*DHP*-min and (S<sub>1</sub>)*twist*1-min ((S<sub>1</sub>)*TS*). In the two *trans*-minima in the S<sub>0</sub> state, the two phenyl rings are deviated slightly from the central ethylenic plane by steric repulsion with the methyl ring, as is the case for dmSB [39], and the difference between these two structures is shown by the *d*<sub>C1C2C4C6</sub> dihedral angles (31.8° for (S<sub>0</sub>)*trans*1-min and -39.8° for (S<sub>0</sub>)*trans*2-min). **Figure 2** shows the optimised geometries of these minima and the TS, as well as the energies of the S<sub>0</sub> and S<sub>1</sub> states relative to (S<sub>0</sub>)*cis*-min, *E*(S<sub>0</sub>) and *E*(S<sub>1</sub>), and four significant geometric parameters: *r*<sub>C5C6</sub>, *d*<sub>C3C1C2C4</sub>, *d*<sub>H7C1C2C3</sub>, and *d*<sub>C8C2C1C4</sub> (*r*<sub>AB</sub> denotes the interatomic distance of *A* and *B* atoms, and *d*<sub>ABCD</sub> denotes the dihedral angle for *A-B-C-D* atoms); *r*<sub>C5C6</sub> is a measure of the cyclisation from the *cis*-form to the *DHP*-form, and *d*<sub>C3C1C2C4</sub> is a measure of the *cis-trans* (or *cis-twist*) structural transformation; *d*<sub>H7C1C2C3</sub> and *d*<sub>C8C2C1C4</sub> represents the deformation of H7 and C8, respectively, from the central ethylenic planar part. The two *trans*-forms have almost the same energies which are very slightly lower than that of the *cis*-form (within 1 kcal/mol).

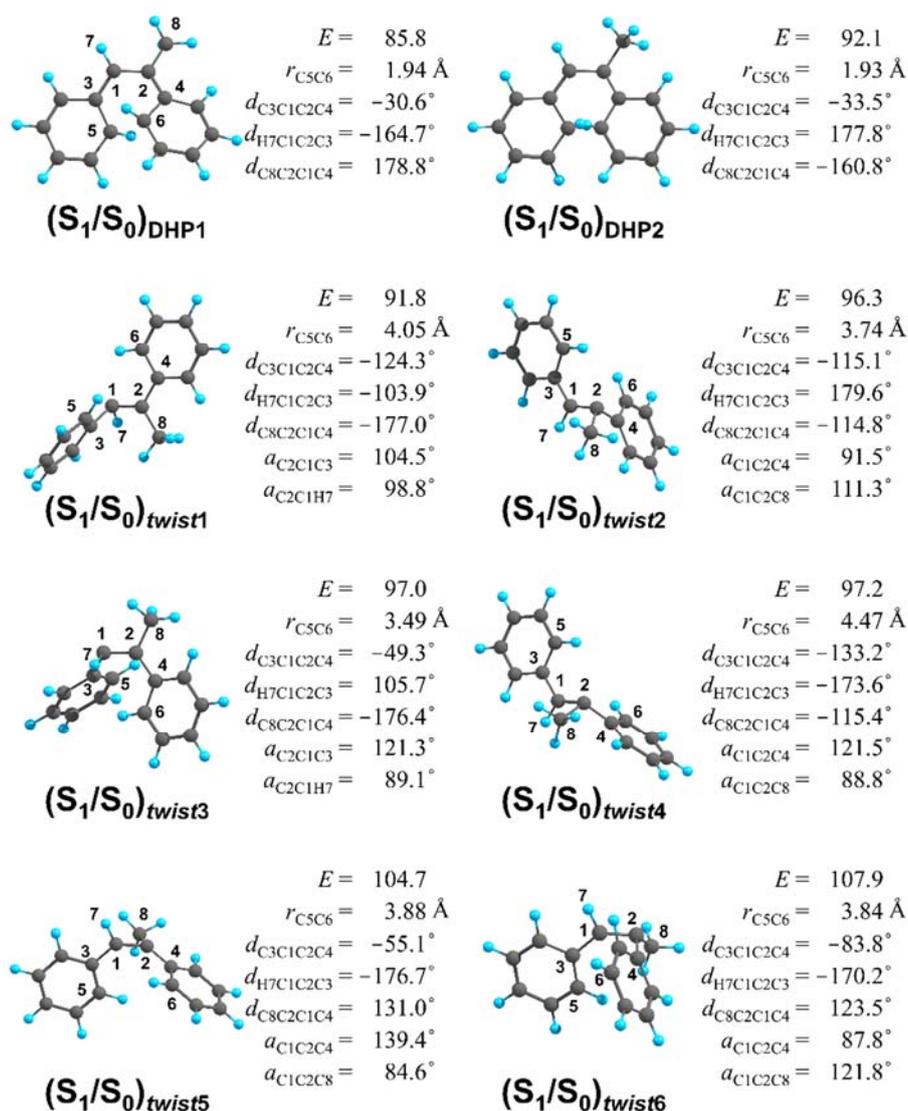
The central C=C part of (S<sub>0</sub>)*DHP*-min shows an almost planar structure (*d*<sub>C3C1C2C4</sub> = -10.1°), whereas the corresponding dihedral angle becomes larger at (S<sub>1</sub>)*DHP*-min (*d*<sub>C3C1C2C4</sub> = -20.6°). In the *twist* region in the S<sub>1</sub> state, two minima, (S<sub>1</sub>)*twist*1-min and (S<sub>1</sub>)*twist*2-min, were found; the former has a pyramidal structure at H-C-phenyl (H-pyramidal) with *d*<sub>H7C1C2C3</sub> = -141.8°, and the latter has a pyramidal structure at CH<sub>3</sub>-C-phenyl (Me-pyramidal) with *d*<sub>C8C2C1C4</sub> = -138.5°. (S<sub>1</sub>)*twist*2-min is energetically higher than (S<sub>1</sub>)*twist*1-min (92.5 vs. 85.5 kcal/mol) because of the repulsion of the methyl ring and phenyl ring. As shown in Fig. 2c, we also located a TS in the S<sub>1</sub> state, and verified that (S<sub>1</sub>)*TS* connects (S<sub>1</sub>)*DHP*-min and (S<sub>1</sub>)*twist*1-min *via* IRC calculations. Thus, the molecule will prefer the H-pyramidal structure to the Me-pyramidal structure in the

