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Measurement of valence-band offsets of InAlN/GaN heterostructures grown by metal-organic vapor phase epitaxy

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The valence band offsets, ΔE_V , of $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ heterostructures grown by metal-organic vapor phase epitaxy were evaluated by using x-ray photoelectron spectroscopy (XPS). The dependence of the energy position and the full width at half maximum of the Al 2p spectrum on the exit angle indicated that there was sharp band bending caused by the polarization-induced electric field combined with surface Fermi-level pinning in each ultrathin InAlN layer. The ΔE_V values evaluated without taking into account band bending indicated large discrepancies from the theoretical estimates for all samples. Erroneous results due to band bending were corrected by applying numerical calculations, which led to acceptable results. The evaluated ΔE_V values were 0.2 ± 0.2 eV for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, 0.1 ± 0.2 eV for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and 0.0 ± 0.2 eV for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$. Despite the large decrease of around 1.0 eV in the band gap of InAlN layers according to the increase in the In molar fraction, the decrease in ΔE_V was as small as 0.2 eV. Therefore, the change in band-gap discontinuity was mainly distributed to that in conduction band offset. © 2011 American Institute of Physics.

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I. INTRODUCTION

Group III-nitrides and their alloys are promising materials that should enable excellent high-power, high-frequency, and optoelectronic devices to be fabricated. Their heterostructures can provide a high-density two-dimensional electron gas (2DEG) with a high saturation velocity in wide-gap semiconductors. Lattice-matched $\text{In}_x\text{Al}_{1-x}\text{N}$ ($x = 0.17\text{--}0.18$)/GaN heterostructures without doping can especially provide a 2DEG with high density exceeding $2 \times 10^{13} \text{ cm}^{-2}$,^{1–8} due to the difference in spontaneous polarization at the interface.⁹ High-performance InAlN/GaN high-electron mobility transistors (HEMTs) have actually been fabricated and reported.^{2,3,5,10} If we expand the scope to include nonlattice-matched InAlN/GaN heterostructures, the band gap, E_g , of InAlN can be greatly changed from 6.2 eV (AlN) to 0.67 eV (InN).¹¹ It has theoretically been predicted that the decrease in the conduction band offset, ΔE_C , according to the increase in the In molar fraction is much larger than decrease in the valence band offset, ΔE_V .¹² This interface property is useful for barrier-height engineering. The large change in ΔE_C according to the molar fraction can be verified by confirming the small change in ΔE_V . However, there have been no reports on measuring the change in ΔE_V by changing the molar fraction.

One of the most useful methods of investigating ΔE_V is by applying x-ray photoelectron spectroscopy (XPS) to the interface between ultrathin InAlN and GaN layers. However, evaluating ΔE_V by using XPS is possibly erroneous due to sharp band bending caused by the spontaneous and piezoelectric polarization in ultrathin InAlN layers. In our prelimi-

nary work,¹³ the ΔE_V of lattice-matched $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$ grown by metal-organic vapor phase epitaxy (MOVPE) was evaluated correctly by taking into account band bending due to the polarization-induced electric field combined with surface Fermi-level pinning. In addition to the lattice-matched $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$ heterostructure, the pseudomorphic heterostructures of $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$ and $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ were also grown and characterized in this work. Erroneous results due to band bending in ultrathin InAlN layers were corrected by applying numerical calculations. Using the measured peak positions of core levels, E_{CL} 's, without correction, the evaluated ΔE_V values deviated from the theoretical estimates according to Refs. 11 and 12. However, the corrected values for all the InAlN/GaN interfaces were in good agreement with the theoretical estimates. Despite the large decrease in E_g of the InAlN layers according to the increase in the In molar fraction, small decrease in ΔE_V were evaluated at the InAlN/GaN interfaces. Therefore the decrease in E_g discontinuity was mostly distributed to that in ΔE_C at the InAlN/GaN interfaces.

