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ON LIMIT STRENGTH OF FROZEN CLAY UNDERGOING TRIAXIAL TENSION Satoshi Nishimura<sup>1, \*</sup>, Hidetoshi Kawasaki<sup>2</sup> and Issei Sato<sup>3</sup> <sup>1</sup>Faculty of Engineering, Hokkaido University, Japan (\*Corresponding author)

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# 10 ABSTRACT

11 A new interpretation of failure of frozen clay in tension is presented, through a unique test program 12 applying 'triaxial tension', in which the axial stress was reduced to negative values while zero or 13 positive radial stress was maintained. The limit strength is determined by a meeting point of two 14 competing mechanisms – strain-hardening in shear, and progressively degrading tensile strength 15 as the shear deformation disturbs the soil structure. In some conditions, the limit strength observed 16 in tension is actually given by the shear resistance limit, not by a true tensile strength. This 17 interpretation is successful in explaining apparently unclear confining stress-dependency of the 18 observed failure envelopes at low stresses. The influence of the pre-freezing effective stress, post-19 freezing total stress and pre-freezing pore water pressure on the strength was also investigated. 20 The tensile strength degradation was formulated by using the specific work input to estimate the 21 operational tensile strength and the strain at which a tensile rupture occurs. Cooling after shear 22 deformation apparently healed the damage and recovered the initial tensile strength, allowing the 23 stress-strain curve to significantly overshoot the proposed tensile strength line. Explaining the 24 combined influence of shear and temperature history on the tensile strength – work relationships 25 requires further study.

26

# 27 INTRODUCTION

28 Recent underground projects involve complex excavation processes at greater depths reaching 29 40-50 m more commonly than ever, often adjacent to existing substructures and/or under soft and 30 unstable ground conditions. Some processes, such as tunnel connection and enlargement, require 31 prior stablisation of surrounding ground. Artificial Ground Freezing (AGF) is an effective technique 32 to stabilise selected localities with high accuracy and reliability. As the recent innovation allows 33 freezing at lower temperatures (for example, circulation of CO<sub>2</sub>, in lieu of conventional brine, can 34 cool the chiller pipes down to -45 °C; Tsuji and Yoshida, 2019), frozen domains are expected to 35 mobilise higher strength. Even colder liquid nitrogen (LN) has long been available for smaller-scale, 36 short-term work, but it is usually produced in remote factory plants, and the freezing was typically 37 one-off and expensive. The above development opens up more possibilities in design, for example 38 by using thinner frozen soil as load-bearing roofs or earth-retaining walls. While keeping the frozen 39 soil thinner is highly cost-effective due to significant reduction in the required heat-conduction time, 40 it risks subjecting the soil to bending in some parts. In this scenario, the design based only on the 41 compressive or shear strength of frozen soil is inappropriate, as the tensile strength limits the 42 bending resistance of the frozen soil. These facts point to importance of an appropriate testing 43 method to characterise the tensile failure in frozen soils, and mechanical theory and insights to 44 interpret the results.

45

46 In the context of deep underground application, it is important to note that considerable total and 47 effective stresses, as well as large pore pressure, exist at the time of freezing, and that they may 48 partially remain even during the excavation processes. As will be reviewed, conventional tension 49 tests on frozen soils have been almost all performed in unconfined states, mostly with soil samples 50 frozen unconfined too. This conventional approach severely limits the possibility of interpreting the 51 frozen soil's failure behaviour in tension by drawing a parallel to the effective-stress-based unfrozen 52 soil mechanics, as will be discussed in the next section. A particularly relevant example of deep 53 AGF application is Trans-Tokyo Bay (Aqua-Line) Expressway in Japan (e.g. Uchida et al., 1993). 54 This project involved construction of submarine tunnels passing through soft alluvial clay at some 55 40-55 m below the sea level (20-35 m below the sea bottom) and vertical shafts connected to them, 56 assisted partially by AGF (e.g. Akagawa, 2021). The pore pressure at the invert level was about 57 600 kPa. A high pore (back) pressure usually prevents tensile failure in unfrozen soft clays by 58 preventing cavitation. It is not clear, however, how this condition affects the tensile failure of frozen 59 clay, nor is it possible to address this issue by the conventional testing methods with no control of 60 confining stress and pore pressure. AGF case histories under similarly deep settings with high 61 water pressure or soft soil conditions have also been reported elsewhere (e.g. Ou et al., 2009; 62 Viggiani et al., 2015).

63

This study adopted a new laboratory test approach by developing 'triaxial tension' testing, in which tensile loading was undertaken with lateral confinement. In contrast to conventional triaxial extension, it allows the minor (total) principal stress  $\sigma_3$  to go negative (i.e. active pulling). Freezing was performed directly after consolidation while the specimen remained under confinement, by improving the method by Wang et al. (2017). To the authors' knowledge, no similar attempt has been reported so far. This study, after considerable trial and error, eventually found a set-up allowing stable temperature control and providing repeatable results.

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This paper will firstly review the basic mechanics of unfrozen soils under tensile loading, which may involve either shear failure or tensile failure. A review of previous studies on tensile strength of

frozen soils will be then given to underline the need for this study's approach. In the authors' view,

investigation on the tensile failure of frozen soils must answer the following questions:

(i) Is the limit strength of frozen soils in tension determined by a 'tensile strength' (i.e.  $\sigma_3$  limit),

- independently from shear strength (deviator stress *q* limit)?
- (ii) Is the tensile strength affected by the initial, pre-freezing pore pressure and the effective stress?
- 79 (iii) Is the peak tensile limit, whatever (i) reveals, to be relied upon in design, given the eventually
- 80 inevitable strain-softening and fracturing?

(iv) How are all the above affected by conditions such as water content *w* (intertwined with the
 above issue of the consolidation effective stress), temperature and strain rate?

The present study provides new, if not comprehensive, insights to these questions through a systematic series of high-quality, repeatable triaxial tension tests in which the pre- and post-freezing stress and the initial pore water pressure were systematically varied. Additional tests adopted complex loading histories involving temperature changes, lateral stress changes and a loading pause.

88

89 This study only considers saturated clay, as initial phase of more thorough investigation. This paper

- 90 uses 'extension' to mean geometric elongation, and 'tension' to mean imposition of an absolute
- 91 negative normal stress.
- 92

#### 93 TENSILE STRENGTH OF UNFROZEN AND FROZEN SATURATED SOILS: A REVIEW

94 Basic mechanics of (unfrozen) saturated soils under tensile loading

95 It may be helpful to briefly review elementary unfrozen soil mechanics concerning tensile failure. 96 The discussion shall be limited to saturated soil in undrained triaxial conditions. In the simplest case 97 (**Figure 1**(a)), when the initial pore water pressure (or the back pressure)  $u_0$  has ample room for 98 reduction, any total stress change, including the average confining total stress decrease due to 99 tensile loading, is compensated by a negative excess pore water pressure under undrained 100 conditions. As a result, the eventual effective stress is unaffected by the total stress changes, and 101 thus the failure occurs when the deviator stress q reaches  $2s_u$ , where  $s_u$  is the undrained shear 102 strength. The pattern is basically a shear failure in triaxial extension, even when the soil is 103 undergoing tension in total stress term (i.e. negative  $\sigma_3$ ).

