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Improved Electron Transport Properties of Zn-rich In-Ga-Zn-O Thin Film Transistors

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Abstract: Amorphous transparent oxide semiconductor InGaZnO₄-based thin film transistors (TFTs) have been practically used as the backplane of flat panel displays. For future higher-definition displays, alternative active materials with a higher field effect mobility (μ_{FE}) are necessary. Although there are a few reports on InGaO₃(ZnO)_m with Zn-rich composition (IGZO_m)-based TFTs, their electron transport properties have not been clarified. Here, we show that a Zn-rich composition enhances the electron transport properties of IGZO_m-TFTs. The best TFT performance was obtained for $m = 7$ ($\mu_{FE} \sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, subthreshold swing $\sim 0.1 \text{ V decade}^{-1}$, and a negligibly small bias stress shift). The carrier effective mass (m^*) of IGZO_m films was found to be $0.16 m_0$, independent of the m -value. We found that μ_{FE} of IGZO_m-TFT increased with the m -value for $m \leq 7$, whereas it decreased for $m > 7$ due to the crystallization. The thermopower modulation analyses revealed that the effective channel thickness increased with m ($m \leq 7$) which resulted in a longer carrier relaxation time. The present results provide an improving strategy toward new material design for next-generation TFTs with higher μ_{FE} values.

Introduction

Amorphous transparent oxide semiconductor InGaZnO₄ (a-IGZO)-based thin film transistors (TFTs) are widely applied as the backplane of flat panel displays.^{1,2} In 2004, Nomura and Hosono *et al.*³ demonstrated that a-IGZO-TFTs exhibit excellent transistor characteristics with a rather high field effect mobility ($\mu_{FE} \sim 7 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$). The reported value is an order of magnitude higher than that of a-Si ($\sim 0.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$). Since then, several display companies have vigorously developed and commercialized large-sized liquid crystal displays (LCDs) and organic light emitting diode displays (OLEDs) using

a-IGZO-TFTs as the backplane. Although a-IGZO-TFTs are widely applied nowadays, μ_{FE} must be increased for higher-definition pixelated and wearable displays. Therefore, alternative active materials with a higher μ_{FE} are still demanded.

In this study, we focus on $\text{InGaO}_3(\text{ZnO})_m$ with higher m -values (IGZO_m hereafter) as a potential alternative active material. IGZO_m crystallizes with a layered structure composed of an InO_2^- layer and a $\text{GaO}(\text{ZnO})_m^+$ block, which are alternately stacked along the c -axis.⁴ Since there is a thermodynamically stable crystalline phase, ceramic targets of IGZO_m can easily be fabricated, good for practical applications that require very large-sized crack-free ceramic targets. Nomura *et al.*⁵ fabricated TFTs in $\text{InGaO}_3(\text{ZnO})_5$ single crystal films using reactive solid-phase epitaxy.⁶ Their TFTs exhibited high mobility $\mu_{FE} \sim 80 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, an order magnitude higher than that of a-IGZO-TFTs ($\mu_{FE} \sim 7 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$). In contrast, Ochi *et al.*⁷ reported that amorphous IGZO_m -TFTs showed a rather low $\mu_{FE} \sim 0.18 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for $m \sim 14$. Thus, the electron transport properties of IGZO_m -TFTs with higher m -values have yet to be clarified.

In 2001, Orita *et al.*⁸ fabricated IGZO_m films with a stable amorphous phase for $m < 4$. The Zn 4s orbitals predominantly contribute to the electron conduction of IGZO_m with a large m -value. In 2006, Nomura *et al.*⁹ precisely investigated the residual carrier concentration (n), which originated from oxygen deficiencies, and Hall mobility (μ_{Hall}) of the IGZO_m system. They found that both n and μ_{Hall} increase as the m -value increases. Since the fabrication of TFTs with normally-off characteristics becomes difficult when n is high, reports on transistor characteristics of IGZO_m -TFTs with a large m -value are rare.^{5, 7}

Here, we show that a Zn-rich composition enhances the electron transport properties of IGZO_m-TFTs. First, we clarified that the carrier effective mass (m^*) of IGZO_m is independent of the m -value ($m^* = 0.16 m_0$, m_0 : electron mass). Then, we found that μ_{FE} of IGZO_m-TFT increased with the m -value for $m \leq 7$. This clearly reveals that the carrier relaxation time (τ) of IGZO_m increases with m ($m \leq 7$). For $m > 7$, μ_{FE} decreases due to crystallization; The carrier electrons scatter at the grain boundaries. The best TFT performance occurs for $m = 7$ ($\mu_{FE} \sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, subthreshold swing $\sim 0.1 \text{ V decade}^{-1}$, and a negligibly small bias stress shift). We attempted thermopower (S) modulation analyses to clarify the electron transport properties of IGZO_m and found that the effective channel thickness increased with m ($m \leq 7$). The present results provide an improving strategy towards new material design for next-generation TFTs with higher μ_{FE} values.

