



Title	Infrared spectra of isoquinolinium (iso-C <sub>9</sub> H <sub>7</sub> NH <sup>+</sup> ) and isoquinolinyl radicals (iso-C <sub>9</sub> H <sub>7</sub> NH and 1-, 3-, 4-, 5-, 6-, 7- and 8-iso-HC <sub>9</sub> H <sub>7</sub> N) isolated in solid para-hydrogen
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**Infrared spectra of isoquinolinium (*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>) and isoquinolinyl radicals (*iso*-C<sub>9</sub>H<sub>7</sub>NH and 1-, 3-, 4-, 5-, 6-, 7-, and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N) isolated in solid *para*-hydrogen**

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## ABSTRACT

Protonated polycyclic aromatic nitrogen heterocycles ( $H^+PANH$ ) are prospective candidates that might contribute to the interstellar unidentified infrared (UIR) emission bands because protonation enhances the relative intensities of bands near 6.2, 7.7, and 8.6  $\mu m$  and the presence of N atom induces a blue shift of the ring-stretching modes so that the spectra of  $H^+PANH$  match better with the 6.2- $\mu m$  feature in the class-A UIR spectra. We report the infrared (IR) spectra of protonated isoquinoline (2-isoquinolinium cation, *iso*- $C_9H_7NH^+$ ), its neutral counterpart (2-isoquinolinyl radical, *iso*- $C_9H_7NH$ ), and another mono-hydrogenated product (6-isoquinolinyl radical, 6-*iso*- $HC_9H_7N$ ), produced on electron-bombardment of a mixture of isoquinoline (*iso*- $C_9H_7N$ ) with excess *para*-hydrogen (*p*- $H_2$ ) during matrix deposition at 3.2 K. To generate additional isomers of hydrogenated isoquinoline, we irradiated *iso*- $C_9H_7N/Cl_2/p$ - $H_2$  matrices at 365 nm to generate Cl atoms, followed by IR irradiation to generate H atoms via  $Cl + H_2 (\nu = 1) \rightarrow HCl + H$ ; the H atoms thus generated reacted with *iso*- $C_9H_7N$ . In addition to *iso*- $C_9H_7NH$  and 6-*iso*- $HC_9H_7N$  observed in electron-bombardment experiments, we identified six additional hydrogenated isoquinoline, 1-, 3-, 4-, 5-, 7-, and 8-*iso*- $HC_9H_7N$  via their IR spectra; hydrogenation on the N atom and all available carbon atoms except two sharing carbon atoms on the fused ring were observed. Spectral groupings were achieved according to their behaviors after maintenance of the matrix in darkness and on secondary photolysis at varied wavelengths. The assignments were supported by comparison of experimental results with vibrational wavenumbers and IR intensities of possible isomers predicted with the B3LYP/6-311++G(d,p) method. The implication in the identification of the UIR band is discussed.

## 1. Introduction

The infrared (IR) emission from many astrophysical objects, showing common major features near 3.3, 6.2, 7.7, 8.6, and 11.2  $\mu\text{m}$ , is referred to as the unidentified infrared (UIR) emission bands.<sup>1,2</sup> Because these emitting wavelengths are characteristic of the CH-stretching, CC-stretching, and CH-bending modes of aromatic compounds, IR fluorescence of polycyclic aromatic hydrocarbons (PAH) upon ultraviolet (UV) excitation has been proposed to be responsible for these UIR bands.<sup>3–6</sup> The slight variations in spectral patterns of UIR features among various interstellar objects are presumably due to distinct distributions of a group of emitters. Although laboratory IR spectra have been acquired for a number of PAH and their derivatives, no PAH has been positively identified as a carrier of the UIR bands. One of the discrepancies between IR spectra of neutral PAH and UIR bands lies in the relative intensities of bands; the intensities for features of neutral PAH near 6.2, 7.7 and 8.6  $\mu\text{m}$  relative to other bands are smaller than those in UIR emission. Theoretical studies suggested that the intensity of these features is enhanced in ions of PAH, such as PAH cations and protonated PAH (designated  $\text{H}^+\text{PAH}$ ).<sup>7,8</sup> We have experimentally demonstrated that IR spectra of relatively large  $\text{H}^+\text{PAH}$ , such as protonated coronene ( $\text{C}_{24}\text{H}_{13}^+$ ),<sup>9</sup> protonated ovalene ( $\text{C}_{32}\text{H}_{15}^+$ ),<sup>10</sup> and non-planar protonated corannulene ( $\text{C}_{20}\text{H}_{11}^+$ ),<sup>11</sup> show such an enhancement and an improved agreement with the UIR emission spectra.<sup>12</sup>

Peeters et al.<sup>13</sup> reported that the UIR emission spectra can be classified as classes A, B, and C according to the band positions and the width of the individual bands in the wavelength region 6–9  $\mu\text{m}$ . The feature near 6.2  $\mu\text{m}$  in class-A spectra is asymmetric, with a maximum near 6.22  $\mu\text{m}$ . This band position does not match the laboratory spectra of neutral PAH with a band near 6.3  $\mu\text{m}$ , especially after accounting for the expected red shift  $\sim 15\text{ cm}^{-1}$  in converting the absorption spectra to UV-induced emission spectra.<sup>7,14,15</sup> The bands of  $\text{H}^+\text{PAH}$  near 6.3  $\mu\text{m}$  associated with the CC-

stretching modes also mismatch with the feature near 6.2  $\mu\text{m}$  in class-A UIR spectra. The substitution of a CH moiety with a nitrogen atom in PAH, known as polycyclic aromatic nitrogen heterocycles (PANH), was suggested to induce a blue shift of the CC-stretching bands of PAH near 6.3  $\mu\text{m}$  so that it matches better with the UIR band near 6.2  $\mu\text{m}$ .<sup>13,15,16</sup> The existence of PANH in interstellar media is plausible because nitrogen is the most abundant elements after hydrogen, helium, and carbon. Nitrogen-containing heterocycles are considered as key chemical species associated with the origin of life.<sup>17</sup> Small PANH, such as quinoline ( $\text{C}_9\text{H}_7\text{N}$ ), isoquinoline (*iso*- $\text{C}_9\text{H}_7\text{N}$ ), and their derivatives, have been detected in meteorites;<sup>18,19</sup> possible mechanisms for the formation of small PANH have been discussed in the literature.<sup>20</sup>

Because of a large proton affinity at the nitrogen site, PANH can be protonated readily by  $\text{H}_3^+$ , which is abundant in interstellar space. For example, the proton affinity of quinoline is 953  $\text{kJ mol}^{-1}$ , which is much greater than a value 803  $\text{kJ mol}^{-1}$  of naphthalene.<sup>21</sup> The protonated PANH (designated  $\text{H}^+\text{PANH}$ ) are considered to be a potentially important class of molecules responsible for UIR bands because protonation enhances the intensity of 6.2, 7.7, and 8.6  $\mu\text{m}$  features and because the spectra of  $\text{H}^+\text{PANH}$  match better with the 6.2- $\mu\text{m}$  feature in class-A UIR spectra.<sup>13,15,16</sup> Infrared multiphoton dissociation (IRMPD) spectra of some small  $\text{PANH}^+$  and  $\text{H}^+\text{PANH}$  have been reported,<sup>22</sup> but the spectra were broad and were typically recorded only in spectral region 600–1600  $\text{cm}^{-1}$ .

We have developed a technique to produce protonated and hydrogenated PAH, of which *para*-hydrogen (*p*- $\text{H}_2$ ) containing a small amount of PAH were bombarded with electrons during deposition at 3.2 K.<sup>12,23,24,25</sup> The ionization of  $\text{H}_2$  molecule led to the formation of  $\text{H}_3^+$  and H atom, which subsequently produced protonated species via proton transfer and their neutral counterparts, i.e., hydrogenated species, via neutralization and hydrogen reactions. Alternatively, for the

production of mono-hydrogenated PAH, we developed an efficient method to study hydrogenation reactions in solid  $p\text{-H}_2$ .<sup>12,23,24,25,26</sup> In this method, a  $\text{Cl}_2$ -doped PAH/ $p\text{-H}_2$  matrix was initially photolyzed at 365 nm to generate Cl atom, followed by IR irradiation to generate H atoms via  $\text{Cl} + \text{H}_2 (\nu = 1) \rightarrow \text{HCl} + \text{H}$ ; H atoms are extremely mobile in quantum-solid  $p\text{-H}_2$  as they can diffuse through the matrix via quantum tunneling. This method produced hydrogenated species with significantly improved efficiency over those from electron-bombardment experiments. Our methods have demonstrated several advantages over other techniques such as the IRMPD and infrared photodissociation (of tagged species).<sup>12,23,24</sup> Our techniques allow the production of protonated and hydrogenated species with little fragmentation. Because we measured IR absorption spectra of species isolated in  $p\text{-H}_2$ , the spectra are characterized by narrow lines with reliable relative IR intensities. A wider spectral coverage, 500–4000  $\text{cm}^{-1}$ , allows us to identify lines associated with the CH- and NH-stretching modes as well. In the investigation of the IR spectra of  $\text{H}^+\text{PANH}$ , information on the NH-stretching and CH-stretching modes provides a definitive identification of isomers. Recently, our laboratory reported the IR spectrum of protonated quinoline (1-quinolinium cation,  $\text{C}_9\text{H}_7\text{NH}^+$ ) in solid  $p\text{-H}_2$ .<sup>25</sup> We observed a line of the NH-stretching mode at 3385  $\text{cm}^{-1}$ , four lines of ring-stretching modes in region 1400–1650  $\text{cm}^{-1}$ , three lines of in-plane CH bending modes in region 1150–1300  $\text{cm}^{-1}$ , and one line of out-of-plane CH bending mode at 790  $\text{cm}^{-1}$ , which is much informative as compared with three broad bands identified with the IRMPD method near 1548, 1354, and 772  $\text{cm}^{-1}$ .<sup>22</sup> The observed lines at 1641.4, 1598.4, and 1562.0  $\text{cm}^{-1}$  were found to be blue-shifted from lines at 1618.7, 1580.8, 1556.7, and 1510.0  $\text{cm}^{-1}$  of the corresponding protonated naphthalene ( $\text{C}_{10}\text{H}_9^+$ ),<sup>27</sup> matching better with the 6.2- $\mu\text{m}$  band of UIR emission. Furthermore, five isomers of quinolinyl radicals ( $\text{C}_9\text{H}_7\text{NH}$  and 3-, 4-, 7-, and 8- $\text{HC}_9\text{H}_7\text{N}$ ) were characterized by their IR spectra.

In this work, we extended our investigation to isoquinoline, which is a structural isomer of quinoline. In isoquinoline, the CH group at site 2 of naphthalene is substituted with a nitrogen atom (Fig. 1); the CH at site 1 is substituted with a nitrogen atom in quinoline. Isoquinoline has 10 possible protonation sites and the numbering of the sites are indicated in Fig. 1. In this paper, the isoquinoline protonated or hydrogenated at the N-atom site is designated as *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> or *iso*-C<sub>9</sub>H<sub>7</sub>NH, respectively, and those protonated or hydrogenated at the C-atom sites of isoquinoline are designated as *n-iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N or *n-iso*-HC<sub>9</sub>H<sub>7</sub>N, in which *n* refers to the carbon number. When no specific isomeric form is referred, we designated the protonated and hydrogenated species as *iso*-C<sub>9</sub>H<sub>8</sub>N<sup>+</sup> and *iso*-C<sub>9</sub>H<sub>8</sub>N, respectively. The IRMPD spectrum of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> has been reported, but, as aforementioned, the spectral region is limited to 600–1600 cm<sup>-1</sup>.<sup>22</sup> The spectral identification of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> was mostly based on an intense band at 774 cm<sup>-1</sup> associated with the out-of-plane CH-bending mode and two other bands at 1360 and 1582 cm<sup>-1</sup>; the characteristic NH-stretching band expected near 3400 cm<sup>-1</sup> was unreported. Electronic excitation spectra of the S<sub>1</sub> ← S<sub>0</sub> and S<sub>3</sub> ← S<sub>0</sub> transitions of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> were reported; the spectra were characterized by band origins near 27500 and 42500 cm<sup>-1</sup>, respectively, and a progression with spacing ~500 cm<sup>-1</sup> associated with the Franck-Condon-active ring-deformation mode.<sup>28,29</sup>

Here, we report the IR identification of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> (2-isoquinolinium cation) and two isomers of isoquinolinyl radicals (*iso*-C<sub>9</sub>H<sub>7</sub>NH and 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N) in solid *p*-H<sub>2</sub> produced in the electron-bombardment experiments. With the hydrogenation technique using Cl<sub>2</sub>, we observed an unprecedented number of isomers of hydrogenated isoquinoline in one experiment. Infrared spectra of eight isomers of isoquinolinyl radicals, including *iso*-C<sub>9</sub>H<sub>7</sub>NH and 1-, 3-, 4-, 5-, 6-, 7-, and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N, were identified; hydrogenation on the N atom and all available C atoms except the two sharing carbon atoms in the fused ring was observed. The implication in the identification

of the UIR band is discussed.

## 2. Experiments

The IR absorption technique applied to species isolated in *p*-H<sub>2</sub> matrices has been described elsewhere.<sup>12,23,24,25</sup> A gold-plated copper substrate for matrix samples was cooled to 3.2 K by a closed-cycle helium refrigerator; it also served as a reflector for IR probing light to have reflective absorption. Absorption spectra were recorded with a Fourier-transform infrared (FTIR) spectrometer (Bruker, VERTEX 80V, KBr beam splitter, Hg-Cd-Te detector cooled to 77 K). Typically, 200 scans at resolution 0.25 cm<sup>-1</sup> were recorded after each stage of the experiment.

The method for converting *normal*-H<sub>2</sub> (*ortho*-H<sub>2</sub> : *p*-H<sub>2</sub> = 3 : 1) to *p*-H<sub>2</sub> has been described previously.<sup>12,23,24</sup> Briefly, *normal*-H<sub>2</sub> was introduced into a copper cell filled with a catalyst (hydrated iron(III) oxide, 30–50 mesh) cooled to 12–13 K with another closed-cycle helium refrigerator. The proportion of *ortho*-H<sub>2</sub> in the resultant *p*-H<sub>2</sub> gas was typically less than 100 ppm. The isoquinoline sample was obtained commercially (97%, Acros organics) and degassed with several freeze–pump–thaw cycles.

Gaseous mixtures of *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> was prepared by flowing *p*-H<sub>2</sub> (flow rate ~10 mmol h<sup>-1</sup>) over a flask containing liquid isoquinoline. The supply of *iso*-C<sub>9</sub>H<sub>7</sub>N vapor was controlled with a needle valve. The mixture was deposited onto the Cu substrate over a period of 7–8 h. During the deposition, the matrix was bombarded with an electron beam (20 μA current and 300 eV kinetic energy) to produce protonated and hydrogenated products. To differentiate various products produced by electron bombardment, the matrix was typically maintained in darkness for ~10 h, followed by secondary photolysis with light at 365 ± 10 nm from a light-emitting diode (375 mW) and at 405 nm from a diode laser (120 mW).



Alternatively, the isoquinolinyl radicals were produced in *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrices on UV photolysis at 365 nm for 1 h, followed by IR irradiation from an external global source for 2 h. Secondary photolysis at 553, 540, 420, 405, and 365 nm (45–60 mW) using an OPO laser (EKSPLA, NT340, 10 Hz) were performed to distinguish various groups of lines according to their photolytic behavior. To avoid the reaction between Cl and *p*-H<sub>2</sub> during the IR spectral acquisition, an IR filter with cutoff wavelength at 2.4 μm was employed.