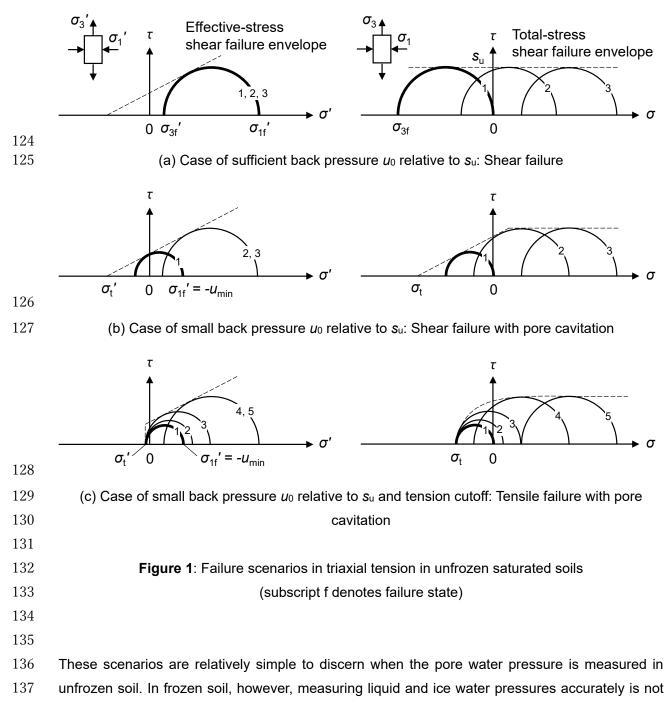
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105 The second pattern, illustrated in **Figure 1**(b), occurs when the pore pressure u decreases to a 106 limit value  $u_{min}$ , and ceases to compensate the total stress changes for any more  $\sigma_3$  reduction. In 107 coarse-grained soils, this may correspond to cavitation pressure of the pore fluid. For finer-grained 108 soils, *u*<sub>min</sub> may be smaller than the vacuum and sustain high suction. If the pore pressure remains 109 at  $u_{\min}$  for further tensile loading, the total stress reduction is directly reflected on the effective stress, 110 and when the latter reaches the failure envelope, failure will occur. It is more likely in reality that, 111 when the pore pressure reaches  $u_{\min}$ , the system cannot maintain undrained conditions, and  $u_{\min}$ 112 rebounds to the atmospheric pressure, causing abrupt tensile failure. For this scenario, the total 113 stress application under undrained conditions before tensile loading influences the eventual tensile

114 strength by raising u, and allowing  $\sigma_3$  to decrease more before u decreases to  $u_{\min}$ .

In fact, there is no a priori reason to assume that the effective-stress failure envelope is linear. The true tensile strength  $\sigma_t$  on the effective stress  $\sigma$ '-axis may not be where extrapolation of the shear failure line would suggest. The third pattern, illustrated in **Figure 1**(c), is a likely response in this case. When the failure is caused by *u* meeting  $u_{min}$ , and  $\sigma'$  meeting  $\sigma_t$ ', this is considered to be a pure tensile failure, with  $\sigma_3$  at failure unaffected by initial confining stress (Mohr's stress circles 1-3 in **Figure 1**(c) right).

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- 123



- 138 an established technique, and no direct manner of observation is possible. The idea in this study
- 139 is therefore to observe how the  $\sigma_3$  limit in tensile loading varies with the major principal stress  $\sigma_1$ .
- 140 If  $\sigma_3$  at failure,  $\sigma_{3f}$ , shifts in parallel with  $\sigma_1$ , what is perceived as 'tensile strength' in frozen soil may
- 141 in reality be governed by the shear strength. If  $\sigma_{3f}$  is unaffected by  $\sigma_1$ , on the contrary, the pore
- 142 materials (ice and liquid) and the soil skeleton may be truly being pulled apart on the perpendicular
- 143 plane. This logic must also be checked with visual observation of the failed specimen.
- 144

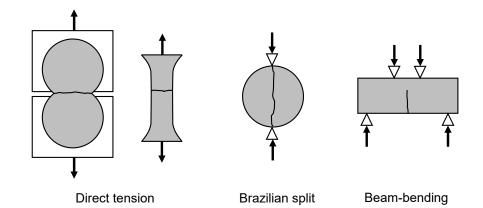
# 145 Testing methods for tensile strength

- 146 The tensile strength of frozen soils has been conventionally investigated by three testing methods; 147 direct tension, Brazilian splitting and beam-bending, as illustrated in Figure 2. Haynes and Karalius 148 (1977), Eckardt (1982) and Zhu and Carbee (1985; 1987), among others, adopted cylindrical 149 specimens with enlarged ends in direct tension. Eckardt (1982) penetrated nails to the specimen 150 ends to ensure the soil-machine coupling. Zhang et al. (2019) tested regular cylindrical specimens, 151 attached to loading plates with epoxy adhesive. In the other studies, the coupling mechanisms are 152 not clearly explained. Akagawa and Nishisato (2009) and Li et al. (2018) adopted '8-shaped' 153 specimens (the leftmost in Figure 2), following Tamrakar et al. (2007). All these studies performed 154 tension tests in unconfined states. One possible exception is a study by 'Jessberger/Ebel' 155 mentioned in Jessberger's (1981) state-of-the-art paper, of which drawing suggests application of 156 radial pressure. However, the source is not cited, and more detail could not be found. Apart from 157 this, there has been no systematic study conducting 'triaxial tension' testing, to the authors' 158 knowledge.
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160 Brazilian split (Bragg and Andersland, 1982; Zhou et al., 2015; You et al., 2021) and beam-bending 161 (Azmatch et al., 2011; Yamamoto and Springman, 2017) tests have also been performed to obtain 162 the tensile strength or observe fracture processes of frozen soils. While these tests are relatively 163 straightforward in principle, the stress field is non-uniform within the specimen. Namikawa and 164 Koseki (2007) compared the above two testing methods with direct tension for cement-treated sand 165 both experimentally and numerically, and concluded that the direct tension test gives the most 166 accurate tensile strength, while the other two are affected by local shear and progressive failure. It 167 would be also difficult to apply confining stress in these two testing methods.

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The literature review suggests that there has been no experimental data or attempt of applying tension from confined states, which, as discussed earlier, would clarify at what conditions the frozen soil fails in shear or reaches a true tensile limit. This background led to development of triaxial tension testing in this study.



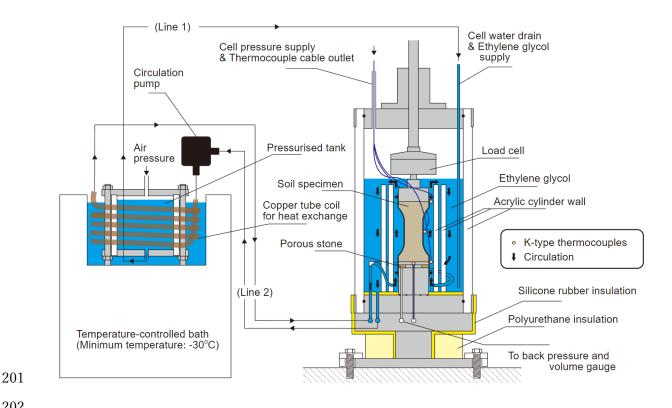
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# Figure 2: Conventional tensile strength tests for soils

# 179 MATERIAL AND TESTING METHOD

180 Apparatus

181 Triaxial apparatus was modified to allow direct tension testing of frozen soils under confining 182 pressure. The system is illustrated in **Figure 3**. The whole system was installed and operated in a 183 room at a temperature of 25 °C. A unique feature of this apparatus is that it allows freezing a soil 184 specimen within the cell, unlike in many other studies in which samples pre-frozen in a separate 185 freezer were transferred to testing machines. This feature permits freezing without releasing the 186 stress after consolidation in unfrozen states, as it happens in-situ. This freezing method prevents 187 unnatural cracking and inhomogenisation due to unconfined freezing (Wang et al., 2017). The 188 freezing is performed by firstly draining out the confining water slowly via an adjustable valve lest 189 the cell pressure drops, and then by pumping in pre-cooled refrigerant into the cell, again slowly 190 not to change the cell pressure. The cooling mechanism was modified from the earlier study by 191 Wang et al. (2017). The refrigerant is controlled via two lines. Line 1 is connected to a pressurised 192 tank in the cold bath, from which the pre-cooled refrigerant (ethylene glycol) is sent to the 193 pressurised triaxial cell. This line was used to adjust the refrigerant level in the triaxial cell. Line 2 194 circulates the refrigerant through the heat exchange coil. This line was used to maintain and adjust 195 the temperature in the triaxial cell. The temperature control of the refrigerant directly in contact with 196 the specimen by Line 2 is much more responsive than more indirect manners, such as controlling 197 the room ambient temperature or placing a heat exchanger within the cell. All through these 198 operations, the cell pressure could be controlled independently (kept constant in this study). 199