II. EXPERIMENTAL

The sample structure used to investigate the interface with XPS is schematically shown in Fig. 1. All the samples were grown on sapphire (0001) substrates by MOVPE without any doping. $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ heterostructures with 2.5 nm-thick InAlN layers were prepared. Here, we have defined the photoelectron exit angle, θ , as the elevation angle indicated in Fig. 1. To measure bulk material constants, thick (around 15 nm) $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$, $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$, and $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layers on GaN and a 2 μm -thick GaN single layer were also prepared. Tri-

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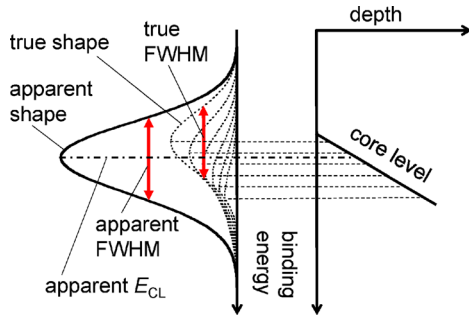


FIG. 4. (Color online) Schematic explaining increases in E_{CL} and FWHM in apparent spectrum.

range of probing depths and observed surface potentials that will be described later, while the nonuniform electric fields were calculated for GaN layers. Similar results have previously been reported.^{6–8} In XPS measurement, only shallow parts of InAlN layers were probed as shown in Fig. 3(b) for the samples with thick InAlN layers with 2DEG, resulting in no detection of the Ga 3d spectrum from the host GaN layer. The samples with ultrathin InAlN layers were without 2DEG, which was confirmed by the contactless eddy current method. Consequently, the band bending in the GaN layers of the samples with ultrathin InAlN layers would be observed as a slight tilt of the potential or a nearly flat band in the range of the probing depth. For example, Fig. 3(c) indicates the observed result, for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ (2.5 nm)/GaN sample, that will be described in Sec. IV. A normally-off type GaN-based HEMT has recently been fabricated by using an ultrathin $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ (1 nm)/AlN (1 nm) barrier,¹⁰ which supports our results because the barrier thickness and the sum of interface charges⁸ are close to those of the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ (2.5 nm)/GaN sample.

In the situations shown in Figs. 3(b) and 3(c), the core levels in the InAlN layer are also bent along with the valence band. As schematically outlined in Fig. 4, the observed spectrum is given by integrating the true spectrum from each depth point along the bent core level. Therefore, sharp band bending results in apparent increases in E_{CL} and the full width at half maximum (FWHM) of core-level spectra from the InAlN layer. This phenomenon becomes more remarkable by increasing θ to extend the probing depth. Furthermore, the E_{CL} of the core-level spectra from the GaN layer increases due to the total potential drop in the InAlN layer for the InAlN/GaN heterostructure sample. With the band bending in Fig. 3(c), ΔE_V might be evaluated to have a negative value despite the type-I alignment. Therefore, we have to derive the actual E_{CL} 's by subtracting increases due to band bending so that we can correctly evaluate ΔE_{CL} .

The sequence of correction for ΔE_V at the InAlN/GaN heterointerfaces involves three steps

- (1) The internal electric field in InAlN is extracted. Assuming a uniform internal electric field in the InAlN layer, the band bending for the entire heterostructure including the GaN layer is calculated by solving the Poisson equation using the measured surface Fermi level, $E_{F\text{surf}}$, position as the boundary condition for strong pinning. Then, the apparent E_{CL} and FWHM due to band bending

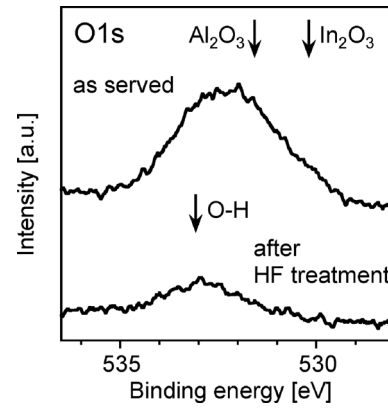


FIG. 5. XPS O 1s spectra observed for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer.

are calculated for Al 2p and Ga 3d. This procedure is repeated to find the most suitable internal electric field in InAlN that achieves the apparent E_{CL} and FWHM observed for Al 2p and Ga 3d.