Because of the low vapor pressure of isoquinoline and the mixing method employed, we were unable to determine accurately the mixing ratio of *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> based on partial pressures. Instead, we estimated the mixing ratios according to the method developed by Tam and Fajardo.<sup>30</sup> The mixing ratio  $x$  (in ppm) of a species in solid *p*-H<sub>2</sub> is represented as

$$x = \frac{2.303 \int \log_{10}(I/I_0) d\nu}{\epsilon l} \times V_m \times 10^6, \quad (1)$$

in which  $\epsilon$  is the absorption coefficient in cm mol<sup>-1</sup>,  $\int \log_{10}(I/I_0) d\nu$  is the observed integrated absorbance of a specific line,  $l$  (in cm) is the IR absorption path length through solid *p*-H<sub>2</sub>, and  $V_m$  (= 23.16 cm<sup>3</sup> mol<sup>-1</sup>)<sup>31</sup> is the molar volume of solid *p*-H<sub>2</sub>. Infrared absorption path length  $l$  was estimated from the integrated band area of transition S<sub>1</sub>(0) + S<sub>0</sub>(0) in 4828–4853 cm<sup>-1</sup> and pure rotational transition at 1167 cm<sup>-1</sup> of the *p*-H<sub>2</sub> matrix.<sup>32</sup> Infrared intensities of a specific vibrational transition calculated with the B3LYP/6-311++G(d,p) method were used for  $\epsilon$ ; the error in calculated  $\epsilon$ , hence  $x$ , might be as large as a factor of 2. By using several lines of the same species reduced the error in the estimated mixing ratio.

### 3. Quantum-chemical calculations

Quantum-chemical calculations were performed using the Gaussian 16 (Rev. B.01) program package.<sup>33</sup> Geometrical optimizations and vibrational analyses were performed with the B3LYP hybrid functionals<sup>34, 35</sup> and standard basis set 6-311++G(d,p).<sup>36</sup> Anharmonic vibrational

wavenumbers and IR intensities were calculated by employing a second-order vibrational perturbation theory (VPT2) implemented in Gaussian 16.<sup>37</sup> Furthermore, calculations on single-point electronic energies with the method coupled cluster with single and double and perturbative triple excitations, CCSD(T),<sup>38</sup> were performed on geometries obtained with the B3LYP/6-311++G(d,p) method; zero-point vibrational energies (ZPVE) were corrected according to the harmonic vibrational wavenumbers calculated with the B3LYP method. Calculations on the electronic excitation and oscillator strength were performed with time-dependent density-functional theory (TD-DFT) using the B3LYP/6-311++G(d,p) method.

### 3.1. Protonated isoquinoline (isoquinolinium cation)

Isoquinoline has 10 possible sites for protonation and the numbering of carbon and nitrogen atoms of isoquinoline is indicated in Fig. 1; N atom is numbered as 2 and the two sharing carbon atoms in the fused ring as 4a and 8a. The geometries of all isomers of protonated isoquinoline, with relative energies and representative geometrical parameters predicted with the CCSD(T)/6-311++G(d,p)//B3LYP/6-311++G(d,p) method, are presented in Fig. S1, ESI†. The isomer protonated at the N-atom site of isoquinoline, 2-isoquinolinium (*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>), was predicted to have the least energy; the energy of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> + H<sub>2</sub> is smaller than *iso*-C<sub>9</sub>H<sub>7</sub>NH + H<sub>3</sub><sup>+</sup> by 538 and 534 kJ mol<sup>-1</sup>, implying a calculated proton affinity of 957 and 989 kJ mol<sup>-1</sup> for *iso*-C<sub>9</sub>H<sub>7</sub>N on the N-atom site at the B3LYP and CCSD(T) levels, respectively. The energies of isomers protonated at the C-atom site of isoquinoline, 1-, 3-, 4-, 5-, 6-, 7-, and 8-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N, are predicted to be 200, 190, 169, 171, 209, 183, and 190 kJ mol<sup>-1</sup> greater than *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>, respectively, much smaller than values 267 and 240 kJ mol<sup>-1</sup> predicted for 4a- and 8a-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N, respectively (Table 1). In all isomers except 4a- and 8a-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N, all carbon and nitrogen atoms remain in a plane and have symmetry of point-group C<sub>s</sub>; the ring skeleton of 4a- and 8a-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N are

non-planar so that these two species belong to point group  $C_1$ . The proton-transfer reactions from  $H_3^+$  to isoquinoline,  $H_3^+ + iso-C_9H_7N \rightarrow H_2 + iso-C_9H_8N^+$ , are barrierless. The potential energy scheme for the proton transfer among these isomers via moving toward the neighboring atom is presented in Fig. S2, ESI†. Although the energies of transition states for isomerization, 231–297  $\text{kJ mol}^{-1}$ , are smaller than  $H_3^+ + iso-C_9H_7N$  (534  $\text{kJ mol}^{-1}$ ), isomerization to other isomers is unlikely to occur once a product is stabilized in solid  $p\text{-H}_2$ .

The scaled harmonic and anharmonic vibrational wavenumbers and IR intensities for  $iso-C_9H_7NH^+$  (Table S1, ESI†), 1-, 3-, and 4- $iso-H^+C_9H_7N$  (Table S2, ESI†), 5-, 6-, and 7- $iso-H^+C_9H_7N$  (Table S3, ESI†), and 8-, 4a-, and 8a- $iso-H^+C_9H_7N$  (Table S4, ESI†) are listed. The scaling of harmonic vibrational wavenumbers is discussed in Section 4.1. The stick spectra in region 1800–600  $\text{cm}^{-1}$  for all isomers of protonated isoquinoline, simulated according to scaled harmonic wavenumbers and IR intensities, are presented in Fig. S3, ESI†.

### 3.2. Hydrogenated isoquinoline (isoquinolinyl radical)

The geometries of all isomers of mono-hydrogenated isoquinoline ( $iso-C_9H_8N$ , isoquinolinyl radical), with key geometrical parameters and relative energies predicted with the CCSD(T)/6-311++G(d,p)//B3LYP/6-311++G(d,p) method, are presented in Fig. S4, ESI†. The isomer hydrogenated at the N-atom site, 2-isoquinolinyl radical ( $iso-C_9H_7NH$ ), was predicted to have the least energy,  $-111 \text{ kJ mol}^{-1}$  relative to  $H + iso-C_9H_7N$ . The energies of other isomers, 1-, 3-, 4-, 5-, 6-, 7-, and 8- $iso-HC_9H_7N$ , are predicted to be 6, 22, 3, 2, 24, 24, and 3  $\text{kJ mol}^{-1}$ , respectively, relative to  $iso-C_9H_7NH$ . The corresponding values for 4a- and 8a- $iso-HC_9H_7N$  are much greater, 100 and 91  $\text{kJ mol}^{-1}$ , respectively. In all isomers except 4a- and 8a- $iso-HC_9H_7N$ , carbon and nitrogen atoms remain planar and have symmetry of point-group  $C_s$ ; 4a- and 8a- $iso-HC_9H_7N$  have non-planar ring skeleton and belong to point group  $C_1$ . The calculated barriers for H-atom addition,

$\text{H} + \text{iso-C}_9\text{H}_7\text{N} \rightarrow \text{iso-C}_9\text{H}_8\text{N}$ , and those for hydrogen transfer to the neighboring site are presented in Fig. S5, ESI†. All barrier heights for reactions  $\text{H} + \text{iso-C}_9\text{H}_7\text{N} \rightarrow \text{iso-C}_9\text{H}_8\text{N}$  are smaller than  $27 \text{ kJ mol}^{-1}$ , except those for the formation of 4a- and 8a-*iso*- $\text{HC}_9\text{H}_7\text{N}$ , which have barriers  $> 46 \text{ kJ mol}^{-1}$  (Table 1). Because the energies of TS for isomerization are greater than that of  $\text{H} + \text{iso-C}_9\text{H}_7\text{N}$ , isomerization reactions will not occur during hydrogenation of isoquinoline in the low-temperature solid *p*- $\text{H}_2$  environment.

The harmonic and anharmonic vibrational wavenumbers and IR intensities for *iso*- $\text{C}_9\text{H}_7\text{NH}$ , 1-, and 3-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (Table S5, ESI†), 4-, 5-, and 6-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (Table S6, ESI†), 7- and 8-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (Table S7, ESI†), and 4a- and 8a-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (Table S8, ESI†) are listed. The IR stick spectra in region  $1800\text{--}600 \text{ cm}^{-1}$  for all isomers of hydrogenated isoquinoline, simulated according to the scaled harmonic vibrational wavenumbers and IR intensities, are presented in Fig. S6, ESI†.

## 4. Experimental results

### 4.1 IR spectrum of *iso*- $\text{C}_9\text{H}_7\text{N}/p\text{-H}_2$ matrices

An IR spectrum recorded after deposition of a *iso*- $\text{C}_9\text{H}_7\text{N}/p\text{-H}_2$  mixture for 5 h is presented in Fig. S7a, ESI†. Intense lines were observed at  $3069.6$ ,  $1632.8$ ,  $1595.0$ ,  $1273.0$ ,  $826.3$ ,  $740.1$ , and  $481.0 \text{ cm}^{-1}$ , with the most intense line at  $826.3 \text{ cm}^{-1}$ . On comparison of observed vibrational wavenumbers and relative intensities with calculations, we assigned 24 observed lines to the fundamental vibrational transitions of *iso*- $\text{C}_9\text{H}_7\text{N}$ . The observed vibrational wavenumbers are plotted against the harmonic vibrational wavenumbers predicted with the B3LYP/6-311++G(d,p) method in Fig. S8, ESI†. A linear regression yielded satisfactory linear relation with a scaling equation  $y = (0.9823 \pm 0.0025)x + (0.6 \pm 2.8)$  for wavenumbers  $600\text{--}2000 \text{ cm}^{-1}$ , in which  $y$  is the scaled vibrational wavenumber and  $x$  is the calculated harmonic vibrational wavenumber; the errors present one standard deviation in fitting. For lines with wavenumbers above  $2000 \text{ cm}^{-1}$ ,

assignments for *iso*-C<sub>9</sub>H<sub>7</sub>N are difficult because of severe interactions of the CH-stretching modes with some overtone or combination modes. We hence adopted the equation for quinoline in *p*-H<sub>2</sub>,  $y = 0.9122 x + 174.3$ ;<sup>25</sup> the error in this region is hence expected to be greater. We used these equations to scale the harmonic vibrational wavenumbers predicted for *iso*-C<sub>9</sub>H<sub>7</sub>N, *iso*-C<sub>9</sub>H<sub>8</sub>N<sup>+</sup>, and *iso*-C<sub>9</sub>H<sub>8</sub>N.

Simulated stick spectra according to scaled harmonic and anharmonic vibrational wavenumbers are presented in Figs. S7b and S7c, ESI†, respectively. The observed and predicted wavenumbers and relative intensities are summarized in Table S9, ESI†. The mean absolute deviation between observed and scaled harmonic vibrational wavenumbers is  $6.5 \pm 9.4 \text{ cm}^{-1}$ ; that between observed and anharmonic vibrational wavenumbers is  $8.8 \pm 10.9 \text{ cm}^{-1}$ , indicating a satisfactory agreement between experimental and predicted vibrational wavenumbers.

The absorption path length of the matrix ( $l$  in Eq. 1) was estimated to be  $\sim 1.0 \text{ mm}$ . Using the integrated absorbance of lines at 1632.8, 1595.0, 826.3, and 740.1  $\text{cm}^{-1}$ , we derived a mixing ratio  $54 \pm 2 \text{ ppm}$  for *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> in Fig. S7a, ESI†; listed error represents one standard deviation among mixing ratios estimated from each absorption line.

## 4.2 IR spectra of electron-bombarded *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> matrices

Figure 2 shows IR spectra of an electron-bombarded *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> matrix in spectral regions 3550–3370 and 1670–600  $\text{cm}^{-1}$ . Figure 2a presents the spectrum recorded after deposition of electron-bombarded matrix for 7 h, with lines of *iso*-C<sub>9</sub>H<sub>7</sub>N stripped using Fig. S7a, ESI†. Therefore, the positive lines in this spectrum are induced by electron bombardment. Figure 2b shows the difference spectrum after maintaining the matrix in darkness for 10 h; positive lines indicate generation and negative lines indicate destruction. After this experimental step, the matrix was irradiated at 365 nm for 25 min followed by 405 nm for 3 h. The difference spectra after each

secondary irradiation step are shown in Figs. 2c and 2d.

It has been demonstrated that electron bombardment during deposition of a  $p\text{-H}_2$  matrix containing the guest molecules produces protonated and hydrogenated products.<sup>12,23,24</sup> During maintenance of the deposited matrix in darkness for a prolonged period, electrons trapped in solid  $p\text{-H}_2$  slowly diffused to neutralize protonated products and produce their neutral counterparts. The trapped hydrogen atoms (i.e., those unreacted during deposition) might further react with parent molecules to produce hydrogenated products. Consequently, during maintenance in darkness, the intensities of lines of protonated species and the parent molecules are expected to decrease, while those of hydrogenated species increase. In the spectra shown in Fig. 2, we identified several groups of lines (groups  $A^+$ , A, and B) according to their characteristic behavior after each experimental step. The intensities of lines in group  $A^+$  decreased after maintaining the matrix sample in darkness for 10 h (Fig. 2b). Those increased at this step was further classified into two groups according to their behavior upon secondary irradiation at 365 and 405 nm. The intensities of lines in groups A and B decreased by similar proportions, 40–50 %, upon irradiation at 365 nm (Fig. 2c), whereas those of lines in group B further decreased by 35 % upon irradiation at 405 nm (Fig. 2d); the percentage is based on the spectrum measured after maintaining the matrix sample in darkness for 10 h.

Lines in group  $A^+$  include intense features observed at 3397.2, 1643.3, 1400.1, 1378.6, and 791.4  $\text{cm}^{-1}$  and weaker ones at 1554.7 and 1497.3  $\text{cm}^{-1}$ ; a line at 720.0  $\text{cm}^{-1}$  might also belong to this group, but it was severely interfered with by the absorption of  $p\text{-H}_2$ . The vibrational wavenumbers and integrated relative intensities of these lines are listed in Table 2. To be discussed in Section 5.1, these lines in group  $A^+$  are assigned to  $iso\text{-C}_9\text{H}_7\text{NH}^+$  (2-isoquinolinium cation). Lines in group A include intense lines at 3495.2, 3491.0, 1613.9, 1476.6, 1440.5, 1180.9, 988.3,

763.1, and 737.3  $\text{cm}^{-1}$  and seven weaker lines, with the most intense line at 988.3  $\text{cm}^{-1}$ . These lines in group A are assigned to *iso*-C<sub>9</sub>H<sub>7</sub>NH (2-isoquinoliny radical), to be discussed in Section 5.2. Lines in group B include intense lines at 1557.5, 1446.8, 1284.5, 890.9, 810.2, and 697.3  $\text{cm}^{-1}$ , with the most intense line at 810.2  $\text{cm}^{-1}$ , and several weaker ones. To be discussed in Section 5.2, these lines in group B are assigned to 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N (6-isoquinoliny radical). Lines in groups A and B were also observed in hydrogenation experiments, to be described in Section 4.3.

### 4.3 IR spectra of UV/IR irradiated *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrices

Figure 3a depicts the IR spectrum in region 1200–600  $\text{cm}^{-1}$ , recorded after photolysis at 365 nm for 1 h of an *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix deposited for 8 h. The corresponding IR spectra covering 3600–3470, 3100–2700, and 1700–1200  $\text{cm}^{-1}$  are presented in Fig. S9, ESI†. Upon photolysis at 365 nm, no new features were observed, except weak lines at 1441.1  $\text{cm}^{-1}$  (assigned to ClOO due to the reaction of Cl with trace O<sub>2</sub> impurity)<sup>39</sup> and at 943.8  $\text{cm}^{-1}$  (the spin-orbit transition of Cl),<sup>40</sup> indicating the production of Cl atoms. This matrix was further exposed to IR irradiation for 2 h; the difference spectrum is shown in Fig. 3b. In addition to those described in the previous section for groups A and B, many new lines were observed. To differentiate these newly observed features, secondary irradiation at 553, 540, 420, 405, and 360 nm was performed sequentially for 30 min at each wavelength; the difference spectra after each step are presented in Figs. 3c–3g, respectively.

