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Production cross sections of ⁴⁷Sc via alpha-particle-induced reactions on natural calcium up to 29 MeV

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Abstract

Production cross sections of a medical radioisotope ⁴⁷Sc via alpha-particle-induced reactions on natural calcium were measured up to 29 MeV. The stacked-foil activation technique and gamma-ray spectrometry were adopted for the experiment. Cross sections of simultaneously formed radionuclides ^{46,44m,44g,43}Sc and ⁴⁷Ca were also determined. Physical yields of the alpha-particle-induced reactions on ^{nat}Ca and ⁴⁴Ca were deduced from the measured cross sections.

Keyword

Scandium-47; Calcium target; Alpha-particle irradiation; Cross section; Excitation function

1. Introduction

Scandium radionuclides are attractive in the field of medical applications due to their properties and availability, such as ^{43,44}Sc for PET [1], ⁴⁷Sc for therapy [2] and their combination for theranostics [3]. These scandium radionuclides can be produced via charged-particle-induced reactions on calcium, scandium, and titanium [3]. The possible reactions should be investigated in detail in order to optimize the radionuclide production for practical use.

In this paper, we focus on production of ⁴⁷Sc ($T_{1/2} = 3.3492$ d) via alpha-particle-induced reactions on natural calcium, of which isotopic composition is ⁴⁰Ca: 96.941%, ⁴²Ca: 0.647%, ⁴³Ca: 0.135%, ⁴⁴Ca: 2.086%, ⁴⁶Ca: 0.004%, and ⁴⁸Ca: 0.187% [4]. Among the calcium isotopes, the alpha-particle-induced reactions on ⁴⁴Ca, ⁴⁶Ca and ⁴⁸Ca can provide ⁴⁷Sc, and its parent ⁴⁷Ca ($T_{1/2} = 4.536$ d) and grandparent ⁴⁷K ($T_{1/2} = 17.50$ s). Taking into account the relatively small abundance of ⁴⁶Ca and ⁴⁸Ca, the most probable reaction for

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production of ⁴⁷Sc in a larger quantity is the ⁴⁴Ca(α ,p)⁴⁷Sc reaction. We found only one experimental study for the ⁴⁷Sc production [5] in a literature survey and the related data in the EXFOR library [6]. Therefore, we performed an experiment to obtain cross sections of the ^{nat}Ca(α ,x)⁴⁷Sc reaction up to 29 MeV. The cross sections on ^{nat}Ca can be normalized to those on enriched ⁴⁴Ca under assumption of negligible contributions from ⁴⁶Ca and ⁴⁸Ca. In addition, production cross sections of ^{46,44m,44g,43}Sc and ⁴⁷Ca were measured to estimate the amount of the co-produced radioactive impurities. The physical yields of the alpha-particle-induced reactions on ^{nat}Ca and ⁴⁴Ca were deduced from the measured cross sections.

2. Experimental details and data analysis

The experiment was conducted with a 29-MeV alpha-particle beam at the AVF cyclotron in RIKEN. The stacked-foil activation technique and high-resolution off-line gamma-ray spectrometry were adopted in the measurement.

