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Title	Mass deformation of twisted super Yang-Mills theory with fuzzy sphere solution
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Citation	Journal of High Energy Physics, 2011(9), 19 https://doi.org/10.1007/JHEP09(2011)019
Issue Date	2011-09
Doc URL	https://hdl.handle.net/2115/47956
Type	journal article
File Information	JHEP2011-9_19.pdf



Mass deformation of twisted super Yang-Mills theory with fuzzy sphere solution

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Abstract

We investigate mass deformation of twisted superalgebra of $U(N)$ super Yang-Mills (SYM) theories in several models and in several dimensions, motivated by the method formulated in [1]. We show that there are several ways to perform the deformation, if a model possesses four scalar supercharges except for two dimensional B-model. We also evaluate classical vacuum solutions of the potentials of scalar fields for each model. We then find that it is always possible to find fuzzy sphere solution in the theories.

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1 Introduction

Presently supersymmetry (SUSY) is believed to be one of the most prospective candidates as the ingredients beyond the standard model and applied to wide variety of models. Since our world is not supersymmetric in fact, SUSY should be broken by some non-perturbative mechanisms. One promising way to deal with such a non-perturbative aspect is to make use of a lattice simulation. Indeed for quantitative evaluation, the lattice simulation is widely used at present for example in search of non-perturbative effects in QCD and SUSY breaking. In addition, it is also interesting to investigate AdS/CFT correspondence [2, 3] from the point of view at lattice or with non-lattice techniques [4, 5].

Even though there are the famous doubling problem and the breakdown of the Leibniz rule on the lattice [6], some approaches address the formulation of supersymmetric theory on a lattice (See reviews in [7–9]). Some insist that the model is supersymmetric invariant up to one or two supersymmetry on the lattice. Full SUSY will, then, be recovered at continuum limit without any fine-tuning [10–13]. While there is another ambitious approach, called link approach [14, 15], which claims the model can be exactly full supersymmetric even on the lattice. Besides it suggests the relationship between the noncommutativity in the link approach and the nonlocality in a new formulation [16].

For the simulation of supersymmetric theory on the lattice, there is a technical problem that a flat direction appears as a classical configuration in the potential of scalar fields [17]. The presence of the flat direction indeed leads to the breakdown of iteration of computation because of over-counting of field configurations. Therefore it should be resolved by some methods, such as addition of mass terms to lift up the potential of scalar fields. In [18, 19] a soft SUSY-breaking mass is introduced to scalar fields in two dimensional models in order to control the flat direction. As a result, one can obtain the correct continuum limit.

In [1], they introduced the mass deformation method [20] to $N = 8, D = 2$ twisted SYM theory to remove the flat direction and consequently showed some steps to obtain $N = 4, D = 4$ SYM theory formulated on the lattice*. In consequence, they succeeded in introducing mass parameter into superalgebra by deforming supercharges to keep some of supersymmetries up to $SU(2)_R$ symmetry. Then, once the theory is deformed to be nilpotent up to $SU(2)_R$ symmetry[†], the action receives mass terms for each field. In addition, they focused on the fact that there appears fuzzy sphere configuration for the potential minimum of scalar field. Utilizing the technique, which is well-known in string theoretical point of view [21, 22], they extended the fuzzy sphere configuration to a noncommutative space. By careful tuning of a lattice spacing, the mass parameter and a parameter of non-commutativity, the model eventually leads to noncommutative $N = 4, D = 4$ SYM theory, where two dimensions correspond to original spacetime and the rest of two dimensions to non-commutative ones. It is also possible to formulate $N = 8, D = 2$ SYM models, which can be obtained without embedding on the fuzzy sphere lattice and simply with tuning the mass parameter to be zero to keep the potential well. They also commented on the case for $N = 4, D = 2$ in the same procedure. In [23], this mass deformation method is also applied to a model combined with other lattice SUSY approach.

This paper is motivated by the mass deformation technique in [1]. We investigated whether it can be applicable to other SYM models and possible to find the fuzzy sphere solution for each model. Throughout this paper, we focus on topological twisted SYM

*In private communication, they said they had also applied the same deformation method to $N = (4, 4), D = 2$ A-model SYM theory. This paper is not yet published.

[†]In this setup, models come to respect only two SUSYs regardless of the fact that the original model possesses full SUSY, say, sixteen supercharges (eight in other models).

theories and choose gauge group as $U(N)$. Indeed the gauge group is taken freely as far as gauge group is large enough to contain $SU(2)$ as a subgroup. The models investigate in this paper are $N = 4, D = 4$ SYM, $N = 8, D = 3$ SYM, $N = 4, D = 3$ SYM theories. More precisely we study A-twisted SYM theories with 16 SUSY in four and three dimensions and B-twisted SYM theories with 8 SUSY in three and two dimensions.

In particular, models have to possess enough number of scalar supercharges to proceed with the mass deformation so that these are defined in four dimensions or less. This is due to the fact that the key element of the mass deformation of superalgebra is to use the model which conserve at least two scalar supercharges with opposite ghost number. Therefore only theories defined in four dimensions or less can be capable of topological twisting and of possessing appropriate number of twisted scalar supercharges simultaneously.

This paper is organized in the following way: Sec. 2 describes a mass deformation method in detail in $N = 4, D = 4$ SYM model and how fuzzy sphere solution arises. From Sec. 3 to Sec. 5, we apply the mass deformation to each SYM model. In Sec. 6, we summarize the results and mention some points.

2 $N = 4, D = 4$ $U(N)$ SYM

Since we study twisted SYM theories as the models which can be formulated on a lattice, we summarize a topological twisting of SYM theories. $N = 4, D = 4$ $U(N)$ SYM theory is obtained through dimensional reduction of $N = 1, D = 10$ $U(N)$ SYM theory. Through the dimensional reduction the isometry decomposes as $SO(9, 1) \rightarrow SO(3, 1) \otimes SO(6)$ where $SO(3, 1)$ is Lorentz group and $SO(6)$ is the internal symmetry group. We here change the isometry $SO(3, 1)$ into $SO(4)$, because we should consider the lattice action in Euclidean spacetime. Furthermore, Euclidean symmetry and the internal R -symmetry are decomposed into subgroups as,

$$SO(4)_E \sim SU(2)_{left} \otimes SU(2)_{right}, \quad (2.1)$$

$$SO(6)_R \rightarrow SU(2)_A \otimes SU(2)_B. \quad (2.2)$$

We define A-Twist as a diagonal sum of $SU(2)_{left}$ and $SU(2)_A$ in mathematical notation namely identifying $SU(2)_{left}$ and $SU(2)_A$. Resulting symmetries are

$$SU(2)'_{left} \otimes SU(2)_{right} \otimes SU(2)_B, \quad (2.3)$$

where $SU(2)_{left} \otimes SU(2)_A \rightarrow SU(2)'_{left}$. As we will see later, a subgroup of $SU(2)_B$ symmetry behaves as ghost number symmetry. We would rename $SU(2)_B$ as $SU(2)_R$ in this section. This twisted SYM theory is called A-model or Vafa-Witten theory [24–27]. The theory possesses twisted supersymmetries whose generators are not spinors but scalars, vectors and tensors. In mass deformation technique we concentrate only on scalar supercharges.

2.1 Action

The twisted action can be off-shell invariant for two scalar supercharges by introducing two kinds of auxiliary fields, h_μ and h_A^+ , and can be rewritten into supercharge exact form. The

action is,

$$\begin{aligned}
S_0^{N=4} = \int d^4x \text{Tr} \Big(& -\mathcal{D}_\mu v \mathcal{D}_\mu v - \frac{1}{4} \mathcal{D}_\mu v_A \mathcal{D}_\mu v_A - \mathcal{D}_\mu \phi \mathcal{D}_\mu \bar{\phi} + (F_{\mu\nu}^+)^2 \\
& - i\psi_\mu (\mathcal{D}_\mu \lambda - \mathcal{D}_\nu \lambda_{\mu\nu}^+) + \frac{i}{2} \phi \{\psi_\mu, \psi_\mu\} - \frac{i}{2} \bar{\phi} \{\lambda, \lambda\} - \frac{i}{8} \bar{\phi} \{\lambda_A^+, \lambda_A^+\} \\
& - iC_\mu (\mathcal{D}_\mu \chi - \mathcal{D}_\nu \chi_{\mu\nu}^+) - \frac{i}{2} \bar{\phi} \{C_\mu, C_\mu\} + \frac{i}{2} \phi \{\chi, \chi\} + \frac{i}{8} \phi \{\chi_A^+, \chi_A^+\} \\
& - \frac{i}{4} v \{\chi_A^+, \lambda_A^+\} + \frac{i}{64} \Gamma_{ABC}^+ v_A^+ \{\chi_B^+, \lambda_C^+\} - i v_{\mu\nu}^+ \{C_\mu, \psi_\nu\} \\
& + i v \{C_\mu, \psi_\mu\} - \frac{i}{4} v_A^+ \{\chi_A^+, \lambda\} + \frac{i}{4} v_A^+ \{\chi, \lambda_A^+\} - i v \{\chi, \lambda\} \\
& - \frac{1}{4} [\phi, v_A^+] [\bar{\phi}, v_A^+] - \frac{1}{32} [v_A^+, v_B^+] [v_A^+, v_B^+] - \frac{1}{4} [v_A^+, v] [v_A^+, v] \\
& - [\phi, v] [\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 - h_\mu h_\mu - \frac{1}{4} h_A^+ h_A^+ \Big), \tag{2.4}
\end{aligned}$$

$$= \int d^4x \text{Tr} \ s \bar{s} \left(\lambda \chi + C_\mu \psi_\mu + \frac{1}{4} \chi_A^+ \lambda_A^+ + \frac{i}{96} \Gamma_{ABC}^+ v_A^+ [v_B^+, v_C^+] - 2 v_A^+ F_A^+ \right). \tag{2.5}$$

Here, we take $\mu, \nu, \dots = 1, 2, 3, 4$. $\mathcal{D}_\mu$ is covariant derivative, and is defined as $\mathcal{D}_\mu \equiv \partial_\mu - i[A_\mu, \dots]$. Then, the superscript, in λ_B^+ and χ_B^+ , denotes antiself-dual two forms and we use the following notation, $\phi_A^+ = \phi_{\mu\nu}^+ = -\frac{1}{2} \epsilon_{\mu\nu\rho\sigma} \phi_{\rho\sigma}^+$. Also, we define a self-dual part of the curvature $F_{\mu\nu}^+ \equiv \frac{1}{4} \delta_{\mu\nu, \rho\sigma}^+ F_{\rho\sigma}$, where $\delta_{\mu\nu, \rho\sigma}^+ = \delta_{\mu\rho} \delta_{\nu\sigma} - \delta_{\mu\sigma} \delta_{\nu\rho} - \epsilon_{\mu\nu\rho\sigma}$. Furthermore, Γ^{+ABC} is antisymmetric in suffixes A, B, C .

In the action, $(\phi, v, \bar{\phi})$ are scalar bosonic fields which transform as triplet under $SU(2)_R$ symmetry and A_μ is gauge boson field which transform as singlet, then (λ, χ) , (C_μ, ψ_μ) and (λ_B^+, χ_B^+) transform as doublet fermionic fields under $SU(2)_R$. Finally, auxiliary fields are singlet under $SU(2)_R$.

This action possesses the following discrete symmetry as,

$$\begin{aligned}
A_\mu &\rightarrow A_\mu, & v &\rightarrow v, \\
\phi &\rightarrow -\bar{\phi}, & \bar{\phi} &\rightarrow -\phi, & v_A^+ &\rightarrow -v_A^+, \\
\chi &\leftrightarrow \lambda, & \chi_A^+ &\leftrightarrow \lambda_A^+, & C_\mu &\leftrightarrow \psi_\mu.
\end{aligned} \tag{2.6}$$

This is the reflection of $SU(2)_R$ symmetry and the presence of s and $\bar{s}$ charges.

Classical configuration of this model is trivial solution for fermionic fields, v_A^+ and gauge fields, but the model possesses flat direction for scalar fields ϕ, v and $\bar{\phi}$ due to the potential terms which appear in the last line in (2.4).

$$\phi = v = \bar{\phi} \propto \mathbf{1}, \quad (\text{other fields}) = 0, \tag{2.7}$$

where $\mathbf{1}$ is $N \times N$ identity matrix.