According to the distinct behavior on secondary photolysis at varied wavelengths, we categorized newly observed features after IR irradiation into eight groups, designated as groups A to H; their photolytic behaviors are summarized in Table 3; the listed percentage variations are based on the spectrum measured after IR irradiation. Results from secondary irradiation at 420 and 405 nm are the most informative because intensity variations of lines were observed in almost all

groups except groups G and H at 420 nm and group E at 405 nm, but information from irradiation at other wavelengths was also required to differentiate each group. At 420 nm, lines in groups A, B, C, and E decreased, whereas those in groups D and F increased. At 405 nm, those in groups A, B, C, and H decreased, whereas those in groups D, F, and G increased. The distinct features of eight groups are indicated with color-coded arrows and group labels in Fig. 3 and Fig. S9, ESI†. These features in groups A–H are assigned to *iso*-C<sub>9</sub>H<sub>7</sub>NH and 6-, 7-, 8-, 4-, 5-, 1-, and 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N, respectively, to be discussed in Sections 5.2–5.5.

We integrated several intense lines in each group, as listed in Table S10, ESI†, to estimate the mixing ratios after UV/IR irradiation of the *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrices. The estimated mixing ratios are [*iso*-C<sub>9</sub>H<sub>7</sub>N] = (90 ± 12) ppm, [*iso*-C<sub>9</sub>H<sub>7</sub>NH] = (2.8 ± 0.5) ppm, [6-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (2.1 ± 1.2) ppm, [7-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (1.1 ± 0.4) ppm, [8-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (1.6 ± 0.1) ppm, [4-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (2.5 ± 0.3) ppm, [5-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (1.5 ± 0.3) ppm, [1-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (1.5 ± 0.3) ppm, and [3-*iso*-HC<sub>9</sub>H<sub>7</sub>N] = (0.42 ± 0.01) ppm; listed errors represent one standard deviation among various lines.

## 5. Discussion

### 5.1 Assignment of lines in group A<sup>+</sup> to 2-isoquinolinium cation (*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>)

As aforementioned, during maintenance of electron-bombarded matrix in darkness, electrons trapped in solid *p*-H<sub>2</sub> will slowly diffuse and neutralize cationic species. The intensities of the lines in group A<sup>+</sup> decayed during this process, indicating that the most likely carrier of lines in this group is isoquinolinium cation, i.e., protonated isoquinoline. In Fig. 4a, we reproduced Fig. 2a, the IR spectrum of the electron-bombarded matrix with absorption lines of *iso*-C<sub>9</sub>H<sub>7</sub>N stripped, and indicated lines in group A<sup>+</sup> with arrows and labels; the regions subjected to interference by intense absorption of the precursor (*iso*-C<sub>9</sub>H<sub>7</sub>N) and other products are shaded gray. The stick spectrum of



*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> simulated according to scaled harmonic vibrational wavenumbers and IR intensities is presented in Fig. 4b.

The observed wavenumbers and relative integrated intensities of lines in group A<sup>+</sup> agree the best with those predicted for *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>, as compared in Table 2. The most intense line observed at 3397.2 cm<sup>-1</sup> is characteristic of the N–H stretching mode; the observed wavenumber agrees satisfactorily with the scaled harmonic vibrational wavenumber 3430 cm<sup>-1</sup> (with intensity 208 km mol<sup>-1</sup>) and anharmonic vibrational wavenumbers at 3407 cm<sup>-1</sup> for the N–H stretching mode of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>. The C–H stretching modes of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> have IR intensities less than 8 km mol<sup>-1</sup> and were severely interfered by precursor absorption, so they could not be definitively identified. Similarly, for all C-protonated isomers except 4a- and 8a-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N, the C–H stretching modes of *iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N are weak and only the CH<sub>2</sub> symmetric stretching mode, predicted in the region 2800–2900 cm<sup>-1</sup> has intensity 32–91 km mol<sup>-1</sup>. For 4a- and 8a-*iso*-H<sup>+</sup>C<sub>9</sub>H<sub>7</sub>N, the new C–H stretching mode on the fused ring is predicted near 2650 cm<sup>-1</sup> with intensity ~50 km mol<sup>-1</sup>. However, no new lines that were induced by electron bombardment and decayed after maintenance in darkness were observed in the 2650–2900 cm<sup>-1</sup> region.

In the region below 1700 cm<sup>-1</sup>, four intense lines were observed at 1643.3, 1400.1, 1378.6, and 791.4 cm<sup>-1</sup>, near the scaled harmonic wavenumbers predicted for *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> at 1639 (ν<sub>9</sub>, N-containing ring stretch), 1394 (ν<sub>16</sub>, ring stretch), 1382 (ν<sub>17</sub>, ring stretch), and 788 (ν<sub>40</sub>, out-of-plane C–H and N–H bend) cm<sup>-1</sup>, respectively. Weaker lines observed at 1554.7 and 1497.3 cm<sup>-1</sup> correspond to 1559 (ν<sub>12</sub>, ring stretch) and 1500 (ν<sub>13</sub>, ring stretch) cm<sup>-1</sup>, respectively. Other relatively intense lines of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> at 1620 (ν<sub>10</sub>), 1587 (ν<sub>11</sub>), 1285 (ν<sub>18</sub>), and 471 (ν<sub>44</sub>) cm<sup>-1</sup> were predicted to have IR intensities 27, 15, 13, and 29 km mol<sup>-1</sup>, respectively. These lines could not be identified definitively because of the spectral interference from the intense absorption of

the parent. The mean absolute deviation between observed and scaled harmonic vibrational wavenumbers is  $8.3 \pm 10.0 \text{ cm}^{-1}$ , which improves to  $4.2 \pm 1.0 \text{ cm}^{-1}$  on excluding the N–H stretching mode ( $\nu_1$ ). The deviation between observed and anharmonic wavenumbers is  $9.8 \pm 6.6 \text{ cm}^{-1}$ , which does not change significantly ( $9.7 \pm 7.2 \text{ cm}^{-1}$ ) after excluding  $\nu_1$ . These results indicate a satisfactory agreement between experiments and predicted vibrational wavenumbers of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>. The comparison of lines in group A<sup>+</sup> with the predicted spectra of all isomers of protonated isoquinoline (*iso*-C<sub>9</sub>H<sub>8</sub>N<sup>+</sup>) is presented in Figs. S10 and S11, ESI†; the observed spectrum of group A<sup>+</sup> agree poorly with the stick IR spectra simulated for isomers of isoquinoline protonated at a carbon site, i.e., 1-, 3-, 4-, 5-, 6-, 7-, 8-, 4a-, and 8a-*iso*-HC<sub>9</sub>H<sub>7</sub>N<sup>+</sup> (Tables S2–S4, ESI†).

Considering the observed chemical behavior, thermodynamic stability, the satisfactory agreement of vibrational wavenumber and relative IR intensity between lines in group A<sup>+</sup> and those predicted for the N-protonated isoquinoline, and the absence of some unique features characteristic for C-protonated isoquinolines, we assign lines in group A<sup>+</sup> to 2-isoquinolium (*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>).

The observed vibrational wavenumbers are also compared with those determined with the IRMPD method<sup>22</sup> in Table 2. In the IRMPD spectrum, three bands were identified at 1582, 1360, and 774 cm<sup>-1</sup>, which are red-shifted from corresponding lines at 1643.3 cm<sup>-1</sup> (and an unobserved line predicted near 1620 cm<sup>-1</sup>), 1400.1 and 1378.6 cm<sup>-1</sup>, and 791.4 cm<sup>-1</sup> observed in solid *p*-H<sub>2</sub>. The red shift is presumably due to the anharmonic excitation process characteristic in the IRMPD experiments. We identified 7 lines, including the characteristic N–H stretching mode that clearly indicates the protonation site is at the N atom.

## 5.2 Assignments of lines in groups A and B to *iso*-C<sub>9</sub>H<sub>7</sub>NH and 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N radicals

As discussed in Section 4.2, the IR lines that were induced by electron bombardment and

gained intensities during maintenance of the matrix in darkness might be due to hydrogenated isoquinoline. These lines were further classified into two groups A and B according to their behaviors on secondary photolysis at 365 nm and 405 nm. Moreover, these lines were also observed in experiments with *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrices irradiated with UV and IR light, in which H atoms were expected to be produced and react with *iso*-C<sub>9</sub>H<sub>7</sub>N to produce hydrogenated species, but not protonated species.

The experimental spectra of lines in groups A and B obtained from electron-bombarded experiment in regions 3600–2400 and 1800–600 cm<sup>-1</sup> are compared with IR stick spectra of all 10 isomers of mono-hydrogenated isoquinoline radicals, simulated according to the scaled harmonic vibrational wavenumbers and IR intensities predicted with the B3LYP/6-311++G(d,p) method in Figs. S12 and S13, ESI†. The difference spectrum after secondary photolysis at 365 nm, Fig. 2c, is inverted and presented in the bottom trace of Figs. S12a and S13a, ESI†; that after secondary photolysis at 405 nm, Fig. 2d, is presented in the top trace. Lines in groups A and B are indicated with blue and red labels and arrows, respectively. Alternatively, lines in groups A and B were also observed in the UV/IR-irradiated *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix. Because of an improved ratio of signal to noise and less interference, we discuss the assignments using the results of the UV/IR-irradiated *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix, presented as the difference spectrum after irradiation at 420 nm from the experiments on *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub>, inverted Fig. 3e, in Fig. 5a to compare with predicted IR stick spectra of possible candidates *iso*-C<sub>9</sub>H<sub>7</sub>NH, 6-, 1-, 3-, and 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N in Figs. 5b–5f, respectively; corresponding comparison with 5-, 7-, 8-, 4a-, and 8a-*iso*-HC<sub>9</sub>H<sub>7</sub>N are presented in Fig. S14, ESI†.

Assignment of lines in group A to the *iso*-C<sub>9</sub>H<sub>7</sub>NH radical is described as follows. In the spectral region presented in Fig. 5, intense lines in group A were observed at 1613.9, 1476.6,

1440.5, 1180.9, 988.3, 763.1, and 737.3  $\text{cm}^{-1}$ , in satisfactory agreement with the predicted spectrum of *iso*-C<sub>9</sub>H<sub>7</sub>NH, which has intense lines at 1614 ( $\nu_9$ , ring stretch), 1478 ( $\nu_{13}$ , C=N stretch), 1440 ( $\nu_{14}$ , ring stretch), 1181 ( $\nu_{21}$ , ring breath + in-plane CH bend), 982 ( $\nu_{26}$ , N-containing ring breath), 761, and 735  $\text{cm}^{-1}$  ( $\nu_{38}$  and  $\nu_{39}$ , both out-of-plane CH bend), shown in Fig. 5b and Fig. S14, ESI†. Weaker lines of group A, which have intensities > 6 % that of the most intense line at 988.3  $\text{cm}^{-1}$ , were observed at 1574.4, 1490.2, 1420.7, 1115.3, 1022.0, 939.2, and 577.2  $\text{cm}^{-1}$ ; these lines are near those predicted at 1577, 1489, 1411 ( $\nu_{10}$ ,  $\nu_{12}$ , and,  $\nu_{15}$ , all ring stretch), 1115 ( $\nu_{24}$ , in-plane CH bend), 1026 ( $\nu_{25}$ , ring breath), 932 ( $\nu_{27}$ , ring deform), and 571 ( $\nu_{42}$ , out-of-plane ring deform)  $\text{cm}^{-1}$ . The *iso*-C<sub>9</sub>H<sub>7</sub>NH isomer has a characteristic N–H stretching mode ( $\nu_1$ ), which is predicted to be 3518  $\text{cm}^{-1}$  (scaled harmonic) or 3476  $\text{cm}^{-1}$  (anharmonic); two intense lines at 3495.2 and 3491.0  $\text{cm}^{-1}$  were observed, similar to the predicted wavenumber; matrix-site effect due to varied interactions of *iso*-C<sub>9</sub>H<sub>7</sub>NH with H<sub>2</sub> could be the reason for this observed doublet.

The simulated spectra for other isomers of isoquinolinyl radicals hydrogenated at the carbon site cannot reproduce satisfactorily the spectrum in group A (Figs. 5 and S12–S14, ESI†). As summarized in Table 4, in which all modes with calculated IR intensity greater than 10 km mol<sup>−1</sup> are listed, we identified 15 modes of *iso*-C<sub>9</sub>H<sub>7</sub>NH. All listed modes except four CH-stretching modes have been observed. The mean absolute deviation between observed and scaled harmonic vibrational wavenumbers is  $4.2 \pm 4.8 \text{ cm}^{-1}$  and that between observed and anharmonic vibrational wavenumbers is  $7.9 \pm 6.9 \text{ cm}^{-1}$ , indicating a satisfactory agreement between experimental and predicted vibrational wavenumbers of *iso*-C<sub>9</sub>H<sub>7</sub>NH.

Assignment of lines in group B to the 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N radical is discussed as follows. In the spectral region below 1650  $\text{cm}^{-1}$ , intense lines in group B were observed at 1557.5, 1446.8, 1284.5, 921.1, 890.9, 810.2, 697.0, and 688.8  $\text{cm}^{-1}$ , in satisfactory agreement with lines of 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N

predicted at 1560, 1450, and 1282 ( $\nu_9$ ,  $\nu_{11}$ , and  $\nu_{17}$ , respectively, all ring stretch), 920 ( $\nu_{37}$ , out-of-plane CH bend), 884 ( $\nu_{26}$ , ring deform), 812, 698, and 686  $\text{cm}^{-1}$  ( $\nu_{39}$ ,  $\nu_{41}$ , and  $\nu_{42}$ , respectively, all out-of-plane CH bend). Weaker lines, which have intensities  $<25\%$  that of the most intense line at 810.2  $\text{cm}^{-1}$ , were observed at 1334.1 and 1184.4  $\text{cm}^{-1}$ ; these lines are near the lines of 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N predicted at 1337 ( $\nu_{16}$ , CH<sub>2</sub> wag) and 1185 ( $\nu_{20}$ , ring stretch)  $\text{cm}^{-1}$ , respectively. Another weaker line, observed at 1041.5  $\text{cm}^{-1}$  and predicted at 1048  $\text{cm}^{-1}$  for the ring breathing ( $\nu_{23}$ ) mode, is tentatively assigned to 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N because the line was only observed in the electron-gun experiment.

In the CH-stretching region, relatively intense lines were observed at 3059.0, 3045.1, 3025.3 in experiments with *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> (Fig. 5a), but definitive mode assignments of these lines were difficult, partly because of severe interference from intense absorptions lines of the parent and partly because of anharmonic interactions. Two lines at 2826.4, and 2802.5  $\text{cm}^{-1}$  (Figs. S12 and S13) were observed in electron bombardment experiments; they were assigned to the characteristic symmetric and antisymmetric CH<sub>2</sub>-stretching ( $\nu_7$  and  $\nu_{33}$ ) modes of 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N. These two lines could not be definitively identified in experiments with *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> (Fig. 5a) because they were interfered with by the intense absorption of HCl and its complexes.