- (2) E_{CL} values are corrected for Al 2p and Ga 3d by subtracting the apparent increases due to band bending to correctly evaluate the ΔE_{CL} value.
- (3) ΔE_V is derived using the corrected ΔE_{CL} value in Eq. (1).

The details on the calculations to estimate the apparent E_{CL} and FWHM of core-level spectra due to band bending are as follows. A core-level spectrum as a function of the binding energy, E , from a layer of thickness d is given by¹⁵

$$I(E) = \int_0^d I_0(E, z) \exp\left(-\frac{z}{\lambda}\right) dz, \quad (2)$$

where z , λ , and $I_0(E, z)$ correspond to the depth from the surface, the escape depth of the photoelectrons, and the spectrum generated at each depth point. Here, to do the calculations within a practical time frame and to avoid heavy computation, the $I_0(E, z)$ for one spin orbital is represented by the pseudo-Voigt function given by

$$V(E, z) = I_{00} \left[\alpha \exp\left\{-\ln 2 \frac{(E - E_0)^2}{(F/2)^2}\right\} + (1 - \alpha) \frac{1}{1 + \left\{-\frac{(E - E_0)^2}{(F/2)^2}\right\}} \right], \quad (3)$$

where I_{00} , α , E_0 , and F are the intensity, the ratio of the Gaussian function, the binding energy of the core level, and the actual FWHM. When surface band bending is not negligible on the scale of the escape depth of photoelectrons, E_0 should be treated as a function of depth z to be $E_0(z)$. The F was assumed to be a sample dependent constant. The α was fixed to 0.73 for both Al 2p and Ga 3d spectra. For Al 2p and Ga 3d spectra, $I_0(E, z)$ is given by a weighted combination of two pseudo-Voigt-function components taking into account the spin-orbit splitting of 0.41 eV for Al 2p and

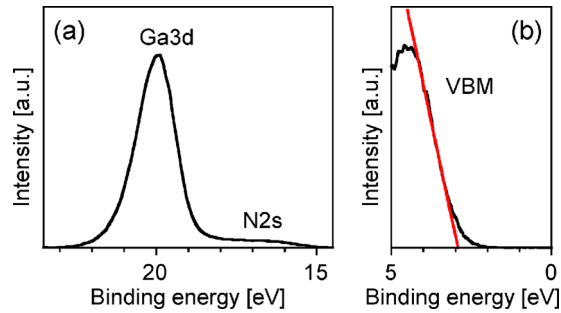


FIG. 6. (Color online) XPS Ga 3d and VBM spectra obtained for thick GaN sample.

0.44 eV for Ga 3d.¹⁶ Furthermore, λ can be changed by changing θ according to

$$\lambda = \lambda_0 \sin \theta, \quad (4)$$

where λ_0 indicates the inelastic mean free path of a photoelectron. The value of λ_0 was calculated using a previously reported equation^{17,18} that was demonstrated to fit the experimental data for several materials including GaN.¹⁹ The λ_0 of Al 2p spectrum was calculated to be 2.87 nm for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$, 2.79 nm for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$, and 2.76 nm for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$. The λ_0 of Ga 3d spectrum was calculated to be 2.58 nm for GaN, 2.95 nm for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$, 2.88 nm for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$, and 2.84 nm for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$.

IV. RESULTS AND DISCUSSION

A. Measurement of bulk material constants

We found that HF treatment efficiently removed oxides from the InAlN and GaN surfaces. Figure 5 shows the O 1s spectra before and after HF treatment for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ at $\theta = 45^\circ$. In_2O_3 component appear at 530.2 eV (Ref. 20) and that for Al_2O_3 appear at 531.6 eV.²¹ The spectral intensity at these binding energies drastically decreased after HF treatment leaving only a faint component assigned to O–H bonding.²² Similar results were also obtained for the $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$, $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$, and GaN surfaces.