Experimental Section

Fabrication of IGZO_m Films. IGZO_m films ($m = 1, 2, 3, 4, 5, 7, 9, 10,$ and 30) were deposited on alkali-free glass substrates (EAGLE XG[®], Corning[®]) by a pulsed laser deposition (PLD) technique at room temperature. A KrF excimer laser (COMPex Pro 102, Coherent) was used to ablate the ceramic targets of IGZO_m. The laser fluence, repetition rate, and the number of pulses were $\sim 2.3 \text{ J cm}^{-2} \text{ pulse}^{-1}$, 10 Hz, and 1800 pulses, respectively. A constant oxygen pressure (1–5 Pa) was used during film deposition to control n . Afterwards, the films were annealed at several temperatures (50–150 °C) in air for 5 min to further control n .

Structure Analyses, Optical and Electron Transport Property Measurements of the IGZO_m Films. The film thickness, relative density, and surface roughness of the resultant IGZO_m films were measured by the X-ray reflectivity (XRR, Cu K α ₁, ATX-G, Rigaku). We also measured the X-ray diffraction (XRD) patterns of the resultant films. The incident angle of the X-ray was fixed at 0.5° while 2 θ scanning was performed. The optical transmission and reflection spectra of the IGZO_m films were recorded using a UV-Vis-NIR spectrophotometer (SolidSpec-3700, Shimadzu). We drew the absorption spectra (Tauc plot) and extracted the optical bandgap of the IGZO_m films. The thermopower (S) was measured by a conventional steady-state method at room temperature. The Hall mobility (μ_{Hall}) and the volume carrier concentration ($n_{3\text{D}}$) were measured by the dc four-probe method with a van der Pauw electrode configuration at room temperature, and ± 7600 G was applied for the Hall effect measurements.

Fabrication of IGZO_m-TFTs. We fabricated IGZO_m-based TFTs (bottom-gate-top-contact type) on alkaline-free glass substrates (EAGLE XG[®], Corning[®]). First, a 106-nm-thick AlO_x film (dielectric permittivity $\epsilon_r = 8$), which served as the gate insulator, was deposited by atomic layer deposition (ALD) onto an ITO-coated glass substrate. The capacitance (C_i) of the AlO_x films was 67 nF cm⁻². Then ~20-nm-thick IGZO_m film was deposited by PLD under an oxygen pressure of 5 Pa at room temperature through a stencil mask. The multilayer film was annealed at 350 °C in air for 5 min. Finally, 20-nm-thick Ti source/drain electrodes were evaporated through the stencil mask at room temperature.

Transistor Characterization and Electric Field Thermopower Modulation. Standard

transistor characteristics of the IGZO_m-TFTs were measured by a semiconductor device analyzer (B1500A, Agilent) at room temperature. We calculated the field effect mobility (μ_{FE}) by using following equation (1),

$$\mu_{FE} = \frac{g_m}{C_i \frac{W}{L} V_D} \quad (1)$$

where g_m is the transconductance, C_i is the capacitance per unit area of the gate insulator (67 nF cm⁻² for AlO_x), W is the channel width, L is the channel length, and V_D is the drain voltage. We obtained μ_{FE} using the maximum value of g_m . Thermopower modulation analyses of the IGZO_m-TFTs were performed at room temperature in air. Details of the thermopower modulation analyses are published elsewhere (SrTiO₃-TFTs,¹⁰ InGaZnO₄-TFTs,¹¹ BaSnO₃-TFTs,¹² SnO₂-TFTs,¹³ InSnZnO_x-TFTs,¹⁴ and GaN-TFTs¹⁵).