The simulated spectra for other isomers of isoquinoliny radicals, *iso*-C<sub>9</sub>H<sub>7</sub>NH, 1-, 3-, 4-, 5-, 7-, 8-, 4a-, and 8a-*iso*-HC<sub>9</sub>H<sub>7</sub>N, cannot reproduce satisfactorily the wavenumbers and relative intensities of lines in group B (Figs. 5 and S12 to S14, ESI<sup>†</sup>). As listed in Table 4 for lines, we assigned 16 lines in group B to 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N. All listed modes with IR intensity greater than 10  $\text{km mol}^{-1}$  except two CH-stretching modes and a line predicted near 1419  $\text{cm}^{-1}$  ( $\nu_{13}$ ) have been observed; the latter line was predicted to have IR intensity of only 1  $\text{km mol}^{-1}$  with the anharmonic calculations. The mean absolute deviation between observed and scaled harmonic vibrational

wavenumbers is  $6.5 \pm 8.3 \text{ cm}^{-1}$  and that between observed and anharmonic vibrational wavenumbers is  $9.1 \pm 9.5 \text{ cm}^{-1}$ , indicating a satisfactory agreement between experimental and predicted wavenumbers.

### 5.3 Assignments of lines in groups C and D to 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N radicals

Intensities of lines in group C decreased upon irradiation at 553, 540, 420, and 405 nm and remained unchanged upon irradiation at 360 nm; in contrast, intensities of lines in group D increased upon irradiation at 553, 420, and 405 nm, decreased upon irradiation at 360 nm, and remained unaltered upon irradiation at 540 nm (Table 3). The difference spectrum in Fig. 3b, recorded after UV/IR irradiation of the *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix, was reproduced in the lower trace of Fig. 6a after truncating negative lines, and the difference spectrum upon secondary irradiation of the matrix at 553 nm, Fig. 3c, was reproduced in the upper trace of Fig. 6a; for clarity, only the region 1000–600  $\text{cm}^{-1}$  is shown. Lines in groups C and D are marked with green and brown labels and arrows, respectively. The stick spectra of isomers of 7- and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N radicals simulated according to scaled harmonic vibrational wavenumbers and IR intensities are depicted in Figs. 6b and 6c, respectively, and compared with the observed lines in groups C and D. Figure S15, ESI†, compares lines in groups C and D with all 10 isomers of isoquinoliny radical in region 1700–1000  $\text{cm}^{-1}$ ; those in regions 3600–3450 and 3100–2700  $\text{cm}^{-1}$  are shown in Fig. S16, ESI†.

As shown in these figures, the spectral pattern of observed lines in group C agrees satisfactorily with IR stick spectrum of the 7-isoquinoliny radical, 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N, but disagree with IR stick spectra of other mono-hydrogenated radicals. Table 4 compares the observed wavenumber and relative IR intensities of lines in group C with both scaled harmonic and anharmonic vibrational wavenumbers and relative IR intensities predicted for 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N.

Intense lines observed at 1551.2, 1331.5, 1284.3, 892.0, 908.4, 817.5, and 682.2  $\text{cm}^{-1}$  agree with those predicted at 1553 ( $\nu_9$ , ring stretch), 1334 ( $\nu_{16}$ ,  $\text{CH}_2$  wag coupled with ring stretch), 1286 ( $\nu_{17}$ , ring stretch), 883 ( $\nu_{26}$ , C6C7C8 symmetric stretch), 912 ( $\nu_{37}$ ,  $\text{CH}_2$  rock), 821 ( $\nu_{39}$ , out-of-plane CH bend), and 681 ( $\nu_{42}$ , out-of-plane C5H bend)  $\text{cm}^{-1}$ . Other weak lines with intensity <12 % of the most intense line at 817.5  $\text{cm}^{-1}$  were observed at 1617.2, 1205.1, 1017.9, 935.4(?), 907.6, and 720.9(?)  $\text{cm}^{-1}$ ; two lines were marked with question marks because they were not clearly visible after IR irradiation, likely due to small intensity. These lines agree with predicted vibrational wavenumbers at 1641 ( $\nu_8$ , ring stretch), 1208 ( $\nu_{20}$ , in-plane CH bend), 1016 ( $\nu_{24}$ , ring stretch), 935 ( $\nu_{36}$ , out-of-plane CH bend), 914 ( $\nu_{25}$ , ring deform), and 694 ( $\nu_{41}$ , 715 from anharmonic calculations, out-of-plane CH bend)  $\text{cm}^{-1}$ .

In the CH-stretching region, lines at 3057.8 and 3025.6  $\text{cm}^{-1}$  in group C agree with predicted vibrational wavenumbers at 3065/3059 (3046/3044,  $\nu_2/\nu_3$ ) and 3028 (3000,  $\nu_6$ )  $\text{cm}^{-1}$ ; values in parentheses are anharmonic vibrational wavenumbers. Lines in the region 2910–2790  $\text{cm}^{-1}$  corresponding to symmetric and antisymmetric  $\text{CH}_2$  stretching modes cannot be assigned definitively because of severe interference of HCl and its complexes. Table 4 lists all vibrational modes with IR intensity greater than 10  $\text{km mol}^{-1}$ . A total of 15 lines in group C were assigned to 7-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (Table 4). All listed modes (with harmonic IR intensity >10  $\text{km mol}^{-1}$ ) except two CH-stretching mode and a line predicted near 1435  $\text{cm}^{-1}$  have been observed; the latter line was predicted to have IR intensity of only 3  $\text{km mol}^{-1}$  in anharmonic calculations. The average absolute deviation between observed and scaled harmonic (anharmonic) vibrational wavenumbers is  $3.2 \pm 2.3$  ( $9.8 \pm 7.3$ )  $\text{cm}^{-1}$ , indicating a satisfactory agreement between experimental and predicted wavenumbers.

Similarly, lines in group D can be assigned to the 8-isoquinolinyl radical, 8-*iso*- $\text{HC}_9\text{H}_7\text{N}$ .

The five most intense lines observed at 1576.2, 1467.6, 1432.2, 830.5, and 653.7  $\text{cm}^{-1}$  are predicted at 1585 ( $\nu_8$ , ring stretch), 1472 ( $\nu_{11}$ , ring stretch), 1431 ( $\nu_{12}$ ,  $\text{CH}_2$  wag coupled with ring stretch), 834, and 654  $\text{cm}^{-1}$  ( $\nu_{39}$  and  $\nu_{42}$ , both out-of-plane CH bend), respectively. Additional two weak lines were observed at 1245.5 and 941.6  $\text{cm}^{-1}$ , near scaled harmonic vibrational wavenumber predicted at 1259 ( $\nu_{18}$ , ring deform) and 935 ( $\nu_{25}$ , ring deform)  $\text{cm}^{-1}$ , respectively. The simulated spectra for other isomers of isoquinolinyl radicals cannot reproduce the wavenumbers and relative intensities of lines in group D (Figs. S15 and S16, ESI†).

Moreover, two lines at 3083.9 and 3076.6  $\text{cm}^{-1}$  in the CH-stretching region agree with the predicted vibrational wavenumbers of  $\nu_1$  and  $\nu_2$  modes at 3080 (3092) and 3070 (3062)  $\text{cm}^{-1}$ ; values in parentheses are anharmonic vibrational wavenumbers (Fig. S16, ESI†). However, lines in region 2910–2790  $\text{cm}^{-1}$  cannot be assigned definitively because of interference from HCl and its complexes. Table 5 compares the observed wavenumber and relative IR intensities for 9 lines in group D with both scaled harmonic and anharmonic vibrational wavenumbers and relative IR intensities predicted for 8-*iso*- $\text{HC}_9\text{H}_7\text{N}$ . All listed modes (with calculated IR intensity greater than 10  $\text{km mol}^{-1}$ ) have been observed except four CH-stretching modes. The average absolute deviation between observed and scaled harmonic (anharmonic) vibrational wavenumbers is  $4.8 \pm 3.8$  ( $6.5 \pm 3.3$ )  $\text{cm}^{-1}$ .

#### 5.4 Assignments of lines in groups E and F to 4-*iso*- $\text{HC}_9\text{H}_7\text{N}$ and 5-*iso*- $\text{HC}_9\text{H}_7\text{N}$ radicals

Intensities of lines in group E decreased upon secondary photolysis at 540, 420, and 360 nm and remained unchanged on photolysis at 553 and 405 nm; in contrast, intensities of lines in group F increased upon secondary photolysis at 420 and 405 nm, but decreased upon photolysis at 360 nm, and remained unaltered on photolysis at 553 and 540 nm. The difference spectrum in Fig. 3b, recorded after UV/IR irradiation of the matrix, was reproduced in the lower trace of Fig. 6d for the



region 1000–600  $\text{cm}^{-1}$  after truncating negative lines, whereas the difference spectrum upon secondary irradiation at 420 nm, Fig. 3e, was reproduced in the upper trace of Fig. 6d. Lines in groups E and F are indicated with light blue and purple labels and arrows, respectively. Figs. 6e and 6f depict IR stick spectra of 4- and 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N according to predicted scaled vibrational harmonic wavenumbers and IR intensities. Observed lines in groups E and F are compared with the IR stick spectra of all 10 isomers of mono-hydrogenated quinoline in regions 1700–1000 and 3600–3450 and 3100–2700 $\text{cm}^{-1}$  in Figs. S17 and S18, ESI†, respectively.

The spectral pattern of observed lines in group E agrees satisfactorily with the IR stick spectrum of 4-isoquinolinyl radical, 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N, but disagrees with predicted spectra of other mono-hydrogenated radicals. Table 5 compares the observed wavenumber and relative IR intensities with both scaled harmonic and anharmonic vibrational wavenumbers and relative IR intensities predicted for 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N. The two most intense lines observed at 1245.9 and 748.1  $\text{cm}^{-1}$  are predicted at 1249 ( $\nu_{18}$ , ring stretch) and 746 ( $\nu_{40}$ , out-of-plane CH bend). Two weak lines were observed at 1483.7 and 1269.4  $\text{cm}^{-1}$ , near scaled harmonic wavenumber predicted at 1485 ( $\nu_{11}$ , ring stretch) and 1274 ( $\nu_{17}$ , ring stretch)  $\text{cm}^{-1}$ . In the CH-stretching region, one line observed at 3087.9  $\text{cm}^{-1}$  agrees with the predicted scaled harmonic vibrational wavenumber at 3083  $\text{cm}^{-1}$  ( $\nu_1$ ) or the anharmonic vibrational wavenumber 3068  $\text{cm}^{-1}$  ( $\nu_2$ ), as shown in Fig. S18, ESI†. This line is tentatively assigned to  $\nu_1$  of 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N, but identification of other intense lines 3070 (3068,  $\nu_2$ ), 3026 (3012,  $\nu_6$ ), and 2874 (2832,  $\nu_7$ )  $\text{cm}^{-1}$  with IR intensities 21, 25, and 29  $\text{km mol}^{-1}$ , respectively, are difficult because of the severe interference with the parent or products. These five lines in group E are hence assigned to 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N. All lines predicted to have IR intensity greater than 10  $\text{km mol}^{-1}$  except four have been observed. However, the missing lines predicted near 1511 and 782  $\text{cm}^{-1}$  have intensities 5 and 2  $\text{km mol}^{-1}$  in anharmonic calculations. Two

additional lines predicted at 782 and 1417  $\text{cm}^{-1}$  ( $\nu_{39}$  and  $\nu_{13}$ ) with IR intensities 21 and 11  $\text{km mol}^{-1}$ , respectively, could not be positively identified because of the interference with the parent and with 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N, respectively. The average absolute deviation between observed and scaled harmonic (anharmonic) vibrational wavenumbers is  $2.8 \pm 1.2$  ( $3.6 \pm 3.2$ )  $\text{cm}^{-1}$ .

Observed lines in group F agree satisfactorily with IR stick spectrum of 5-isoquinoliny radical, 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N, but disagree with those of other mono-hydrogenated radicals. Table 5 compares the observed wavenumber and relative IR intensities predicted for 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N. Three most intense lines observed at 821.9, 708.6, and 652.3  $\text{cm}^{-1}$  are predicted at 823 ( $\nu_{39}$ , out-of-plane CH bend), 710 ( $\nu_{41}$ , out-of-plane ring deform), and 651  $\text{cm}^{-1}$  ( $\nu_{42}$ , out-of-plane CH bend). Two additional lines were observed at 1402.0 and 1201.3  $\text{cm}^{-1}$ , near scaled harmonic vibrational wavenumber predicted at 1418 ( $\nu_{13}$ , CH<sub>2</sub> scissor) and 1202 ( $\nu_{20}$ , CH<sub>2</sub> wag coupled with ring stretch). Only two lines of 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N with predicted IR intensities greater than 10  $\text{km mol}^{-1}$  in the region 1700–600  $\text{cm}^{-1}$ , 1582 ( $\nu_8$ , 18  $\text{km mol}^{-1}$ ) and 1426  $\text{cm}^{-1}$  ( $\nu_{12}$ , 10  $\text{km mol}^{-1}$ ) remain unobserved because of interference; the former might be interfered with by the line of group A at 1574.4  $\text{cm}^{-1}$  and the latter one has intensity of 5  $\text{km mol}^{-1}$  according to anharmonic calculations. In the CH-stretching region, one line observed at 3083.2  $\text{cm}^{-1}$  agrees with the predicted scaled harmonic (anharmonic) vibrational wavenumbers at 3080 (3082,  $\nu_1$ )  $\text{cm}^{-1}$ , as shown in Fig. S18, ESI†. This line is tentatively assigned to the  $\nu_1$  mode of 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N, but identification of other intense lines 3070 (3054,  $\nu_2$ ), 3064 (3038,  $\nu_3$ ), 3031 (2995,  $\nu_6$ ), and 2874 (2841,  $\nu_7$ )  $\text{cm}^{-1}$  is difficult because of the severe interference of parent or products in the CH-stretching region. These six lines in group F are hence assigned to 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N; the average absolute deviation between observed and scaled harmonic (anharmonic) vibrational wavenumbers is  $4.1 \pm 5.9$  ( $2.3 \pm 2.1$ )  $\text{cm}^{-1}$ .

### 5.5 Assignments of lines in group G and H to 1-*iso*-HC<sub>9</sub>H<sub>7</sub>N and 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N radicals

Apart from those lines assigned to groups A–E, we observed four lines in group G with intensities that increased upon secondary irradiation at 405 nm, decreased upon secondary irradiation at 360 nm, and remained unchanged on irradiation at 553, 540, and 420 nm. We also observed two weak lines with intensities decreased upon secondary irradiation at 553 and 405 nm and remained unaltered upon irradiation at other wavelengths; these lines are categorized as group H. The difference spectrum in Fig. 3c, recorded after UV/IR irradiation of the matrix, was reproduced in the lower trace of Fig. 6g after truncating negative lines; the difference spectrum recorded upon secondary irradiation of the matrix at 360 nm, Fig. 3g, was reproduced in the upper trace of Fig. 6g for region 1000–600  $\text{cm}^{-1}$ . Lines in groups G and H are denoted with orange and olive labels and arrows, respectively. Figures 6h and 6i depict IR stick spectra of 1- and 3-*iso*- $\text{HC}_9\text{H}_7\text{N}$ , simulated according to scaled vibrational harmonic wavenumbers, for comparison with observed lines. Observed lines in groups G and H in region 1700–1000  $\text{cm}^{-1}$  are compared with the IR stick spectra of all 10 isomers of mono-hydrogenated quinoline, as illustrated in Fig. S19, ESI†.