We define twisted supercharges as s and $\bar{s}$ and they satisfy the following subalgebra up to gauge transformation at off-shell.

$$s^2 = \bar{s}^2 = \delta_g, \quad \{s, \bar{s}\} = \delta_g. \tag{2.8}$$

Supercharges s and $\bar{s}$ possess ghost number $+1$ and -1 respectively. Therefore s and $\bar{s}$ also transform as doublet under $SU(2)_R$. By integrating out the auxiliary fields, twisted $N = 4$ supersymmetry with sixteen supercharges are recovered. s and $\bar{s}$ transformations of each field are shown in the table in Appendix B.1.

2.2 Mass deformation of superalgebra

Next, we introduce deformation terms into supersymmetric charges as,

$$Q = s + \Delta s, \quad \bar{Q} = \bar{s} + \Delta \bar{s}, \quad (2.9)$$

and the deformation terms Δs and $\Delta \bar{s}$ satisfy the following algebras,

$$(\Delta s)^2 = (\Delta \bar{s})^2 = 0, \quad \{\Delta s, \Delta \bar{s}\} = 0. \quad (2.10)$$

We then define the following nontrivial mass deformed transformation for auxiliary fields and fermionic fields. These fermionic fields with nontrivial transformation are selected from the ones which do not appear in the transformation of auxiliary fields.

$$\Delta s(h_\mu) = MC_\mu, \quad \Delta \bar{s}(h_\mu) = M\psi_\mu, \quad (2.11)$$

$$\Delta s(h_A^+) = -M\lambda_A^+, \quad \Delta \bar{s}(h_A^+) = M\chi_A^+, \quad (2.12)$$

$$\Delta s(\lambda) = -2M\phi, \quad \Delta \bar{s}(\lambda) = 2Mv, \quad (2.13)$$

$$\Delta s(\chi) = -2Mv, \quad \Delta \bar{s}(\chi) = -2M\bar{\phi}, \quad (2.14)$$

where M is a mass parameter. By adding deformed terms, new supercharges, Q and $\bar{Q}$, satisfy the following algebra involving $SU(2)$ generators.

$$Q^2 = MJ_{++}, \quad \bar{Q}^2 = -MJ_{--}, \quad \{Q, \bar{Q}\} = -MJ_0 = -2MJ'_0, \quad (2.15)$$

where J_{++} and J_{--} are ladder operators, and eigenvalue of J_0 corresponds to the ghost number of each field. Transformations by J_{++} , J_{--} and J_0 are shown in Table 1. As seen in Table 1, fermionic fields, (C_ν, ψ_ν) , (λ_B^+, χ_B^+) and (λ, χ) , form doublet and bosons $(\phi, v, \bar{\phi})$ form triplet.

Each coefficient appearing on the right hand side of (2.11)-(2.14) is determined to keep the algebra (2.15). A relation corresponding to (2.15) frequently appear in various models in this paper. We would therefore keep the same form of the relation (2.15) in the following sections, so that the general form of resulting action and other algebra do not receive any changes by the dimensional reduction.

2.3 Deformed action

Since Q and $\bar{Q}$ also form $SU(2)$ doublet, they satisfy following algebra,

$$[J_{++}, Q] = 0, \quad [J_{--}, \bar{Q}] = 0, \quad (2.16)$$

$$[J_{++}, \bar{Q}] = Q, \quad [J_{--}, Q] = \bar{Q}, \quad (2.17)$$

$$[J_0, Q] = Q, \quad [J_0, \bar{Q}] = -\bar{Q}. \quad (2.18)$$

The action can be rewritten with new supercharges Q and $\bar{Q}$ as,

$$\mathcal{S}^{N=4} = (Q\bar{Q} - M)\mathcal{F}_0 \quad (2.19)$$

$$= \mathcal{S}_0^{N=4} + (s\Delta \bar{s} + \Delta s\bar{s} + \Delta s\Delta \bar{s} - M)\mathcal{F}_0, \quad (2.20)$$

where

$$\mathcal{F}_0 = \int d^4x \text{Tr} \left(\lambda\chi + C_\mu\psi_\mu + \frac{1}{4}\chi_A^+\lambda_A^+ + \frac{i}{96}\Gamma_{ABC}^+v_A^+[v_B^+, v_C^+] - 2v_A^+F_A^+ \right). \quad (2.21)$$

	gh#	J_{++}	J_{--}	J_0
C_ν	1	0	ψ_ν	C_ν
ψ_ν	-1	C_ν	0	$-\psi_\nu$
λ_B^+	1	0	$-\chi_B^+$	λ_B^+
χ_B^+	-1	$-\lambda_B^+$	0	$-\chi_B^+$
λ	1	0	$-\chi$	λ
χ	-1	$-\lambda$	0	$-\chi$
ϕ	2	0	$-2v$	2ϕ
v	0	$-\phi$	$\bar{\phi}$	0
$\bar{\phi}$	-2	$2v$	0	$-2\bar{\phi}$
v_B^+	0	0	0	0
A_ν	0	0	0	0
h_ν	0	0	0	0
h_B^+	0	0	0	0

Table 1: Transformation by $J_{\pm\pm}$ and J_0

The action (2.19) is invariant under Q and $\bar{Q}$ transformation as long as $\mathcal{F}_0$ has $SU(2)_R$ invariance. For example, applying the operator Q to the action, we obtain $M\bar{Q}J_{++}\mathcal{F}_0$ term from the commutation relation of J_{++} and $\bar{Q}$. Since this term must vanish, $SU(2)_R$ invariance of $\mathcal{F}_0$ is required. $SU(2)_R$ invariance of $\mathcal{F}_0$ can be easily confirmed because $SU(2)_R$ doublet fermions are paired together and the other terms are composed of $SU(2)_R$ singlet fields, and therefore each term in $\mathcal{F}_0$ possesses zero ghost number.

Regardless of the mass deformation in superalgebra, two supersymmetries are still conserved. Consequently, the action obtains several terms as,

$$\begin{aligned} \mathcal{S}^{N=4} = \mathcal{S}_0^{N=4} + 6Miv[\phi, \bar{\phi}] - 4M^2(v^2 + \phi\bar{\phi}) \\ + M \left(2\lambda\chi - 2C_\mu\psi_\mu - \frac{1}{2}\chi_A^+\lambda_A^+ - \frac{i}{96}\Gamma_{ABC}^+v_A^+[v_B^+, v_C^+] + 2v_A^+F_A^+ \right). \end{aligned} \quad (2.22)$$

Because of the second term in (2.19), $\mathcal{F}_0$ itself appears in the action as mass terms of the fermions, besides there are new combinations of scalar fields.

The potential term including v, ϕ and $\bar{\phi}$ which is analogous with the original work [1] is

$$\mathcal{V}_1 = -[\phi, v][\bar{\phi}, v] + \frac{1}{4}[\phi, \bar{\phi}]^2 + 6Miv[\phi, \bar{\phi}] - 4M^2(v^2 + \phi\bar{\phi}). \quad (2.23)$$

In the next subsection, we will focus on this scalar potential and investigate how classical configuration changes.

2.4 Fuzzy sphere solution

In the previous sections, we supposed that ϕ and $\bar{\phi}$ are real-valued independent fields, but here we would like to consider ϕ and $\bar{\phi}$ are hermitian conjugate each other, namely $\phi^\dagger = \bar{\phi}$. Although there are some ambiguities in interpretation, we would leave this problem for later discussion since it indeed does not affect degrees of freedom in this model. In addition to this reinterpretation, for simplicity, we replace anti-hermitian fields ϕ, v , and $\bar{\phi}$ into hermitian fields, $\phi \rightarrow i\phi, v \rightarrow iv$ and $\bar{\phi} \rightarrow i\bar{\phi}$.

Subsequently, the scalar potential (2.23) becomes,

$$\mathcal{V}_1 = -[\phi, v][\bar{\phi}, v] + \frac{1}{4}[\phi, \bar{\phi}]^2 + 6Mv[\phi, \bar{\phi}] + 4M^2(v^2 + \phi\bar{\phi}). \quad (2.24)$$

This potential is semi-positive definite because it can be written as sum of perfect squares as,

$$\begin{aligned}\mathcal{V}_1 &= \frac{1}{4} ([\phi, \bar{\phi}] + 4Mv)^2 + ([\phi, v] - 2M\phi) (-[\bar{\phi}, v] - 2M\bar{\phi}) \\ &= \frac{1}{4} ([\phi, \bar{\phi}] + 4Mv)^2 + ([\phi, v] - 2M\phi) ([\phi, v] - 2M\phi)^\dagger \geq 0.\end{aligned}\quad (2.25)$$

This potential has a nontrivial vacuum configuration due to the deformation. We then obtain the conditions:

$$[\phi, \bar{\phi}] + 4Mv = 0, \quad (2.26)$$

$$[\phi, v] - 2M\phi = 0, \quad (2.27)$$

$$-[\bar{\phi}, v] - 2M\bar{\phi} = 0, \quad (2.28)$$

such that the solution is analogous to Lie algebraic relations as,

$$v = -2ML_3, \quad (2.29)$$

$$\phi = \pm 2ML_+, \quad (2.30)$$

$$\bar{\phi} = \pm 2ML_-, \quad (2.31)$$

where L_3 and $L_\pm$ are N -dimensional representation of $SU(2)$ generators, which is subgroup of $U(N)$ gauge group. We find fuzzy sphere solution for four dimensional $N = 4$ $U(N)$ SYM theory, and the flat direction is resolved consequently.

2.5 Adding $\Delta\mathcal{F}$ terms

In mass deformation process, the action has further possibility to add extra terms $\Delta\mathcal{F}$ as long as it is $SU(2)_R$ invariant.

$$\mathcal{S} = (Q\bar{Q} - M)(\mathcal{F}_0 + \Delta\mathcal{F}), \quad (2.32)$$

where

$$\mathcal{F}_0 = \int d^4x \text{Tr} \left(\lambda\chi + C_\mu\psi_\mu + \frac{1}{4}\chi_A^+\lambda_A^+ + \frac{i}{96}\Gamma_{ABC}^+v_A^+[v_B^+, v_C^+] - 2v_A^+F_A^+ \right), \quad (2.33)$$

$$\Delta\mathcal{F} = \tilde{m}(v_B^+v_B^+). \quad (2.34)$$

In this particular model, the previous deformation does not produce mass term of v_B^+ , so it is reasonable to include mass term of v_B^+ , where $\tilde{m}$ is some arbitrary parameters with mass dimension one.

The action becomes,

$$\begin{aligned}\mathcal{S} &= \mathcal{S}_0^{N=4} + (6Miv[\phi, \bar{\phi}] - (2M)^2(v^2 + \phi\bar{\phi})) \\ &\quad + M \left(2\lambda\chi - 2C^\mu\psi_\mu - \frac{1}{2}\chi_A^+\lambda_A^+ - \frac{i}{96}\Gamma^{+ABC}v_A^+[v_B^+, v_C^+] + 2v_A^+F_A^+ \right) \\ &\quad + \tilde{m} \left(\frac{i}{32}\Gamma_{BCD}^+v_B^+[v_C^+, v_D^+] - iv_B^+[v_B^+, v] - 2v_B^+F_B^+ + v_B^+h_B^+ \right) + \frac{1}{2}\tilde{m}\chi_B^+\lambda_B^+ - \tilde{m}Mv_B^+v_B^+.\end{aligned}\quad (2.35)$$

The third line includes new terms due to $\Delta\mathcal{F}$.

After eliminating auxiliary fields by solving equation of motion, we can also focus on the potential terms with respect to v_B^+ .

$$\mathcal{V}_2 = -\frac{1}{32}[v_A^+, v_B^+][v_A^+, v_B^+] + \left(\frac{i}{32}\tilde{m} - \frac{i}{96}M\right)\Gamma_{ABC}^+ v_A^+[v_B^+, v_C^+] + (\tilde{m}^2 - \tilde{m}M)v_B^+ v_B^+. \quad (2.36)$$

By inserting appropriate parameter $\tilde{m} = -\frac{M}{3}$, the action also becomes semi-positive definite as follows.

$$\mathcal{V}_2 = \frac{1}{32} \left([v_A^+, v_B^+] + \frac{i}{3}M\Gamma_{ABC}^+ v_C^+ \right)^\dagger \left([v_A^+, v_B^+] + \frac{i}{3}M\Gamma_{ABC}^+ v_C^+ \right), \quad (2.37)$$

where a formula, $\Gamma_{ABC}^+ \Gamma_{ABD}^+ = 2 \cdot 4^2 \delta_{CD}^+$, is used.