Calcium-fluoride (CaF₂) was used as a target material. The stacked target consisted of thin CaF₂ layers, two types of Al foils (Al_H, purity: 99.999%, thickness: 10 µm, size: 100 × 100 mm, Goodfellow Cambridge Ltd., UK and Al_L, purity: >99%, thickness: 5 µm, size: 100 × 100 mm, The Nilaco Corporation, Japan) and ^{nat}Ti foil (purity: 99.5%, thickness: 5 μ m, size: 50 × 100 mm, The Nilaco Corporation, Japan). By measuring the lateral size and mass of the purchased Al and Ti foils, their average thicknesses were determined for the data analysis. The measured thicknesses of Al_H, Al_L and ^{nat}Ti foils were 2.57, 1.50 and 2.30 mg/cm², respectively. The ^{nat}Ti foil was used for monitoring the beam parameters by using the ^{nat}Ti(α, x)⁵¹Cr monitor reaction and to check the energy loss of the alpha-particles through the stacked target. The AlL foil was used to catch the recoiled reaction products from the ^{nat}Ti foil in the stack. The Al_H foil was used as a backing of the CaF_2 layer. The CaF_2 layer was deposited on the Al_H foil using a vacuum evaporation method. The thickness of the CaF₂ layer was determined to be 0.135 mg/cm² from the deposited area (90 \times 90 mm) and weight of CaF₂. All foils were cut into 10×10 mm pieces to construct the stacked target and to fit a target holder served as a Faraday cup. The CaF₂ layers were sandwiched with the Al_H foils as Al_H(CaF₂)-(CaF₂)Al_H to collect the recoiled reaction products from the thin CaF₂ layer and to increase the target thickness. Twelve sets of the Al_H(CaF₂)-(CaF₂)Al_H foils and seven sets of the ^{nat}Ti-Al_L foils were stacked together in the target holder. The configuration of the stacked target is listed in Table 1.

The stacked target was irradiated for 30 min with an alpha-particle beam collimated to 3 mm in diameter. The average beam intensity measured by charge integration on the Faraday cup was 175 nA. The primary beam energy measured using the time-of-flight (TOF) method [7] was 29.0 \pm 0.2 MeV. Energy degradation in the stacked target was calculated using stopping powers obtained from the SRIM code [8]. The calculation of the energy degradation was checked by using the ^{nat}Ti(α ,x)⁵¹Cr monitor reaction.

The high-resolution gamma-ray spectrometry of the activated foils was done without chemical separation. A high-purity germanium detector (ORTEC GEM30P4-70) and a dedicated software (SEIKO EG&G Gamma Studio) were used for the gamma-ray measurement and spectrum analysis. The efficiency of the detector was calibrated with a multiple gamma-ray emitting point-like source (Eckert & Ziegler Isotope Products). The point-like source consisted of eight radionuclides, ^{57,60}Co, ⁸⁸Y, ¹⁰⁹Cd, ¹¹³Sn ¹³⁷Cs, ¹³⁹Ce and ²⁴¹Am to cover the energy range between 60 and 1836 keV. Gamma rays emitted from each activated foil were measured using the energy- and geometry-dependent, and efficiency-calibrated detector. Each set of the Al_H(CaF₂)-(CaF₂)Al_H foils was measured 5 times with cooling times from 3.2 h to 77.0 d. The associated dead times were kept below 2.1%. The experimental parameters are summarized in Table 2. Nuclear data required to determine cross sections collected in Table 3 were retrieved from the online databases, NuDat 2.8 [9], LiveChart [10], Lund/LBNL Nuclear Data Search [11], and QCalc [12].

Cross sections σ were deduced using the following activation formula

$$\sigma = \frac{N_{\gamma}\lambda}{\varepsilon_d \varepsilon_\gamma \varepsilon_t N_t N_b (1 - e^{-\lambda t_b}) e^{-\lambda t_c} (1 - e^{-\lambda t_m})}$$
(1)

where N_{γ} is the measured net count of the peak area, λ is the decay constant (s⁻¹), ε_d is the detector efficiency, ε_{γ} is the gamma-ray intensity, ε_t is the dead time correction, N_t is the surface density of target atoms (cm⁻²), N_b is the number of projectiles per unit time (s⁻¹), t_b is the irradiation time (s), t_c is the cooling time (s), and t_m is the acquisition time (s).

Cross sections of the ^{nat}Ti(α ,x)⁵¹Cr monitor reaction were derived to assess the beam parameters and the target thicknesses. The gamma line at 320.08 keV ($I_{\gamma} = 9.91\%$) emitted from each pair of the ^{nat}Ti-Al_L foils was measured. Recoiled products from the ^{nat}Ti foils were caught by the next Al_L foils. The cross sections derived using the spectra with dead times less than 0.6% were compared with the IAEA recommended values [13]. According to the comparison, the beam intensity was corrected by +5.6% to have agreement with the amplitude of the measured and recommended values. The corrected beam intensity 185 nA was accepted and used in the data analysis. The peak position was found to be slightly shifted to the lower energy side. As the beam energy was measured by the TOF method the primary beam energy was accepted and the energy loss of the alpha particle in the stack was adjusted by changing the thicknesses of the Al_H and Al_L foils in the stack by -1% within the uncertainty of the given foil thicknesses. The measured thicknesses of the ^{nat}Ti foil and the CaF₂ layer were adopted without any correction. The cross sections using the corrected parameters agree with the recommended values as shown in Fig. 1. The adopted parameters for deduction of cross sections are also listed in Table 2.