twist region.



**Figure 2.** The optimised geometries of (a) the four minima in the  $S_0$  state, (b) three minima in the  $S_1$  state, and (c) one TS in the  $S_1$  state for mSB. The energies of the  $S_0$  and  $S_1$  states relative to ( $S_0$ )*cis*-min (in kcal/mol) and significant geometrical parameters are also given.

Concerning the S<sub>1</sub>/S<sub>0</sub>-MECI structures, we located two structures in the DHP region and six structures in the *twist* region, as shown in **Fig. 3**. In the DHP region, (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP1</sub> corresponds to the H-pyramidal structure ( $d_{\text{H7C1C2C3}} = -164.7^\circ$  and  $d_{\text{C8C2C1C4}} = 178.8^\circ$ ), while (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP2</sub> corresponds to the Me-pyramidal structure ( $d_{\text{H7C1C2C3}} = 177.8^\circ$  and  $d_{\text{C8C2C1C4}} = -160.8^\circ$ ), although the pyramidal parts are close to the planar structure compared to the S<sub>1</sub>-minima. (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP1</sub> is more stable than (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP2</sub> (85.9 vs. 92.1 kcal/mol) because of the steric repulsion between the methyl group and the phenyl ring in the Me-pyramidal structure. The (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP1</sub> is energetically close to (S<sub>1</sub>)<sub>DHP-min</sub> (83.6 kcal/mol), indicating that the molecule could easily relax to the ground state through (S<sub>1</sub>/S<sub>0</sub>)<sub>DHP1</sub>. The six S<sub>1</sub>/S<sub>0</sub>-MECI structures in the *twist* region are given in the order of the relative energies. (S<sub>1</sub>/S<sub>0</sub>)<sub>twist1</sub> and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist3</sub> correspond to the H-pyramidal structure, whereas (S<sub>1</sub>/S<sub>0</sub>)<sub>twist2</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist4</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist5</sub>, and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist6</sub> correspond to the Me-pyramidal structure. These structures can be further classified into the *trans*-type and *cis*-type according to the relative positions of two phenyl rings; (S<sub>1</sub>/S<sub>0</sub>)<sub>twist1</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist2</sub>, and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist4</sub> correspond to the twisted-*trans*-pyramidal structure, and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist3</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist5</sub>, and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist6</sub> correspond to the twisted-*cis*-pyramidal structure. The latter twisted-*cis*-pyramidal-type MECI has not been reported previously. Taking into account these two classifications, the six S<sub>1</sub>/S<sub>0</sub>-MECI structures in the *twist* region can be classified into four types, *i.e.*, H-twisted-*trans*-pyramidal ((S<sub>1</sub>/S<sub>0</sub>)<sub>twist1</sub>), H-twisted-*cis*-pyramidal ((S<sub>1</sub>/S<sub>0</sub>)<sub>twist3</sub>), Me-twisted-*trans*-pyramidal ((S<sub>1</sub>/S<sub>0</sub>)<sub>twist2</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist4</sub>), and Me-twisted-*cis*-pyramidal ((S<sub>1</sub>/S<sub>0</sub>)<sub>twist5</sub>, (S<sub>1</sub>/S<sub>0</sub>)<sub>twist6</sub>). The difference between (S<sub>1</sub>/S<sub>0</sub>)<sub>twist2</sub> and (S<sub>1</sub>/S<sub>0</sub>)<sub>twist4</sub> is the bond angle in the pyramidal part:  $a_{\text{C1C2C4}}$  is close to  $90^\circ$  in (S<sub>1</sub>/S<sub>0</sub>)<sub>twist2</sub>, and  $a_{\text{C1C2C8}}$  is close to  $90^\circ$  in (S<sub>1</sub>/S<sub>0</sub>)<sub>twist4</sub>. Similarly,  $a_{\text{C1C2C4}}$  is close to  $90^\circ$  in (S<sub>1</sub>/S<sub>0</sub>)<sub>twist6</sub>, and  $a_{\text{C1C2C8}}$  is close to  $90^\circ$  in (S<sub>1</sub>/S<sub>0</sub>)<sub>twist5</sub>. Similar geometrical features of the MECI structures were observed for SB and dmSB [31,39]. (S<sub>1</sub>/S<sub>0</sub>)<sub>twist3</sub> can also be classified by this viewpoint ( $a_{\text{C2C1C7}}$  is close to  $90^\circ$ ), unlike (S<sub>1</sub>/S<sub>0</sub>)<sub>twist1</sub> (the smallest bond angle is  $98.8^\circ$ ).