Since v_B^+ is self-dual fields, all the v_B^+ are rewritten into v_{i4}^+ through the relation $v_B^+ = v_{\mu\nu}^+ = -\frac{1}{2}\epsilon_{\mu\nu\rho\sigma}v_{\rho\sigma}^+$. In the end, we can find another fuzzy sphere solution with the formula $\Gamma_{i4,j4,k4}^+ = -\epsilon_{ijk}$,

$$v_{i4}^+ = \frac{4M}{3}L_i \quad (2.38)$$

where L_i is N -dimensional representation of $SU(2)$.

Before ending this section it should be noted that choices of fuzzy sphere solution leads to a problem concerning to positivity of kinetic terms. In section 2.4, we redefined some fields from anti-hermitian into hermitian ones. Since one fuzzy sphere solution is only related to scalar fields v, ϕ and $\bar{\phi}$ and the other solution is only to the fields v_A^+ , we independently select hermiticity of these fields to obtain the fuzzy sphere solution. However the kinetic term of v_A^+ is not positive definite. Indeed the negative sign of kinetic term causes the instability of the action itself, and therefore it would be not appropriate to adopt two solutions simultaneously.

3 $N = 8, D = 3$ $U(N)$ SYM

By dimensional reduction of $N = 4, D = 4$ SYM theory, new two scalar supercharges, s_4 and $\bar{s}_4$, appear so that it is possible to consider new pairs of supercharges to perform mass deformation of superalgebra as we have shown in $N = 4, D = 4$ action.

To analyze isometry of this model, we firstly consider naive dimensional reduction of ten-dimensional $N = 1$ SYM theory down to $D = 3$. Global symmetry is,

$$SU(2)_E \otimes SO(7). \quad (3.1)$$

This internal symmetry can be further decomposed as,

$$SO(7) \rightarrow SU(2)_R \otimes SU(2)_1 \otimes SU(2)_2. \quad (3.2)$$

To obtain desired twisted model, we define a twist as a diagonal sum of $SU(2)_E$ and $SU(2)_2$, so that the model possesses isometry of,

$$SU(2)'_E \otimes SU(2)_R \otimes SU(2)_1, \quad (3.3)$$

where $SU(2)_E \otimes SU(2)_2 \rightarrow SU(2)'_E$. This internal symmetry contains one extra $SU(2)$ symmetry compared to $N = 4, D = 4$ SYM model. It suggests that there would be other topological supercharges which correspond to this $SU(2)$ symmetry, indeed it corresponds to two scalar supercharges, s_4 and $\bar{s}_4$, as noted in the beginning of this section. In this model, we expect $U(1)$ subgroup of either $SU(2)_R$ or $SU(2)_1$ corresponds to ghost number symmetry.

3.1 Action

Three dimensional $N = 8$ twisted $U(N)$ SYM action is,

$$\begin{aligned}
\mathcal{S}_0^{N=8} = \int d^3x \text{Tr} \Big[& -\mathcal{D}_i v \mathcal{D}_i v - \mathcal{D}_i v_j \mathcal{D}_i v_j + [A_4, v_j][A_4, v_j] \\
& - \mathcal{D}_i \phi \mathcal{D}_i \bar{\phi} + \mathcal{D}_i A_4 \mathcal{D}_i A_4 + \frac{1}{2} F_{ij} F_{ij} \\
& - i\psi_i (\mathcal{D}_i \lambda + \epsilon_{ijk} \mathcal{D}_j \lambda_k + i[A_4, \lambda_i]) - i\psi_4 (\mathcal{D}_i \lambda_i - i[A_4, \lambda]) \\
& - iC_i (\mathcal{D}_i \chi + \epsilon_{ijk} \mathcal{D}_j \chi_k + i[A_4, \chi_i]) - iC_4 (\mathcal{D}_i \chi_i - i[A_4, \chi]) \\
& + \frac{i}{2} \phi \{\psi_i, \psi_i\} + \frac{i}{2} \phi \{\psi_4, \psi_4\} - \frac{i}{2} \bar{\phi} \{\lambda, \lambda\} - \frac{i}{2} \bar{\phi} \{\lambda_i, \lambda_i\} \\
& - \frac{i}{2} \bar{\phi} \{C_i, C_i\} - \frac{i}{2} \bar{\phi} \{C_4, C_4\} + \frac{i}{2} \phi \{\chi, \chi\} + \frac{i}{2} \phi \{\chi_i, \chi_i\} \\
& - iv \{\chi_i, \lambda_i\} - i\epsilon_{ijk} v_i \{\chi_j, \lambda_k\} + i\epsilon_{ijk} v_i \{C_j, \psi_k\} - iv_i \{C_i, \psi_4\} + iv_i \{C_4, \psi_i\} \\
& + iv \{C_i, \psi_i\} + iv \{C_4, \psi_4\} - iv_i \{\chi_i, \lambda\} + iv_i \{\lambda_i, \chi\} - iv \{\chi, \lambda\} \\
& + [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 \\
& - [\phi, v_i][\bar{\phi}, v_i] - \frac{1}{2} [v_i, v_j][v_i, v_j] - [v_i, v][v_i, v] - h_i h_i - h_4 h_4 - h_{i4}^+ h_{i4}^+ \Big],
\end{aligned} \tag{3.4}$$

where Lorentz indices run, $i, j = 1, 2, 3$. Most of the situations correspond to $N = 4, D = 4$ case, while there appear new characters. New scalar fields appear from the four dimensional vector fields, and new vector fields are induced by the selfdual two-forms in $N = 4, D = 4$ action, such as λ_A^+ and χ_A^+ , because of the self-dual condition: for example, $\lambda_{ij}^+ = -\frac{1}{2} \epsilon_{ijk4} \lambda_{k4}^+$.

In this models, bosonic scalar fields $(\phi, v, \bar{\phi}, A_4)$ possess the ghost number as $(+2, 0, -2, 0)$ respectively. Fermionic scalar fields $(\lambda, \chi, C_4, \psi_4)$ possess ghost number as $(+1, -1, +1, -1)$ respectively, and the ghost number of vector fields $(C_i, \psi_i, \lambda_i, \chi_i)$ are assigned as $(+1, -1, +1, -1)$ respectively.

3.2 Combinations of supercharges

In this subsection, we consider pairings of scalar supercharges. As in the previous section, the action is SUSY invariant with two scalar supercharges even after the mass deformation. Supertransformations of scalar supercharges are shown in Appendix B.2. Since ghost number of $(s, \bar{s}, s_4, \bar{s}_4)$ is assigned as $(+1, -1, -1, +1)$ respectively, it is possible to consider four different ways of choosing two supercharges among them. In the following, the action can be written as supercharge exact form with these possible combination of supercharges.

$$\mathcal{S}_0^{N=8} = s\bar{s} \int d^3x \text{Tr} \left(\lambda\chi + C_4\psi_4 + C_i\psi_i + \chi_i\lambda_i - \frac{2i}{3} \epsilon_{ijk} v_i [v_j, v_k] - 4v_i \mathcal{D}_i A_4 + 2\epsilon_{ijk} v_i F_{jk} \right) \tag{3.5}$$

$$= \bar{s}_4 s_4 \int d^3x \text{Tr} \left(\lambda\chi + C_4\psi_4 - C_i\psi_i - \chi_i\lambda_i - \frac{2i}{3} \epsilon_{ijk} v_i [v_j, v_k] + 4v_i \mathcal{D}_i A_4 + 2\epsilon_{ijk} v_i F_{jk} \right) \tag{3.6}$$

$$= ss_4 \int d^3x \text{Tr} (\psi_4 \lambda + C_4 \chi + \lambda_i \psi_i + C_i \chi_i + 2i\mathcal{T}_{CS} - 4iv_i \mathcal{D}_i v + 2i\epsilon_{ijk} v_i \mathcal{D}_j v_k) \tag{3.7}$$

$$= \bar{s}_4 \bar{s} \int d^3x \text{Tr} (\lambda\psi_4 - C_4 \chi + \lambda_i \psi_i + C_i \chi_i - 2i\mathcal{T}_{CS} - 4iv_i \mathcal{D}_i v - 2i\epsilon_{ijk} v_i \mathcal{D}_j v_k) \tag{3.8}$$

where, in the last two lines, $\mathcal{T}_{CS}$ denotes Chern-Simons term.

$$\mathcal{T}_{CS} = \epsilon_{ijk} \left(A_i \partial_j A_k - \frac{i}{3} A_i [A_j, A_k] \right). \quad (3.9)$$

Thus, we can perform the mass deformation about these actions. There are some comments about the above combinations: the first combinations is expected results, since it can be derived straightforwardly from dimensional reduction of $N = 4, D = 4$ model (2.5). The second model is composed with $\mu = 4$ component of vector supercharges, s_4 and $\bar{s}_4$. On the other hand, the last two are totally new results. One remarkable point is Chern-Simons term appears in the integrands, so this term would play a role to give a mass to the gauge field [28] in a deformed action in the end.

However, it is not necessary to study all these deformations. The action (3.4) is invariant under some discrete symmetries, which may be the consequence of the fact that the action possesses two $SU(2)$ symmetries in internal symmetry. There are indeed three kinds of discrete symmetries and its details are shown in Appendix A.1. With this useful tool, we can show equivalence between the deformation of $(s, \bar{s})$ and $(s_4, \bar{s}_4)$ and between (s, s_4) and $(\bar{s}, \bar{s}_4)$ [‡].

3.3 Mass deformation of superalgebra in $(s, \bar{s})$ combination

We first consider $(s, \bar{s})$ pair because the combination of $(s_4, \bar{s}_4)$ is equivalent with it through discrete symmetries.

Let define the following deformation of algebra,

$$Q = s + \Delta s, \quad \bar{Q} = \bar{s} + \Delta \bar{s}, \quad (3.10)$$

and impose the following algebra in the same manner as the previous section:

$$Q^2 = MJ_{++}, \quad \bar{Q}^2 = -MJ_{--}, \quad \{Q, \bar{Q}\} = -MJ_0. \quad (3.11)$$

We can define the following nontrivial transformations of deformed terms to satisfy the above algebra.

$$\Delta s(h_4) = MC_4, \quad \Delta \bar{s}(h_4) = M\psi_4, \quad (3.12)$$

$$\Delta s(h_i) = MC_i, \quad \Delta \bar{s}(h_i) = M\psi_i, \quad (3.13)$$

$$\Delta s(h_{i4}^+) = -M\lambda_i, \quad \Delta \bar{s}(h_{i4}^+) = M\chi_i, \quad (3.14)$$

$$\Delta s(\lambda) = -2M\phi, \quad \Delta \bar{s}(\lambda) = 2Mv, \quad (3.15)$$

$$\Delta s(\chi) = -2Mv, \quad \Delta \bar{s}(\chi) = -2M\bar{\phi}. \quad (3.16)$$

This deformed algebra can also be directly derived by the dimensional reduction of $N = 4, D = 4$ model. The final result is indeed equivalent to the one which is obtained from dimensional reduction.

[‡]One of discrete symmetries makes the model invariant, and other symmetries changes the models to the corresponding partner after the deformation. But these discrete symmetries do not completely show that all models are equivalent.

However, there seems to be nontrivial "discrete symmetry" which connects remaining two, for example (3.5) and (3.7). This *symmetry* corresponds to taking complex conjugation of complexified covariant derivatives, namely $\mathcal{D}_\mu^\pm \equiv \partial_\mu - i(A_\mu \pm i v_\mu)$. Interchanging of $\mathcal{D}_\mu^\pm$ results in the action invariant, and it also makes the relation among supercharge as, $s \rightarrow s_4 \rightarrow \bar{s}_4 \rightarrow \bar{s} \rightarrow -s$. This might be connected with $SO(4)$ symmetry since R-symmetry is $SU(2) \otimes SU(2)$, which is isomorphic to $SO(4)$. But we should carefully examine that there seems to be no reasonable representation of the symmetry just like interchanging of fields, and the origin of the rotational symmetry of supercharges has not been obvious yet.