No.	Foil/(Layer)	No.	Foil/(Layer)	No.	Foil/(Layer)	No.	Foil/(Layer)
1	Al_H	9	Ti	19	Al_H	29	Ti
	(CaF ₂)	10	$\operatorname{Al}_{\operatorname{L}}$		(CaF ₂)	30	Al_L
2	(CaF_2)	11	Al_{H}	20	(CaF ₂)	31	Al_H
	Al_H	11	(CaF ₂)		Al_{H}		(CaF ₂)
3	Al_H	10	(CaF_2)	21	Ti	20	(CaF_2)
	(CaF ₂)	12	$Al_{\rm H}$	22	Al_L	32	$Al_{\rm H}$
4	(CaF_2)	13	Al_H	23	Al_H	33	Ti
	Al_{H}		(CaF ₂)		(CaF ₂)	34	Al_{L}
_	Al_{H}	14	(CaF_2)	24	(CaF ₂)	35	Al_H
Ð	(CaF_2)		Al_{H}		Al_{H}		(CaF ₂)
6	(CaF ₂)	15	Ti	25	Ti	0.0	(CaF_2)
	$Al_{ m H}$	16	Al_L	26	Al_L	36	Al_{H}
7	Al_H	17	Al _H	27	Al_{H}	37	Ti
	(CaF ₂)		(CaF ₂)		(CaF_2)	38	Al_L
8	(CaF ₂)		(CaF ₂)		(CaF ₂)		
	Al_H	18	Al_H	28	Al_H		

Table 1. Configuration of the stacked target. The CaF_2 layers were sandwiched with the Al_H foils.

Table 2. Experimental parameters.

Target				
Measured (corrected) thickness (mg/cm ²)	CaF ₂ : 0.135			
	²⁷ Al (Al _H , purity: 99.999%): 2.57 (2.55)			
	²⁷ Al (Al _L , purity: >99%): 1.50 (1.49)			
	^{nat} Ti: 2.30			
Stack composition	12 sets of $Al_H(CaF_2)$ -(CaF ₂) Al_H foils			
	7 sets of ^{nat} Ti-Al _L foils			
Beam				
Measured (corrected) beam current (nA)	175 (185) ±11			
Primary energy (MeV)	29.0 ± 0.2			
Irradiation period (min)	30.0			
Measurement				
Series: Cooling time (distance, dead time)	Ser. 1: 3.2-5.4 h (10 cm, 0.8-1.9%)			
	Ser. 2: 5.6-21.6 h (5 cm, 0.1-1.8%)			
	Ser. 3: 1.2-2.8 d (1 cm, 0.1%)			
	Ser. 4: 2.9-3.0 d (1 cm, <0.1%)			
	Ser. 5: 43.5-77.0 d (1 cm, <0.1%)			