The Ga 3d and VBM spectra observed at $\theta = 45^\circ$ for the HF-treated thick GaN layer on a sapphire substrate are shown in Fig. 6. The efficiency of HF treatment can also be seen here. The shape of the Ga 3d spectrum was identical to that obtained for the GaN (2×2) clean surface grown by molecular beam epitaxy, which indicated that the surface was nearly oxide free. The VBM position was determined by linear extrapolation of the low binding energy edge of the valence band spectrum. The value of $(E_{\text{Ga } 3d}^{\text{GaN}} - E_{\text{V}}^{\text{GaN}})$ was evalu-

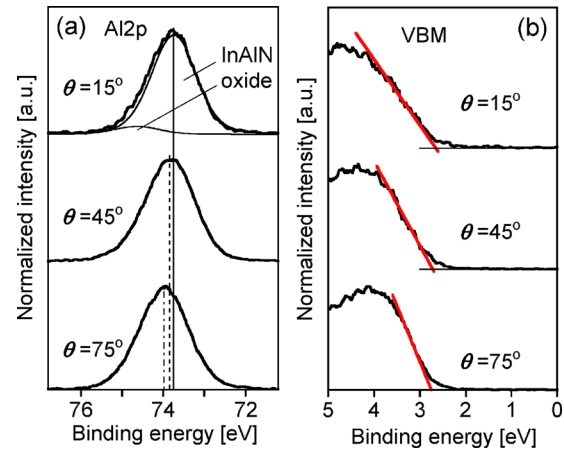


FIG. 7. (Color online) XPS Al 2p and VBM spectra obtained for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ (14.5 nm)/GaN heterostructure.

ated to be 17.10 ± 0.07 eV. The same value was also obtained for the thick GaN layer on an SiC substrate.

Figures 7(a) and 7(b) show the Al 2p and VBM spectra obtained at three different θ 's for the thick $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ samples. No oxide components were detected for Al 2p spectra at $\theta = 45^\circ$ or 75° , even though they were detected at the surface sensitive angle of $\theta = 15^\circ$. It can be seen that the Al 2p spectrum shifts to higher binding energies with the increase in θ , which indicates that there was sharp band bending in the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer. We also found that the FWHM of the Al 2p spectrum increased with θ . The results are summarized in Table I with the values calculated by assuming an internal electric field of 1.2 MV/cm in the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer. Note that the corrected VBM position in Table I is close to that reported for $\text{In}_{0.16}\text{Al}_{0.84}\text{N}$.²³ Consequently, the bulk material constant $(E_{\text{Al } 2p}^{\text{InAlN}} - E_{\text{V}}^{\text{InAlN}})$ was determined to be 71.10 ± 0.07 eV for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$. The same value was also obtained for the thick $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ layer by applying the same method. However, applying the same method to the thick $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layer did not lead to an acceptable result. The thickness needed to investigate the bulk material constant was probably too large for the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layer to maintain good bulk quality. (Precise evaluation of the critical thickness was beyond the scope of this work.) Therefore, the value of $(E_{\text{Al } 2p}^{\text{InAlN}} - E_{\text{V}}^{\text{InAlN}})$ for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ was extrapolated from the results of two other samples to be 71.10 eV.

B. Measurement of band offsets

To determine ΔE_{V} , the ΔE_{CL} between Al 2p and Ga 3d spectra was measured for each heterostructure with the thin

TABLE I. Summary of measured and calculated values of E_{CL} and FWHM for thick $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer.