Even though only four lines were observed in group G, they correspond well with lines of 1-*iso*- $\text{HC}_9\text{H}_7\text{N}$  with predicted IR intensities greater than 10  $\text{km mol}^{-1}$  in region 1700–600  $\text{cm}^{-1}$  except two lines predicted at 945 and 699  $\text{cm}^{-1}$  ( $\nu_{25}$  and  $\nu_{42}$  with IR intensity 10 and 15  $\text{km mol}^{-1}$ , respectively); the absence of these two lines is due to interference from the parent and 6-*iso*- $\text{HC}_9\text{H}_7\text{N}$  (group B), respectively. Table 6 compares the observed wavenumber and relative IR intensities with both scaled harmonic and anharmonic vibrational wavenumbers and relative IR intensities predicted for 1-*iso*- $\text{HC}_9\text{H}_7\text{N}$ . The four lines observed at 1416.6, 918.3, 769.2 and 714.6  $\text{cm}^{-1}$  were predicted at 1414 ( $\nu_{13}$ ,  $\text{CH}_2$  scissor coupled with in-plane C3H bend), 920 ( $\nu_{26}$ , ring deform), 772 and 715  $\text{cm}^{-1}$  ( $\nu_{40}$  and  $\nu_{41}$ , both out-of-plane CH bend). Similarly, the CH-stretching

region could not be characterized because of severe interference from the parent and products. Even though only a limited number of lines were observed, we feel confident to assign these lines in group G to 1-*iso*-HC<sub>9</sub>H<sub>7</sub>N because major lines were observed. The average absolute deviation between observed and scaled harmonic (anharmonic) vibrational wavenumbers is  $2.6 \pm 1.9$  ( $9.6 \pm 7.1$ ) cm<sup>-1</sup>.

Two lines in group H observed at 1642.7 and 876.4 cm<sup>-1</sup>, similar to those predicted at 1644 (ν<sub>8</sub>, CN stretch coupled with CH<sub>2</sub> scissor) and 870 (ν<sub>26</sub>, ring stretch) cm<sup>-1</sup>, which are the second and third most intense lines of 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N in region 1700–600 cm<sup>-1</sup>. However, definitive identification of the most intense line predicted at 746 cm<sup>-1</sup> was unachievable because of interference from the parent and the most intense line in group E (4-*iso*-HC<sub>9</sub>H<sub>7</sub>N) at 748.1 cm<sup>-1</sup>. The four other missing lines in region 600–1700 cm<sup>-1</sup> were predicted to have IR intensities less than 14 km mol<sup>-1</sup>, and all of them were predicted to have IR intensities less than 9 km mol<sup>-1</sup> with the anharmonic calculations. The comparison of observed wavenumbers and those predicted for 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N is listed in Table 6. Because only two lines were observed, we hence tentatively assigned the observed two lines in group H to 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N. The smallest number of observed lines is consistent with the estimated smallest mixing ratio,  $0.42 \pm 0.01$  ppm, among all observed isomers.

## 5.6 Mechanism of formation in *p*-H<sub>2</sub>

The mechanism to form 2-isoquinolinium (*iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>) is expected to be similar to those for the formation of 1-quinolinium cation<sup>25</sup> and protonated cyclic compounds including PAH.<sup>22, 23</sup> During deposition of *p*-H<sub>2</sub> under electron bombardment, H<sub>2</sub> molecules are ionized by electron to produce H<sub>2</sub><sup>+</sup>; the H<sub>2</sub><sup>+</sup> product readily react with another H<sub>2</sub> molecule to produce H and H<sub>3</sub><sup>+</sup>. Because the proton affinity of isoquinoline, 957 kJ mol<sup>-1</sup> for the formation of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> and

$>689 \text{ kJ mol}^{-1}$  for the formation of  $iso\text{-H}^+\text{C}_9\text{H}_7\text{N}$ , is much greater than that of  $\text{H}_2$ ,  $422 \text{ kJ mol}^{-1}$ ,<sup>20</sup> proton transfer from  $\text{H}_3^+$  to isoquinoline occurs readily. The exothermicity of the proton transfer reaction,  $\text{H}_3^+ + iso\text{-C}_9\text{H}_7\text{N} \rightarrow \text{H}_2 + iso\text{-C}_9\text{H}_8\text{N}^+$ , is significantly greater than the energy barriers for proton transfer between isomers of  $iso\text{-C}_9\text{H}_8\text{N}^+$  (Fig. S2, ESI†); all isomers can be produced, in principle, via the proton transfer reaction. However, the observation of only  $iso\text{-C}_9\text{H}_7\text{NH}^+$  in our experiments indicates that either the proton transfer to the N site of isoquinoline is significantly more favorable or the energy dissipation by solid  $p\text{-H}_2$  is efficient enough to quench the isomerization process, so that only the thermodynamically most stable  $iso\text{-C}_9\text{H}_7\text{NH}^+$  was produced. Moreover, no fragmentation or ring-opening products of  $iso\text{-C}_9\text{H}_8\text{N}^+$  could be identified in these experiments. Hansen et al. reported the significant formation of HCN in their UV photodissociation experiments of protonated quinoline and isoquinoline,<sup>41</sup> but we did not observe HCN (expected to be near  $3303$  and  $720 \text{ cm}^{-1}$ )<sup>42</sup> either. These authors calculated an energy threshold of  $\sim 448 \text{ kJ mol}^{-1}$  for the formation of HCN and ring-opening coproducts. Even though this energy is smaller than the exothermicity of  $\text{H}_3^+ + iso\text{-C}_9\text{H}_7\text{NH} \rightarrow iso\text{-C}_9\text{H}_7\text{NH}^+ + \text{H}_2$ ,  $\sim 534 \text{ kJ mol}^{-1}$ , one would expect the excess energy of  $iso\text{-C}_9\text{H}_7\text{NH}^+$  be rapidly dissipated in solid  $p\text{-H}_2$  below its dissociation threshold.

In electron bombardment experiments, the isoquinolinyl radicals are expected to be formed on neutralization of isoquinolinium cation,  $iso\text{-C}_9\text{H}_8\text{N}^+ + \text{e}^- \rightarrow iso\text{-C}_9\text{H}_8\text{N}$ , or by the reaction of  $iso\text{-C}_9\text{H}_7\text{N}$  with an H atom,  $\text{H} + iso\text{-C}_9\text{H}_7\text{N} \rightarrow iso\text{-C}_9\text{H}_8\text{N}$ . The former process tends to produce the neutral counterpart of protonated species; therefore, the formation of  $iso\text{-C}_9\text{H}_7\text{NH}$  is expected, which is in accordance with our experiments that lines in group A (assigned to  $iso\text{-C}_9\text{H}_7\text{NH}$ ) were observed. The mixing ratios of  $iso\text{-C}_9\text{H}_7\text{NH}^+$  and  $iso\text{-C}_9\text{H}_7\text{NH}$  after electron-bombarded matrix deposition were estimated to be  $0.65 \pm 0.09$  and  $2.2 \pm 1.1 \text{ ppm}$ , respectively, in which errors

represent the standard deviations in averaging mixing ratios derived from several line. After maintenance of the matrix in darkness for 10 h, the mixing ratio of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> decreased by  $0.25 \pm 0.16$  ppm and that of *iso*-C<sub>9</sub>H<sub>7</sub>NH increased by  $0.9 \pm 0.5$  ppm, indicating that the observed abundance of *iso*-C<sub>9</sub>H<sub>7</sub>NH might not be explained by considering the neutralization process alone. However, because of the large errors that might be associated with the predicted IR intensities, we cannot exclude the possibility that the mixing ratio of *iso*-C<sub>9</sub>H<sub>7</sub>NH was overestimated as compared with that of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>. In the electron-bombardment experiments, for hydrogenated species, in addition to *iso*-C<sub>9</sub>H<sub>7</sub>NH, we were able to identify clearly only the formation of 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N radical (lines in group B). The mixing ratio of 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N after the matrix deposition was estimated to be  $0.9 \pm 0.3$  ppm and its increase after maintenance of the matrix in darkness for 10 h is  $1.0 \pm 0.4$  ppm, similar to the increase of *iso*-C<sub>9</sub>H<sub>7</sub>NH ( $0.9 \pm 0.5$  ppm). It is unclear why only 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N among all possible carbon sites for hydrogenation was observed.

For H-addition reactions, because they have barriers  $> 20$  kJ mol<sup>-1</sup>, quantum tunneling is required in low-temperature solid *p*-H<sub>2</sub>. The probability of tunneling reaction depends on the width and height of the barrier. The barrier height for the formation of isoquinolinyl radicals are predicted to be 21–27 kJ mol<sup>-1</sup> for *iso*-C<sub>9</sub>H<sub>7</sub>NH and 1-, 3-, 4-, 5-, 6-, 7-, 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N, respectively, but much greater ( $\sim 50$  kJ mol<sup>-1</sup>) for 4a- and 8a-*iso*-HC<sub>9</sub>H<sub>7</sub>N. It is hence conceivable that all isomers of *iso*-C<sub>9</sub>H<sub>8</sub>N except 4a- and 8a-*iso*-HC<sub>9</sub>H<sub>7</sub>N can be produced from the reaction of H with isoquinoline.

In the H-atom reactions, 8 isomers of *iso*-C<sub>9</sub>H<sub>8</sub>N were observed; the estimated mixing ratios for *iso*-C<sub>9</sub>H<sub>7</sub>NH, 1-, 3-, 4-, 5-, 6-, 7-, and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N are  $2.8 \pm 0.5$ ,  $1.5 \pm 0.3$ ,  $0.42 \pm 0.01$ ,  $2.5 \pm 0.3$ ,  $1.5 \pm 0.3$ ,  $2.1 \pm 1.2$ ,  $1.1 \pm 0.4$ , and  $1.6 \pm 0.1$  ppm, respectively. Because of errors in the calculated IR intensities and barriers of hydrogenation reactions, obtaining a definitive relationship

between observed mixing ratios and reaction barriers is difficult. The rate of tunneling reactions depends on the barrier height and the width of the barrier, which requires more sophisticated quantum-chemical calculations. No apparent correlation between observed mixing ratios and barriers for formation calculated with the CCSD(T) method could be derived. Nevertheless, if one compares the barriers calculated with the B3LYP method, formation of *iso*-C<sub>9</sub>H<sub>7</sub>NH is associated with a significantly smaller barrier ( $\sim 6 \text{ kJ mol}^{-1}$  vs.  $> 11 \text{ kJ mol}^{-1}$ ) and the largest mixing ratio ( $2.8 \pm 0.5 \text{ ppm}$ ) for *iso*-C<sub>9</sub>H<sub>7</sub>NH; formation of 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N is associated with the largest barrier ( $15 \text{ kJ mol}^{-1}$  from B3LYP and  $27 \text{ kJ mol}^{-1}$  from CCSD(T)) and the smallest mixing ratio ( $0.42 \pm 0.01 \text{ ppm}$ ).

## 5.7 Secondary irradiation

The UV spectra of *iso*-C<sub>9</sub>H<sub>7</sub>N and all 10 mono-hydrogenated isomers of *iso*-C<sub>9</sub>H<sub>7</sub>N predicted with the TD-B3LYP/6-311++G(d,p) method are depicted in Fig. S20, ESI†; associated wavelengths and oscillator strengths are listed in Tables S11–S13, ESI†. The calculated UV spectra provide satisfactory predictions in the choice of wavelengths used for secondary irradiation. The wavelengths for secondary irradiation were selected to be greater than 300 nm to avoid photolysis of the parent, which was predicted to have UV absorption starting near 300 nm. The significant destruction of 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N and 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N upon photolysis at 553 nm agrees well with the predicted absorption bands near 544 and 520 nm, corresponding to HOMO  $\rightarrow$  LUMO transition for both isomers. The increase of 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N at this wavelength seems to indicate that 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N  $\rightarrow$  8-*iso*-HC<sub>9</sub>H<sub>7</sub>N isomerization occurs. We could hence differentiate these three mono-hydrogenated isomers of *iso*-C<sub>9</sub>H<sub>7</sub>N upon irradiation at 553 nm. At 420 and 405 nm, 1-, 5- and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N either increase in intensity or remained nearly the same; these three isomers are the ones that have no significant absorption for wavelength greater than 320 nm; the increase might

be due to the hydrogen transfer from a neighboring site. For 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N, it has no significant absorption for wavelength greater than 370 nm, so we observed no change at 420 nm and a small decrease at 405 nm. The destruction of 1-, 4-, 5-, and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N upon irradiation at 360 nm is consistent with the predicted absorption bands at 322/347, 324/342, 330, and 322/348 nm, respectively, with oscillator strength greater than 0.11. The destruction of *iso*-C<sub>9</sub>H<sub>7</sub>NH upon irradiation at 540, 420, and 405 nm agrees with the predicted absorption bands near 585/483 and 426 nm.

### 5.8 Implication to astrochemistry

The UIR emission spectrum observed from the Orion Bar PDR<sup>43</sup> is compared with the stick IR spectra of protonated isoquinoline, protonated quinoline,<sup>25</sup> and protonated naphthalene (1-C<sub>10</sub>H<sub>9</sub><sup>+</sup>)<sup>27</sup> in Fig. S21, ESI†. In the region ~6.2 μm, the line of 1-C<sub>10</sub>H<sub>9</sub><sup>+</sup> in solid *p*-H<sub>2</sub> at 1618.7 cm<sup>-1</sup> (~6.18 μm) is blue-shifted to 1643.3 cm<sup>-1</sup> (ν<sub>9</sub>, ~6.09 μm) in protonated isoquinoline and to 1641.4 cm<sup>-1</sup> (~6.09 μm) in protonated quinoline. The line of *iso*-C<sub>9</sub>H<sub>7</sub>NH in solid *p*-H<sub>2</sub> at 1554.7 cm<sup>-1</sup> (ν<sub>12</sub>) is also blue-shifted from that of 1-C<sub>10</sub>H<sub>9</sub><sup>+</sup> at 1510 cm<sup>-1</sup>. Two additional lines of 1-C<sub>10</sub>H<sub>9</sub><sup>+</sup> in solid *p*-H<sub>2</sub> were observed at 1556.7 and 1580.8 cm<sup>-1</sup>, the corresponding lines of *iso*-C<sub>9</sub>H<sub>7</sub>NH were predicted at 1587 and 1620 cm<sup>-1</sup>, but was not observed due to interference. Because we recorded absorption spectra, which should be slightly different from the UIR emission spectra. Typically, the maxima of emission bands are slightly red-shifted from those of absorption bands because of the anharmonicity associated with the vibrationally excited emitting states produced on UV irradiation.<sup>14,44</sup> Thus, the band associated with ring stretching modes of protonated quinoline and isoquinoline match the 6.2-μm UIR band better than those of protonated naphthalene. Relatively intense lines of the out-of-plane CH-bending mode were observed at 791.4 cm<sup>-1</sup> (12.64 μm) in protonated isoquinoline and at 790.2 cm<sup>-1</sup> (12.65 μm) in protonated quinoline, matching



the 12.8- $\mu\text{m}$  band of UIR emission, although the UIR bands in region 12.4–13.3  $\mu\text{m}$  are generally considered to originate from out-of-plane C–H bending modes associated with trio CH that has three adjacent peripheral CH moieties.<sup>45,46,47</sup> The other UIR bands near 7.7, 8.6, and 11.2  $\mu\text{m}$  are unobserved or weak in the absorption spectra of protonated quinoline and protonated isoquinoline; these two protonated species are hence unlikely to be the carrier of the UIR bands. This is consistent with the expectation that quinoline and isoquinoline are too small to survive the intense UV field in interstellar media. Therefore, it is desirable to extend the current investigation to larger  $\text{H}^+\text{PANH}$  to confirm that the blue shifts of ring stretching modes and the appearance of the band  $\sim 12.7$   $\mu\text{m}$  are common in  $\text{H}^+\text{PANH}$ .