Nuclide	Half-life	Decay mode (%)	E _γ (keV)	I _γ (%)	Contributing monthing	Q-value	Threshold energy
					Contributing reaction	(MeV)	(MeV)
⁴⁷ Sc	3.3492 d	β- (100)	159.381	68.3(4)	$^{44}Ca(\alpha,p)$	-2.0	2.2
					$^{46}Ca(\alpha,t)$	-11.3	12.3
					48 Ca(α ,t2n)	-28.6	30.9
					⁴⁷ Ca decay		
⁴⁶ Sc	83.79 d	β- (100)	889.277	99.9840(10)	$^{43}Ca(\alpha,p)$	-1.5	1.7
			1120.545	99.9870(10)	$^{44}Ca(\alpha,d)$	-10.4	11.4
					46 Ca(α ,tn)	-22.0	23.9
^{44m} Sc	58.61 h	$\epsilon + \beta^+ (1.20)$	271.241	86.74(6)	$^{42}Ca(\alpha,d)$	-11.7	12.8
		IT (98.80)			43 Ca(α ,t)	-13.4	14.6
					⁴⁴ Ca(α,tn)	-24.5	26.7
^{44g} Sc	3.97 h	$\epsilon + \beta^+ (100)$	511.0	188.54(10)	$^{42}Ca(\alpha,d)$	-11.4	12.5
			1157.020	99.9(4)	$^{43}Ca(\alpha,t)$	-13.1	14.3
					$^{44}Ca(\alpha,tn)$	-24.2	26.5
					⁴⁴ Ti decay		
⁴³ Sc	3.891 h	$\epsilon + \beta^+ (100)$	372.9	22.5(7)	40 Ca(α ,p)	-3.5	3.9
			511.0	176.2(16)	$^{42}Ca(\alpha,t)$	-14.9	16.3
					43 Ca(α ,tn)	-22.8	24.9
					${}^{40}Ca(\alpha,n){}^{43}Ti$ decay	-11.2	12.3
⁴⁷ Ca	4.536 d	β- (100)	1297.09	67(13)	46 Ca(α , 3 He)	-13.3	14.5
					48 Ca(α , α n)	-10.0	10.8
					${}^{48}Ca(\alpha, \alpha p){}^{47}K$ decay	-15.8	17.1
⁵¹ Cr	27.704 d	ε (100)	320.0824	9.910(10)			

Table 3. Reaction and decay data retrieved from online databases [9-12].

Gamma lines in bold were used for data evaluation.



Fig. 1. Excitation function of the ${}^{nat}Ti(\alpha,x){}^{51}Cr$ monitor reaction with the IAEA recommended values [13].

3. Results and discussion

The production cross sections of ^{47,46,44m,44g,43}Sc and ⁴⁷Ca via the alpha-particle-induced reactions measured on CaF2 targets up to 29 MeV were converted to those on natural calcium targets. The derived cross sections are summarized in Table 4. The results are graphically shown in Figs. 2-7 with experimental data in the literature [5,14,15] and the model based theoretical cross sections available in the TENDL-2019 library [16]. The cross sections on the enriched targets [5,15] were normalized to those of natural targets using the isotopic ratios of natural calcium. Physical yields of ⁴⁷Sc in the alpha-particle-induced reactions on ^{nat}Ca and the enriched ⁴⁴Ca were deduced using the measured cross sections and are shown in Fig. 8.

The median projectile energies in each sandwiched pair of the CaF_2 layers are listed in Table 4. The total energy uncertainties of 0.2-0.8 MeV were propagated from the uncertainties of the primary beam energy (0.2 MeV) and the target thicknesses (1-5%). The total uncertainties of 9.3-35.1% for the cross sections are calculated from the square roots of the quadratic summation of the linearly propagated components. The components taken into account are beam intensity (6%), gamma-ray intensity (<19.4%), detector efficiency (5%), target thickness (5%) and counting statistics (0.3-27.8%).