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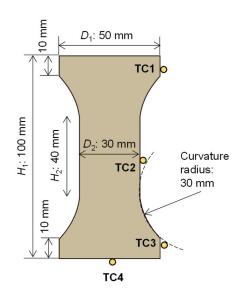
203 204

Figure 3: Triaxial tension test apparatus for frozen soil

206 Higher cooling efficiency was achieved by setting a double acrylic wall around the specimen and 207 restricting the circulation path. This cell-in-a-cell proved very effective in insulating the specimen 208 from the room environment. The refrigerant comes to contact with the specimen first as it enters 209 the triaxial cell, and thus effectively deprives the heat from the specimen before receiving heat from 210 outside. For -30 °C set at the cold bath, -23 °C was achieved at the specimen surface. This is much 211 higher cooling efficiency than in the previous design (Wang et al., 2017). The refrigerant outside 212 the inner double-wall works as a buffer, and was only mildly cold. This helps limiting the outer cell 213 surface frosting that hinders visual inspection. Another improvement was to use a taller pedestal. 214 The specimen bottom was slightly warmer in this apparatus due to the heat leak through the 215 pedestal. A taller pedestal alleviated this problem. The temperature was measured by K-type 216 thermocouples at four locations, three at the specimen's outer surface and one at the bottom end 217 centre (Figure 4). TC1 was at -14~-13 °C when TC2-TC4 were at -15 °C. This difference was 218 smaller for higher temperatures.

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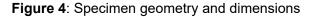
220 The axial ram was driven by a very precise and powerful direct-drive motor capable of applying 221 compression and extension loads equally. Pore water drainage during consolidation was allowed 222 towards the porous stone ring (Figure 3) embedded in the pedestal via filter paper strips attached 223 to the specimen's lateral surface (Figure 4 right).





(a) Dimensions and themocouples (TC)

(b) Specimen, as trimmed, and as enclosed by rubber membrane with filter paper strips



232 The soil specimens were trimmed into a dumbbell shape with enlarged ends, as illustrated in Figure 233 4. The diameters of the middle part and the enlarged ends were 30 mm and 50 mm, respectively. 234 The total height was 100 mm. The ends were trimmed smoothly, and simply put into contact with 235 smooth metal platen surfaces. The soil-platen coupling to transmit a tensile load was established 236 solely by freezing adhesion. It was therefore unnecessary to perform any special treatment during 237 specimen set-up in this study. The bottom platen was polished stainless steel, and the top platen 238 was anodised aluminum. The cross-sectional area ratio between the specimen ends and middle is 239 2.8. The tensile failure is therefore forced to occur in the middle part as long as the platen-specimen 240 adhesion is greater than 1/2.8 of the tensile strength. The transition from the middle to the enlarged 241 parts had a gentle curvature with a radius of 30 mm to avoid stress concentration. This geometry 242 was successful in causing the failure in the middle part. By adopting the quick freezing under 243 confining pressure, the water content ratio of the middle part to the ends, as measured after 244 freezing, was 0.996-0.1011, meaning almost no internal axial water migration. Wang et al. (2017) 245 observed slightly greater radial water migration, with a rim-to-core water content ratio of 1.03 on 246 average for 30 mm-diameter Kasaoka clay specimens. This is still small and accepted in this study.

247

#### 248 Material and specimens

Reconstituted saturated samples of Kasaoka clay were tested in this study. The clay has been extensively tested in frozen states by the first author in the past. The physical properties are shown

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251 in **Table 1**. The gravimetric unfrozen water contents  $w_u$  measured at several temperatures by 252 Nuclear Magnetic Resonance (NMR) are shown in Figure 5, along with the calculated degree of 253 liquid saturation  $S_{I}$ . The clay was made into slurry at the water content *w*=100%, then mixed and 254 poured into a 220mm-diameter consolidometer under vacuum for de-airing, and then one-255 dimensionally pre-consolidated. Thus prepared samples were saturated, showing B value > 0.95. 256 Nine specimens were taken from each consolidated cake. The sample pre-consolidated at  $\sigma_{vp}$ '=100 257 and 200 kPa were used for tests with pre-freezing isotropic effective stress  $p_c$ '=100 and 600 kPa, 258 respectively. Ideally,  $\sigma_{vp}$  should have been smaller than 50 kPa for  $p_c$ '=100 kPa so that the effect 259 of one-dimensional consolidation would be eventually erased at  $p_c$ '. However, such a soft specimen 260 was difficult to be prepared into the desired shape. The specimens were formed by carefully 261 trimming along curved guides in a rotary soil lathe.

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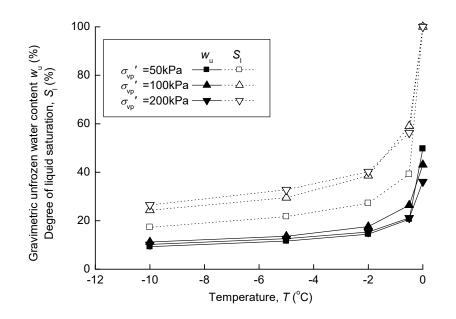
263 Wang et al. (2017) and Nishimura and Wang (2019) presented the triaxial compression strength of 264 Kasaoka clay prepared in a similar manner. At -10, -5 and -2 °C, the clay did not exhibit any change 265 of strength against post-freezing total stress  $p_{\rm f}$  in the range of 200-600 kPa, as shown in **Figure 6** 266 (solid markers; note that the ultimate strength, not the peak strength, was published by Nishimura 267 and Wang, 2019), unlike against the pre-freezing stress (open markers). This insensitivity to  $p_{\rm f}$  is a 268 characteristic of fine-grained soils, seen in other earlier studies too (e.g. Chamberlain et al., 1972; 269 Wang et al., 2004). Inspired by Ladanyi and Morel (1990), Nishimura and Wang (2019) explained 270 this by less dilative nature of fine-grained soils leading to shear failure in the pattern of Figure 1(a) 271 right. Although there has been clear evidence that polycrystalline ice's strength is pressure-272 dependent (Rist and Murrell, 1994; Singh and Jordaan, 1996; Shulson, 2001), the ice in smaller 273 pores in fine-grained soils may have less micro-defects that respond to confining pressure.

#### Table 1: Physical properties of Kasaoka clay

Plastic limit (%)	Liquid limit (%)	Particle density (Mg/m <sup>3</sup> )	Percentage finer than	
			2/5/63/75 µm (%)	
28	62	2.65	40/60/98/100	

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Figure 5: Gravimetric unfrozen water content  $w_u$  and degree of liquid saturation  $S_l$  in frozen Kasaoka clay one-dimensionally consolidated at  $\sigma_{VP}$ '

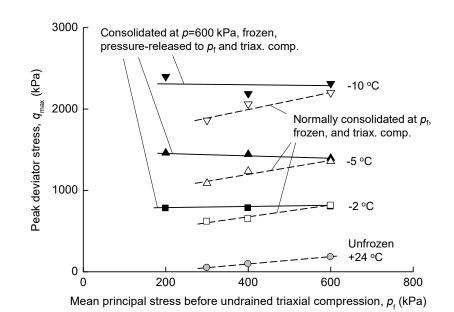




Figure 6: Triaxial compression peak strength of frozen Kasaoka clay: Influence of pre- and postfreezing stress changes (pre-freezing back pressure=200 kPa in all cases)

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#### 287 Test program

30 triaxial tension tests were performed in total, as summarised in **Table 2**. The main variables were the temperature *T* (-1, -4, -9 or -15 °C), pre-freezing effective consolidation stress  $p_c'$  (100 or 600 kPa) and post-freezing total stress  $p_f$  (0, 200 or 650 kPa). A series consisting of Tests 1-23

291 covers all the combinations of these, except  $p_f=0$  kPa at -1 °C. The adhesion at the bottom end

turned out to be too weak in this condition to apply tension. Note that the *T* values cited are nominal,
and the measured values are shown in **Table 2**. They are taken from thermocouple TC2, since the
failure occurred in the specimen middle.