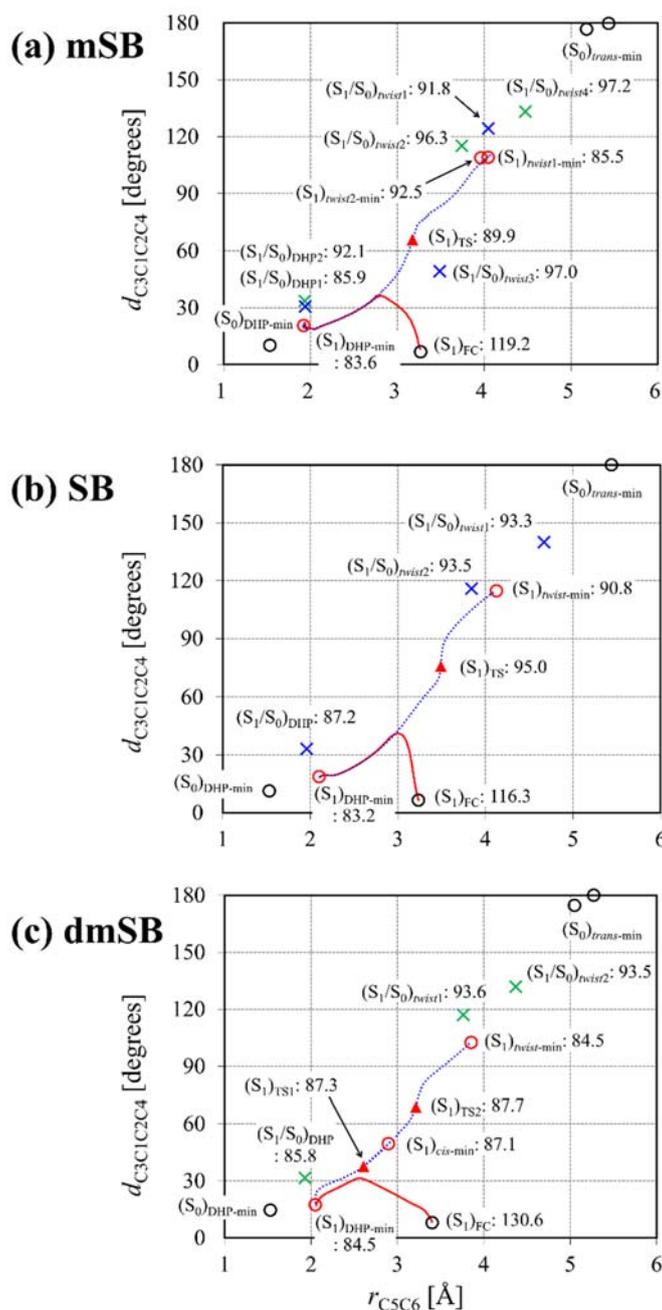
**Figure 3.** The optimized geometries of the calculated eight  $S_1/S_0$ -MECIs for mSB. The energy relative to ( $S_0$ )<sub>cis-min</sub> ( $E$  in kcal/mol) and significant geometrical parameters are also shown.

In the dynamics simulations, the molecule is expected to reach the energetically low-lying CI region. We performed meta-IRC calculations starting from the eight MECIs in the *twist* region, and verified that all the pathways reach ( $S_1$ )<sub>twist1-min</sub> or ( $S_1$ )<sub>twist2-min</sub> without a barrier, indicating that the energy level of each MECI can be used to judge whether the corresponding MECI is accessible by trajectories wandering around the  $S_1$ -minima.