Energy (MeV)	⁴⁷ Sc (mb)	⁴⁶ Sc (mb)	^{44m} Sc (mb)	^{44g} Sc (mb)	⁴³ Sc (mb)	⁴⁷ Ca (mb)
28.6 ±0.2	0.486 ± 0.048	7.47 ± 1.02	3.33 ± 0.31	1.67 ± 0.18	$53.9 \pm \!\! 5.4$	0.0724 ± 0.0254
27.6 ± 0.2	0.530 ± 0.069	$7.70 \pm \! 0.94$	$3.27 \pm \! 0.31$	1.75 ± 0.19	$66.9 \pm \! 6.6$	
26.6 ± 0.2	0.562 ± 0.061	7.32 ± 0.99	$3.13 \pm \! 0.30$	1.74 ± 0.19	81.8 ± 8.1	
25.6 ± 0.2	0.636 ± 0.069	$6.78 \pm \! 0.91$	$3.03 \pm \! 0.29$	$1.87 \pm \! 0.20$	$97.9 \pm \! 9.7$	
23.9 ± 0.3	$0.902 \pm \! 0.092$	6.09 ± 0.74	$2.76 \pm \! 0.26$	1.47 ± 0.17	$150 \pm \! 15$	
22.8 ± 0.3	1.18 ± 0.12	5.63 ± 0.72	2.63 ± 0.25	1.36 ± 0.16	$193 \pm \! 19$	
$20.9 \pm \! 0.3$	1.67 ± 0.16	$4.56 \pm \! 0.67$	1.74 ± 0.17	0.906 ± 0.108	$292 \pm \!\!29$	
19.6 ± 0.4	$2.13 \pm \! 0.20$	$3.33 \pm \! 0.48$	1.03 ± 0.10	0.588 ± 0.082	371 ± 36	
17.5 ± 0.4	$2.82 \pm \! 0.27$	$1.91 \pm \! 0.49$	0.142 ± 0.023	0.116 ± 0.031	$516\pm\!\!51$	
15.1 ±0.5	$2.57 \pm \! 0.25$				557 ± 55	
12.5 ±0.6	1.64 ± 0.16				478 ±47	
9.3 ±0.8	0.752 ± 0.078				232 ± 23	

Table 4. Production cross sections via the alpha-particle-induced reactions on natural calcium.

3.1 The ^{nat}Ca(α ,x)⁴⁷Sc reaction

⁴⁷Sc (T_{1/2} = 3.3492 d) can be produced from three stable isotopes of calcium. The main contribution is expected to be the direct ⁴⁴Ca(α,p)⁴⁷Sc reaction. ⁴⁷Sc can also be formed by decay of the co-produced ⁴⁷Ca (T_{1/2} = 4.536 d) and ⁴⁷K (T_{1/2} = 17.50 s). ⁴⁷Ca can be formed in the alpha-particle-induced reactions on two calcium isotopes (⁴⁶Ca: 0.004% and ⁴⁸Ca: 0.187%), while ⁴⁷K can be formed only on ⁴⁸Ca. The gamma line at 159.381 keV (I_γ = 68.3%) followed by the decay of ⁴⁷Sc was measured after cooling times of 1.2-2.8 d (Ser. 3). Their contributions from the two parent radionuclides were expected to be small because of the very low abundances of the two target isotopes. The measured cross section for ⁴⁷Ca production is only 0.0765 mb at 28.6 MeV (see section 3.6). No experimental information is available for the ⁴⁸Ca(α,αp)⁴⁷K reaction, but according to the TENDL-2019 database the predicted cross sections are in the range of micro barn at a 30 MeV alpha-particle energy.

The cross sections of the ^{nat}Ca(α, x)⁴⁷Sc reaction were derived and shown in Fig. 2 in comparison with the previous experimental data [5] and the theoretical values in the TENDL-2019 library [16]. The previous data of the ⁴⁴Ca(α, p)⁴⁷Sc reaction are normalized using the isotopic ratio of natural calcium. The peak position of the previous study is slightly shifted to the lower energy than that of ours. The previous data are larger below 15 MeV and smaller between 15 and 20 MeV than ours. The theoretically calculated values in the TENDL-2019 library overestimates both experimental results by a factor of 2 around at the peak energy.



Fig. 2. Cross sections of the ^{nat}Ca(α ,x)⁴⁷Sc reaction with the normalized values from the previous data [5] and the theoretical values from the TENDL-2019 library [16].