θ	$E_{\text{Al } 2p}$ (eV)			FWHM of Al 2p (eV)			VBM (eV)		
	Measured	Calculated		Measured	Calculated		Measured	Calculated	
		Apparent	Corrected		Apparent	Corrected		Apparent	Corrected
15°	73.75	73.75	73.66	1.28	1.29	1.28	2.65	2.63	2.54
45°	73.87	73.87	73.66	1.40	1.38	1.28	2.72	2.74	2.54
75°	73.98	73.95	73.66	1.42	1.44	1.28	2.75	2.75	2.54

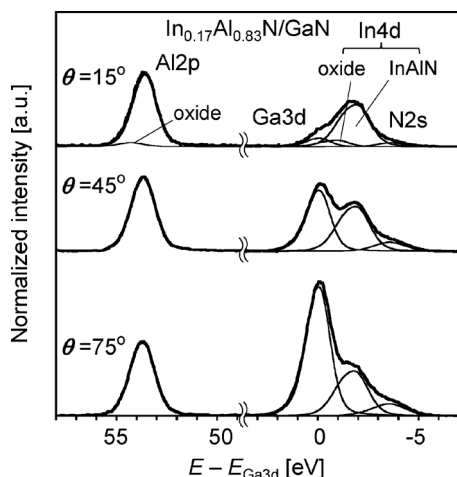


FIG. 8. XPS Al 2*p*, Ga 3*d*, In 4*d*, and N 2*s* core-level spectra obtained for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ (2.5 nm)/GaN heterostructure.

(2.5 nm) InAlN top layer. The spectra we obtained are shown in Fig. 8 for the lattice-matched $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN heterostructure, in Fig. 9 for the pseudomorphic $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ /GaN heterostructure, and in Fig. 10 for the pseudomorphic $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ /GaN heterostructure. The measured values of ΔE_{CL} and ΔE_{V} without correction have been summarized in Table II. Using the previously mentioned bulk constant values, all the ΔE_{V} values were evaluated to be negative and were not independent of θ . The evaluated ΔE_{V} values in Table II are also plotted in Fig. 11 to compare them with the calculated values according to Refs. 11 and 12. The ΔE_{V} values without correction deviated from the calculated values and indicated type-II alignments for all the samples despite the prediction of type-I alignments. In addition, the E_{CL} and FWHM of the Al 2*p* spectra for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN and $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ /GaN interfaces increased with increasing θ as summarized in Table III. These results indicate that there was sharp band bending in the ultrathin InAlN layers again. Here, the detected internal electric field decreased according to the increase in the In molar fraction. The observed trend is in agreement with the theoretical estimates.⁹

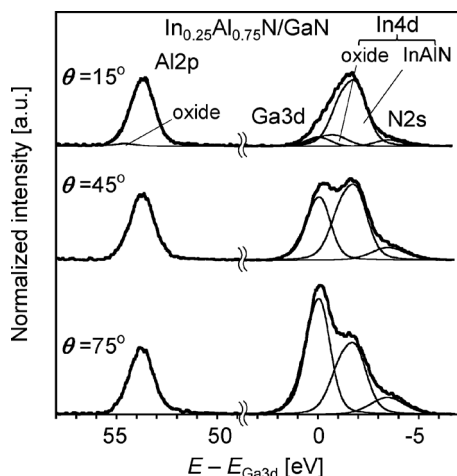


FIG. 9. XPS Al 2*p*, Ga 3*d*, In 4*d*, and N 2*s* core-level spectra obtained for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ (2.5 nm)/GaN heterostructure.

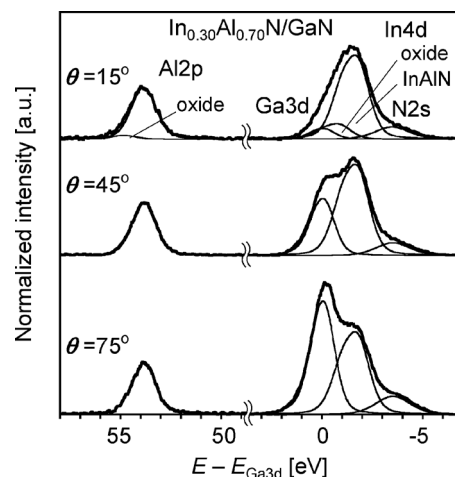


FIG. 10. XPS Al 2*p*, Ga 3*d*, In 4*d*, and N 2*s* core-level spectra obtained for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ (2.5 nm)/GaN heterostructure.