## B. Reaction pathways in the S<sub>1</sub> state

The  $\pi\pi^*$ -excited *cis*-mSB should move down the slope of the S<sub>1</sub>-PES and branch into the DHP and *twist* regions. Following the previous studies on *cis*-SB [31] and *cis*-dmSB [39], we first examined the S<sub>1</sub>-PES in the two-dimensional coordinate space spanned by  $r_{C5C6}$  (which characterises the transformation to the DHP) and  $d_{C3C1C2C4}$  (which characterises the transformation to the *twist* form). **Figure 4** shows the positions of the stationary points in the S<sub>0</sub> and S<sub>1</sub> states and the S<sub>1</sub>/S<sub>0</sub>-MECIs, as well as the reaction pathways of the meta-IRC from the Franck-Condon *cis*-structure (red solid line) and the IRC between the DHP and *twist* minima (blue dotted line), in two-dimensional coordinate space for (a) mSB, (b) SB, and (c) dmSB. The H-pyramidal S<sub>1</sub>/S<sub>0</sub>-MECIs are indicated by blue cross marks, and the Me-pyramidal S<sub>1</sub>/S<sub>0</sub>-MECIs are indicated by green cross marks. As shown in Fig. 4, the DHP and *twist* regions are separated by a single TS for mSB and SB, whereas there are DHP, *cis*, and *twist* regions separated by two TSs, for dmSB. The meta-IRC path from the FC structure reaches the IRC pathway between (S<sub>1</sub>)<sub>DHP-min</sub> and (S<sub>1</sub>)<sub>twist1-min</sub>, indicating a branched reaction pathway.

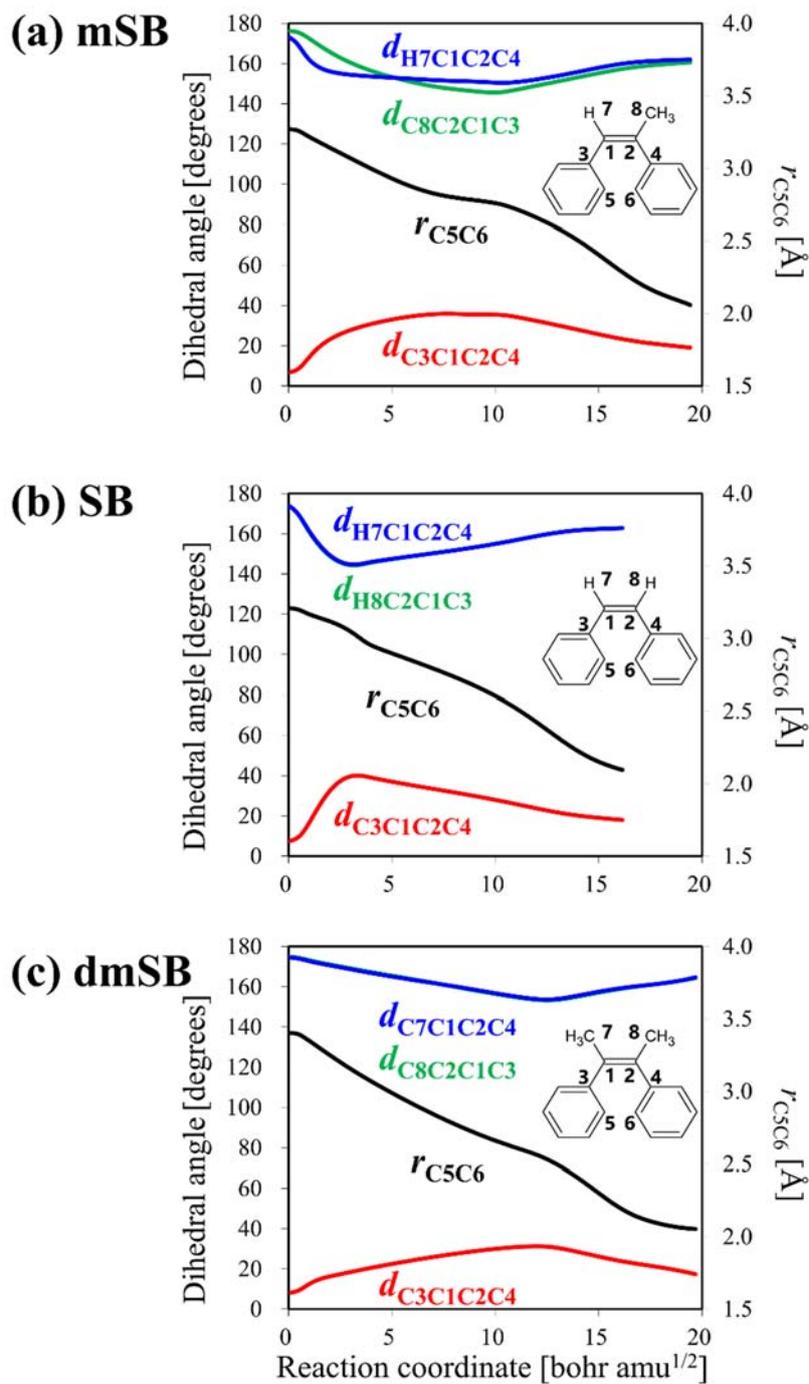
The profile of the meta-IRC route suggests that there is a preference for the DHP region over the *twist* region in all cases. In the case of SB, however, the direction of the meta-IRC path is toward the *twist* region before reaching the IRC between the DHP and *twist* minima as shown in Fig. 4b, and thus, the inertial force will push the molecule to the *twist* region. On the other hand, in the case of dmSB, the direction of the meta-IRC path is toward the DHP region from the start of the pathway, as shown in Fig. 4c, indicating that the DHP-form is expected as a major product. The case of mSB is between those of SB and dmSB, and the branching ratio over the DHP and *twist* regions will be sensitive to the dynamic effects in the same way as in the case of SB [31].