3.2 The ^{nat}Ca(α ,x)⁴⁶Sc reaction

The radionuclide ⁴⁶Sc can be produced in the alpha-particle-induced reactions on the ^{43,44,46,48}Ca target isotopes. However, contribution from the reactions on ⁴⁸Ca is not possible in the energy range of this work. ⁴⁶Sc has two isomeric states, the ground state ⁴⁶gSc ($T_{1/2} = 83.79$ d) and the short-lived meta-stable state ⁴⁶mSc ($T_{1/2} = 18.75$ s) decaying 100% to the ground state by the IT decay mode. The co-produced isomer ⁴⁶mSc completely decayed to the ground state before the measurement started. The cross sections of the ^{nat}Ca(α ,x)^{46g+m}Sc reaction were determined using the gamma line at 889.277 keV ($I_{\gamma} = 99.9840\%$). The spectra collected after cooling times of 43.5-77.0 d (Ser. 5) were used to determine the cross sections. The derived cross sections are compared with the previous experimental data [5] and the theoretically calculated values in the TENDL-2019 library [16] as shown in Fig. 3. The experimental data on the enriched targets of ⁴³Ca and ⁴⁴Ca in the previous study [3] were normalized using the isotopic ratio of natural calcium. The dominant contribution comes from ⁴⁴Ca due to approximately 15 times larger abundance of ⁴⁴Ca than ⁴³Ca. The normalized values for the ⁴⁴Ca(α ,x)⁴⁶Sc reaction in the previous study are slightly larger than ours below 25 MeV (see Fig 3.). The theoretical values of the TENDL-2019 library overestimates both experimental results above 20 MeV.



Fig. 3. Cross sections of the $^{nat}Ca(\alpha,x)^{46}Sc$ reaction with the normalized values from the previous data [5] and the theoretical values from the TENDL-2019 library [16].

3.3 The ^{nat}Ca(α ,x)^{44m}Sc reaction

Contribution for the formation of ^{44m}Sc ($T_{1/2} = 58.61$ h) is possible from ^{42,42,44}Ca up to 30-MeV alphaparticle energy. The gamma line at 271.241 keV ($I_{\gamma} = 86.74\%$) following the decay of ^{44m}Sc was used to assess the activity of the activated target foils. The spectra measured after cooling times of 1.2-2.8 d (Ser. 3) were adopted for cross section determination. The measured cross sections of the ^{nat}Ca(α,x)^{44m}Sc reaction are shown in Fig. 4, in which the results are compared with the previous experimental data [5] and the theoretical predication in the TENDL-2019 library [16]. The previous data on the enriched targets of ⁴²Ca and ⁴³Ca were normalized using the isotopic ratio of natural calcium. The normalized values from the ⁴²Ca target agree with our results above 23 MeV, however the literature data below 20 MeV are higher than our deduced cross sections. The theoretical prediction in the TENDL-2019 library is almost consistent with our data below 25 MeV, but is lower than ours at the higher energies.



Fig. 4. Cross sections of the ^{nat}Ca(α ,x)^{44m}Sc reaction with the normalized values from the previous data [5] and the theoretical values from the TENDL-2019 library [16].

3.4 The ^{nat}Ca(α ,x)^{44g}Sc reaction

Beside the direct production of ^{44g}Sc ($T_{1/2} = 3.97$ h) it can be populated by decay of the much longerlived ^{44m}Sc ($T_{1/2} = 58.61$ h) and ⁴⁴Ti ($T_{1/2} = 59.1$ y). The gamma line at 1157.020 keV ($I_{\gamma} = 99.9\%$) from the decay of ^{44g}Sc ($T_{1/2} = 3.97$ h) was used to assess the activity of the target foils. The spectra including the partial contribution of ^{44m}Sc were measured after cooling times of 5.6-21.6 h (Ser. 2). To subtract the contribution of ^{44m}Sc from the measured net counts, the activities of ^{44g}Sc (A_g) and ^{44m}Sc (A_m) at the end-ofbombardment were calculated using the cross sections (σ_m) derived in section 3.3 as:

$$A_{g} = N_{t}N_{b}\sigma_{g}\left(1 - e^{-\lambda_{g}t_{b}}\right) + \varepsilon_{IT}N_{t}N_{b}\sigma_{m}\frac{\lambda_{g}\lambda_{m}}{(\lambda_{g} - \lambda_{m})}\left(\frac{\left(1 - e^{-\lambda_{m}t_{b}}\right)}{\lambda_{m}} - \frac{\left(1 - e^{-\lambda_{g}t_{b}}\right)}{\lambda_{g}}\right)$$
(2)