## 6. Conclusion

Electron bombardment during matrix deposition of a mixture of isoquinoline (*iso*- $\text{C}_9\text{H}_7\text{N}$ ) and *p*- $\text{H}_2$  produced *iso*- $\text{C}_9\text{H}_7\text{NH}^+$ , *iso*- $\text{C}_9\text{H}_7\text{NH}$ , and 6-*iso*- $\text{HC}_9\text{H}_7\text{N}$ ; their IR lines were classified according to their behavior after maintenance of the matrix in darkness for prolonged time and upon secondary photolysis at 365 and 405 nm. To enhance the hydrogenation reactions, we irradiated an *iso*- $\text{C}_9\text{H}_7\text{N}/\text{Cl}_2/p\text{-H}_2$  matrix at 365 nm, followed by IR light to generate H atoms for the reaction with *iso*- $\text{C}_9\text{H}_7\text{N}$ ; the formation of all 8 isomers of *iso*- $\text{C}_9\text{H}_8\text{N}$  radicals singly hydrogenated on non-ring-sharing atoms were observed, indicating the  $\text{Cl}$ -induced hydrogenation reaction is more effective than electron bombardment for hydrogenation. The IR lines observed in  $\text{Cl}_2$ -doped experiments were classified into eight groups and assigned to *iso*- $\text{C}_9\text{H}_7\text{NH}$ , 6-, 7-, 8-, 4-, 5-, 1-, and 3-*iso*- $\text{HC}_9\text{H}_7\text{N}$  according to their behavior upon secondary photolysis at 553, 540, 420, 405, and 360 nm and comparison with quantum-chemical calculations at the B3LYP/6-311++G(d,p) level of theory. Nearly all lines predicted to have IR intensities greater than 10  $\text{km mol}^{-1}$  have been observed for *iso*- $\text{C}_9\text{H}_7\text{NH}$ , 6-, 7-, 8-, 4-, and 5-*iso*- $\text{HC}_9\text{H}_7\text{N}$ ; assignment of lines

in group H to 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N is tentative, because only two lines were unambiguously assigned.

Three bands of gaseous *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> were observed previously with the IRMPD method.<sup>22</sup> In addition to these three features, we observed additional 4 lines of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> to provide improved IR characterization of this species. The IR spectra of all isomers of mono-hydrogenated isoquinoline are previously unreported; they provide important information for future identification of these species in combustion or astronomical environments. Observation of all 8 isomers of the feasible mono-hydrogenated PAH in a single experiment has never been demonstrated previously; the high efficiency of H-tunneling reaction, absence of multiple hydrogenation, and simplified spectra with narrow lines in matrix isolation are unique characters of this type of H reactions in solid *p*-H<sub>2</sub> using Cl atoms and IR light to induce hydrogenation reactions and allow identification of various isomers.

Lines of protonated naphthalene (1-C<sub>10</sub>H<sub>9</sub><sup>+</sup>) at 1618.7, 1580.8/1556.7, and 1510.0 cm<sup>-1</sup> (Ref. 27) are blue-shifted to 1641.4, 1598.4, and 1562.0 cm<sup>-1</sup> in protonated quinoline, C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>,<sup>25</sup> and 1643.3, (~1620/1587), and 1554.7 cm<sup>-1</sup> in protonated isoquinoline, *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup>; the predicted lines at 1620/1587 cm<sup>-1</sup> were unobserved because of interference. C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> and *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> are not expected to be candidates for the carrier of the UIR band because they are unlikely to survive the strong UV field in the interstellar media, considering the small sizes. Nevertheless, this observation experimentally supports the proposed blue shift of the ring stretching bands upon substitution of a CH moiety of H<sup>+</sup>PAH with a nitrogen atom to form H<sup>+</sup>PANH; these blue-shifted band matches the UIR bands near 6.2 μm better than those of H<sup>+</sup>PAH.

### Supplementary information

Electronic supplementary information (ESI) available: Tables: scaled harmonic and anharmonic vibrational wavenumbers and IR intensities of 10 isomers of isoquinolinium cation and 10 isomers

of isoquinolinyl radical. Comparison of observed vibrational wavenumbers of isoquinoline with calculated values. Estimated mixing ratios of isoquinolinyl radical. Vertical excitation wavelengths and oscillator strengths of all isomers of isoquinolinyl radicals. Figures: calculated geometries and relative energies of isomers of isoquinolinium cations and isoquinolinyl radicals. Potential energy diagrams for various isomers of isoquinolinium cations and transition states for isomerization and for various isomers of isoquinolinyl radicals and transition states for their formation and isomerization. Predicted IR spectra for 10 isomers of isoquinolinium cations and isoquinolinyl radicals. Comparison of experimental and predicted IR spectra of *iso*-C<sub>9</sub>H<sub>7</sub>N. Calibration curve for observed and calculated harmonic vibrational wavenumbers of *iso*-C<sub>9</sub>H<sub>7</sub>N. Spectra of *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrices at various stages of experiments. Comparison of experimental spectrum of group A<sup>+</sup> with simulated spectra of 10 isomers of isoquinolinium cations. Comparison of experimental spectra of groups A and B, groups C and D, groups E and F, and groups G and H with simulated spectra of 10 isomers of isoquinolinyl radicals. Predicted UV spectra of isomers of hydrogenated isoquinoline. Comparison of experimental spectra of protonated isoquinoline, isoquinoline, and naphthalene with UIR bands observed from Orion Bar PDR. See <https://doi.org/10.1039/d2cpxxxxxxx>.

### **Conflict of interest**

There are no conflicts to declare.

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**Table 1** Relative energies of isomers of protonated isoquinoline and relative energies and barrier heights for the formation of various isomers of mono-hydrogenated isoquinoline from hydrogenation of isoquinoline calculated with the CCSD(T)/6-311++G(d,p)//B3LYP/6-311++G(d,p) method

Site <sup>a</sup>	<i>iso</i> -C <sub>9</sub> H <sub>8</sub> N <sup>+</sup>	<i>iso</i> -C <sub>9</sub> H <sub>8</sub> N	
	Relative energy <sup>b</sup>	Relative energy <sup>b</sup>	Barrier for formation <sup>c</sup>
	/kJ mol <sup>-1</sup>	/kJ mol <sup>-1</sup>	/kJ mol <sup>-1</sup>
1	200 (198) <sup>d</sup>	6 (23) <sup>d</sup>	24 (14) <sup>d</sup>
2	0 <sup>e</sup>	0 <sup>f</sup>	23 (6)
3	190 (186)	22 (39)	27 (15)
4	169 (166)	3 (19)	22 (11)
5	171 (167)	2 (20)	21 (11)
6	209 (204)	24 (41)	26 (14)
7	183 (180)	24 (41)	26 (14)
8	190 (186)	3 (20)	22 (12)
4a	267 (271)	100 (125)	46 (39)
8a	240 (245)	91 (118)	52 (36)

<sup>a</sup> See Fig. 1 for protonation or hydrogenation sites. <sup>b</sup> Zero-point vibrational energies are corrected using unscaled harmonic vibrational wavenumbers calculated with the B3LYP/6-311++G(d,p) method. <sup>c</sup> Energy of the transition state relative to that of H + *iso*-C<sub>9</sub>H<sub>7</sub>N. <sup>d</sup> Energies calculated with the B3LYP//6-311++G(d,p) are listed in parentheses for comparison. <sup>e</sup> The energy of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> + H<sub>2</sub> is smaller than *iso*-C<sub>9</sub>H<sub>7</sub>NH + H<sub>3</sub><sup>+</sup> by 534 (538) kJ mol<sup>-1</sup>, implying a proton affinity of 957 (989) kJ mol<sup>-1</sup> for *iso*-C<sub>9</sub>H<sub>7</sub>N on the N-atom site. <sup>f</sup> The energy of *iso*-C<sub>9</sub>H<sub>7</sub>NH is smaller than H + *iso*-C<sub>9</sub>H<sub>7</sub>N by 111 (148) kJ mol<sup>-1</sup>.

**Table 2** Comparison of observed vibrational wavenumbers ( $\text{cm}^{-1}$ ) and relative IR intensities of lines in group A<sup>+</sup> with scaled harmonic and anharmonic vibrational wavenumbers ( $\text{cm}^{-1}$ ) and IR intensities ( $\text{km mol}^{-1}$ ) of *iso*-C<sub>9</sub>H<sub>7</sub>NH<sup>+</sup> predicted with the B3LYP/6-311++G(d, p) method

Mode	Sym.	B3LYP/6-311++G(d, p)				Experiments	
		Scaled harmonic		Anharmonic		IRMPD <sup>a</sup>	<i>p</i> -H <sub>2</sub>
$\nu_1$	a'	3430 <sup>b</sup>	(208) <sup>c</sup>	3407	(180) <sup>c</sup>		3397.2 (250) <sup>d</sup>
$\nu_9$	a'	1639	(83)	1625	(60)	1582	1643.3 (40)
$\nu_{10}$	a'	1620	(27)	1610	(19)		
$\nu_{11}$	a'	1587	(15)	1576	(11)		
$\nu_{12}$	a'	1559	(7)	1550	(10)		1554.7 (10)
$\nu_{13}$	a'	1500	(14)	1494	(12)		1497.3 (7)
$\nu_{16}$	a'	1394	(43)	1382	(40)	1360	1400.1 (69)
$\nu_{17}$	a'	1382	(49)	1378	(25)		1378.6 (31)
$\nu_{18}$	a'	1285	(13)	1285	(12)		
$\nu_{40}$	a''	788	(150)	806	(39)	774	791.4 (100)
$\nu_{42}$	a''	724	(23)	715	(58)		<sup>e</sup>
$\nu_{44}$	a''	471	(29)	471	(32)		

<sup>a</sup> Infrared multiphoton dissociation (IRMPD) experiments by Galu   *et al.*<sup>22</sup> <sup>b</sup> Harmonic vibrational wavenumbers scaled according to  $0.9122 x + 174.3$  for wavenumbers  $>2000 \text{ cm}^{-1}$  and  $0.9823 x + 0.6$  for wavenumbers  $<2000 \text{ cm}^{-1}$ . <sup>c</sup> The IR intensities in  $\text{km mol}^{-1}$  are listed in parentheses; only modes with IR intensities  $> 10 \text{ km mol}^{-1}$  are listed, unless a mode was observed. <sup>d</sup> Percentage IR intensities with respect to that of the line at  $791.4 \text{ cm}^{-1}$  are listed in parentheses. <sup>e</sup> Interference by absorption of *p*-H<sub>2</sub>.

**Table 3** Summary of behaviors on secondary irradiation of various groups of lines of *iso*-C<sub>9</sub>H<sub>8</sub>N

Group	Assignments	553 nm	540 nm	420 nm	405 nm	360 nm
A	<i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	no change	decrease (~ 6 %) <sup>b</sup>	decrease* <sup>a</sup> (~ 11 %)	decrease (~ 9 %)	no change
B	6- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N	no change	no change	decrease* (~ 12 %)	decrease (~ 8 %)	no change
C	7- <i>iso</i> HC <sub>9</sub> H <sub>7</sub> N	decrease* (~ 13 %)	decrease (~ 3 %)	decrease (~ 8 %)	decrease (~ 6 %)	no change
D	8- <i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	increase* (~ 15 %)	no change	increase (~ 3%)	increase (~ 1%)	decrease (~ 2 %)
E	4- <i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	no change	decrease* (~ 7 %)	decrease (~ 2%)	no change	decrease (~ 1 %)
F	5- <i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	no change	no change	increase* (~ 6%)	increase (~ 2 %)	decrease (~ 2 %)
G	1- <i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	no change	no change	no change	increase (~ 3 %)	decrease* (~ 5 %)
H	3- <i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH	decrease* (~ 4 %)	no change	no change	decrease (~ 2 %)	no change

<sup>a</sup>The most significant variation in each species among all wavelengths is marked with \*. <sup>b</sup>The percentage variation is based on the spectrum measured after IR irradiation.

**Table 4** Comparison of observed vibrational wavenumbers (in  $\text{cm}^{-1}$ ) and relative IR intensities of lines in groups A–C with scaled harmonic and anharmonic vibrational wavenumbers and IR intensities of *iso*-C<sub>9</sub>H<sub>7</sub>NH, 6-*iso*-HC<sub>9</sub>H<sub>7</sub>N, and 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N predicted with the B3LYP/6-311++G(d, p) metho

Mode	Sym.	<i>iso</i> -C <sub>9</sub> H <sub>7</sub> NH (group A)			6- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N (group B)			7- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N (group C)		
		Scaled harm. <sup>a</sup>	Anharmonic	Experiment	Scaled harm. <sup>a</sup>	Anharmonic	Experiment	Scaled harm. <sup>a</sup>	Anharmonic	Experiment
v <sub>1</sub>	a'	3518 (107) <sup>b</sup>	3476 (58) <sup>b</sup>	3495.2 (105) <sup>c</sup> 3491.0 (45)	3071 (20)	3054 (17)	3059.0 <sup>d</sup>	3069 (27)	3028 (4)	<sup>f</sup>
v <sub>2</sub>	a'				3067 (27)	3041 (29)	3045.1 <sup>d</sup>	3065 (20)	3046 (19)	<sup>f</sup>
v <sub>3</sub>	a'	3101 (12)	3071 (9)	<sup>f</sup>	3059 (10)	3032 (32)	3025.3 <sup>d</sup>	3059 (17)	3044 (41)	3057.8 <sup>d</sup>
v <sub>4</sub>	a'	3083 (15)	3074 (13)	<sup>f</sup>						
v <sub>5</sub>	a'				3048 (14)	3017 (9)	<sup>f</sup>			
v <sub>6</sub>	a'	3066 (27)	3051 (43)	<sup>f</sup>	3030 (25)	3014 (10)	<sup>f</sup>	3028 (20)	3000 (12)	3025.6 <sup>d</sup>
v <sub>7</sub>	a'				2850 (39)	2808 (4)	2826.4 (69) <sup>c</sup>	2849 (48)	2759 (6)	(HCl) <sup>f</sup>
v <sub>8</sub>	a'	3050 (10)	3006 (1)	<sup>f</sup>				1641 (2)	1624 (0)	1617.2 (11) <sup>c</sup>
v <sub>9</sub>	a'	1614 (48)	1603 (31)	1613.9 (81)	1560 (25)	1550 (19)	1557.5 (35)	1553 (35)	1542 (16)	1551.2 (69)
v <sub>10</sub>	a'	1577 (20)	1565 (6)	1574.4 (18)						
v <sub>11</sub>	a'				1450 (22)	1440 (16)	1446.8 (46)			
v <sub>12</sub>	a'	1489 (11)	1484 (10)	1490.2 (13)				1435 (10)	1430 (3)	(D) <sup>f</sup>
v <sub>13</sub>	a'	1478 (25)	1465 (2)	1476.6 (58)	1419 (11)	1420 (1)	(A) <sup>f</sup>			
v <sub>14</sub>	a'	1440 (17)	1436 (10)	1440.5 (44)						
v <sub>15</sub>	a'	1411 (14)	1400 (7)	1420.7 (20)						
v <sub>16</sub>	a'				1337 (8)	1322 (8)	1334.1 (18)	1334 (4)	1319 (4)	1331.5 (29)
v <sub>17</sub>	a'				1282 (13)	1282 (10)	1284.5 (52)	1286 (10)	1283 (1)	1284.3 (51)
v <sub>20</sub>	a'				1185 (7)	1186 (4)	1184.4 (22)	1208 (11)	1201 (5)	1205.1 (14)
v <sub>21</sub>	a'	1181 (22)	1165 (14)	1180.9 (33)						
v <sub>23</sub>	a'				1048 (5)	1047 (4)	1041.5? (8)			
v <sub>24</sub>	a'	1115 (10)	1114 (16)	1115.3 (13)				1016 (4)	1017 (3)	1017.9 (11)
v <sub>25</sub>	a'	1026 (6)	1026 (5)	1022.0 (8)				914 (6)	914 (5)	907.6 (8)
v <sub>26</sub>	a'	982 (90)	980 (61)	988.3 (100)	884 (21)	885 (19)	890.9 (54)	883 (12)	884 (11)	892.0 (60)
v <sub>27</sub>	a'	932 (5)	937 (14)	939.2 (17)						
v <sub>33</sub>	a'				2840 (13)	2755 (18)	2802.5 (49)	2838 (15)	2753 (21)	(HCl) <sup>f</sup>
v <sub>36</sub>	a''							935 (3)	958 (3)	935.4? (11)
v <sub>37</sub>	a''				920 (9)	931 (2)	921.1 (25)	912 (12)	911 (13)	908.4 (51)
v <sub>38</sub>	a''	761 (49)	760 (67)	763.1 (66)						
v <sub>39</sub>	a''	735 (77)	737 (17)	737.3 (78)	812 (36)	815 (30)	810.2 (100)	821 (45)	828 (38)	817.5 (100)
v <sub>41</sub>	a''				698 (24)	700 (26)	697.0 (36)	694 (0)	715 (4)	720.9?
v <sub>42</sub>	a''	571 (17)	550 (15)	577.2 (6)	686 (6)	689 (4)	688.8 (29)	681 (13)	667 (7)	682.2 (46)

<sup>a</sup> Harmonic vibrational wavenumbers scaled according to  $0.9122x + 174.3$  for wavenumbers  $>2000 \text{ cm}^{-1}$  and  $0.9823x + 0.6$  for wavenumbers  $<2000 \text{ cm}^{-1}$ . <sup>b</sup> IR intensities in  $\text{km mol}^{-1}$ . <sup>c</sup> IR intensities as per cent of that of the most intense line in each species. <sup>d</sup> Severe interference of parent absorption lines intricate the estimation of observed IR intensities. <sup>f</sup> Interference due to absorption of parent or products; in the latter case, the group number of the interfering line or HCl are indicated in parentheses.