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Table 2: List of performed triaxial tension tests

ID 7		Г (°С)	) <i>p</i> c	p <sub>c</sub> ′	pc' Uc	$p_{\mathrm{f}}$	$e_{c}^{2)}$	Note
	Nominal	Measured <sup>1)</sup>	(kPa)	(kPa)	(kPa)	(kPa)		
1	-1	-0.71	200	100	100	200	1.133	
2		-0.63	200	100	100	650	1.129	
3		-0.78	650	600	50	200	0.819	
4		-0.75	650	600	50	650	0.802	
5		-0.63	650	600	50	200	0.802	(Duplicate of Test 3)
6		-4.2	200	100	100	0	1.123	
7		-4.0	200	100	100	200	1.131	
8	-4	-4.4	200	100	100	650	1.126	
9	-4	-4.3	650	600	50	0	0.796	
10		-4.4	650	600	50	200	0.797	
11		-4.3	650	600	50	650	_3)	
12		-8.3	200	100	100	0	1.109	(Slightly warm)
13		-9.0	200	100	100	200	1.129	
14	0	-9.2	200	100	100	650	1.114	
15	-9	-8.9	650	600	50	0	0.798	
16		-9.2	650	600	50	200	0.801	
17		-9.0	650	600	50	650	_3)	
18		-14.4	200	100	100	0	1.104	
19		-14.7	200	100	100	200	1.099	
20		-14.5	200	100	100	650	1.105	
21	-15	-14.6	650	600	50	0	0.794	
22		-14.6	650	600	50	200	0.798	
23		-14.6	650	600	50	650	0.795	
24		-14.7	650	100	550	0	1.116	
25	-15	-14.7	650	100	550	200	1.102	
26		-14.7	650	100	550	650	1.112	
27	-4 to -15	-4.2 to -14.7	650	600	50	650	0.792	Pause and <i>T</i> change
28	-4	-4.3	650	600	50	650	_3)	Pause
29	-15	-14.7	650	600	50	650	0.799	Pause
30	-4	-4.4	650	600	50	σ <sub>1</sub> : 650 to 0	0.791	$\sigma_1$ change

<sup>297</sup> 298 299

<sup>1)</sup> Mid-height thermocouple (TC2), average over 2,000 sec around the tensile peak

299 <sup>2)</sup>  $e_c$ : Void ratio after consolidation

<sup>3)</sup> Final water content could not be measured accurately due to membrane rupture when removing the specimen

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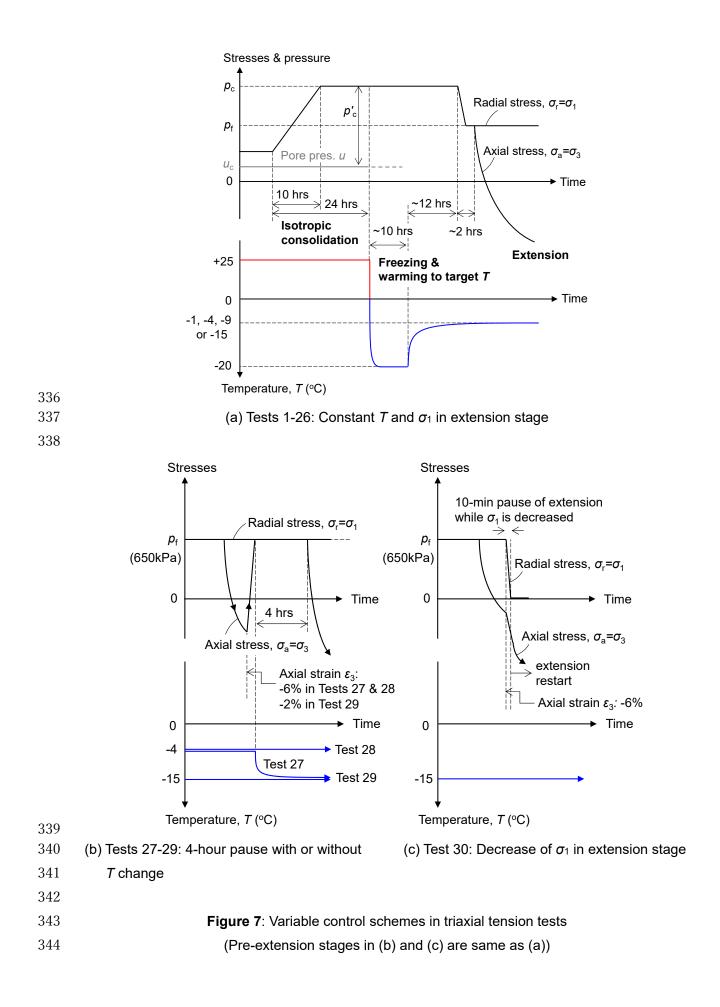
The back pressure after consolidation  $u_c$  was basically 100 kPa for  $p_c$ '=100 kPa, and 50 kPa for  $p_c$ '=600 kPa. Ideally, the latter would have been set at 100 kPa too – however, the slightly smaller value had to be accepted due to the cell capacity of 700 kPa. In Tests 24-26,  $u_c$ =550 kPa was set for  $p_c$ '=100 kPa. Test 5 is a duplicate test of Test 3 performed early in the program, out of a concern about test repeatability, as a slight temperature fluctuation has relatively large influence on the result at such high temperatures as -1 °C. Tests 18-26 constitute an interesting comparative set, allowing observation of influence of  $p_c$ ',  $p_f$  and  $u_c$  independently while keeping the other two constant ( $u_c$ =50 kPa and 100 kPa are assumed nearly equal here).

311

312 The above basic series (Tests 1-26) were conducted in the following steps. Firstly, the pre-313 consolidated, unfrozen specimen was set up in the apparatus after being covered with a rubber 314 membrane with a diameter of 30 mm. The enlarged part therefore underwent greater membrane 315 force, as the membrane size was chosen for the middle part. This influence was not deemed 316 important, as the failure always occurred at the specimen middle. After being isotropically 317 consolidated to  $p_c'$ , the specimen was frozen under constant cell pressure, as described earlier. 318 This process was undrained. The convergence of axial displacement can be considered to mark 319 temperature equilibrium, and the freezing was complete in 20 minutes. At this moment, the 320 specimen was frozen at -20 °C, and this state was kept at least for 10 hours. The temperature was 321 then increased to a target value. The change was fully transferred to the specimen core within an 322 hour (as inferred from TC4 and numerical analysis), but about 12 hours were allowed for ensuring 323 equilibration. The total stress was then changed to  $p_{\rm f}$  in an hour, and after the axial strain rate 324 became smaller than 10<sup>-4</sup> %/min, the extension stage was initiated. These stages are illustrated in 325 Figure 7(a).