**Figure 4.** The minima, TSs, and MECI structures and the reaction pathways in the  $S_1$  state for (a) mSB, (b) SB [31], and (c) dmSB [39] projected onto two-dimensional coordinate space in terms of  $r_{C_5C_6}$  and  $d_{C_3C_1C_2C_4}$ . The  $S_1$ -minima and  $S_1$ -TS are depicted by red circles and red triangles, respectively; the H-pyramidal  $S_1/S_0$ -MECIs are given by blue cross marks, and the Me-pyramidal  $S_1/S_0$ -MECIs are given by green cross marks. The  $S_0$ -minima are also given by black circles. The meta-IRC path from the FC point is denoted by a solid red line, and the IRC paths are denoted by a blue dotted line. The  $S_1$  energies (in kcal/mol) relative to  $(S_0)_{cis-min}$  are given for the respective structures.

**Figure 5** shows the variations in the dihedral angles  $d_{C3C1C2C4}$ ,  $d_{R7C1C2C4}$ , and  $d_{R8C2C1C3}$  ( $R = H$  or  $C$ ) and  $r_{C5C6}$  along the meta-IRC from the FC point of *cis*-form for (a) mSB, (b) SB, and (c) dmSB, where  $d_{C3C1C2C4}$  (in red) corresponds to the rotational angle of two phenyl groups about the C=C bond, and  $d_{R7C1C2C4}$  (in blue) and  $d_{R8C2C1C3}$  (in green) correspond to the deviations in R7 and R8, respectively, with respect to the central ethylenic plane. In SB and dmSB,  $d_{R7C1C2C4}$  and  $d_{R8C2C1C3}$  have the same values along the meta-IRC pathway because of the  $C_2$  symmetry. The  $d_{C3C1C2C4}$  and a pair of  $d_{R7C1C2C4}$  and  $d_{R8C2C1C3}$  show a synchronous change where the former shows an increase from  $10^\circ$  to  $40^\circ$ , whereas the latter shows a decrease from  $175^\circ$  to  $150^\circ$ . The gradients of the change in these dihedral angles in the initial stage indicate that  $\pi\pi^*$ -excited *cis*-SB is promptly accelerated toward the twisted-form, whereas, in the case of  $\pi\pi^*$ -excited *cis*-dmSB, the corresponding acceleration is much smaller, and the  $\pi\pi^*$ -excited *cis*-mSB is the middle of *cis*-SB and *cis*-dmSB, which is consistent with the reaction route profiles in Fig. 4. The difference in the reaction path profile between *cis*-SB and *cis*-dmSB can be ascribed to the difference in the mass of the H and CH<sub>3</sub> groups in the central ethylenic part [39]. Actually, in the case of mSB,  $d_{H7C1C2C4}$  decreases more rapidly than  $d_{C8C2C1C3}$  as shown in Fig. 5a. The existence of H and CH<sub>3</sub> group in mSB will introduce mixed dynamics effects.

In Fig. 4, we note the energy difference between the S<sub>1</sub>-minima and S<sub>1</sub>/S<sub>0</sub>-MECIs in the *twist* region. The most interesting feature of the  $\pi\pi^*$ -excited dmSB is the long lifetime of the P\* state, and this was ascribed to the difference in the S<sub>1</sub>-minimum and S<sub>1</sub>/S<sub>0</sub>-MECI structures, as well as a relatively large energy difference between these two structures [39]. In the case of mSB, all S<sub>1</sub>-minima and S<sub>1</sub>/S<sub>0</sub>-MECI have a pyramidal structure, and the energy difference of the lowest S<sub>1</sub>-minimum and S<sub>1</sub>/S<sub>0</sub>-MECI is 6.3 kcal/mol which is smaller than the dmSB case (9 kcal/mol) and larger than the SB case (3 kcal/mol). Therefore, the  $\pi\pi^*$ -excited *cis*-mSB will reach the S<sub>1</sub>/S<sub>0</sub>-CI structures in the *twist* region much faster than the dmSB case, although its excited-state lifetime should be longer than the SB case.