**Table 5** Comparison of observed vibrational wavenumbers (in  $\text{cm}^{-1}$ ) and relative IR intensities of lines in groups D–F with scaled harmonic and anharmonic vibrational wavenumbers and IR intensities of 8-, 4-, and 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N predicted with the B3LYP/6-311++G(d, p) method

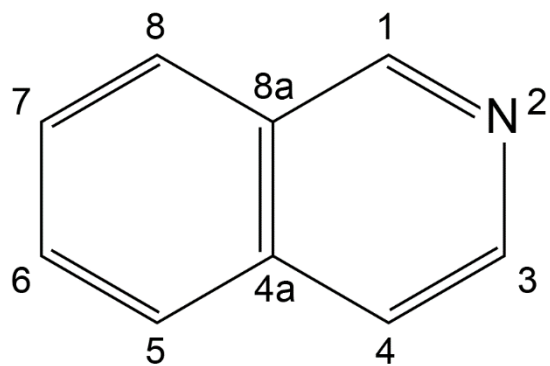
Mode Sym.		8- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N (group D)			4- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N (group E)			5- <i>iso</i> -HC <sub>9</sub> H <sub>7</sub> N (group F)		
		Scaled harm. <sup>a</sup>	Anharmonic	Experiment	Scaled harm. <sup>a</sup>	Anharmonic	Experiment	Scaled harm. <sup>a</sup>	Anharmonic	Experiment
$\nu_1$	a'	3080 (16) <sup>b</sup>	3092 (21) <sup>b</sup>	3083.9 <sup>c</sup>	3083 (18)	3041 (21)	3087.9 <sup>c</sup>	3080 (16)	3082 (10)	3083.2 <sup>c</sup>
$\nu_2$	a'	3070 (35)	3062 (4)	3076.6 <sup>c</sup>	3070 (21)	3068 (24)	<sup>e</sup>	3070 (17)	3054 (15)	<sup>e</sup>
$\nu_3$	a'							3064 (26)	3038 (36)	<sup>e</sup>
$\nu_5$	a'	3046 (17)	3007 (23)	<sup>e</sup>						
$\nu_6$	a'	3024 (25)	2976 (12)	<sup>e</sup>	3026 (25)	3012 (24)	<sup>e</sup>	3031 (22)	2995 (10)	<sup>e</sup>
$\nu_7$	a'	2874 (20)	2838 (5)	(HCl) <sup>e</sup>	2874 (29)	2832 (7)	(HCl) <sup>e</sup>	2874 (26)	2841 (7)	(HCl) <sup>e</sup>
$\nu_8$	a'	1585 (52)	1569 (32)	1576.2 <sup>c</sup>				1582 (18)	1568 (10)	<sup>e</sup>
$\nu_{10}$	a'				1511 (10)	1493 (5)	<sup>e</sup>			
$\nu_{11}$	a'	1472 (14)	1464 (13)	1467.6 (85) <sup>d</sup>	1485 (11)	1477 (8)	1483.7 (20) <sup>d</sup>			
$\nu_{12}$	a'	1431 (19)	1425 (4)	1432.2 (43)				1426 (10)	1415 (5)	(A) <sup>e</sup>
$\nu_{13}$	a'				1417 (11)	1398 (10)	(A) <sup>e</sup>	1418 (14)	1396 (6)	1402.0 (45) <sup>d</sup>
$\nu_{17}$	a'				1274 (10)	1270 (2)	1269.4 (14)			
$\nu_{18}$	a'	1259 (10)	1249 (7)	1245.5 (24)	1249 (25)	1239 (21)	1245.9 (34)			
$\nu_{20}$	a'							1202 (10)	1200 (9)	1201.3 (42)
$\nu_{25}$	a'	935 (17)	937 (15)	941.6 (29)						
$\nu_{33}$	a''	2875 (10)	2796 (14)	(HCl) <sup>e</sup>	2877 (11)	2794 (15)	(HCl) <sup>e</sup>	2875 (11)	2797 (15)	(HCl) <sup>e</sup>
$\nu_{39}$	a''	834 (39)	835 (31)	830.5 (62)	782 (21)	815 (2)	<sup>e</sup>	823 (30)	825 (23)	821.9 (67)
$\nu_{40}$	a''				746 (64)	748 (65)	748.1 (100)			
$\nu_{41}$	a''							710 (15)	709 (9)	708.6 (50)
$\nu_{42}$	a''	654 (43)	657 (38)	653.7 (100)				651 (45)	653 (39)	652.3 (100)

<sup>a</sup> Harmonic vibrational wavenumbers scaled according to  $0.9122x + 174.3$  for wavenumbers  $>2000 \text{ cm}^{-1}$  and  $0.9823x + 0.6$  for wavenumbers  $<2000 \text{ cm}^{-1}$ . <sup>b</sup> IR intensities in  $\text{km mol}^{-1}$ . <sup>c</sup> Severe interference of parent absorption lines intricate the estimation of observed IR intensities. <sup>d</sup> IR intensities as per cent of that of the most intense line in each species. <sup>e</sup> Interference due to absorption of parent or products; in the latter case, the group number of the interfering line or HCl are indicated in parentheses.

**Table 6** Comparison of observed vibrational wavenumbers (in  $\text{cm}^{-1}$ ) and relative IR intensities of lines in groups G and H with scaled harmonic and anharmonic vibrational wavenumbers and IR intensities of  $\text{C}_9\text{H}_8\text{N}$  predicted with the B3LYP/6-311++G(d, p) method

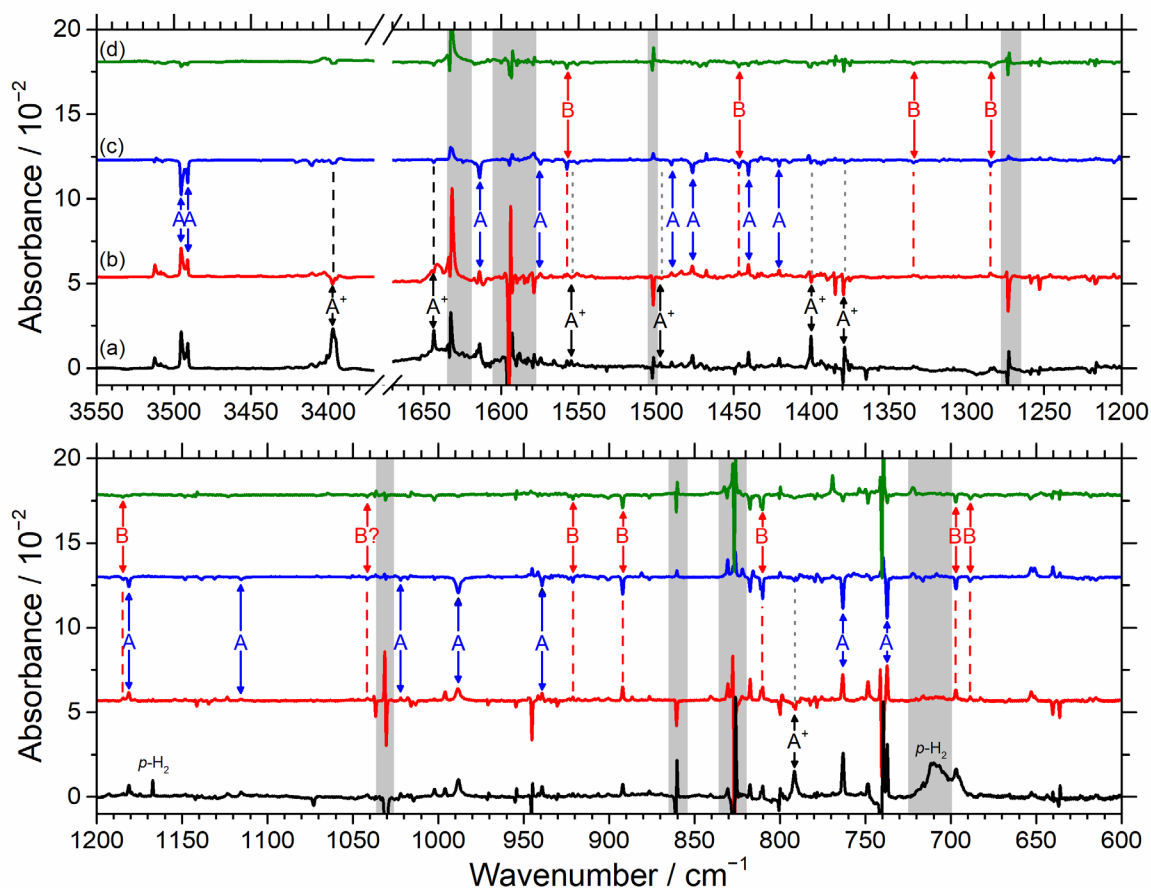
Mode	Sym.	1- <i>iso</i> - $\text{HC}_9\text{H}_7\text{N}$ (group G)			3- <i>iso</i> - $\text{HC}_9\text{H}_7\text{N}$ (group H)		
		Scaled harm. <sup>a</sup>	Anharmonic	Experiment	Scaled harm. <sup>a</sup>	Anharmonic	Experiment
$\nu_1$	a'	3083 (17) <sup>b</sup>	3068 (6) <sup>b</sup>	<sup>d</sup>	3085 (16) <sup>b</sup>	3072 (22) <sup>b</sup>	<sup>d</sup>
$\nu_2$	a'	3076 (21)	3052 (8)	<sup>d</sup>	3072 (21)	3047 (21)	<sup>d</sup>
$\nu_3$	a'	3070 (15)	3076 (44)	<sup>d</sup>			<sup>d</sup>
$\nu_4$	a'				3054 (20)	3004 (12)	
$\nu_6$	a'	3012 (37)	2988 (53)	<sup>d</sup>	2985 (35)	2922 (26)	<sup>d</sup>
$\nu_7$	a'	2884 (18)	2827 (15)	(HCl) <sup>d</sup>	2864 (42)	2789 (23)	(HCl) <sup>d</sup>
$\nu_8$	a'				1644 (57)	1623 (35)	1642.7 (100) <sup>c</sup>
$\nu_{10}$	a'				1538 (11)	1530 (7)	<sup>d</sup>
$\nu_{13}$	a'	1414 (14)	1398 (31)	1416.6 (19) <sup>c</sup>			
$\nu_{19}$	a'				1227 (14)	1223 (7)	<sup>d</sup>
$\nu_{23}$	a'				1034 (10)	1032 (9)	(A) <sup>d</sup>
$\nu_{25}$	a'	945 (10)	942 (12)	(D) <sup>d</sup>			
$\nu_{26}$	a'	920 (16)	920 (61)	918.3 (14)	870 (32)	871 (26)	876.4 (60)
$\nu_{27}$	a'						
$\nu_{35}$	a''				970 (11)	968 (8)	<sup>d</sup>
$\nu_{40}$	a''	772 (63)	773 (236)	769.2 (100)	748 (64)	752 (41)	(E) <sup>d</sup>
$\nu_{41}$	a''	715 (7)	729 (122)	714.6 (33)	724 (10)	717 (16)	(C) <sup>d</sup>
$\nu_{42}$	a''	699 (15)	698 (38)	(B) <sup>d</sup>			

<sup>a</sup> Harmonic vibrational wavenumbers scaled according to  $0.9122x + 174.3$  for wavenumbers  $>2000 \text{ cm}^{-1}$  and  $0.9823x + 0.6$  for wavenumbers  $<2000 \text{ cm}^{-1}$ . <sup>b</sup> IR intensities in  $\text{km mol}^{-1}$ . <sup>c</sup> IR intensities as per cent of that of the most intense line in each species. <sup>d</sup> Interference due to absorption of parent or products; in the latter case, the group number of the interfering line or HCl are indicated in parentheses.



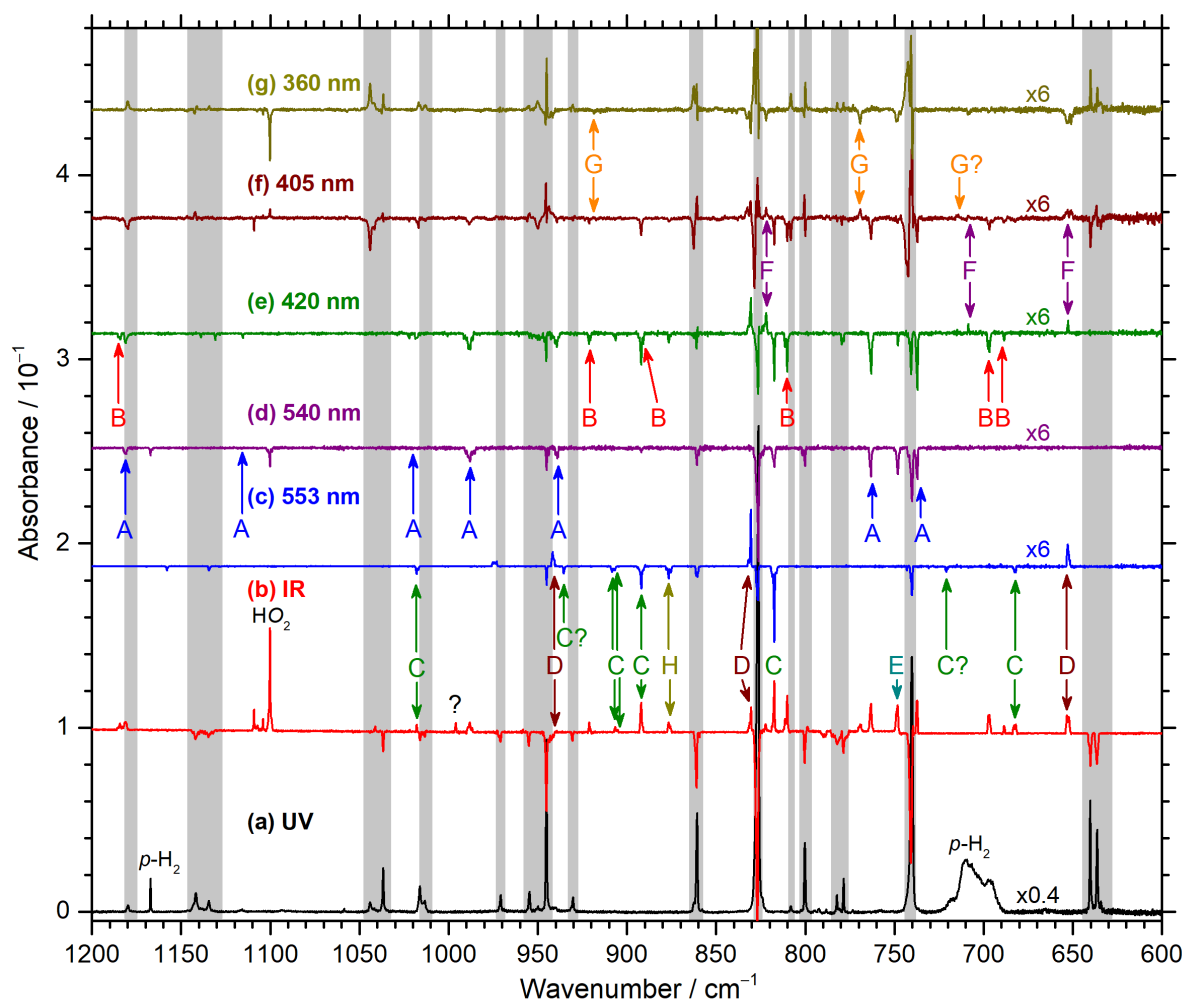
**Fig. 1** Numbering of carbon and nitrogen atoms in isoquinoline (*iso*-C<sub>9</sub>H<sub>7</sub>N) according to IUPAC nomenclature.