326

327 The extension was applied at a constant axial strain rate of -0.05 %/min. This rate was sufficiently 328 slow for the frozen clay to remain well in the static, ductile regime (e.g. Parameswaran 1980) but 329 allows a failure to be reached within the day under the operator's constant watch. In this paper, the 330 nominal axial strain  $\varepsilon_a$  (=minor principal strain  $\varepsilon_3$ ) is defined by the axial displacement d divided by 331 the total specimen height  $H_1$  (see **Figure 4**). Given the enlarged ends, this definition may 332 underestimate the true strain in the middle, and is only nominal. The axial stress  $\sigma_a$  (= $\sigma_3$ ) is defined 333 by the axial load divided by the mid-height cross-sectional area  $A_2$  (= $\pi D_2^2/4$ ) after consolidation, 334 without update for horizontal thinning due to extension. The deviator stress, q, is defined as the 335 axial stress minus the radial stress (i.e. negative in axial tension).



345 Tests 27-30 were designed to investigate the influence of more complex state history on the failure. 346 They followed the same steps except in the extension stage. Tests 27-29 involved unloading and 347 reloading of deviator stress, with a 4-hour pause in between with or without temperature decrease. 348 as illustrated in **Figure 6**(b). The purpose of the pause was to allow temperature equilibration in 349 Test 27, and to have the same time to failure as this in Tests 28 and 29. The intension will be 350 discussed more later. The unloading in these tests started at  $\varepsilon_a$ =-6% in Tests 27 and 28 (*T*=-4 °C) 351 and  $\varepsilon_a$ =-2% in Test 29 (*T*=-15 °C), approximately 75% of the tensile rupture strain observed in Tests 352 11 and 23, respectively. In Test 30, extension was paused at  $\varepsilon_a$ =-6%, the radial stress  $\sigma_1$  was 353 reduced from 650 to 0 kPa in 10 minutes, then the extension was resumed. This is illustrated in 354 Figure 6(c).

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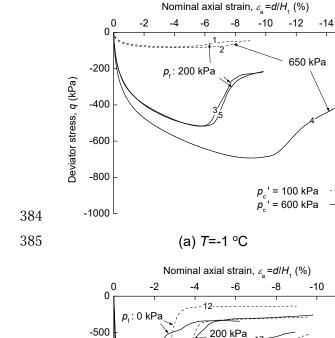
#### 356 PEAK STRENGTH OBSERVED IN EXTENSION

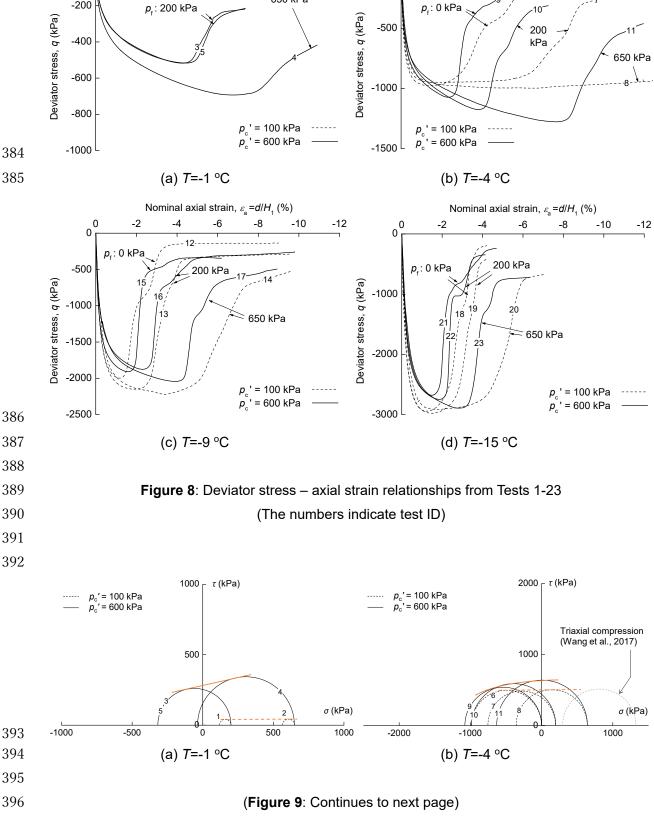
357 Stress-strain relationships and failure patterns

358 The observed relationships between the deviator stress q and the nominal axial strain  $\varepsilon_a$  from Tests 359 1-23 are shown in **Figure 8** for different temperatures *T*. Consistency of Tests 3 and 5 assures the 360 repeatability, even in sensitive temperature ranges. For T=-4, -9 and -15 °C, the pre-peak parts of 361 the curves are unique to each  $p_{c'}$ , unaffected by the post-freezing stress  $p_{f}$ . This is consistent with 362 the earlier results in **Figure 6** showing no influence of the post-freeze stress change on the triaxial 363 compression strength at moderate stress levels. Beyond the initial part, the influence of  $p_{\rm f}$  becomes 364 clearer, as the initially same stress-strain relationships start branching out one after another, 365 starting from lower  $p_{\rm f}$ , due to failure. In some cases (*T*=-4 and -15 °C,  $p_{\rm c}$ '=100 kPa), this occurred 366 after common q peaks (-950 kPa and -2950 kPa, respectively) were reached, meaning that the 367 minimum q value is unaffected by  $p_{f}$ , and only the post-failure strength reduction behaviour is 368 different. For the rest of conditions, lower  $p_f$  led to slightly lower strength (i.e. higher q) as the stress-369 strain curves fell short of travelling up to the common *q* peak.

370

371 The peak strength (minimum q) is shown as Mohr's stress circles in **Figure 9**, along with earlier 372 triaxial compression data obtained by the authors' team where the conditions are same (the 373 difference in the strain rate, 0.01 %/min vs 0.05 %/min, is corrected by the observed rate-374 dependency; Wang et al., 2017). The failure envelopes are ambiguous when viewed with Figure 375 1(a)-(c) right as possible templates. As discussed above, some cases showed no influence of  $p_f$  on 376 the strength, leading to a flat envelope in line with **Figure 1**(a) right. The others show only gentle 377 slopes – far gentler than expected from the effective stress-based shear resistance angle  $\Phi'$ . 378 Testing with negative  $\sigma_1$  would lead to a clearer picture, but this all-round negative stress is an 379 unlikely regime in field, and hard to achieve in tests too. Whilst it is probably reasonable to 380 approximate  $\phi$  as 0, except at T=-1 °C, from **Figure 8** in practice, this 'as-plotted' Mohr's stress 381 circle leaves ambiguity.





Nominal axial strain,  $\varepsilon_{a} = d/H_{1}$  (%)