The mass deformed action becomes,

$$\mathcal{S}_{s,\bar{s}} = (Q\bar{Q} - M)\mathcal{F}_{s,\bar{s}}, \quad (3.17)$$

where

$$\mathcal{F}_{s,\bar{s}} = \int d^3x \text{Tr} \left(\lambda\chi + C_4\psi_4 + C_i\psi_i + \chi_i\lambda_i - \frac{2i}{3}\epsilon_{ijk}v_i[v_j, v_k] - 4v_i\mathcal{D}_iA_4 + 2\epsilon_{ijk}v_iF_{jk} \right). \quad (3.18)$$

The explicit form of the action is,

$$\begin{aligned} \mathcal{S}_{s,\bar{s}} = \mathcal{S}_0^{N=8} + 6iMv[\phi, \bar{\phi}] - 4M^2(v^2 + \phi\bar{\phi}) \\ - 2M \left(-\lambda\chi + C_4\psi_4 + C_i\psi_i + \chi_i\lambda_i - \frac{i}{3}\epsilon_{ijk}v_i[v_j, v_k] - 2v_i\mathcal{D}_iA_4 + \epsilon_{ijk}v_iF_{jk} \right). \end{aligned} \quad (3.19)$$

The potential terms containing scalar fields are,

$$\mathcal{V}_{s,\bar{s}} = [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4}[\phi, \bar{\phi}]^2 - 4M^2(v^2 + \phi\bar{\phi}) + 6iMv[\phi, \bar{\phi}] \quad (3.20)$$

In this case, we would focus on $(\phi, v, \bar{\phi})$ which are triplet under $SU(2)_R$. Similarly we can find the fuzzy sphere solution, since in the following potential the last two terms are the same as the potential (2.25).

$$\mathcal{V}_{s,\bar{s}} = -[A_4, v][A_4, v] - [A_4, \phi][A_4, \bar{\phi}] + \frac{1}{4}([\phi, \bar{\phi}] + 4Mv)^2 + ([\phi, v] - 2M\phi)([\phi, v] - 2M\phi)^\dagger, \quad (3.21)$$

where redefinition of the fields $(\phi, v, \bar{\phi})$ are already done as in the previous section: $\Phi \rightarrow i\Phi, \Phi \in \{\phi, v, \bar{\phi}\}$. In this potential, $(\phi, v, \bar{\phi})$ contribute to the fuzzy sphere configuration, while classical configuration of A_4 is trivial.

In the following subsections, we only consider (s, s_4) combination in the deformation. As mentioned above, $(s, \bar{s})$ pair is equivalent to $(s_4, \bar{s}_4)$, and (s, s_4) pair is also equivalent to $(\bar{s}, \bar{s}_4)$ through the discrete symmetries. A procedure of the deformation of (s, s_4) combination is the same as the previous one, so only the results are shown.

3.4 Mass deformation of superalgebra in (s, s_4) combination

In this subsection, we focus on the cases where Chern-Simons term appear in the integrand. Similar to the previous procedure, let define the deformed supercharges as,

$$Q^+ = s + \Delta s, \quad Q^- = s_4 + \Delta s_4. \quad (3.22)$$

We then impose the following algebra,

$$(Q^+)^2 = MJ_{++}, \quad (Q^-)^2 = -MJ_{--}, \quad \{Q^+, Q^-\} = -MJ_0, \quad (3.23)$$

so that we define the transformations as,

$$\Delta s(L) = -M\lambda, \quad \Delta s_4(L) = M\psi_4, \quad (3.24)$$

$$\Delta s(N_i) = M\lambda_i, \quad \Delta s_4(N_i) = M\psi_i, \quad (3.25)$$

$$\Delta s(M_i) = MC_i, \quad \Delta s_4(M_i) = M\chi_i, \quad (3.26)$$

$$\Delta s(C_4) = -2M\phi, \quad \Delta s_4(C_4) = 2iMA_4, \quad (3.27)$$

$$\Delta s(\chi) = -2iMA_4, \quad \Delta s_4(\chi) = -2M\bar{\phi}. \quad (3.28)$$

where N_i, M_i and L are auxiliary fields. In this case, A_4 takes part in the transformation instead of v .

The deformed action takes the form,

$$\mathcal{S}_{s,s_4} = (Q^+ Q^- - M) \mathcal{F}_{s,s_4}, \quad (3.29)$$

where,

$$\mathcal{F}_{s,s_4} = \int d^3x \text{Tr} (C_4 \chi + \psi_4 \lambda + \lambda_i \psi_i + C_i \chi_i + 2i \mathcal{T}_{CS} - 4iv_i \mathcal{D}_i v + 2i \epsilon_{ijk} v_i \mathcal{D}_j v_k), \quad (3.30)$$

$$\mathcal{T}_{CS} = \epsilon_{ijk} \left(A_i \partial_j A_k - \frac{i}{3} A_i [A_j, A_k] \right). \quad (3.31)$$

The final result of the action is,

$$\begin{aligned} \mathcal{S}_{s,s_4} = \mathcal{S}_0^{N=8} - 6M A_4 [\phi, \bar{\phi}] + 4M^2 (A_4^2 - \phi \bar{\phi}) \\ - 2M (-C_4 \chi + \psi_4 \lambda + \lambda_i \psi_i + C_i \chi_i + i \mathcal{T}_{CS} - 2iv_i \mathcal{D}_i v + i \epsilon_{ijk} v_i \mathcal{D}_j v_k). \end{aligned} \quad (3.32)$$

The potential terms containing scalar fields are,

$$\mathcal{V}_{s,s_4} = [A_4, v][A_4, v] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} ([\phi, \bar{\phi}] + 4M A_4)^2 + ([\phi, A_4] - 2M \phi)([\phi, A_4] - 2M \phi)^\dagger. \quad (3.33)$$

The fields ϕ and $\bar{\phi}$ are already redefined as hermitian. In this model, $(\phi, A_4, \bar{\phi})$ transform as triplet under $SU(2)_R$ and contribute to the fuzzy sphere solution. On the other hand, the classical solution of v is trivial.

The followings are the potential terms including scalar fields in each combination of supercharges:

$$\begin{aligned} \mathcal{V}_{s,\bar{s}} &= [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 + 6Miv[\phi, \bar{\phi}] - 4M^2(v^2 + \phi \bar{\phi}), \\ \mathcal{V}_{\bar{s}_4, s_4} &= [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 - 6Miv[\phi, \bar{\phi}] - 4M^2(v^2 + \phi \bar{\phi}), \\ \mathcal{V}_{s, s_4} &= [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 - 6M A_4 [\phi, \bar{\phi}] + 4M^2(A_4^2 - \phi \bar{\phi}), \\ \mathcal{V}_{\bar{s}_4, \bar{s}} &= [A_4, v][A_4, v] + [A_4, \phi][A_4, \bar{\phi}] - [\phi, v][\bar{\phi}, v] + \frac{1}{4} [\phi, \bar{\phi}]^2 - 6M A_4 [\phi, \bar{\phi}] + 4M^2(A_4^2 - \phi \bar{\phi}). \end{aligned}$$

Subscripts of $\mathcal{V}$ denote the combination of the supercharges. As seen in the list of the potential terms, these potentials lead to the same solution except that there is an interchange of v and A_4 . There are signs difference in some terms but they are absorbed in a redefinition of the parameter M . It is then concluded that the models considered in this section have the same fuzzy sphere solution consequently.

4 $N = 4, D = 3$ $U(N)$ SYM

From here mass deformation procedure for three dimensional $N = 4$ SYM is treated. Since $N = 2, D = 4$ SYM model does not possess two scalar supercharges, it is not possible to perform the deformation. We then consider dimensional reduction of $N = 2, D = 4$ SYM model down to three dimensions so that one scalar supercharge arises from vector supercharge.

Three dimensional $N = 4$ SYM model is obtained through dimensional reduction either directly from $N = 1, D = 6$ SYM or via $N = 2, D = 4$ SYM. This theory is invariant under eight supercharges, namely half SUSY of the previous sections. And the isometry is given as, from dimensional reduction of $N = 1, D = 6$ theory,

$$\begin{aligned} SO(6)_E \otimes SU(2)_R &\rightarrow SO(3)_E \otimes SO(3)_N \otimes SU(2)_R \\ &\sim SU(2)_E \otimes (SU(2)_N \otimes SU(2)_R) \end{aligned} \quad (4.1)$$

where $SU(2)_R$ is R-symmetry in $N = 1, D = 6$ SYM.

Since there are two $SU(2)$ internal symmetries, there are two possible twists, in other words diagonal sum of two subgroups, choosing either $SU(2)_R$ or $SU(2)_N$ with $SU(2)_E$.

- A-model (or so called Super-BF [29–31]): Twist $SU(2)_E$ with $SU(2)_R$

$$SU(2)_E \otimes SU(2)_N \otimes SU(2)_R \rightarrow SU(2)_{E'} \otimes SU(2)_N \quad (4.2)$$

Field content consists of $SU(2)_N$ doublet scalar fermions, $SU(2)_N$ doublet vector fermions and $SU(2)_N$ triplet scalar bosons with appropriate numbers of auxiliary fields to adjust the degree of freedom. There are two scalar supercharges which transform as doublets of $SU(2)_N$.

- B-model [32]: Twist $SU(2)_E$ with $SU(2)_N$

$$SU(2)_E \otimes SU(2)_N \otimes SU(2)_R \rightarrow SU(2)_{E'} \otimes SU(2)_R \quad (4.3)$$

This model contains of $SU(2)_R$ doublet scalar fermions, $SU(2)_R$ doublet vector fermions and $SU(2)_R$ singlet vector bosons with appropriate numbers of auxiliary fields. There are two scalar supercharges which transform as doublets of $SU(2)_R$.

In each twisting the model possesses one $SU(2)$ symmetry (either $SU(2)_R$ or $SU(2)_N$) as internal symmetry, so that the models commonly possess two scalar supercharges, transforming as doublet under its $SU(2)$ symmetry. Therefore, it is expected to perform mass deformation and find a fuzzy sphere solution in each model.

4.1 A-model

The action is shown as,

$$\begin{aligned} \mathcal{S}_0^A = \int d^3x \text{Tr} &\left(\frac{1}{2} F_{\mu\nu} F_{\mu\nu} + \mathcal{D}_\mu N \mathcal{D}_\mu N + \mathcal{D}_\mu \phi \mathcal{D}_\mu \bar{\phi} - i\psi_\mu \mathcal{D}_\mu \psi - i\chi_\mu \mathcal{D}_\mu \chi - i\epsilon_{\mu\nu\rho} \psi_\mu \mathcal{D}_\nu \chi_\rho \right. \\ &- H_\mu H_\mu + \frac{i}{2} \bar{\phi} \{ \psi_\mu, \psi_\mu \} + \frac{i}{2} \bar{\phi} \{ \chi, \chi \} + \frac{i}{2} \phi \{ \chi_\mu, \chi_\mu \} + \frac{i}{2} \phi \{ \psi, \psi \} \\ &\left. + N \{ \psi, \chi \} - N \{ \psi_\mu, \chi_\mu \} - [N, \bar{\phi}] [N, \phi] + \frac{1}{4} [\phi, \bar{\phi}]^2 \right) \end{aligned} \quad (4.4)$$

$$= \int d^3x \text{Tr} \, s\bar{s} \left(\chi\psi + \psi_\mu \chi_\mu + 2i\epsilon_{\mu\nu\rho} \left(A_\mu \partial_\nu A_\rho - \frac{i}{3} A_\mu [A_\nu, A_\rho] \right) \right), \quad (4.5)$$

where $(\phi, N, \bar{\phi})$ are bosonic scalar fields transforming as triplet under $SU(2)_R$, besides A_μ and H_μ are gauge fields and auxiliary fields respectively, both transforming as singlet under $SU(2)_R$. While (χ, ψ) and (ψ_μ, χ_μ) are fermionic scalar and vector fields respectively, transforming as doublet under $SU(2)_R$.

This model possesses two scalar supercharges, out of eight supercharges, $s, \bar{s}, s_\mu$ and $\bar{s}_\mu$. Similar to the previous section, new supercharges are defined as,

$$Q = s + \Delta s, \quad \bar{Q} = \bar{s} + \Delta \bar{s}. \quad (4.6)$$

Supercharges and deformation terms satisfy the following algebra;

$$\begin{aligned} s^2 = \bar{s}^2 = \{s, \bar{s}\} &= \delta_g, \\ \Delta s^2 = \Delta \bar{s}^2 = \{\Delta s, \Delta \bar{s}\} &= 0. \end{aligned} \quad (4.7)$$

In this model, supercharges are nilpotent up to gauge transformation.