(single column)



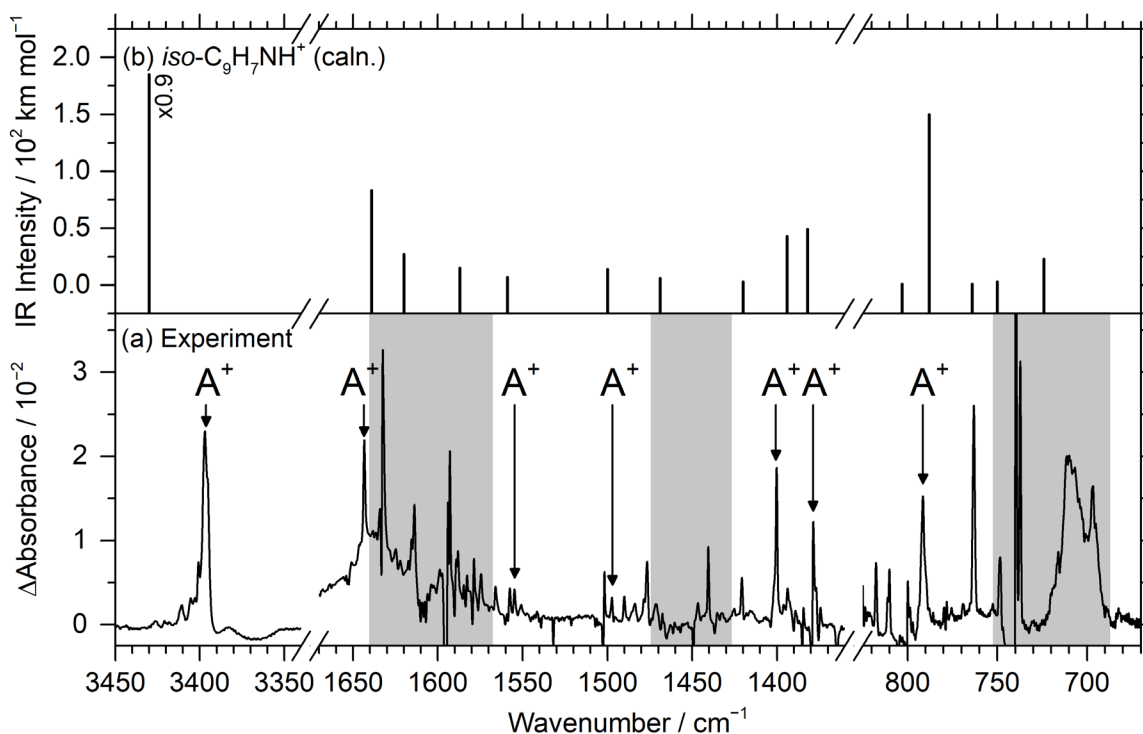
**Fig. 2** Infrared spectra of an electron-bombarded *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> matrix after each experimental step. (a) After deposition at 3.2 K for 7 h. The lines of *iso*-C<sub>9</sub>H<sub>7</sub>N were stripped using a spectrum of a *iso*-C<sub>9</sub>H<sub>7</sub>N/*p*-H<sub>2</sub> matrix without electron bombardment. (b) Difference spectrum after maintenance of the matrix in darkness for 10 h. (c) Difference spectrum after secondary irradiation at 365 nm for 25 min. (d) Difference spectrum after further irradiation at 405 nm for 3 h. The lines in groups A<sup>+</sup>, A, and B are indicated with labels and arrows. Spectral regions subjected to interference from intense absorption of *iso*-C<sub>9</sub>H<sub>7</sub>N are shaded gray. Baselines were shifted for clarity.

(double column)



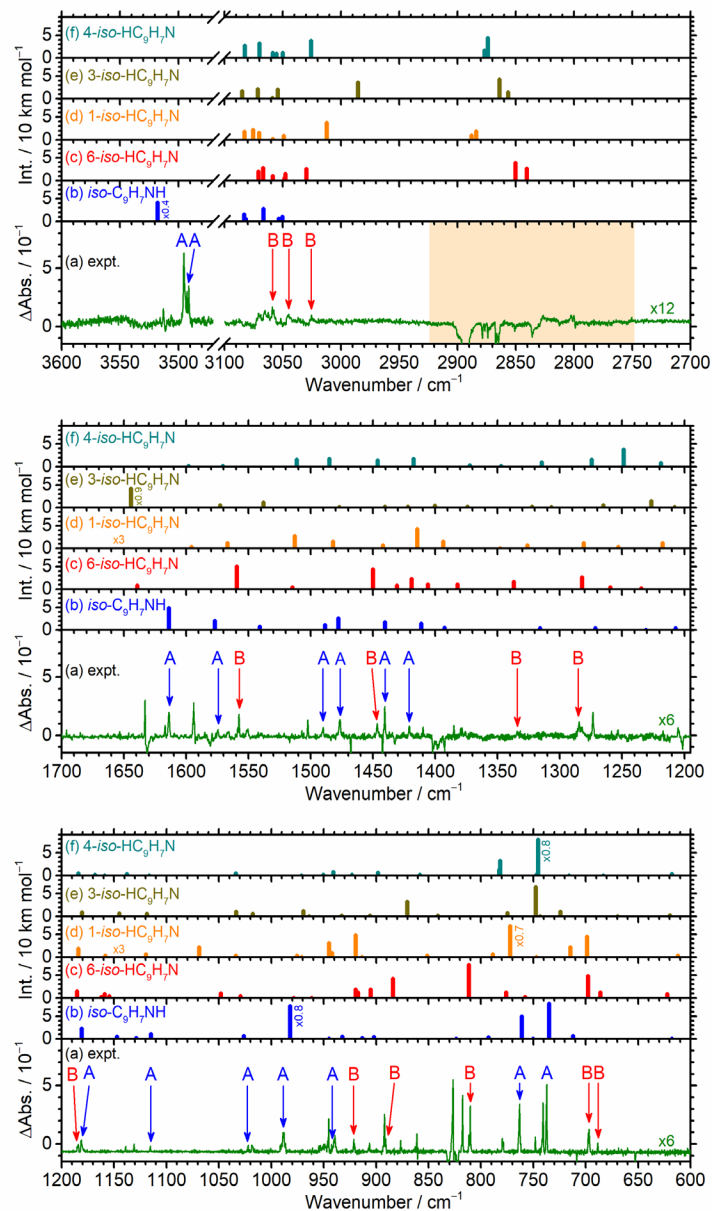
**Fig. 3** IR spectra of an *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix in region 1200–600 cm<sup>−1</sup> at various stages of experiments. (a) Spectrum of an *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix deposited at 3.2 K for 8 h and irradiated at 365 nm for 1 h. (b) Difference spectrum of the matrix after subsequent irradiation with IR light for 2 h. (c) Difference spectra of the matrix upon secondary irradiation at 553 nm for 30 min. (d) Difference spectra of the sample upon secondary irradiation at 540 nm for 30 min. (e) Difference spectra of the sample upon secondary irradiation at 420 nm for 30 min. (f) Difference spectra of the sample upon secondary irradiation at 405 nm for 30 min. (g) Difference spectra of the sample upon secondary irradiation at 360 nm for 30 min. Lines in groups A (*iso*-C<sub>9</sub>H<sub>7</sub>NH), B (6-*iso*-HC<sub>9</sub>H<sub>7</sub>N), C (7-*iso*-HC<sub>9</sub>H<sub>7</sub>N), D (8-*iso*-HC<sub>9</sub>H<sub>7</sub>N), E (4-*iso*-HC<sub>9</sub>H<sub>7</sub>N), F (5-*iso*-HC<sub>9</sub>H<sub>7</sub>N), G (1-*iso*-HC<sub>9</sub>H<sub>7</sub>N), and H (3-*iso*-HC<sub>9</sub>H<sub>7</sub>N) are indicated. Spectral regions subjected to interference from intense absorption of *iso*-C<sub>9</sub>H<sub>7</sub>N are shaded gray. Baselines were shifted for clarity.

(double column)



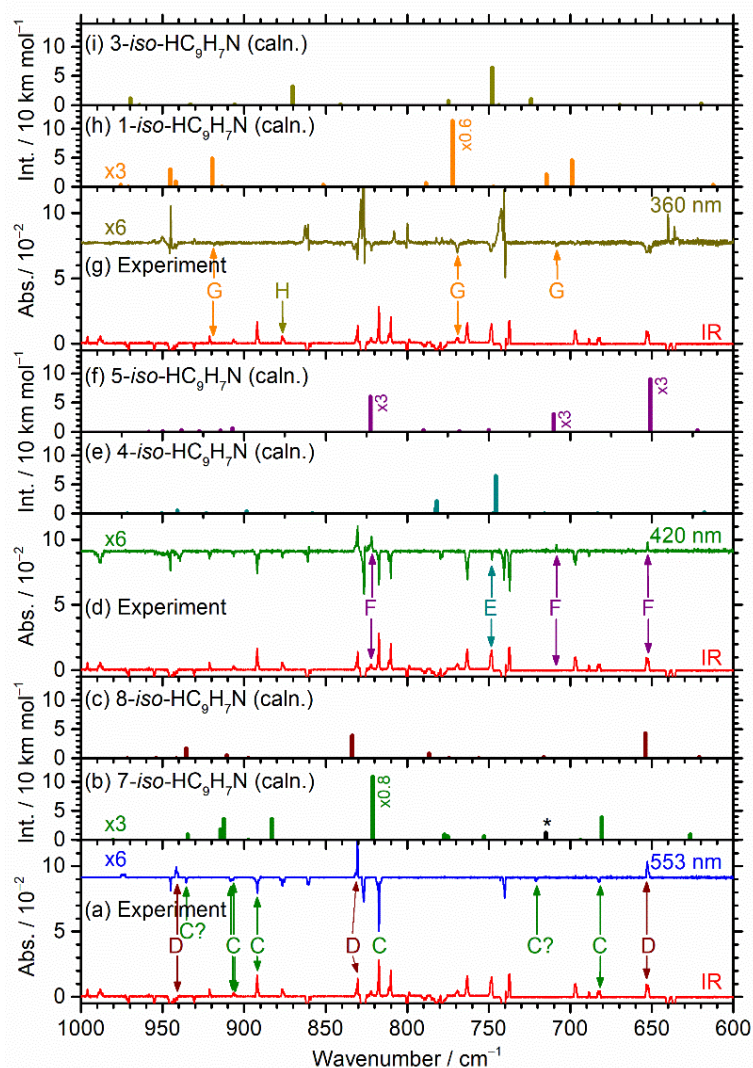
**Fig. 4** Comparison of observed lines of group  $A^+$  with predicted IR stick spectrum of  $iso\text{-C}_9\text{H}_7\text{NH}^+$ . (a) The spectrum after deposition; lines of  $iso\text{-C}_9\text{H}_7\text{N}$  are stripped. Lines in group  $A^+$  are indicated with labels and arrows. Gray rectangles indicate spectral regions suffering from severe interference from parent ( $iso\text{-C}_9\text{H}_7\text{N}$ ). (b) Simulated IR stick spectrum of  $iso\text{-C}_9\text{H}_7\text{NH}^+$  according to scaled harmonic wavenumbers and IR intensities predicted with the B3LYP/6-311++G(d,p) method.

(double column)



**Fig. 5** Comparison of observed lines of groups A and B in  $iso-C_9H_7N/Cl_2/p-H_2$  experiments with IR stick spectra of  $iso-C_9H_8N$  predicted with theory. (a) Difference spectrum after irradiation at 420 nm (inverted Fig. 3e); lines belonging to groups A ( $iso-C_9H_7NH$ ) and B (6- $iso-HC_9H_7N$ ) are indicated with blue and red labels and arrows. IR stick spectra of  $iso-C_9H_7NH$  (b), 6- $iso-HC_9H_7N$  (c), 1- $iso-HC_9H_7N$  (d), 3- $iso-HC_9H_7N$  (e), and 4- $iso-HC_9H_7N$  (f) simulated according to scaled harmonic vibrational wavenumbers and harmonic IR intensities predicted with the B3LYP/6-311++G(d,p) method are presented. Spectral regions that suffer severe interference from absorption of HCl and HCl complexed with water are shaded with light orange.

(double column)

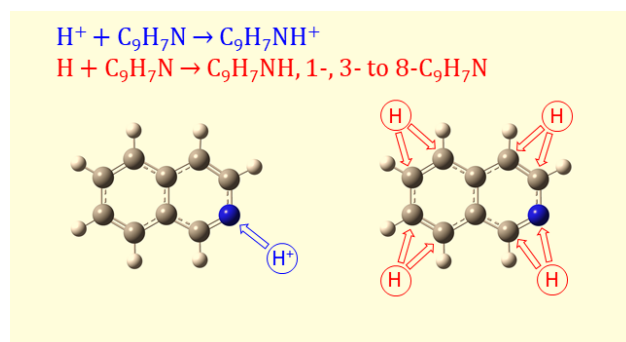


**Fig. 6** Comparison of observed lines of groups C to H produced from IR irradiation of an *iso*-C<sub>9</sub>H<sub>7</sub>N/Cl<sub>2</sub>/*p*-H<sub>2</sub> matrix with theoretically predicted stick spectra of representative isomers of hydrogenated isoquinoline (*iso*-HC<sub>9</sub>H<sub>7</sub>N). (a) Difference spectra after IR irradiation (lower trace, from Fig. 3c) and after secondary photolysis at 553 nm (upper trace, from Fig. 3d); lines of groups C and D are indicated. IR stick spectrum of 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N (b) and 8-*iso*-HC<sub>9</sub>H<sub>7</sub>N (c). (d) Same as in (a) except that the upper trace is from secondary photolysis at 420 nm, Fig. 3e; lines of groups E and F are indicated. IR stick spectrum of 4-*iso*-HC<sub>9</sub>H<sub>7</sub>N (e) and 5-*iso*-HC<sub>9</sub>H<sub>7</sub>N (f). (g) Same as in (a) except that the upper trace is from secondary photolysis at 360 nm, Fig. 3g, but is inverted; lines of groups G and H are indicated. IR stick spectrum of 1-*iso*-HC<sub>9</sub>H<sub>7</sub>N (e) and 3-*iso*-HC<sub>9</sub>H<sub>7</sub>N (f). All IR stick spectra are shown according to scaled harmonic vibrational wavenumbers and IR intensities predicted with the B3LYP/6-311++G(d,p) method. The mark \* indicates anharmonic vibrational wavenumber of 7-*iso*-HC<sub>9</sub>H<sub>7</sub>N.

(single column)



## Graphic abstract



Protonation of isoquinoline was observed at the N-atom site, whereas hydrogenation of isoquinoline was observed at the N-atom site and all C-atom sites except the sharing C atoms on the fused ring.

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