For the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN and $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ /GaN samples, increases in the E_{CL} and FWHM of the Al 2*p* spectra were sufficiently large to extract the internal electric field in the InAlN layers. The previously mentioned calculation reproduced the observed phenomenon for Al 2*p* in the range of experimental error as is also summarized in Table III. The assumed internal electric fields were 3.2 MV/cm for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer and 2.4 MV/cm for the $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ layer. For Ga 3*d* spectra, we found that the apparent shift in $E_{\text{Ga 3d}}$ due to band bending in GaN was a few tens of millielectron volts for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN sample and negligible for the other samples. The corrected ΔE_{V} values were 0.2 eV for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN and 0.1 eV for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ /GaN, as listed in Table IV with the corrected values of ΔE_{CL} . The measured band alignments for the quasiflat-band situation using the reported E_{g} values¹¹ are summarized in Figs. 12(a) and 12(b) for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ /GaN and $\text{In}_{0.25}\text{Al}_{0.75}\text{N}$ /GaN interfaces. Note that the resulting E_{Fsurf} position for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}$ layer in Fig. 12(a) is close to the experimental result previously reported for $\text{In}_{0.16}\text{Al}_{0.84}\text{N}$.²³

For the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ /GaN sample, however, no large increases depending on θ were observed in either the $E_{\text{Al 2p}}$ or FWHM as summarized in Table III. Nevertheless, piezoelectric polarization did not completely compensate for spontaneous polarization in the pseudomorphic $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layer

TABLE II. Summary of measured results without correction for InAlN/GaN interfaces.

x	θ	ΔE_{CL} (eV)	ΔE_{V} (eV)
0.17	15°	53.61	−0.4
	45°	53.71	−0.3
	75°	53.74	−0.3
0.25	15°	53.70	−0.3
	45°	53.75	−0.3
	75°	53.78	−0.2
0.30	15°	53.85	−0.2
	45°	53.85	−0.2
	75°	53.85	−0.2

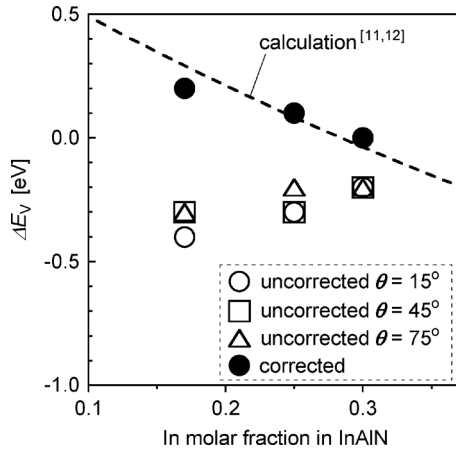


FIG. 11. Measured ΔE_V values with and without correction vs. In molar fraction in comparison with theoretical calculation according to Refs. 11 and 12.

according to the theoretical estimate in Ref. 9. Therefore, it is likely there was an internal electric field in the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layer. The numerical calculations indicated that the increases in apparent $E_{\text{Al } 2p}$ and FWHM were negligible within experimental error with internal electric fields up to 1.2 MV/cm in the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ layer. This maximum value for the electric field is in good agreement with the theoretical estimates in Ref. 9 and leads to the correction of ΔE_V at the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ interface to be 0.0 eV as shown in Fig. 12(c) and Table IV. The rationality of this correction is given by resultant E_{Fsurf} positions as follows.