Let define the transformation of deformation terms as,

$$\Delta s(H_\mu) = -M\psi_\mu \quad \Delta \bar{s}(H_\mu) = -M\chi_\mu, \quad (4.8)$$

$$\Delta s(\chi) = -2M\phi \quad \Delta \bar{s}(\chi) = 2iMN, \quad (4.9)$$

$$\Delta s(\psi) = -2iMN \quad \Delta \bar{s}(\psi) = 2M\bar{\phi}. \quad (4.10)$$

The action becomes, as far as Q and $\bar{Q}$ satisfy the algebra consistently used in the previous sections,

$$\mathcal{S} = (Q\bar{Q} - M)\mathcal{F}_0^A \quad (4.11)$$

$$\begin{aligned} &= \mathcal{S}_0^A + 6MN[\phi, \bar{\phi}] + 4M^2(N^2 + \phi\bar{\phi}) \\ &\quad - 2M \left(-\chi\psi + \psi_\mu\chi_\mu + i\epsilon_{\mu\nu\rho} \left(A_\mu\partial_\nu A_\rho - \frac{i}{3}A_\mu[A_\nu, A_\rho] \right) \right), \end{aligned} \quad (4.12)$$

where

$$\mathcal{F}_0^A = \chi\psi + \psi_\mu\chi_\mu + 2i\epsilon_{\mu\nu\rho} \left(A_\mu\partial_\nu A_\rho - \frac{i}{3}A_\mu[A_\nu, A_\rho] \right). \quad (4.13)$$

The potential terms including scalar fields are,

$$\mathcal{V}_A = -[N, \bar{\phi}][N, \phi] + \frac{1}{4}[\phi, \bar{\phi}]^2 + 6MN[\phi, \bar{\phi}] + 4M^2(N^2 + \phi\bar{\phi}). \quad (4.14)$$

However the notation is different to $N = 4, D = 4$ model, this potential is the same as that in $N = 4, D = 4$ model and $N = 8, D = 3$ model, and the fuzzy sphere solution is also obtained.

4.2 B-model

Next, we consider another twisted model (B-twisted model). $N = 2, D = 4$ B-model possesses two scalar supercharges, but these supercharges have the same ghost number. Thus exact part of the action corresponding to $\mathcal{F}_0$ possesses nonzero ghost number. However $D = 3$ B-twisted model possesses two scalar supercharges whose ghost number is opposite. Therefore it is possible to perform the mass deformation.

The B-model action is,

$$\begin{aligned} \mathcal{S}_0^B &= \int d^3x \text{Tr} \left(\frac{1}{4}F_{\mu\nu}F_{\mu\nu} + \frac{1}{2}\mathcal{D}_\mu V_\nu \mathcal{D}_\mu V_\nu - 2i\lambda_\mu \mathcal{D}_\mu \tilde{\lambda} - 2i\tilde{\lambda}_\mu \mathcal{D}_\mu \lambda \right. \\ &\quad - 2i\epsilon_{\mu\nu\rho}\lambda_\mu \mathcal{D}_\nu \tilde{\lambda}_\sigma + 2iV_\mu \{\lambda_\mu, \tilde{\lambda}\} + 2iV_\mu \{\tilde{\lambda}_\mu, \lambda\} \\ &\quad \left. + 2i\epsilon_{\mu\nu\rho}V_\mu \{\lambda_\nu, \tilde{\lambda}_\rho\} - \frac{1}{4}[V_\mu, V_\nu]^2 + 2K^2 + 2G\tilde{G} \right), \quad (4.15) \\ &= \int d^3x \text{Tr} \, s\bar{s} \left(-2\lambda\tilde{\lambda} - 2\lambda_\mu \tilde{\lambda}_\mu + \epsilon_{\mu\nu\rho}F_{\mu\nu}V_\rho - i\epsilon_{\mu\nu\rho} \left(A_\mu\partial_\nu A_\rho - \frac{1}{3}A_\mu[A_\nu, A_\rho] \right) \right). \quad (4.16) \end{aligned}$$

This model includes gauge field A_μ , bosonic vector field V_μ , fermionic fields $\lambda, \tilde{\lambda}, \lambda_\mu$ and $\tilde{\lambda}_\mu$ and three scalar auxiliary fields $G, \tilde{G}$ and K . These fields are also classified to $SU(2)_N$ doublet scalar fermions, $SU(2)_N$ doublet vector fermions, $SU(2)_N$ singlet vector bosons. What is special in this model is that auxiliary fields possess ghost numbers, and they transform as triplet under $SU(2)_N$.

Classical configuration of this model is trivial and there exists flat direction in the solution of $-\frac{1}{4}[V_\mu, V_\nu]^2$ term, so that we should apply mass deformation method. A special feature of this model is that there are vector bosons and no scalar bosonic fields, which is slightly different to the previous one, especially in treatment of fuzzy sphere solution.

Supercharges satisfy the following algebra,

$$s^2 = \bar{s}^2 = \{s, \bar{s}\} = 0. \quad (4.17)$$

The action and fields are strictly nilpotent without any gauge transformation, and transformation is shown in Appendix B.3. Due to internal symmetry of $SU(2)_N$, these scalar supercharges s and $\bar{s}$ possess ghost number $+1$ and -1 respectively.

Then, we define the deformed supercharges as,

$$Q = s + \Delta s, \quad \bar{Q} = \bar{s} + \Delta \bar{s}, \quad (4.18)$$

and consider the deformed transformation respecting the following algebra.

$$Q^2 = MJ_{++}, \quad \bar{Q} = -MJ_{--}, \quad \{Q, \bar{Q}\} = -MJ_0. \quad (4.19)$$

As a consequence, we define the following transformations.

$$\Delta s(G) = 0, \quad \Delta \bar{s}(G) = 2M\lambda, \quad (4.20)$$

$$\Delta s(\tilde{G}) = -2M\tilde{\lambda}, \quad \Delta \bar{s}(\tilde{G}) = 0, \quad (4.21)$$

$$\Delta s(K) = M\lambda, \quad \Delta \bar{s}(K) = M\tilde{\lambda}, \quad (4.22)$$

$$\Delta s(\lambda_\mu) = 0, \quad \Delta \bar{s}(\lambda_\mu) = -MV_\mu, \quad (4.23)$$

$$\Delta s(\tilde{\lambda}_\mu) = MV_\mu, \quad \Delta \bar{s}(\tilde{\lambda}_\mu) = 0. \quad (4.24)$$

The transformation of auxiliary fields should vanish because auxiliary fields G and $\tilde{G}$ possess ghost number $+2$ and -2 and there is no field possessing ghost number ± 3 in this model. Even though there are some differences, the mass deformation can be performed.

The deformed action is,

$$\mathcal{S} = (Q\bar{Q} - M) \mathcal{F}_0, \quad (4.25)$$

where,

$$\mathcal{F}_0 = -2\lambda\tilde{\lambda} - 2\lambda_\mu\tilde{\lambda}_\mu + \epsilon_{\mu\nu\rho}F_{\mu\nu}V_\rho - i\epsilon_{\mu\nu\rho}\left(A_\mu\partial_\nu A_\rho - \frac{1}{3}A_\mu[A_\nu, A_\rho]\right). \quad (4.26)$$

In the end, we obtain,

$$\begin{aligned} \mathcal{S} = & \mathcal{S}_0^B + 2M^2V_\mu V_\mu + iM\epsilon_{\mu\nu\rho}V_\mu[V_\nu, V_\rho] \\ & - 4M\lambda\tilde{\lambda} + 4M\lambda_\mu\tilde{\lambda}_\mu + iM\epsilon_{\mu\nu\rho}\left(A_\mu\partial_\nu A_\rho - \frac{1}{3}A_\mu[A_\nu, A_\rho]\right) \end{aligned} \quad (4.27)$$

Potential terms are,

$$\mathcal{V}_B = -\frac{1}{4}[V_\mu, V_\nu][V_\mu, V_\nu] + 2M^2V_\mu V_\mu + iM\epsilon_{\mu\nu\rho}V_\mu[V_\nu, V_\rho]. \quad (4.28)$$

This is totally different to the previous potentials. We redefine the mass parameter as $M \rightarrow iM$. Then, the potential is further written as,

$$\begin{aligned}\mathcal{V}_B &= \left(\frac{1}{2}[V_\mu, V_\nu] - iM\epsilon_{\mu\nu\rho}V_\rho \right) \left(-\frac{1}{2}[V_\mu, V_\nu] + iM\epsilon_{\mu\nu\sigma}V_\sigma \right) \\ &= \left(\frac{1}{2}[V_\mu, V_\nu] - iM\epsilon_{\mu\nu\rho}V_\rho \right) \left(\frac{1}{2}[V_\mu, V_\nu] - iM\epsilon_{\mu\nu\rho}V_\rho \right)^\dagger \geq 0.\end{aligned}\quad (4.29)$$

Hence a classical configuration of the potential satisfies,

$$\frac{1}{2}[V_\mu, V_\nu] - iM\epsilon_{\mu\nu\rho}V_\rho = 0. \quad (4.30)$$

This takes the same form as the commutation relation of Lie algebra, so solution can be chosen as,

$$V_\mu = 2ML_\mu, \quad (4.31)$$

where L_μ ($\mu = 1, 2, 3$) is N -dimensional representation of $SU(2)$, which is a subgroup of gauge group $U(N)$.

In contrast to A-models in the previous sections, bosonic vector fields form fuzzy sphere solution in this model and they are proportional to generators of $SU(2)$.

5 $N = 4, D = 2$ $U(N)$ SYM

We can obtain two dimensional models by dimensional reduction of three dimensional models. There appear two new scalar supercharges from vector supercharges in three dimensional model, so that it is similarly possible to consider several combinations of supercharges.

Before continuing to perform mass deformation, we would evaluate the isometry of two dimensional model. The isometry of the $N = 4, D = 2$ model can be obtained dimensional reduction from $N = 1, D = 6$ model. In this section we consider the following decomposition of the isometry:

$$SO(6) \times SU(2)_R \rightarrow SO(2)_E \times (SU(2)_N \times SU(2)_R), \quad (5.1)$$

where $SU(2)_R$ is an internal symmetry in $N = 1, D = 6$ SYM[§]. Euclidean symmetry is taken as $SO(2)_E$ which is subgroup of $SO(3)_E$.

With the above decomposition, it is possible to consider two ways of a twist.

- A-model: Taking diagonal sum of $SO(2)_E$ and $SO(2)_R$ which is subgroup of $SU(2)_R$

$$SO(2)_E \times (SU(2)_N \times SU(2)_R) \rightarrow SO(2)_{E'} \times SU(2)_N \quad (5.2)$$

In this model, we consider two scalar supercharges transforming as a doublet under $SU(2)_N$.

- B-model: Taking diagonal sum of $SO(2)_E$ with $SO(2)_N$ which is subgroup of $SU(2)_N$

$$SO(2)_E \times (SU(2)_N \times SU(2)_R) \rightarrow SO(2)_{E'} \times SU(2)_R \quad (5.3)$$

In this model, we consider two supercharges transforming as a doublet under $SU(2)_R$.

From these consideration we conclude that the internal symmetries in each model are the same one in three dimensional models, even though new scalar supercharges appear.

[§] $SO(6)$ can take maximal subgroups as $SO(6) \rightarrow SO(4) \times SO(2)$ or $SO(6) \rightarrow SO(3) \otimes SO(3)$. In the present model we consider the latter decomposition and further take $SO(2)_E$ subgroup of the $SO(3)$

5.1 A-model

Two dimensional A-twisted action is,

$$\begin{aligned}
\mathcal{S}_0 = \int d^2x \text{Tr} & \left(\frac{1}{2} F_{\mu\nu} F_{\mu\nu} + \mathcal{D}_\mu A_3 \mathcal{D}_\mu A_3 + \mathcal{D}_\mu N \mathcal{D}_\mu N - [A_3, N][A_3, N] \right. \\
& + \mathcal{D}_\mu \phi \mathcal{D}_\mu \bar{\phi} - [A_3, \phi][A_3, \bar{\phi}] - [N, \bar{\phi}][N, \phi] + \frac{1}{4} [\phi, \bar{\phi}]^2 \\
& - i\psi_\mu \mathcal{D}_\mu \psi - \psi_3 [A_3, \psi] - i\chi_\mu \mathcal{D}_\mu \chi - \chi_3 [A_3, \chi] \\
& - i\epsilon_{\mu\nu} \psi_\mu \mathcal{D}_\nu \chi_3 - i\epsilon_{\mu\nu} \psi_3 \mathcal{D}_\mu \chi_\nu + \epsilon_{\mu\nu} \psi_\mu [A_3, \chi_\nu] \\
& - H_\mu H_\mu - H_3 H_3 + \frac{i}{2} \bar{\phi} \{ \psi_\mu, \psi_\mu \} + \frac{i}{2} \bar{\phi} \{ \psi_3, \psi_3 \} \\
& + \frac{i}{2} \phi \{ \chi_\mu, \chi_\mu \} + \frac{i}{2} \phi \{ \chi_3, \chi_3 \} + \frac{i}{2} \bar{\phi} \{ \chi, \chi \} + \frac{i}{2} \phi \{ \psi, \psi \} \\
& \left. + N \{ \psi, \chi \} + N \{ \psi_\mu, \chi_\mu \} + N \{ \psi_3, \chi_3 \} \right), \tag{5.4}
\end{aligned}$$

where $\mu = 1, 2$.