Results and Discussion

First, we fabricated IGZO_m films and analyzed their structure and optical properties. **Figure 1a** shows the XRD patterns of the resultant IGZO_m films. For $m \leq 7$, a halo of IGZO_m appears at $2\theta \sim 34^\circ$, whereas sharp diffraction peaks appear for $m > 7$. These results indicate that films with $m \leq 7$ are amorphous whereas those with $m > 7$ are randomly-oriented polycrystalline. These observations are consistent with the report by Orita *et al.* in which IGZO_m crystallized at room temperature for $m > 5$.⁸ The optical bandgap (**Figure 1b**) gradually decreases as the m -value increases (3.25 eV for $m = 1$ and 3.15 eV for $m = 5$). This tendency is similar to that of polycrystalline IGZO_m.¹⁶ These observations demonstrate that IGZO_m ($m = 1, 2, 3, 4, 5, 7, 9, \text{ and } 30$) films are successfully fabricated.

Then, we measured the electron transport properties of the IGZO_m films. **Figure 2a** plots the thermopower (S) as a function of the volume carrier concentration (n_{3D}) for the resultant IGZO_m films at room temperature in air. The S -values are always negative, confirming that the films are n-type semiconductors. The absolute value of S decreases with n_{3D} in all cases. The density of states effective mass (m^*) was extracted using the relationship between n_{3D} and S .¹⁷ Although the data is a bit scattered due to the difficulty of accurate n_{3D} measurement when the mobility is small, all the S - n_{3D} relationships follow a simulated line with a carrier effective mass (m^*) of $0.16 m_0$, similar to the reported m^* of a-IGZO ($m_d^*/2 = 0.17 m_0$)¹⁸⁻²⁰, indicating that m^* does not depend on the m -values (Note we extracted averaged value since the data was scattered, though extraction of m^* requires accurate carrier concentration values). This indicates that the Zn 4s orbital dominated conduction band bottom exhibits a similar band dispersion to that of the In 5s orbital dominated conduction band bottom. In contrast, μ_{Hall} does not show a clear trend as a function of n_{3D} (**Figure 2b**). Since all IGZO_m films have similar m^* values, the difference in the observed μ_{Hall} is due to the variation in the carrier relaxation time ($\tau = e \cdot (\mu \cdot m^*)^{-1}$). Since the IGZO_m films were deposited with different oxygen pressure as well as with different annealing temperatures in order to tune the carrier concentration, the film quality of the IGZO_m films were different each other. The difference in the film quality largely affects carrier relaxation time. Therefore, the Hall effect measurement was a bit difficult when the film quality was poor.

Next, we fabricated IGZO_m-TFTs on ITO-AlO_x glass-insulator substrates and measured their transistor characteristics. Compared to the other high- k dielectrics, AlO_x has a

relatively larger band offset with IGZO_m.^{21, 22} Therefore, we expected that the gate leakage current and bias instabilities would be suppressed by the use of AlO_x gate insulator for IGZO_m-TFTs. **Figure 3a** shows the typical transfer characteristics of IGZO_m-TFTs at room temperature. The channel length (L), width (W), and applied drain voltage (V_d) are 100 μm , 400 μm , and +5 V in all cases, respectively. The on current gradually increases with m for $m \leq 7$. However, the off current increases and the on current gradually decreases with m for $m > 7$. **Figure 3b** plots μ_{FE} of the IGZO_m-TFTs. For $m \leq 7$, μ_{FE} gradually increases with m , but it decreases with m for $m > 7$. The highest μ_{FE} is $\sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ for $m = 7$. It is noteworthy that the low μ_{FE} values of $m = 10$ and 30 should originate from the grain boundaries since IGZO_m crystallizes for $m > 7$. This is also reflected by the subthreshold swing ($S.S.$) value. The $S.S.$ value remains almost constant ($\sim 0.1 \text{ V decade}^{-1}$) for $m = 1-7$ but it increases from $0.1-0.3 \text{ V decade}^{-1}$ for higher values of m . This would be attributed to the increase in oxygen deficiency-related defects at the interface of the IGZO_m channel layer and the gate dielectric.