$$A_m = N_t N_b \sigma_m (1 - e^{-\lambda_m \iota_b}) \tag{3}$$

where σ_g is the cross section for ^{44g}Sc, λ_g and λ_m are the decay constants of ^{44g}Sc and ^{44m}Sc (s⁻¹), ε_{IT} is the decay branching ratio (IT: 98.8%) from ^{44m}Sc to ^{44g}Sc. The expected net count of the gamma line N_{γ} was:

$$N_{\gamma} = \varepsilon_d \varepsilon_{\gamma} \varepsilon_t \left(A_g \frac{e^{-\lambda_g t_c} \left(1 - e^{-\lambda_g t_m}\right)}{\lambda_g} + \varepsilon_{IT} A_m \frac{\lambda_g}{(\lambda_g - \lambda_m)} \left(\frac{e^{-\lambda_m t_c} \left(1 - e^{-\lambda_m t_m}\right)}{\lambda_m} - \frac{e^{-\lambda_g t_c} \left(1 - e^{-\lambda_g t_m}\right)}{\lambda_g} \right) \right).$$
(4)

From Eqs. (2)-(4) and the measured net counts, we could determine σ_g . The contribution of ⁴⁴Ti was negligibly small due to its long half-life (T_{1/2} = 59.1 y) and no peaks at 67.87 (I_γ = 93.0%) and 78.32 keV (I_γ = 96.4%) from the decay of ⁴⁴Ti were found in any measured spectra.

The independent cross sections of the ^{nat}Ca(α, x)^{44g}Sc reaction are thus determined and shown in Fig. 5 in comparison with the previous experimental data [5] and the theoretical predication in the TENDL-2019 library [16]. The previous data on the enriched targets of ⁴²Ca and ⁴³Ca were normalized using the isotopic ratio of natural calcium. The normalized values of Levkovski (1991) for the ⁴²Ca(α, x)^{44g}Sc reaction alone (not including the contribution of the ⁴³Ca(α, x)^{44g}Sc reaction) are slightly larger than our measured cross sections. The prediction of the model based theoretical values in the TENDL-2019 library agrees well with our result.



Fig. 5. Cross sections of the $^{nat}Ca(\alpha,x)^{44g}Sc$ reaction with the normalized values from the previous data [5] and the theoretical values from the TENDL-2019 library [16].

3.5 The ^{nat}Ca(α ,x)⁴³Sc reaction

Reactions on the ^{40,42,43}Ca target isotopes can contribute to the direct formation of ⁴³Sc ($T_{1/2} = 3.891$ h) in the investigated energy range. Decay of co-produced ⁴³Ti ($T_{1/2} = 509$ ms) may also provide considerable contribution as the ⁴⁰Ca(α ,n)⁴³Ti reaction takes place on the most abundant ⁴⁰Ca (96.941%) isotope. To assess the ^{nat}Ca(α ,x)⁴³Sc process the gamma line at 372.9 keV ($I_{\gamma} = 22.5\%$) from decay of ⁴³Sc was used. Spectra collected after cooling times of 5.6-21.6 h (Ser. 2) were used. The applied cooling times assured complete decay of the co-produced ⁴³Ti to ⁴³Sc. The cumulative cross sections of the ^{nat}Ca(α ,x)⁴³Sc reaction were derived. The result shown in Fig. 6 is compared with the previous experimental data [5,14,15] and the theoretical predication in the TENDL-2019 library [16]. The previous data of the ⁴⁰Ca(α ,p)⁴³Sc reaction [5,15] are normalized using the isotopic ratio of natural calcium. The two experimental datasets by Levkovski (1991) and Alabyad et al. (2018) agree with our data above 12 MeV within uncertainties. Their data at the lower energy region below 12 MeV are larger than ours. The data by Howard et al. (1974) show similar shape but higher amplitude above 10 MeV. The peak position of the TENDL-2019 values is almost consistent with the experimental data, although the peak amplitude is higher than our data.