In this model, there are four scalar supercharges ($s, \bar{s}, s_3, \bar{s}_3$) possessing ghost numbers as $(+, -, -, +)$ respectively. Then, the action can be expressed as supercharge exact form in four different ways:

$$\mathcal{S}_0 = \int d^2x \text{Tr} s \bar{s} (\chi\psi + \psi_\mu \chi_\mu + \psi_3 \chi_3 + 2i\epsilon_{\mu\nu} A_3 F_{\mu\nu}) \tag{5.5}$$

$$= \int d^2x \text{Tr} \bar{s}_3 s_3 (\chi\psi - \psi_\mu \chi_\mu + \psi_3 \chi_3 - 2i\epsilon_{\mu\nu} A_3 F_{\mu\nu}) \tag{5.6}$$

$$= \int d^2x \text{Tr} s s_3 (\psi_3 \psi + \chi_3 \chi + \epsilon_{\mu\nu} \psi_\mu \chi_\nu - 2i\epsilon_{\mu\nu} N F_{\mu\nu}) \tag{5.7}$$

$$= \int d^2x \text{Tr} \bar{s}_3 \bar{s} (-\psi_3 \psi - \chi_3 \chi + \epsilon_{\mu\nu} \psi_\mu \chi_\nu - 2i\epsilon_{\mu\nu} N F_{\mu\nu}). \tag{5.8}$$

In this model all the integrands include BF-term which corresponds to Chern-Simons term in three dimensions. Performing the mass deformation, we can take four different choices of scalar supercharges. It can be shown that these procedures are all equivalent by using some discrete symmetries shown in Appendix A.2. As in the case of $N = 8, D = 3$ SYM model, there are some discrete symmetries in the above action. The model, however, possesses a cyclic symmetry, for example supercharges transform as $s \rightarrow s_3, s_3 \rightarrow -\bar{s}_3, \bar{s}_3 \rightarrow \bar{s}$ and $\bar{s} \rightarrow s$. From the discussion in the beginning of this section, the internal symmetry is $SU(2)$ after the twist. However the internal symmetry might be enhanced, because the cyclic symmetry implies that these supercharges belong to the same multiplet in a large group. As we mention in the previous footnote, maximal subgroup of $SO(6)$ is $SO(2) \otimes SU(2) \otimes SU(2)$. If this $SO(2)$ symmetry is identified with two dimensional Euclidean one, the internal symmetry becomes $SO(4)$ after the twist.

Since these procedures of the mass deformation are equivalent, it is sufficient to consider the case of $(s, \bar{s})$ pair. The simplest way to obtain the result is to perform dimensional reduction of deformed $N = 4, D = 3$ SYM model. The results are similar to that of the three dimensional model, and therefore we omit details here.

5.2 B-model

Two dimensional B-model action is

$$\begin{aligned}
\mathcal{S}_0 = \int d^2x \text{Tr} & \left(\frac{1}{4} F_{\mu\nu} F_{\mu\nu} + \frac{1}{2} \mathcal{D}_\mu A_3 \mathcal{D}_\mu A_3 + \frac{1}{2} \mathcal{D}_\mu V_\nu \mathcal{D}_\mu V_\nu + \frac{1}{2} \mathcal{D}_\mu V_3 \mathcal{D}_\mu V_3 \right. \\
& - \frac{1}{2} [A_3, V_\mu] [A_3, V_\mu] - \frac{1}{2} [A_3, V_3] [A_3, V_3] \\
& - 2i\lambda_\mu \mathcal{D}_\mu \tilde{\lambda} - 2\lambda_3 [A_3, \tilde{\lambda}] - 2i\tilde{\lambda}_\mu \mathcal{D}_\mu \lambda - 2\tilde{\lambda}_3 [A_3, \lambda] \\
& - 2i\epsilon_{\mu\nu} \lambda_\mu \mathcal{D}_\nu \tilde{\lambda}_3 - 2i\epsilon_{\mu\nu} \lambda_3 \mathcal{D}_\mu \tilde{\lambda}_\nu + 2i\epsilon_{\mu\nu} \lambda_\mu [A_3, \tilde{\lambda}_\nu] \\
& + 2iV_\mu \{\lambda_\mu, \tilde{\lambda}\} + 2iV_3 \{\lambda_3, \tilde{\lambda}\} + 2iV_\mu \{\tilde{\lambda}_\mu, \lambda\} + 2iV_3 \{\tilde{\lambda}_3, \lambda\} \\
& + 2i\epsilon_{\mu\nu} V_\mu \{\lambda_\nu, \tilde{\lambda}_3\} - 2i\epsilon_{\mu\nu} V_\mu \{\lambda_3, \tilde{\lambda}_\nu\} + 2i\epsilon_{\mu\nu} V_3 \{\lambda_\mu, \tilde{\lambda}_\nu\} \\
& \left. - \frac{1}{4} [V_\mu, V_\nu]^2 - \frac{1}{2} [V_\mu, V_3] [V_\mu, V_3] + 2K^2 + 2G\tilde{G} \right) \quad (5.9)
\end{aligned}$$

Then, the exact form of the action is

$$\mathcal{S}_0 = \int d^2x \text{Tr} s \bar{s} \left(2\tilde{\lambda}\lambda + 2\tilde{\lambda}_\mu \lambda_\mu + 2\tilde{\lambda}_3 \lambda_3 + \epsilon_{\mu\nu} (V_3 F_{\mu\nu} - iA_3 F_{\mu\nu} + 2V_\mu \mathcal{D}_\nu A_3) \right) \quad (5.10)$$

$$= \int d^2x \text{Tr} \bar{s}_3 s_3 \left(2\tilde{\lambda}\lambda - 2\tilde{\lambda}_\mu \lambda_\mu + 2\tilde{\lambda}_3 \lambda_3 + \epsilon_{\mu\nu} (-V_3 F_{\mu\nu} + iA_3 F_{\mu\nu} + 2V_\mu \mathcal{D}_\nu A_3) \right) \quad (5.11)$$

$$= \int d^2x \text{Tr} s s_3 \left(2\lambda\lambda_3 + 2\tilde{\lambda}\lambda_3 + 2\epsilon_{\mu\nu} \tilde{\lambda}_\mu \lambda_\nu + 2iV_3 \mathcal{D}_\mu V_\mu \right) \quad (5.12)$$

$$= \int d^2x \text{Tr} \bar{s}_3 \bar{s} \left(-2\lambda\tilde{\lambda}_3 - 2\tilde{\lambda}\lambda_3 + 2\epsilon_{\mu\nu} \tilde{\lambda}_\mu \lambda_\nu - 2iV_3 \mathcal{D}_\mu V_\mu \right) \quad (5.13)$$

In contrast to the A-model there is no cyclic symmetry. Not four all different procedures of the mass deformation are not equivalent There are however some discrete symmetries in this model, so that it is sufficient to consider $(s, \bar{s})$ and (s, s_3) cases in the deformation. A characteristic of this model is that the integrand of $(s, \bar{s})$ includes BF-term and that of (s, s_3) does not include it. We will try to perform the deformation of each model.

5.2.1 $s, \bar{s}$ action

We define

$$Q = s + \Delta s, \quad \bar{Q} = \bar{s} + \Delta \bar{s}, \quad (5.14)$$

and assume the following algebra,

$$Q^2 = MJ_{++}, \quad \bar{Q}^2 = -MJ_{--}, \quad \{Q, \bar{Q}\} = -MJ_0. \quad (5.15)$$

We obtain the following transformation law of Δs and $\Delta \bar{s}$,

$$\Delta s(G) = 0, \quad \Delta \bar{s}(G) = 2M\lambda, \quad (5.16)$$

$$\Delta s(\tilde{G}) = -2M\tilde{\lambda}, \quad \Delta \bar{s}(\tilde{G}) = 0, \quad (5.17)$$

$$\Delta s(K) = M\lambda, \quad \Delta \bar{s}(K) = M\tilde{\lambda}, \quad (5.18)$$

$$\Delta s(\lambda_3) = 0, \quad \Delta \bar{s}(\lambda_3) = -MV_3, \quad (5.19)$$

$$\Delta s(\tilde{\lambda}_3) = MV_3, \quad \Delta \bar{s}(\tilde{\lambda}_3) = 0, \quad (5.20)$$

$$\Delta s(\lambda_\mu) = 0, \quad \Delta \bar{s}(\lambda_\mu) = -MV_\mu, \quad (5.21)$$

$$\Delta s(\tilde{\lambda}_\mu) = MV_\mu, \quad \Delta \bar{s}(\tilde{\lambda}_\mu) = 0. \quad (5.22)$$

The action is now deformed as,

$$\mathcal{S} = (Q\bar{Q} - M)\mathcal{F}_0, \quad (5.23)$$

where

$$\mathcal{F}_0 = 2\tilde{\lambda}\lambda + 2\tilde{\lambda}_\mu\lambda_\mu + 2\tilde{\lambda}_3\lambda_3 + \epsilon_{\mu\nu}(V_3F_{\mu\nu} - iA_3F_{\mu\nu} + 2V_\mu\mathcal{D}_\nu A_3). \quad (5.24)$$

Then, the deformed action is,

$$\begin{aligned} \mathcal{S} = \mathcal{S}_0 + 2M \Big(& \tilde{\lambda}\lambda - 2\tilde{\lambda}_\mu\lambda_\mu - 2\tilde{\lambda}_3\lambda_3 \Big) \\ & + 3iM\epsilon_{\mu\nu}V_\mu[V_\nu, A_3] + iM\epsilon_{\mu\nu}A_3F_{\mu\nu} + 2M^2(V_\mu V_\mu + V_3V_3). \end{aligned} \quad (5.25)$$

The potential terms which involve V_μ and V_3 are,

$$V = -\frac{1}{4}[V_\mu, V_\nu]^2 - \frac{1}{2}[V_\mu, V_3][V_\mu, V_3] + 3iM\epsilon_{\mu\nu}V_\mu[V_\nu, V_3] + 2M^2(V_\mu V_\mu + V_3V_3). \quad (5.26)$$

These results are exactly the same as those which are obtained through the dimensional reduction of three dimensional B-model. A vacuum solution can be obtained in the similar way, and therefore we have to identify V_μ and V_3 with $SU(2)$ generators to form the fuzzy sphere solution.

5.2.2 s, s_3 action (same for $\bar{s}_3, \bar{s}$)

In this combination, we can not find any descriptions of supertransformation of deformed charges. The cause of this problem has not been cleared yet. This problem might be related to the internal symmetry in these models. From the discussion of the discrete symmetry A-model has larger internal symmetry than that of B-model. However, it is not clear that this fact affects whether the mass deformation is possible.

6 Summary and discussion

In this paper, we have evaluated whether the mass deformation method, introduced in [1] is applicable to any SYM models and also focused on the fuzzy sphere solution to resolve the flat direction which typically appear in higher N -extended SYM models.

We have found that the key ingredients of mass deformation are that the theory should have at least two topological supercharges and their ghost numbers should be opposite. Therefore we have focused on the models which satisfy both characteristics: maximal SUSY or half maximal SUSY models in four dimensions or less.

The mass deformation procedure is as follow:

1. Find the combination of supercharges whose ghost numbers are opposite, and write the action in exact form in their supercharges.
2. Deform their supercharges, and define deformed algebra.
3. The action is exact with deformed supercharges, and is invariant under its two super-symmetries.
4. Through this process, the action receives mass terms, and the potential terms are changed.

5. Classical configuration constrains scalar fields to satisfy the algebra of $SU(2)$. Thus the fuzzy sphere solution appear.

We have studied the following models: $N = 4, D = 4$ SYM, $N = 8, D = 3$ SYM, two different twisted models of $N = 4, D = 3$ SYM and of $N = (4, 4), D = 2$ SYM. We investigate the classical solution of the potential in each case. Some of them are connected by the dimensional reduction, so that the result is sometimes straightforward. But dimensional reduction produces new scalar supercharges derived from vector supercharge in higher dimensions, and it contributes to new combinations of supercharges. Even though there appear new supercharges and new combinations of them, they are connected each other through the discrete symmetries. It is therefore sometimes enough to evaluate only one model.