The Fermi-level stabilization energy, E_{Fstab} , has been reported to be an energy position where E_{F} stabilizes for sufficiently high damage density.²⁴ The E_{Fstab} lies at 4.9 eV from the vacuum level in all semiconductors.²⁴ This energy level can therefore be used as a common benchmark point for E_{Fsurf} position. In the quasiflat-band situation of InAlN/GaN heterointerface, the value of $(E_{\text{Fsurf}} - E_{\text{Fstab}})$ should be identical for both InAlN and GaN layers. The measured values of $(E_{\text{Fsurf}} - E_{\text{Fstab}})$ are summarized in Table V. Note that there was 0.2 eV of difference in $(E_{\text{Fsurf}} - E_{\text{Fstab}})$ between $\text{In}_{0.3}\text{Al}_{0.7}\text{N}$ and GaN without correction for the $\text{In}_{0.3}\text{Al}_{0.7}\text{N}/\text{GaN}$ interface, whereas no such differences were

TABLE IV. Summary of corrected ΔE_{CL} and ΔE_V values for InAlN/GaN interfaces.

x	ΔE_{CL} (eV)	ΔE_V (eV)
0.17	54.20	0.2
0.25	54.13	0.1
0.30	54.04	0.0

seen in the corrected data for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$ and $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$ interfaces. The difference in uncorrected values of $(E_{\text{Fsurf}} - E_{\text{Fstab}})$ for the $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ interface must have been caused by the internal electric field. After correction assuming 1.2 MV/cm in $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$, the difference disappeared, as seen in Table V. Moreover, the dispersion of E_{Fsurf} for the same surface pretreatment for InAlN/GaN heterostructures should be discussed. Even though the E_{Fsurf} positions of the HF-treated samples may disperse statistically depending on the sample, the dispersion for the same treatment is likely to be within a few 0.1 eV. As can be seen from Table V, the $(E_{\text{Fsurf}} - E_{\text{Fstab}})$ was evaluated to be 0.0 eV for the $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$ interface and -0.2 eV for the $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$ interface. For the $\text{In}_{0.3}\text{Al}_{0.7}\text{N}/\text{GaN}$ sample, however, the value was estimated to be 0.3 eV for the GaN surface without correction resulting in a wide dispersion of 0.5 eV. This discrepancy dependent on the sample is not acceptable. After correction assuming 1.2 MV/cm in $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$, the dispersion width reduced to be an acceptable value of 0.2 eV. Therefore, correction resulting in a ΔE_V value of 0.0 eV for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}$ was appropriate.

As can be seen in Fig. 11, the corrected ΔE_V values are in good agreement with the calculated values for all the interfaces. The experimental error was evaluated by considering the fitting errors and the discrepancy between repeated measurements to be ± 40 meV for ΔE_{CL} and ± 70 meV for the two $(E_{\text{CL}} - E_V)$ components in Eq. (1). Without correction by using numerical calculations, the total error for ΔE_V is ± 180 meV. Fitting by numerical calculations introduces an additional error of ± 60 meV resulting in a total error of ± 240 meV. Since the final results were approximated values in units of 0.1 eV, the errors in both cases should be ± 0.2 eV,

TABLE III. Comparison of measured and calculated values for $E_{\text{Al } 2p}$ and FWHM of $\text{Al } 2p$ spectra from ultrathin InAlN layers on GaN .

x	θ	$E_{\text{Al } 2p}$ (eV)			FWHM of $\text{Al } 2p$ (eV)		
		Measured	Calculated		Measured	Calculated	
			Apparent	Corrected		Apparent	Corrected
0.17	15°	74.01	74.00	73.80	1.36	1.35	1.28
	45°	74.11	74.11	73.80	1.40	1.40	1.28
	75°	74.13	74.14	73.80	1.41	1.41	1.28
0.25	15°	73.69	73.68	73.52	1.40	1.41	1.36
	45°	73.74	73.76	73.52	1.43	1.43	1.36
	75°	73.77	73.77	73.52	1.44	1.44	1.36
0.30	15°	73.70	73.69	73.61	1.42	1.42	1.41
	45°	73.72	73.73	73.61	1.42	1.42	1.41
	75°	73.73	73.74	73.61	1.42	1.42	1.41