Fig. 6. Cross sections of the $^{nat}Ca(\alpha,x)^{43}Sc$ reaction with the previous data [14], the normalized values from the previous data [5,15] and the theoretical values from the TENDL-2019 library

[16].

3.6 The ^{nat}Ca(α ,x)⁴⁷Ca reaction

Since ⁴⁷Ca can be produced only in the reactions from the low abundant ^{46,48}Ca target isotopes, the expected cross section is low. The cross sections of the ^{nat}Ca(α,x)⁴⁷Ca reaction were determined using the gamma line at 1297.09 keV (I₇ = 67%) from decay of ⁴⁷Ca (T_{1/2} = 4.536 d). Decay contribution from the energetically possible co-produced parent ⁴⁷K (T_{1/2} = 17.50 s), which completely decayed to ⁴⁷Ca during the cooling times of 1.2-2.8 d (Ser. 3), is also included. Therefore, the deduced cross section is considered to be cumulative. Only one cross section of the ^{nat}Ca(α,x)⁴⁷Ca reaction was deduced at the highest bombarding particle energy of 28.6 MeV with a statistical uncertainty below 30%. Due to the large relative uncertainty (19.4%) of the gamma-ray intensity and the poor counting statistics (27.8%), the cross section has a large total relative uncertainty of 35.0%. The result is shown in Fig. 7 and compared with the theoretical predication in the TENDL-2019 library [16]. The TENDL-2019 predicted value is a factor of two higher than our experimental data. No experimental data were found in the literature for this reaction.



Fig. 7. Cross sections of the $^{nat}Ca(\alpha,x)^{47}Ca$ reaction with the theoretical values from the TENDL-2019 library [16].

3.7 Physical yield of ⁴⁷Sc

Physical yields [17] for production of ⁴⁷Sc in the alpha-particle-induced reactions on ^{nat}Ca and the enriched ⁴⁴Ca were deduced from the measured cross sections of the ^{nat}Ca(α ,x)⁴⁷Sc reaction presented in section 3.1. Assuming that the contributions from reactions on low abundant target isotopes ^{46,48}Ca were negligibly small, the yield was transformed to a 100% ⁴⁴Ca target. The deduced physical yields using ^{nat}Ca and ⁴⁴Ca can reach 34.5 and 1650 kBq/µAh at 20 MeV, respectively. Above 20 MeV, the amount of the longer-lived ⁴⁶Sc impurity is getting too high. For practical use of ⁴⁷Sc, the amount of the co-produced longer-lived impurity of ⁴⁶Sc should be estimated. The reaction threshold energy of the ⁴⁴Ca(α ,d)⁴⁶Sc reaction is E_{th} = 11.4 MeV. However, the lowest energy cross section point measured by us for the ^{nat}Ca(α ,d)⁴⁶Sc reaction is 17.5 MeV. Therefore, the experimental data were extended down to the threshold energy of the reaction by using the TENDL-2019 prediction. Using this assumption, the estimated yields of the co-produced ⁴⁶Sc impurity using the ^{nat}Ca and ⁴⁴Ca targets are 0.74 and 35.5 kBq/µAh, respectively.



Fig. 8. Physical yields of ⁴⁷Sc in the alpha-particle-induced reactions on ^{nat}Ca and ⁴⁴Ca.

4. Summary

We measured the activation cross sections of the alpha-particle-induced reactions on natural calcium. The experiment was performed at the RIKEN AVF cyclotron using the stacked-foil activation technique and high-resolution gamma-ray spectrometry. The production cross sections of ^{47,46,44m,44g,43}Sc and ⁴⁷Ca using CaF₂ target material were measured and converted to those on natural calcium. The derived cross sections are compared with the literature data and the model based theoretically calculated values of the TENDL-2019 library. Physical yields of ⁴⁷Sc via the alpha-particle-induced reactions on ^{nat}Ca and ⁴⁴Ca were deduced. The experimental result is expected to contribute to research and development of nuclear medicine using the promising therapeutic radionuclide ⁴⁷Sc.

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Declarations of interest

None

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