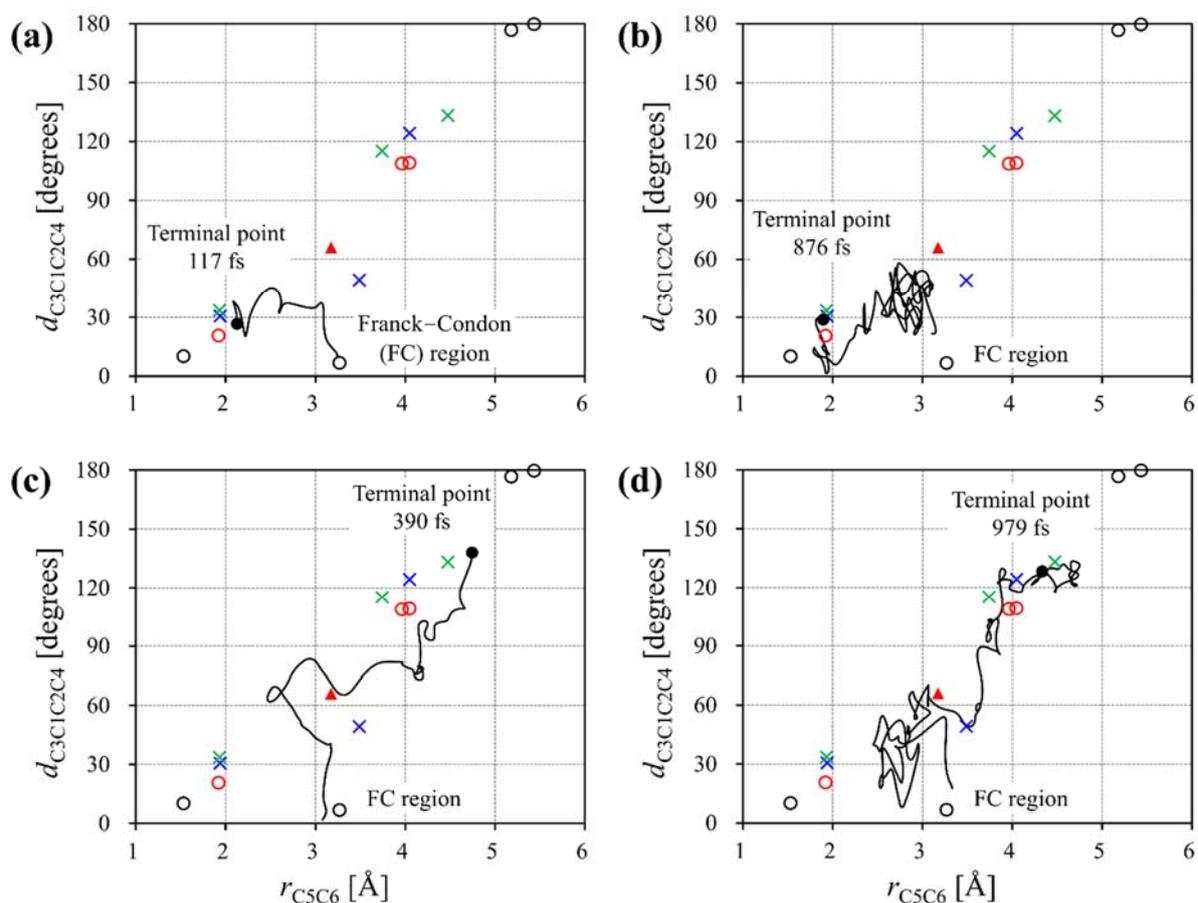
**Figure 5.** Variations in the dihedral angles and  $r_{C5C6}$  along the meta-IRC from the Franck-Condon structures for (a) *cis*-mSB, (b) *cis*-SB, and (c) *cis*-dmSB;  $d_{C3C1C2C4}$  (in red) corresponds to the opening of two phenyl rings from *cis*- to *twist*-structure,  $d_{R7C1C2C4}$  (in blue) corresponds to the motion of R7, and  $d_{R8C2C1C3}$  (in green) corresponds to the motion of R8.