-6

6

-8

-10

-12

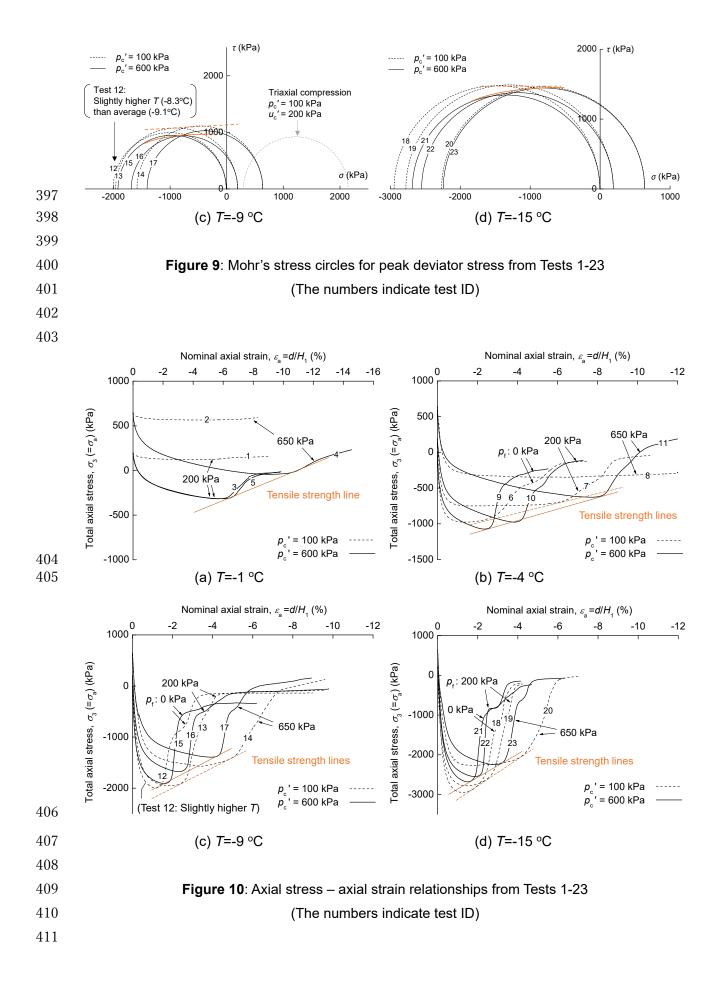
-2

0

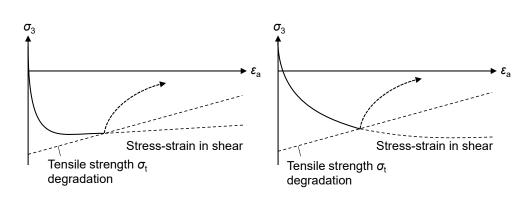
0

-16

-4



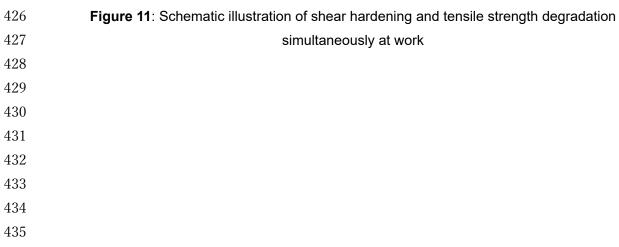
412 Another perspective offers a clearer view of the possible process of failure. Figure 10 replots 413 **Figure 8** in terms of axial stress  $\sigma_a = \sigma_3$ . For each  $p_f$  and T condition, the failure points are neatly 414 aligned along a single line. A possible mechanism hypothesised from this observation is that the 415 tensile strength is not constant but decreases as the shear strain accumulates. If the stress-strain 416 relationship shows comparatively higher initial stiffness and smaller strain to peak shear resistance, 417 the peak *q* value is dictated by the shear resistance, and only the post-peak behaviour is affected by tensile failure, as illustrated in **Figure 11**(a). This was the case for T=-4 and -15 °C,  $p_c$ '=100 kPa. 418 419 In the other cases, the intersection of the decreasing tensile strength line with the stress-strain 420 curve occurred before the latter reached the peak, as shown in Figure 11(b). In such cases, lower 421  $p_{\rm f}$  (and hence lower  $\sigma_3$  for a given q since  $\sigma_3 = p_{\rm f} + q$ ) leads to higher q (lower |q|) as it meets the 422 tensile strength line earlier. 423



425

(b) Larger  $p_c'$  and smaller ice content: (a) Smaller  $p_c'$  and larger ice content: Smaller shear strength and larger initial stiffness stiffness

Larger shear strength and smaller initial



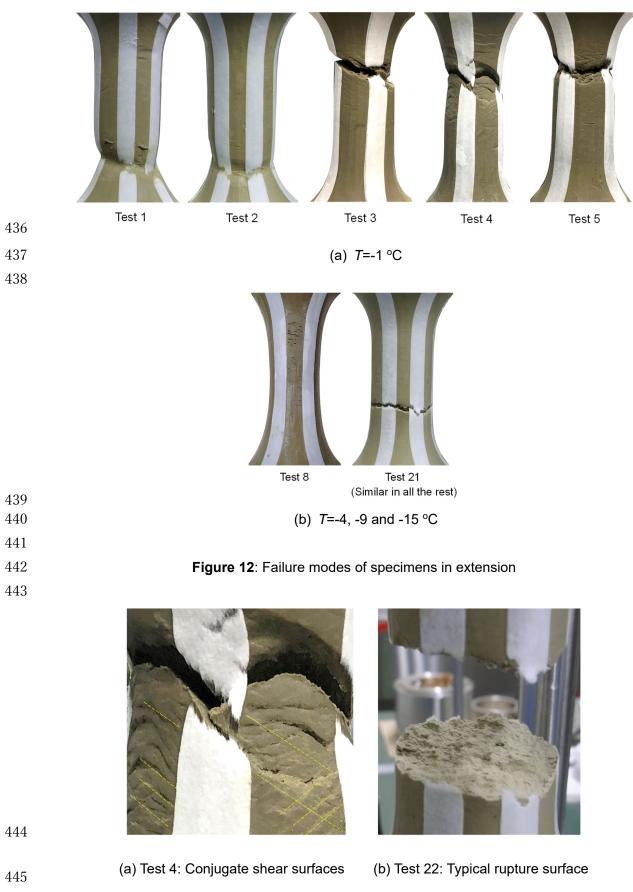


Figure 13: Close-up of failure

447 Almost all the tests exhibited a clear tensile failure surface at the end, as shown in Figure 12(b, 448 right) and **Figure 13**(b). The surface was perpendicular to the axis, and the surface texture was 449 rough with no trace of shear. The exceptions include Test 8, which did not reach clear failure despite 450 significant extension beyond -12%. The other exceptions are Test 1-5, performed at T=-1 °C (more 451 accurately, around -0.7 °C). These tests exhibited only marginally higher strength than unfrozen 452 states, due to the significant unfrozen water content. It is probably also relevant that the sample 453 had been frozen initially to -20 °C and thawed back to -0.7 °C. This freezing and undrained 'half-454 thawing' might have led to significant softening by generating excess pore water pressure 455 (Nishimura et al., 2020). As a result, the strength was very low, and  $\sigma_3$  remained positive at failure 456 in Tests 1 and 2. The failure in these tests was therefore triaxial extension shear without involving 457 eventual tensile failure, as shown in **Figure 12**(a). In the tests with  $p_c$ '=600 kPa (Tests 3-5),  $\sigma_3$  did 458 decrease to negative values (i.e. tensile stress). While these tests eventually showed tensile failure 459 in form of rupture, diagonal traces of triaxial extension shear had developed as seen in Figure 460 **12**(a) due to the large shear strain before reaching the tensile failure line;  $\varepsilon_a$ =-6.2%, -10.8% and -461 6.5% in Tests 3, 4 and 5, respectively. A close-up view of the shear traces is shown in Figure 13(a). 462 Conjugate shear surfaces with approximately 30° dip were formed, along which the tensile rupture 463 occurred, unlike at lower temperatures where the tensile rupture was almost horizontal. These tests 464 almost failed in triaxial extension shear, only engaging the decreasing tensile strength after shear 465 localisation occurred.

466

#### 467 Initial tensile strength and degradation

468 The above observations suggest that the rather ambiguous nature of the failure envelopes in 469 Figure 9 – neither completely flat nor obviously stress-dependent – can be explained by a fine 470 balance of two competing factors; the rate of strain-hardening in shear deformation, and the rate 471 of tensile strength degradation, as illustrated in **Figure 11**. In general, smaller  $p_c'$  led to faster initial 472 strain-hardening. For same Kasaoka clay, Wang et al. (2017; 2019) found that the initial stiffness 473 was higher for lower  $p_c$  due to larger w and hence larger ice content, while the peak shear strength 474 was lower for lower  $p_c$ ' due to less inter-soil particle friction. The pre-peak strain in shear therefore 475 tended to be smaller for lower  $p_c$ , as seen in **Figure 8**. Mobilising the shear strength, which was 476 not post-freezing stress-dependent, before engaging the tensile strength led to generally flatter 477 peak strength envelope for  $p_c$ '=100 kPa.