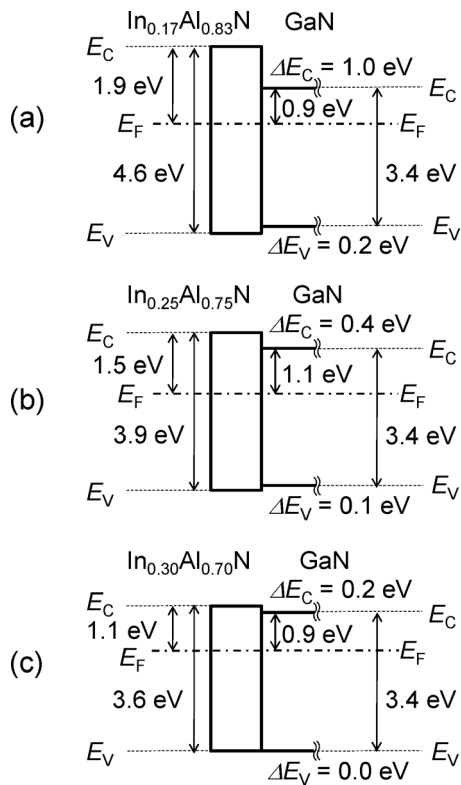


FIG. 12. Schematic band alignments using measured ΔE_V values and reported E_g values (Ref. 11), for (a) $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, (b) $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and (c) $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ interfaces in quasiflat-band situation.

which is a typical value for the XPS measurements of ΔE_V .^{25–28} Even though the evaluated ΔE_V values were small, the E_g of semiconductors on both sides was much larger than the total error. Therefore, the present results provide important information that ΔE_V values are small despite the large band gaps of all samples. The measurement of small ΔE_V values with similar error values can also be found in other reports on TiO_2/GaN (0.09 ± 0.25 eV) (Ref. 27) and HfO_2/GaN (-0.1 ± 0.1 eV) (Ref. 29) interfaces. As can be seen in Fig. 12, the decrease in ΔE_V according to the increase in the In molar fraction should have been much smaller than that in ΔE_C . In other words, the change in E_g discontinuity was mostly distributed to that in ΔE_C at the InAlN/GaN interfaces.

V. CONCLUSION

The ΔE_V of $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$ heterostructures grown by MOVPE were

TABLE V. ($E_{\text{Fsurf}} - E_{\text{Fstab}}$) values for the InAlN(2.5 nm)/GaN heterostructures in quasiflat-band situation.

	$E_{\text{Fsurf}} - E_{\text{Fstab}}$ (eV)			
	$\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$ (corrected)	$\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$ (corrected)	$\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$	
			Uncorrected	Corrected
InAlN surface	0.0	-0.2	0.1	0.0
GaN surface	0.0	-0.2	0.3	0.0

evaluated by using XPS. The E_{CL} and FWHM of the Al 2p increased according to the exit angle, which was explained by band bending due to the polarization-induced electric field in ultrathin InAlN layers combined with surface Fermi-level pinning. The ΔE_V values evaluated by the conventional method were in disagreement with the theoretical estimates for all samples. The evaluations were therefore corrected by numerical calculations taking into account band bending, which led to the ΔE_V values being in good agreement with the theoretical estimates for all the interfaces. The evaluated ΔE_V values were 0.2 ± 0.2 eV for $\text{In}_{0.17}\text{Al}_{0.83}\text{N}/\text{GaN}$, 0.1 ± 0.2 eV for $\text{In}_{0.25}\text{Al}_{0.75}\text{N}/\text{GaN}$, and 0.0 ± 0.2 eV for $\text{In}_{0.30}\text{Al}_{0.70}\text{N}/\text{GaN}$, which indicated that the large change in band-gap discontinuity was mostly distributed to that in ΔE_C .

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