### C. On-the-fly MD simulations

The dynamics of the  $\pi\pi^*$ -excited *cis*-mSB, showing the branched nature toward the DHP and *twist* regions on the  $S_1$ -PES, was examined by on-the-fly MD simulations at the SF-TDDFT level. In total, 40 trajectories were run from the FC region of the *cis*-form, in which 11 trajectories reach the DHP region while 29 trajectories reach the *twist* region, resulting in a DHP:*twist* branching ratio of 0.275:0.725. There is one exceptional trajectory that enters the *twist* region first, and then it moves to the DHP region later without reaching  $S_1/S_0$ -CIs in the *twist* region. In our previous studies, the branching ratio for DHP:*twist* was calculated as 0.26:0.74 for *cis*-SB [31] and 0.85:0.15 for *cis*-dmSB [39], and thus, the ratio for *cis*-mSB is very close to that for *cis*-SB. Only in the case of *cis*-dmSB are the terminal points of trajectories scattered mainly in the DHP region, which can be explained by the meta-IRC profile being directed toward the DHP-form from the start of the run [39]. In the case of *cis*-mSB, the meta-IRC profile and the position and energy of  $(S_1)_{TS}$  suggest that there is a preference for the *twist* region over the DHP region in the dynamics.

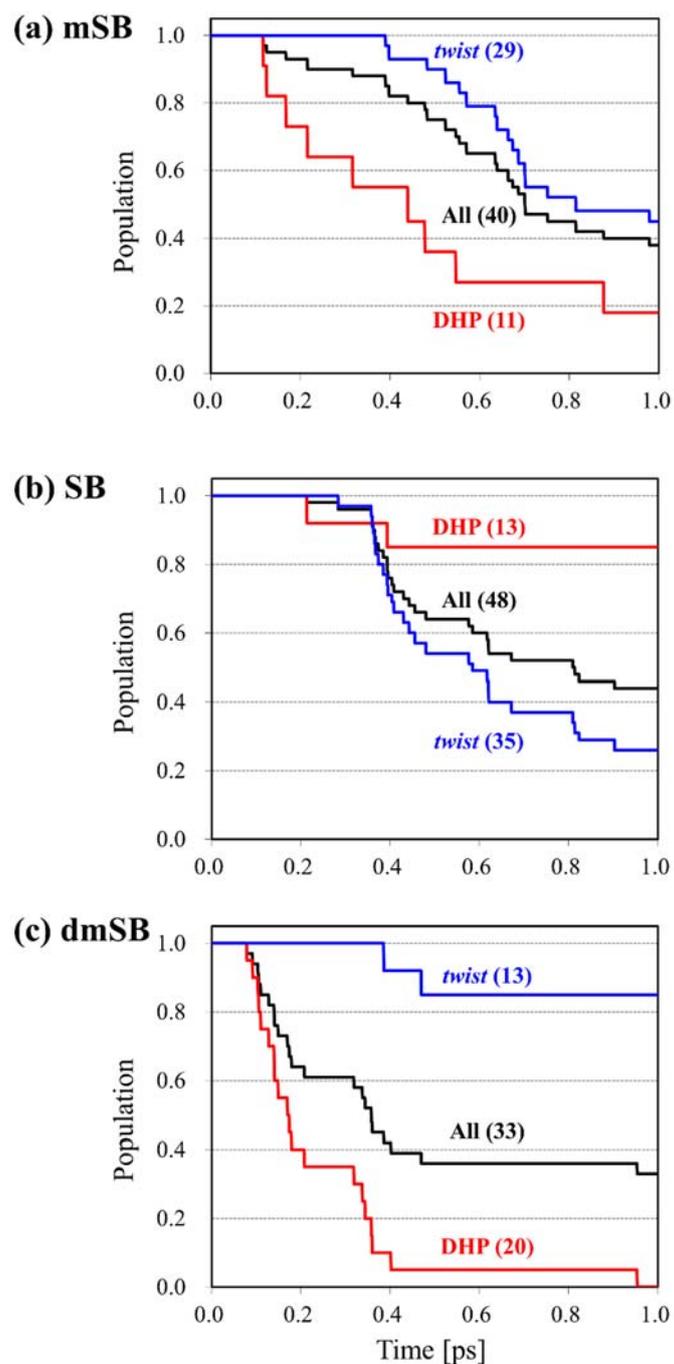
During a total simulation time of 1.0 ps, nine trajectories among the 11 trajectories entering the DHP region reached the CI region in 364 fs on average, while 16 trajectories among the 29 trajectories entering the *twist* region reached the CI region in 635 fs on average. Thus, the trajectories entering the DHP region from the FC region can more easily reach the  $S_1/S_0$ -CI and hop to the ground state than those trajectories entering the *twist* region. **Figure 6** shows projections of four trajectories to the two-dimensional coordinate space, as examples: the trajectories entering the DHP region and reaching the CI region with (a) the shortest time (117 fs) and (b) the longest time (876 fs), and the trajectories entering the *twist* region and reaching the CI region with (c) the shortest time (390 fs) and (d) the longest time (979 fs). As shown in **Fig. 6c** and **6d**, the trajectories which finally entered the *twist* region wandered while in the region between  $(S_1)_{DHP-min}$  and  $(S_1)_{TS}$ , leading to the relatively longer lifetime compared to the trajectories in the DHP region. We also analyzed the terminal CI points of nine trajectories in