478

One of the key issues in understanding the limit strength is the tensile strength degradation during shear. **Figure 14**(a) summarises the tensile strength lines taken from **Figure 10** for all the conditions. Although there is a consistent trend, different slopes make the description difficult. When the same envelopes are replotted against the specific shear work input (i.e. the area between the  $q - \varepsilon_a$  curve and q=0 line in **Figure 10**), the relationships have an almost constant slope, as shown in **Figure 14**(b). This slope can be used for tensile strength degradation modelling. The intercept 485 at  $\varepsilon_a=0$  is the initial tensile strength  $\sigma_{t0}$ , or a hypothetical tensile strength that would be observed if 486 the shear stress-strain relationship were perfectly rigid plastic. The  $\sigma_{t0}$  values are plotted against 487 temperature T in **Figure 15**, along with the actually observed maximum |q| (= $|q|_{max}$ ) in the triaxial 488 tensile Tests 1-2, 6-8, 14, 18-20. These tests are likely to have mobilised the shear strength before 489 reaching the tensile failure. Also shown are  $|q|_{max}$  values from triaxial compression tests (Wang et 490 al., 2017), corrected for the strain rate difference. Although the anisotropy and the intermediate 491 principal stress ( $\sigma_2$ )-effect lead to different trends, these  $|q|_{\text{max}}$  values represent the shear strength. 492 They are seen to be only slightly lower than  $\sigma_{t0}$ . Therefore, whether the limit strength in extension 493 is determined by reaching the potential maximum shear strength or being cut short by degrading 494 tensile strength (i.e. shear-induced reduction from  $\sigma_{t0}$ ) depends on the fine balance of these 495 mechanisms.

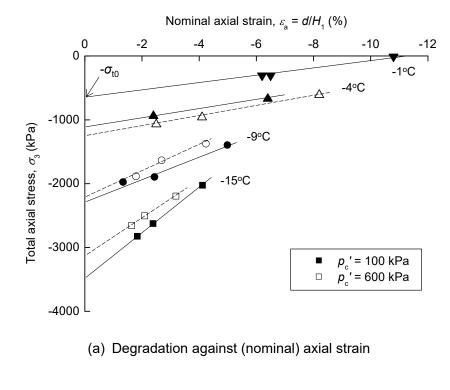
496

497 Interestingly,  $\sigma_{t0}$  is higher than the ice tensile strength compiled by Cuda and Ash (1984), despite

the strain rate slower by one order (0.05%/min against 0.9%/min) and significant unfrozen water

499 content in the clay. The microscopic failure of ice in the clay's small pores is probably harder to

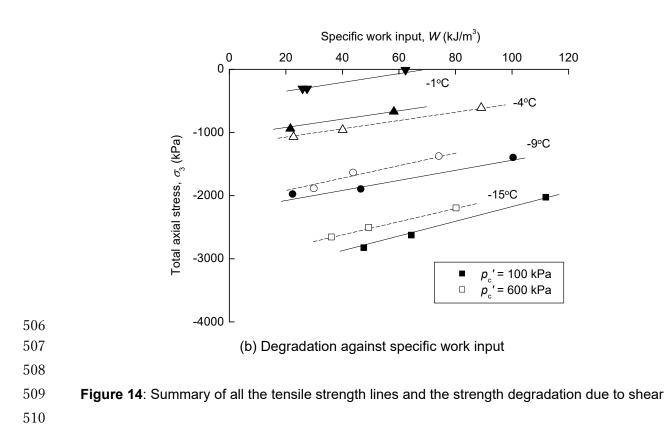
- 500 develop, just as the pore liquid water is harder to cavitate.
- 501

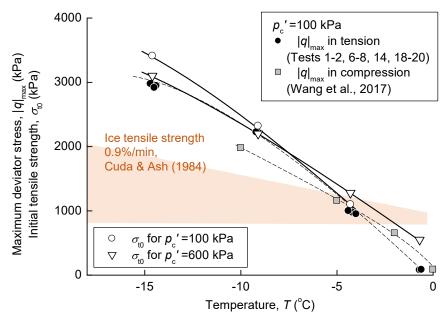


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502

(Figure 14: Continues to next page)



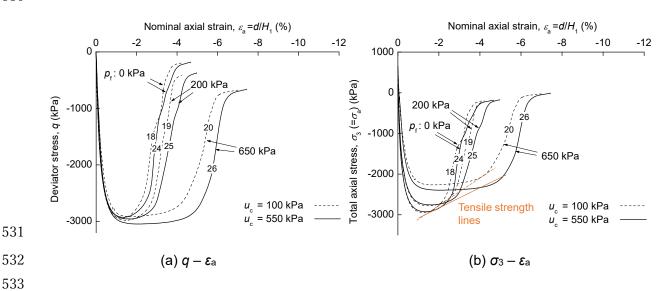


**Figure 15**: Initial tensile strength against temperature, in comparison with shear strength

- 515 TENSILE STRENGTH FOR MORE GENERAL STATES
- *Pore water pressure effects: Tests* 24-26
- 517 Test 24-26 applied a higher pre-freezing back pressure  $u_c$  of 550 kPa, while the samples were
- 518 consolidated at  $p_c$ '=100 kPa. **Figure 16** compares these tests with Tests 18-20, with same  $p_c$ '=100

519 kPa and T=-15 °C but with  $u_c$ =100 kPa. The higher back pressure only marginally affected the 520 failure, providing a small extra strain before the tensile rupture at each  $p_{\rm f}$ . The multi-phase coupled 521 system theory by Nishimura and Wang (2019) does not differentiate these cases for shear 522 behaviour, as a frozen state is essentially an undrained state, and the pore water (ice + liquid) 523 pressure at frozen state will be only dependent on the initial effective stress, just as in the unfrozen 524 soil mechanics. In Test 24, for example, the pore water pressure would decrease by 450 kPa in this 525 theory when eventually pf was reduced from 650 kPa to 200 kPa, bringing it to the same level as 526 in Test 19. The almost same results seen for Tests 18-20 and 24-26 confirm the theory. A possible 527 explanation for the slight difference is that less air escapes from the pore water when it turns ice 528 under higher pressure, giving the pore ice with less defect. Figure 16 suggests, however, that this 529 effect is negligible in practice.

530



- 534 **Figure 16**: Influence of initial (pre-freezing) back pressure on stress-strain curves at -15 °C
- 535 536

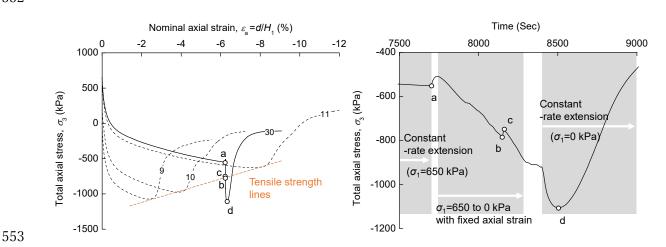
#### 537 Major principal stress $\sigma_1$ changes: Test 30

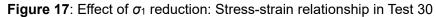
538 The framework of shear-induced degradation of the tensile strength was tested for more general 539 testing conditions. Test 30 involved reduction of the radial stress  $\sigma_{\rm r}$  (= the major principal stress  $\sigma_{\rm 1}$ ) 540 from 650 kPa to 0 kPa after reaching  $\varepsilon_a$ =-6%, or 75% of  $\varepsilon_a$  at tensile rupture. When the extension 541 was resumed with  $\sigma_1=0$  kPa, the  $q - \varepsilon_a$  curve obviously overshot the tensile strength line established 542 from Tests 9-11 with same T=-4 °C and  $p_c$  =600 kPa, as shown in **Figure 17**(a). However, a closer 543 look at the  $\sigma_1$ -reduction stage in **Figure 17**(b) reveals that the failure actually had initiated at Point 544 b. As the axial ram was fixed, the  $\sigma_3$  change is a spontaneous reaction to  $\sigma_1$  reduction. It usually 545 changes parallelly with  $\sigma_1$  to keep q constant (because there is no shear strain change), but at 546 Point b, there was a temporary rebound of  $\sigma_3$  to Point c. These points fall exactly on the expected 547 tensile strength line. The fact that the test reached tensile failure immediately after the overshooting

to point d suggests the strength is still bound by this line, despite the temporary excessive strength.
 This overshoot is similar to TESRA behaviour proposed by Tatsuoka et al. (2002), reflecting the
 temporary effect of strain acceleration due to sudden loading restart.