To clarify the origin of why $m = 7$ TFT exhibited the highest μ_{FE} , various gate voltages were applied to measure the thermopower of the IGZO_m-TFTs (**Figure 4a**, electric field thermopower modulation analyses). A constant V_g was initially applied to the TFTs to accumulate carrier electrons. Then, S of the channel was measured during the V_g application. **Figure 4b** plots the electric field modulated S as a function of n_{2D} . For $m = 1$ and 2 , several S -values around the lower n_{2D} are not on a straight line, which may be due to the non-parabolic shaped tail states²³ just below the original conduction band bottom at a lower m -value.¹¹ It shows a decreasing tendency when n_{2D} is high. On the other hand, the absolute values of S for $m = 3, 5,$ and 7 showed a simple decreasing

tendency. It should be noted that $|S|$ increases as the m -value increases. Then, we compared the n_{2D} (**Figure 4b**) with n_{3D} (**Figure 2a**) for a given $|S|$ (**Figure 4b**) to estimate the effective channel thickness (n_{2D}/n_{3D}). Interestingly, the n_{2D}/n_{3D} increases from ~ 5 nm to ~ 28 nm with m -value, similar to the increasing tendency of μ_{FE} . Thus, an increase of τ with m -value originates from the increase of the effective channel thickness. An increased thickness reduces electron scattering at the interface between dielectric and semiconducting channel. It should be noted that the field effect mobility is obtained with the Fermi energy level above the conduction band edge, at carrier scattering at defect centers becomes negligibly small.

According to the Poisson equation, the carrier accumulation layer becomes thick when the dielectric permittivity of the semiconductor is large or the residual carrier concentration is high. The dielectric constant of amorphous InGaZnO₄ is 16^{24} and that of ZnO is $\sim 9^{25}$. Therefore, it is impossible to explain the origin of why the effective thickness of IGZO_{*m*} increases with m using the dielectric constant scenario. In short, the thickness increase with m is most likely due to an increase of the residual carrier concentration with m . Since the bandgap of IGZO_{*m*} decreased with m , the increasing tendency of the residual carrier concentration with m is reasonable.

Finally, we performed the positive bias stress (PBS) and negative bias stress (NBS) measurements for TFTs with $m \leq 7$. For NBS, TFTs were evaluated in ambient air at $V_g = -20$ V for a total stress time of 10000 s. Except for $m = 1$, the transfer characteristics of the TFTs do not change significantly after NBS (**Figures S1b–S4b**). In contrast, for PBS with $V_g = +20$ V, the stress conditions affect the positive ΔV_{th} (**Figures S1a–S4a**).

This would be attributed to charge trapping in gate insulators and the creation of electron traps in the channel. **Figure S5** shows the change in the normalized V_{th} shift of the IGZO_m-TFTs under both PBS and NBS as a function of the stress time. Large V_{th} shift occurs after PBS for $m = 1$ (**Figure S1a**), indicating that n_{2D} should be lower than that assumed as $C_1 \cdot (V_g - V_{th})$. In contrast, the V_{th} shift for $m = 3, 5,$ and 7 (**Figures S2a, S3a, and S4a**) is negligible, indicating that the electric field thermopower measurements provide an accurate and effective channel thickness. These results reveal that $m = 7$ IGZO_m-TFT exhibits the best TFT characteristics ($\mu_{FE} \sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, $S.S. \sim 0.1 \text{ V decade}^{-1}$ and a negligibly small bias stress shift) because the wider effective channel thickness provides a longer τ . The present results provide an improving strategy towards new material design for next-generation TFTs with higher μ_{FE} values.

Conclusion

We studied the electron transport properties of InGaO₃(ZnO)_m ($m = 1-30$) with higher m -values (IGZO_m) thin films and TFTs. The best TFT performance was obtained for $m = 7$ ($\mu_{FE} \sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, subthreshold swing $\sim 0.1 \text{ V decade}^{-1}$, and a negligibly small bias stress shift). First, we clarified that the m^* of IGZO_m films is $0.16 m_0$, independent of the m -value. Then, we found that μ_{FE} of IGZO_m-TFT increased with the m -value for $m \leq 7$, whereas it decreased for $m > 7$ due to the crystallization. The thermopower modulation analyses revealed that the effective channel thickness increased with m ($m \leq 7$) that resulted in the longer τ . The present results provide an improving strategy towards new material design for next-generation TFTs with higher μ_{FE} values.

ASSOCIATED CONTENT

The Supporting Information is available free of charge at

<https://pubs.acs.org/doi/10.1021/acs.jpcc.xxxxxxx>.