the DHP region. Four terminal structures can be categorised as H-pyramidal CI (close to  $(S_1/S_0)_{DHP1}$ ), and one terminal structure can be categorised as Me-pyramidal CI (close to  $(S_1/S_0)_{DHP2}$ ). However, it is difficult to classify the other four structures into H-pyramidal or Me-pyramidal structures:  $(d_{H7C1C2C3}, d_{C8C2C1C4}) = (-177.7^\circ, -173.2^\circ)$ ,  $(-174.3^\circ, -175.0^\circ)$ ,  $(-170.0^\circ, -173.2^\circ)$ , and  $(164.5^\circ, -164.0^\circ)$ . In the case of the 16 trajectories that terminated in the *twist* region, the terminal CIs are all characterised as an H-twisted-*trans*-pyramidal structures corresponding to  $(S_1/S_0)_{twist1}$ .



**Figure 6.** Four trajectories that start from the FC region of the *cis*-form and run on the  $S_1$ -PES: (a) one terminated at the CI in the DHP region with the shortest time; (b) one terminated at the CI in the DHP region with the longest time; (c) one terminated at the CI in the *twist* region with the shortest time; and (d) one terminated at the CI in the *twist* region with the longest time.

**Figure 7** shows the decay of the  $S_1$ -population as a function of time calculated from on-the-fly MD simulations for (a) *cis*-mSB, (b) *cis*-SB [31], and (c) *cis*-dmSB [39], where all  $S_1$  trajectories, those terminated in the DHP region, and those terminated in the *twist* region are plotted in black, red, and blue, respectively. Here, the " $S_1$ -population" is defined as a ratio of the trajectories before reaching the crossing points of the  $S_0$  and  $S_1$  states. The trajectories in the *twist* region reached the CI more quickly in SB (Fig. 7b), whereas the trajectories in the DHP region reached the CI more quickly in dmSB (Fig. 7c). These results can be explained by the directions of the meta-IRC path from the FC region in the initial stages, as well as the geometrical and energy differences of the  $S_1$ -minimum and  $S_1/S_0$ -CI structures in the *twist* region for dmSB [39]. In the case of mSB, the trajectories in the DHP region reach the CI more quickly than in the *twist* region, as for dmSB, but the trajectories in the *twist* region reach the CI smoothly. In other words, *cis*-mSB has no long-lifetime component in the  $S_1$  state, unlike the *cis*-SB and *cis*-dmSB cases. This feature of *cis*-mSB is ascribed to the initial direction of the meta-IRC, which is oriented between the *twist* (*cis*-SB) and DHP (*cis*-dmSB) regions. In the *twist* region, the decay rate for mSB is slightly longer than that for SB because the energy difference of  $S_1$ -minima and  $S_1/S_0$ -MECIs is larger for mSB than for SB. The present on-the-fly MD simulations demonstrate that the introduction of light (H) and relatively heavy ( $\text{CH}_3$ ) fragments to the central part of the molecule tunes the lifetime and results in the disappearance of the slow component in the excited-state dynamics.



**Figure 7.** The decay of the S<sub>1</sub> population as a function of time derived from on-the-fly MD simulations for (a) *cis*-mSB, (b) *cis*-SB [31], and (c) *cis*-dmSB [39]. The S<sub>1</sub> population for all the trajectories is plotted in black, while the decays for the trajectories remaining in the DHP region and for those in the *twist* region are plotted in red and in blue, respectively.

## IV. Conclusions

In the present study, the reaction pathways and branching dynamics on the excited-state PES for  $\pi\pi^*$ -excited *cis*-mSB were investigated by SF-TDDFT calculations and on-the-fly MD simulations, and a comparison with the previous theoretical studies on *cis*-SB [31] and *cis*-dmSB [39] has been made. In all the systems, the meta-IRC pathway starting from the FC structure of the *cis*-form is connected to the minimum in the DHP region, but the reaction path profile suggests that dynamics effects push the molecule partly to the *twist* region. In the case of *cis*-SB, all trajectories enter the *twist* region once because of the inertial force; subsequently, some trajectories enter the DHP region. In the case of *cis*-dmSB, most trajectories enter the DHP region along the meta-IRC path, and some trajectories move to the *twist* region later. The case of *cis*-mSB is intermediate between those of *cis*-SB and *cis*-dmSB, and, as the result of dynamics simulations, the branching ratio was calculated to be  $\text{DHP:twist} = 0.275:0.725$ , which is similar to the case of *cis*-SB. In the *twist* region, mSB hops to the ground state through the H-twisted-*trans*-pyramidal  $S_1/S_0$ -CI. Concerning  $S_1$  population decay, *cis*-SB has a slow component assigned to trajectories terminating in the DHP region, whereas *cis*-dmSB has a slow component assigned to trajectories terminating in the *twist* region. On the other hand, *cis*-mSB has no slow component because its reaction path profile is intermediate between those of *cis*-SB and *cis*-dmSB.

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