Since A-twisted models always possess at least one $SU(2)$ R -symmetry, it results in the fact that the theories possess two nilpotent supercharges which transform as doublet under $SU(2)$ R -symmetry. Indeed, the models which satisfy above conditions are always possible to perform mass deformation and no exception arises with respect to the model we treat in this paper.

The difference appears in B-twisted models. The mass deformation is applicable to $N = 4, D = 3$ model, and the problem does not arise. Through the dimensional reduction we obtain $N = (4, 4), D = 2$ B-model. But the deformations using some new combinations of supercharges are not allowed, because there is no reasonable definition of supertransformation which is consistent with deformed algebra. This might be caused by insufficiency of R -symmetry, since the cyclic discrete symmetry is not found in this model.

There are some points which are not fully clarified, and they are left for future work:

- The difference of the results in A- and B-models, from the point of internal symmetry.
- Treatment for complex terms which should be excluded for lattice simulation.

Acknowledgment

We would like to thank N. Kawamoto, F. Sugino and I. Kanamori for fruitful discussion and comments. A. M. would like to thank K. Ohta for useful comments.

A Discrete symmetry

A.1 Discrete symmetries in $N = 8, D = 3$ model

Symmetry obtained through 4D (ghost number flipping): $(s \leftrightarrow \bar{s}), (\bar{s}_4 \leftrightarrow s_4)$

$$\begin{array}{lll} A_i \rightarrow A_i, & A_4 \rightarrow A_4, & v \rightarrow v, \\ \phi \rightarrow -\bar{\phi}, & \bar{\phi} \rightarrow -\phi, & v_i \rightarrow -v_i \\ \chi \leftrightarrow \lambda, & \chi_i \leftrightarrow \lambda_i, & C_i \leftrightarrow \psi_i, \quad C_4 \leftrightarrow \psi_4. \end{array}$$

ghost number flipping: $(s \rightarrow s_4, s_4 \rightarrow -s), (\bar{s}_4 \rightarrow \bar{s}, \bar{s} \rightarrow -\bar{s}_4)$

$$\begin{array}{lll} A_i \rightarrow A_i, & v_i \rightarrow v_i, & v \rightarrow v, \\ \phi \rightarrow -\bar{\phi}, & \bar{\phi} \rightarrow -\phi, & A_4 \rightarrow -A_4, \\ \lambda \rightarrow -\psi_4, & \psi_4 \rightarrow \lambda, & \chi \rightarrow C_4, \quad C_4 \rightarrow -\chi, \\ \lambda_i \rightarrow \psi_i, & \psi_i \rightarrow -\lambda_i, & \chi_i \rightarrow -C_i, \quad C_i \rightarrow \chi_i. \end{array}$$

ghost number conserving: $(s \rightarrow \bar{s}_4, \bar{s}_4 \rightarrow -s), (\bar{s} \rightarrow s_4, s_4 \rightarrow -\bar{s})$

$$\begin{array}{llll}
A_i \rightarrow A_i, & v_i \rightarrow v_i, & \phi \rightarrow \phi, & \bar{\phi} \rightarrow \bar{\phi}, \\
A_4 \rightarrow -A_4, & v \rightarrow -v, & & \\
\lambda \rightarrow C_4, & C_4 \rightarrow -\lambda, & \chi \rightarrow \psi_4, & \psi_4 \rightarrow -\chi, \\
\lambda_i \rightarrow -C_i, & C_i \rightarrow \lambda_i, & \chi_i \rightarrow -\psi_i, & \psi_i \rightarrow \chi_i.
\end{array}$$

A.2 Discrete symmetries in $N = 4, D = 2$ A-model SYM

ghost number flipping: $(s \rightarrow \bar{s}, \bar{s} \rightarrow -s), (s_3 \rightarrow \bar{s}_3, \bar{s}_3 \rightarrow -s_3)$

$$\begin{array}{llll}
A_\mu \rightarrow A_\mu, & A_3 \rightarrow A_3, & H_\mu \rightarrow H_\mu, & H_3 \rightarrow H_3, \\
\phi \leftrightarrow \bar{\phi}, & N \rightarrow -N & & \\
\psi_\mu \rightarrow \chi_\mu, & \psi \rightarrow \chi, & \psi_3 \rightarrow \chi_3, & \\
\chi_\mu \rightarrow -\psi_\mu, & \chi \rightarrow -\psi, & \chi_3 \rightarrow -\psi_3. &
\end{array}$$

ghost number flipping: $(s \rightarrow s_3, s_3 \rightarrow -s), (\bar{s}_3 \rightarrow \bar{s}, \bar{s} \rightarrow -\bar{s}_3)$

$$\begin{array}{llll}
A_\mu \rightarrow A_\mu, & N \rightarrow N, & H_\mu \rightarrow H_\mu, & H_3 \rightarrow H_3 \\
\phi \leftrightarrow \bar{\phi}, & A_3 \rightarrow -A_3, & & \\
\psi_\mu \rightarrow \epsilon_{\mu\nu} \chi_\nu, & \psi \rightarrow \psi_3, & \chi_3 \rightarrow \chi, & \\
\chi_\mu \rightarrow \epsilon_{\mu\nu} \psi_\nu, & \psi_3 \rightarrow -\psi, & \chi \rightarrow -\chi_3. &
\end{array}$$

ghost number conserving: $(s \rightarrow \bar{s}_3, \bar{s}_3 \rightarrow -s), (s_3 \rightarrow \bar{s}, \bar{s} \rightarrow -s_3)$

$$\begin{array}{llll}
A_\mu \rightarrow A_\mu, & A_3 \rightarrow A_3, & N \rightarrow N, & H_3 \rightarrow H_3, \\
\phi \rightarrow \phi, & \bar{\phi} \rightarrow \bar{\phi}, & H_\mu \rightarrow -H_\mu, & \\
\psi_\mu \rightarrow -\epsilon_{\mu\nu} \psi_\nu, & \psi \rightarrow \chi_3, & \psi_3 \rightarrow \chi, & \\
\chi_\mu \rightarrow -\epsilon_{\mu\nu} \chi_\nu, & \chi_3 \rightarrow -\psi, & \chi \rightarrow -\psi_3. &
\end{array}$$

ghost number flipping forward cyclic: $(\bar{s}_3 \rightarrow \bar{s} \rightarrow s \rightarrow s_3 \rightarrow -\bar{s}_3)$

$$\begin{array}{llll}
A_\mu \rightarrow A_\mu, & H_3 \rightarrow H_3, & & \\
\phi \leftrightarrow \bar{\phi}, & A_3 \leftrightarrow N, & H_\mu \rightarrow -\epsilon_{\mu\nu} H_\nu, & \\
\psi \rightarrow \psi_3, & \psi_3 \rightarrow -\chi_3, & \chi_3 \rightarrow \chi, & \chi \rightarrow \psi, \\
\psi_\mu \rightarrow \epsilon_{\mu\nu} \chi_\nu, & \chi_\mu \rightarrow \psi_\mu. & &
\end{array}$$

ghost number flipping backward cyclic: $(s_3 \rightarrow s \rightarrow \bar{s} \rightarrow \bar{s}_3 \rightarrow -s_3)$

$$\begin{array}{llll}
A_\mu \rightarrow A_\mu, & H_3 \rightarrow H_3, & & \\
\phi \leftrightarrow \bar{\phi}, & A_3 \leftrightarrow N, & H_\mu \rightarrow \epsilon_{\mu\nu} H_\nu, & \\
\psi \rightarrow \chi, & \chi \rightarrow \chi_3, & \chi_3 \rightarrow -\psi_3, & \psi_3 \rightarrow \psi, \\
\psi_\mu \rightarrow \chi_\nu, & \chi_\mu \rightarrow -\epsilon_{\mu\nu} \psi_\mu. & &
\end{array}$$

A.3 Discrete symmetries in $N = 4, D = 2$ B-model SYM

ghost number flipping: $(s \rightarrow \bar{s}, \bar{s} \rightarrow -s), (s_3 \rightarrow \bar{s}_3, \bar{s}_3 \rightarrow -s_3)$

$$\begin{aligned} V_\mu &\rightarrow V_\mu, & V_3 &\rightarrow V_3, & A_\mu &\rightarrow A_\mu, & A_3 &\rightarrow A_3, \\ G &\rightarrow -\tilde{G}, & \tilde{G} &\rightarrow -G, & K &\rightarrow -K, \\ \lambda &\rightarrow -\tilde{\lambda}, & \lambda_\mu &\rightarrow \tilde{\lambda}_\mu, & \lambda_3 &\rightarrow \tilde{\lambda}_3, \\ \tilde{\lambda} &\rightarrow \lambda, & \tilde{\lambda}_\mu &\rightarrow -\lambda_\mu, & \tilde{\lambda}_3 &\rightarrow -\lambda_3. \end{aligned}$$

ghost number flipping: $(s \rightarrow s_3, s_3 \rightarrow -s), (\bar{s}_3 \rightarrow \bar{s}, \bar{s} \rightarrow -\bar{s}_3)$

$$\begin{aligned} V_\mu &\rightarrow -V_\mu, & V_3 &\rightarrow -V_3, & A_\mu &\rightarrow A_\mu, & A_3 &\rightarrow -A_3, \\ G &\rightarrow \tilde{G}, & \tilde{G} &\rightarrow G, & K &\rightarrow -K, \\ \lambda &\rightarrow -\tilde{\lambda}_3, & \tilde{\lambda} &\rightarrow \lambda_3, & \lambda_\mu &\rightarrow \epsilon_{\mu\nu} \tilde{\lambda}_\nu, \\ \tilde{\lambda}_3 &\rightarrow \lambda, & \lambda_3 &\rightarrow -\tilde{\lambda}, & \tilde{\lambda}_\mu &\rightarrow \epsilon_{\mu\nu} \lambda_\nu. \end{aligned}$$

ghost number conserving: $(s \rightarrow \bar{s}_3, \bar{s}_3 \rightarrow -s), (s_3 \rightarrow \bar{s}, \bar{s} \rightarrow -s_3)$

$$\begin{aligned} V_\mu &\rightarrow -V_\mu, & V_3 &\rightarrow V_3, & A_\mu &\rightarrow A_\mu, & A_3 &\rightarrow A_3, \\ G &\rightarrow G, & \tilde{G} &\rightarrow \tilde{G}, & K &\rightarrow K, \\ \lambda &\rightarrow -\lambda_3, & \tilde{\lambda} &\rightarrow \tilde{\lambda}_3, & \lambda_\mu &\rightarrow -\epsilon_{\mu\nu} \lambda_\nu, \\ \lambda_3 &\rightarrow \lambda, & \tilde{\lambda}_3 &\rightarrow -\tilde{\lambda}, & \tilde{\lambda}_\mu &\rightarrow -\epsilon_{\mu\nu} \tilde{\lambda}_\nu. \end{aligned}$$

B Tables of supertransformation

B.1 Supertransformation in $N = 4, D = 4$

	gh#	s	$\bar{s}$
v	0	$\frac{1}{2}\lambda$	$\frac{1}{2}\chi$
v_B^+	0	$-\frac{1}{2}\lambda_B^+$	$\frac{1}{2}\chi_B^+$
λ	1	$-i[\phi, v]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
λ_B^+	1	$i[\phi, v_B^+]$	$\frac{i}{32}\Gamma_{BCD}^+[v_C^+, v_D^+] + i[v_B^+, v] - 2F_B^+ + h_B^+$
ψ_μ	-1	$i\mathcal{D}_\mu v + i\mathcal{D}_\rho v_{\mu\rho}^+ + h_\nu$	$i\mathcal{D}_\mu \bar{\phi}$
ϕ	2	0	λ
$\bar{\phi}$	-2	$-\chi$	0
χ	-1	$\frac{i}{2}[\phi, \bar{\phi}]$	$i[\bar{\phi}, v]$
χ_B^+	-1	$\frac{i}{32}\Gamma_{BCD}^+[v_C^+, v_D^+] - i[v_B^+, v] - 2F_B^+ + h_B^+$	$i[\bar{\phi}, v_B^+]$
C_μ	1	$-i\mathcal{D}_\mu \phi$	$i\mathcal{D}_\mu v - i\mathcal{D}_\rho v_{\mu\rho}^+ - h_\mu$
A_μ	0	$-\frac{i}{2}C_\mu$	$-\frac{i}{2}\psi_\mu$
h_μ	0	$-\frac{i}{2}\mathcal{D}_\mu \lambda + \frac{i}{2}\mathcal{D}_\rho \lambda_{\mu\rho}^+ + \frac{i}{2}[C_\rho, v_{\mu\rho}^+] - \frac{i}{2}[v, C_\mu] - \frac{i}{2}[\phi, \psi_\mu]$	$\frac{i}{2}\mathcal{D}_\mu \chi - \frac{i}{2}\mathcal{D}_\rho \chi_{\mu\rho}^+ + \frac{i}{2}[\psi_\rho, v_{\mu\rho}^+] + \frac{i}{2}[v, \psi_\mu] - \frac{i}{2}[\bar{\phi}, C_\mu]$
h_B^+	0	$-\frac{i}{2}[\phi, \chi_B^+] + \frac{i}{2}[v, \lambda_B^+] + \frac{i}{2}[v_B^+, \lambda] - \frac{i}{2}\delta_{B,\mu\nu}^+ \mathcal{D}_\mu C_\nu + \frac{i}{32}\Gamma_{BCD}^+[\lambda_C^+, v_D^+]$	$\frac{i}{2}[\bar{\phi}, \lambda_B^+] + \frac{i}{2}[v, \chi_B^+] - \frac{i}{2}[v_B^+, \chi] - \frac{i}{2}\delta_{B,\mu\nu}^+ \mathcal{D}_\mu \psi_\nu - \frac{i}{32}\Gamma_{BCD}^+[\chi_C^+, v_D^+]$