Bias stress characteristics of IGZO₁-TFT, Bias stress characteristics of IGZO₃-TFT, Bias stress characteristics of IGZO₅-TFT, Bias stress characteristics of IGZO₇-TFT,

Normalized V_{th} shift (ΔV) as a function of stress time.

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Author Contributions

P.G., H.Y., T.F., and H.O. fabricated the samples and measured the transistor characteristics and thermopower. Y.Z. analyzed the thermopower. Y. Matsuo fabricated the AlO_x films. Y. Magari analyzed the transistor characteristics. H.O. planned and supervised the project. All authors discussed the results and commented on the manuscript.

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Notes

The authors declare no competing financial interest.

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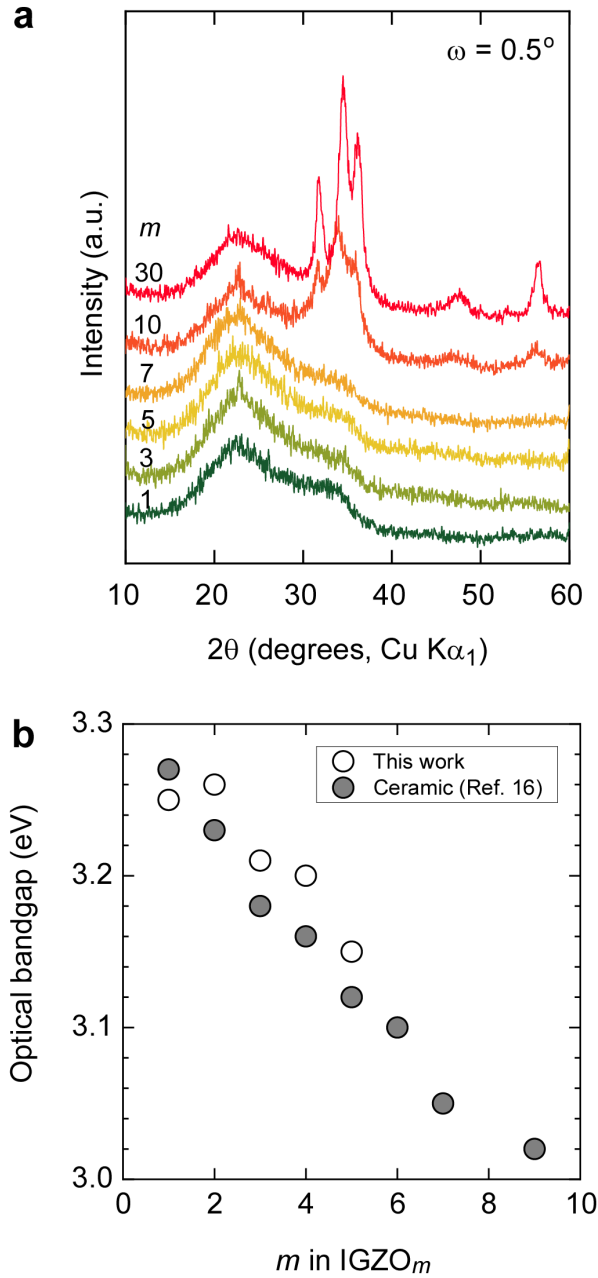


Figure 1. Structural and optical characterization of the resultant IGZO_m films. (a) XRD patterns with a fixed X-ray incident angle of 0.5°. For $m \leq 7$, a halo of IGZO_m is seen at $2\theta \sim 34^\circ$, whereas sharp diffraction peaks appear for $m > 7$. For $m > 7$, crystallization occurs. (b) Optical bandgap. As m increases, the optical bandgap of IGZO_m decreases, which is a similar overall tendency as that of IGZO_m ceramics.¹⁶