- 551
- 552

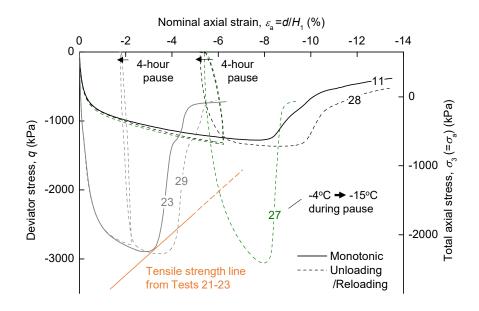
554 555 556





#### 557 Time to failure and temperature changes: Tests 27-29

558 A temperature change from T=-4 °C to -15 °C was conducted during extension in Test 27. As 559 described earlier, a pause of 4 hours was set between the unloading and reloading to allow 560 temperature equilibrium at the specimen core. Reloading after cooling to -15 °C resulted in 561 considerably high tensile stress, as shown in Figure 18, unbound by the established tensile 562 strength line. The peak q, observed at around -3000 kPa, is the value expected if Test 23 could 563 continue shear deformation without being bound by the tensile strength line. Obviously, 4-hour 564 cooling 'healed' the tensile strength degradation induced by the previous shear up to  $\varepsilon_a$ =-6%. It was 565 not clear, however, whether the cooling or the pause was responsible for this. Tests 28 and 29 were 566 therefore conducted, in which similar unloading, 4-hour pause and reloading was conducted but 567 without *T* change. Figure 18 shows that the pause alone has limited healing effect, only moderately 568 affecting the tensile strength line. The observed full recovery of the tensile strength by cooling 569 cannot be explained just by the additional freezing of unfrozen pore water, because Figure 5 570 suggests  $S_{\rm I}$  reduction of only 10% from -4 °C to -15 °C. Regelation of the damaged ice must have 571 also occurred. A longer pause, or more generally, a longer time to failure including slower loading 572 rates, may have greater influence on the tensile strength line. The different loading rate conditions 573 affect the shear stress-strain curves too. Obviously more study is necessary in incorporating the 574 combined effects of temperature and deformation rate into the tensile strength framework.



579 580

577 **Figure 18**: Influence of pause and/or temperature decrease (Tests 27-29) on stress-strain 578 relationship: Comparison at  $p_c$ '=600 kPa and  $p_f$ =650 kPa

581 CONCLUSIONS AND IMPLICATIONS

A series of 'triaxial tension' tests was conducted on frozen Kasaoka clay, in which the axial stress was reduced to negative values while zero or positive radial stress was maintained. The samples were consolidated to different stresses in the triaxial cell and frozen without releasing the pressure before the tension tests, as expected in artificial ground freezing. This testing method enabled investigating the influence of the pre-freezing consolidation effective stress, post-freezing total radial stresses as well as the initial pore water pressure – a feature not attained in the conventional uniaxial tensile, split or bending tests.

589

590 Apparently flat to sub-flat failure envelopes, showing undecisive influence of stress, to Mohr's 591 stress circle were obtained. It was shown, however, that they are explicable by considering two 592 competing mechanisms; shear strain-hardening and tensile strength degradation. The former is 593 independent of the confining (or radial) stress in Kasaoka clay for the studied stress range (0-650 594 kPa), while the latter is uniquely correlated to the specific work input as shear deformation 595 progresses. A peak shear strength may be mobilised before tension rupture eventually occurs, or 596 the tensile loading may be cut short by engaging the degrading tensile strength before attaining the 597 peak shear strength. Question (i) raised in the introduction has been thus answered – both tensile 598 and shear strengths play a role.

599

600 The tensile strength is described by two factors; the initial tensile strength, or a hypothetical strength

601 mobilised at zero shear strain, and its degradation rate as the soil structure is disturbed by shearing. 602 They are strongly dependent on the temperature, although the degradation rate can be normalised 603 by using the specific work input. More modest influence comes from the initial pore water pressure 604 and time to failure. The higher the pore water pressure, and the longer the time to failure, the frozen 605 clay endured to slight larger strains before eventually failing in tensile rupture. An interesting 606 phenomenon was observed that cooling from -4 °C to -15 °C just before the tensile failure 607 apparently healed the damage and fully recovered the tensile strength even at large strain. 608 Elucidating more generally the combined effect of temperature changes and loading rate (or time 609 to failure) requires much further investigation. These observations answer questions (ii) and (iv), 610 while leaving a room for further research.

611

612 In a more practical perspective, it may be acceptable to approximate the clay's failure envelope, 613 such as shown in Figure 9, as flat and stress-independent, except at very warm, half-thawing 614 temperatures (above -1 °C, for example). It should be remembered, however, that such envelopes 615 do not reflect the deformability after reaching the limit tensile stress. Given the abrupt loss of 616 resistance in tensile failure, a safety margin in design should take account of the deformability. The 617 ductile behaviour such as observed in Test 8 (-4 °C) in Figure 8 is assuring when exploiting the 618 tensile strength in design, while colder temperatures lead to much more brittle behaviour. The 619 proposed framework involving the work-induced tensile strength degradation model will be useful 620 in estimating not just the peak strength, but also the strain to reach it. This is a tentative answer to 621 question (iii).

622

#### 623 ACKNOWLEDGMENTS

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### 720 LIST OF NOTATIONS

- 721 A<sub>2</sub>: Area of the middle part of the specimen
- $D_1$ ,  $D_2$ : Diameters of the end and middle parts of the specimen
- 723 ec: Void ratio after consolidation
- $H_1$ ,  $H_2$ : Total and middle part heights of the specimen
- $p_c$ ': The effective mean effective stress at consolidation
- *p*<sub>f</sub>: The total mean effective stress after freezing and before extension
- q,  $|q|_{max}$ : Deviator stress (=axial stress radial stress) and its absolute maximum value
- 728 Si: Degree of liquid saturation (=volume of ice / volume of pore)
- 729 su: Undrained shear strength
- *T*: Temperature
- *u*: Pore water pressure
- *u*<sub>0</sub>: Back pressure
- *u*<sub>c</sub>: Pore water pressure after consolidation (= back pressure in the experiment)
- *u*<sub>min</sub>: The minimum sustainable pore water pressure
- 735 w: Water content
- *w*<sub>u</sub>: Gravimetric unfrozen water content
- $\epsilon_1, \epsilon_3$ : Major and minor principal strains
- $\varepsilon_{a}$ ,  $\varepsilon_{r}$ : Axial and radial strains (= $\varepsilon_{3}$  and  $\varepsilon_{1}$ , respectively)
- $\sigma$ ,  $\sigma$ ': Generic expressions for total and effective normal stresses
- $\sigma_1$ ,  $\sigma_2$ ,  $\sigma_3$ : Major, intermediate, and minor principal stresses
- $\sigma_{1f}$ ,  $\sigma_{3f}$ ,  $\sigma_{1f}$ ',  $\sigma_{3f}$ ': Major and minor principal stresses at failure, and their effective values
- $\sigma_a, \sigma_r$ : Axial and radial stresses (= $\sigma_3$  and  $\sigma_1$ , respectively)
- $\sigma_t$ ,  $\sigma_t$ ': Total and effective tensile strength
- $\sigma_{t0}$ : (Total) tensile strength at zero shear strain
- $\sigma_{vp}$ ': One-dimensional pre-consolidation effective stress
- $\tau$ : Shear stress