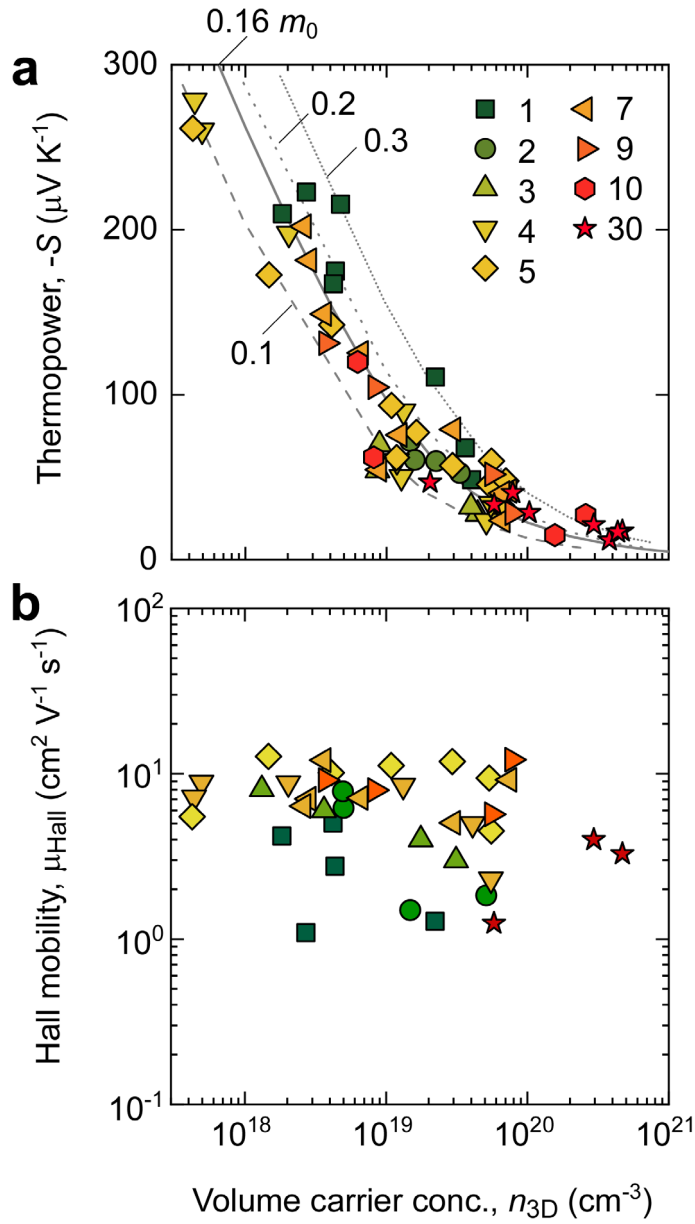


Figure 2. Electron transport properties of the a-IGZO_m films at room temperature. Changes in (a) thermopower (S) and (b) Hall mobility (μ_{Hall}) with the volume carrier concentration (n_{3D}). (a) shows a clear trend, but (b) does not.

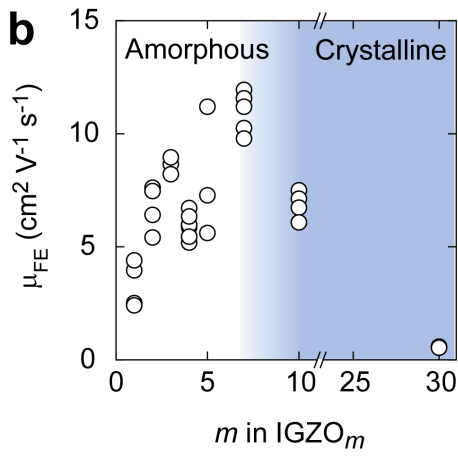
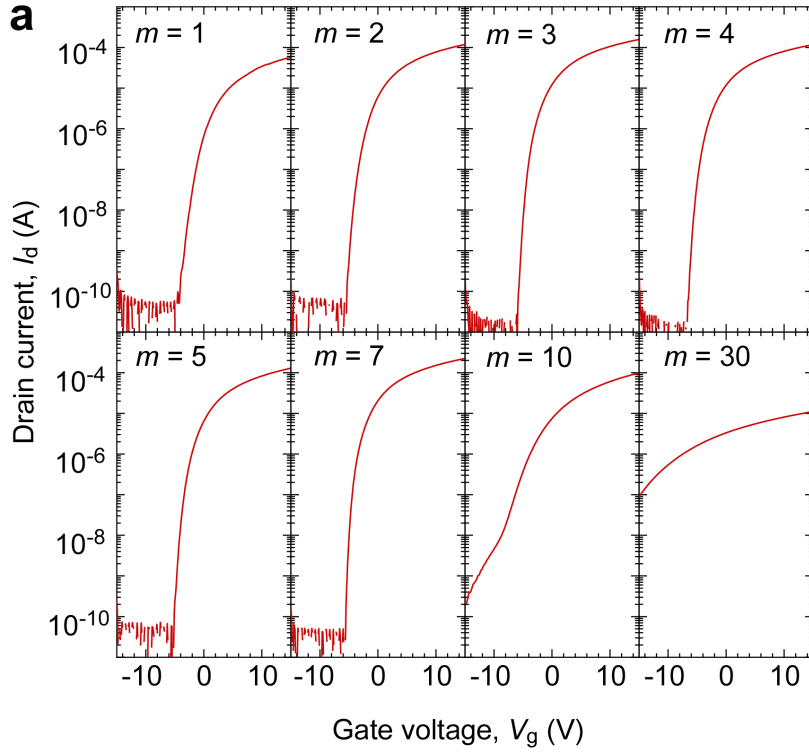


Figure 3. Typical transistor characteristics of IGZO_m-TFTs at room temperature.

(a) Transfer characteristics ($L = 100 \mu\text{m}$, $W = 400 \mu\text{m}$, $V_d = +5 \text{ V}$). On current gradually increases with m for $m \leq 7$. For $m > 7$, the off current increases and the on current gradually decreases with m . (b) Field effect mobility of the IGZO_m-TFTs. For $m \leq 7$, μ_{FE} gradually increases with m , but it drops with m for $m > 7$. Highest μ_{FE} is $\sim 12 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, which is obtained for $m = 7$.

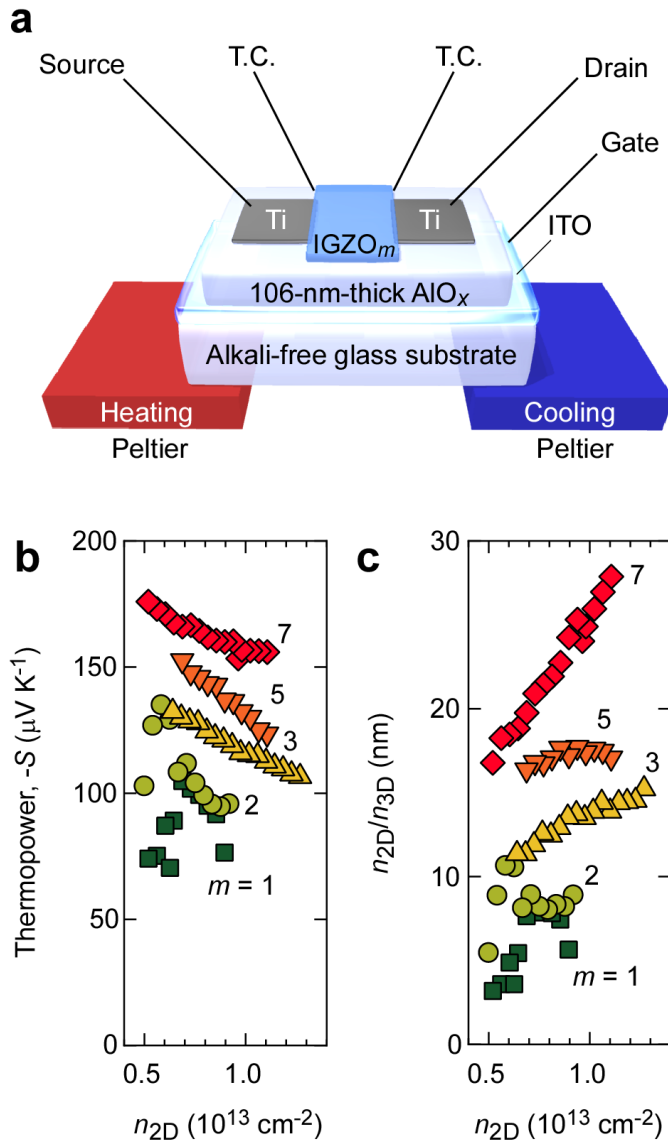


Figure 4. Electric field thermopower modulation analyses of the IGZO_m-TFTs. (a) Schematic of the experimental setup. IGZO_m-TFT is placed between two Peltier devices. Upon applying a gate voltage, the thermopower is measured by applying a temperature difference between the source and drain electrodes. (b) Thermopower and (c) effective thickness of IGZO_m-TFTs as a function of n_{2D} . Note that the effective thickness (n_{2D}/n_{3D}) varies from ~ 2 nm to ~ 28 nm depending on the m -value.