B.2 Supertransformation in $N = 8, D = 3$

$(s, \bar{s})$ supertransformation

	gh#	s (gh# = 1)	$\bar{s}$ (gh# = -1)
v	0	$\frac{1}{2}\lambda$	$\frac{1}{2}\chi$
v_i	0	$-\frac{1}{2}\lambda_i$	$\frac{1}{2}\chi_i$
λ	1	$-i[\phi, v]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
λ_i	1	$i[\phi, v_i]$	$-\frac{i}{2}\epsilon_{ijk}[v_j, v_k] + i[v_i, v]$ $-\mathcal{D}_i A_4 + \frac{1}{2}\epsilon_{ijk}F_{jk} + h_{i4}^+$
ψ_i	-1	$i\mathcal{D}_i v - i\epsilon_{ijk}\mathcal{D}_j v_k + [A_4, v_i] + h_i$	$i\mathcal{D}_i \bar{\phi}$
ψ_4	-1	$[A_4, v] - i\mathcal{D}_j v_j + h_4$	$[A_4, \bar{\phi}]$
ϕ	2	0	λ
$\bar{\phi}$	-2	$-\chi$	0
χ	-1	$\frac{i}{2}[\phi, \bar{\phi}]$	$i[\bar{\phi}, v]$
χ_i	-1	$-\frac{i}{2}\epsilon_{ijk}[v_j, v_k] - i[v_i, v]$ $-\mathcal{D}_i A_4 + \frac{1}{2}\epsilon_{ijk}F_{jk} + h_{i4}^+$	$i[\bar{\phi}, v_i]$
C_i	1	$-i\mathcal{D}_i \phi$	$i\mathcal{D}_i v + i\epsilon_{ijk}\mathcal{D}_j v_k - [A_4, v_i] - h_i$
C_4	1	$-[A_4, \phi]$	$[A_4, v] + i\mathcal{D}_j v_j - h_4$
A_i	0	$-\frac{i}{2}C_i$	$-\frac{i}{2}\psi_i$
A_4	0	$-\frac{i}{2}C_4$	$-\frac{i}{2}\psi_4$

Transformation of auxiliary fields by s and $\bar{s}$

	s	$\bar{s}$
h_{i4}^+	$-\frac{i}{2}[\phi, \chi_i] - \frac{i}{2}[\lambda_i, v] + \frac{i}{2}[v_i, \lambda] - \frac{1}{2}[C_i, A_4]$ $-\frac{i}{2}\epsilon_{ijk}[\lambda_j, v_k] - \frac{i}{2}\mathcal{D}_i C_4 + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j C_k$	$\frac{i}{2}[\bar{\phi}, \lambda_i] - \frac{i}{2}[\chi_i, v] - \frac{i}{2}[v_i, \chi] - \frac{1}{2}[\psi_i, A_4]$ $+\frac{i}{2}\epsilon_{ijk}[\chi_i, v_k] - \frac{i}{2}\mathcal{D}_i \psi_4 + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \psi_k$
h_i	$-\frac{i}{2}[\phi, \psi_i] + \frac{i}{2}[C_i, v] + \frac{i}{2}[C_4, v_i] + \frac{1}{2}[A_4, \lambda_i]$ $-\frac{i}{2}\epsilon_{ijk}[C_j, v_k] - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \lambda_k - \frac{i}{2}\mathcal{D}_i \lambda$	$-\frac{i}{2}[\bar{\phi}, C_i] - \frac{i}{2}[\psi_i, v] + \frac{i}{2}[\psi_4, v_i] - \frac{1}{2}[A_4, \chi_i]$ $-\frac{i}{2}\epsilon_{ijk}[\psi_j, v_k] + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \chi_k + \frac{i}{2}\mathcal{D}_i \chi$
h_4	$-\frac{i}{2}[\phi, \psi_4] + \frac{i}{2}[C_4, v] - \frac{1}{2}[A_4, \lambda] - \frac{i}{2}[C_j, v_j]$ $-\frac{i}{2}\mathcal{D}_j \lambda_j$	$-\frac{i}{2}[\bar{\phi}, C_4] - \frac{i}{2}[\psi_4, v] + \frac{1}{2}[A_4, \chi] - \frac{i}{2}[\psi_j, v_j]$ $+\frac{i}{2}\mathcal{D}_j \chi_j$

$(\bar{s}_4, s_4)$ supertransformation

	gh#	$\bar{s}_4$ (gh# = 1)	s_4 (gh# = -1)
v	0	$-\frac{1}{2}C_4$	$-\frac{1}{2}\psi_4$
v_i	0	$\frac{1}{2}C_i$	$-\frac{1}{2}\psi_i$
λ	1	$-[A_4, \phi]$	$-[A_4, v] - i\mathcal{D}_j v_j + g$
λ_i	1	$-i\mathcal{D}_i \phi$	$-i\mathcal{D}_i v + [A_4, v_i] + i\epsilon_{ijk}\mathcal{D}_j v_k + g_{i4}^+$
ψ_i	-1	$\frac{i}{2}\epsilon_{ijk}[v_j, v_k] - i[v_i, v]$ $-\mathcal{D}_i A_4 - \frac{1}{2}\epsilon_{ijk}F_{jk} + g_i$	$-i[\bar{\phi}, v_i]$
ψ_4	-1	$\frac{i}{2}[\phi, \bar{\phi}]$	$-i[\bar{\phi}, v]$
ϕ	2	0	C_4
$\bar{\phi}$	-2	$-\psi_4$	0
χ	-1	$-[A_4, v] + i\mathcal{D}_j v_j - g$	$[A_4, \bar{\phi}]$
χ_i	-1	$-i\mathcal{D}_i v - [A_4, v_i] - i\epsilon_{ijk}\mathcal{D}_j v_k - g_{i4}^+$	$i\mathcal{D}_i \bar{\phi}$
C_i	1	$-i[\phi, v_i]$	$\frac{i}{2}\epsilon_{ijk}[v_j, v_k] + i[v_i, v]$ $-\mathcal{D}_i A_4 - \frac{1}{2}\epsilon_{ijk}F_{jk} + g_i$
C_4	1	$i[\phi, v]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
A_i	0	$-\frac{i}{2}\lambda_i$	$-\frac{i}{2}\chi_i$
A_4	0	$-\frac{i}{2}\lambda$	$-\frac{i}{2}\chi$

Transformation of auxiliary fields by $\bar{s}_4$ and s_4

	$\bar{s}_4$	s_4
g_{i4}^+	$\frac{i}{2}[\phi, \chi_i] + \frac{i}{2}[\lambda_i, v] + \frac{i}{2}[\lambda, v_i] - \frac{1}{2}[A_4, C_i]$ $+\frac{i}{2}\mathcal{D}_i C_4 + \frac{i}{2}\epsilon_{ijk}[\lambda_j, v_k] - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j C_k$	$\frac{i}{2}[\bar{\phi}, \lambda_i] - \frac{i}{2}[\chi_i, v] + \frac{i}{2}[\chi, v_i] + \frac{1}{2}[A_4, \psi_i]$ $-\frac{i}{2}\mathcal{D}_i \psi_4 + \frac{i}{2}\epsilon_{ijk}[\chi_j, v_k] + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \psi_k$
g_i	$-\frac{i}{2}[\phi, \psi_i] + \frac{i}{2}[C_i, v] - \frac{i}{2}[v_i, C_4] - \frac{1}{2}[\lambda_i, A_4]$ $-\frac{i}{2}\epsilon_{ijk}[C_j, v_k] - \frac{i}{2}\mathcal{D}_i \lambda - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \lambda_k$	$\frac{i}{2}[\bar{\phi}, C_i] + \frac{i}{2}[\psi_i, v] + \frac{i}{2}[v_i, \psi_4] - \frac{1}{2}[\chi_i, A_4]$ $+\frac{i}{2}\epsilon_{ijk}[\psi_j, v_k] - \frac{i}{2}\mathcal{D}_i \chi - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \chi_k$
g	$\frac{i}{2}[\phi, \chi] + \frac{i}{2}[\lambda, v] + \frac{1}{2}[A_4, C_4] - \frac{i}{2}[\lambda_j, v_j]$ $+\frac{i}{2}\mathcal{D}_j C_j$	$\frac{i}{2}[\bar{\phi}, \lambda] - \frac{i}{2}[\chi, v] - \frac{1}{2}[A_4, \psi_4] - \frac{i}{2}[\chi_j, v_j]$ $-\frac{i}{2}\mathcal{D}_j \psi_j$

Supertransformation of (s, s_4)

	gh#	s (gh# = 1)	s_4 (gh# = -1)
v	0	$\frac{1}{2}\lambda$	$-\frac{1}{2}\psi_4$
v_i	0	$-\frac{1}{2}\lambda_i$	$-\frac{1}{2}\psi_i$
λ	1	$-i[\phi, v]$	$-[A_4, v] - i\mathcal{D}_j v_j + L$
λ_i	1	$i[\phi, v_i]$	$-i\mathcal{D}_i v + [A_4, v_i] + i\epsilon_{ijk}\mathcal{D}_j v_k - N_i$
ψ_i	-1	$i\mathcal{D}_i v - i\epsilon_{ijk}\mathcal{D}_j v_k + [A_4, v_i] + N_i$	$-i[\bar{\phi}, v_i]$
ψ_4	-1	$[A_4, v] - i\mathcal{D}_j v_j + L$	$-i[\bar{\phi}, v]$
ϕ	2	0	C_4
$\bar{\phi}$	-2	$-\chi$	0
χ	-1	$\frac{i}{2}[\phi, \bar{\phi}]$	$[A_4, \bar{\phi}]$
χ_i	-1	$-\frac{i}{2}\epsilon_{ijk}[v_j, v_k] - i[v_i, v]$ $-\mathcal{D}_i A_4 + \frac{1}{2}\epsilon_{ijk}F_{jk} + M_i$	$i\mathcal{D}_i \bar{\phi}$
C_i	1	$-i\mathcal{D}_i \phi$	$\frac{i}{2}\epsilon_{ijk}[v_j, v_k] + i[v_i, v]$ $-\mathcal{D}_i A_4 - \frac{1}{2}\epsilon_{ijk}F_{jk} - M_i$
C_4	1	$-[A_4, \phi]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
A_i	0	$-\frac{i}{2}C_i$	$-\frac{i}{2}\chi_i$
A_4	0	$-\frac{i}{2}C_4$	$-\frac{i}{2}\chi$

Transformation of auxiliary fields in (s, s_4) case

	s	s_4
M_i	$-\frac{i}{2}[\phi, \chi_i] - \frac{i}{2}[\lambda_i, v] + \frac{i}{2}[v_i, \lambda] - \frac{1}{2}[C_i, A_4]$ $-\frac{i}{2}\epsilon_{ijk}[\lambda_j, v_k] - \frac{i}{2}\mathcal{D}_i C_4 + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j C_k$	$-\frac{i}{2}[\bar{\phi}, C_i] - \frac{i}{2}[\psi_i, v] - \frac{i}{2}[v_i, \psi_4] + \frac{1}{2}[\chi_i, A_4]$ $-\frac{i}{2}\epsilon_{ijk}[\psi_j, v_k] + \frac{i}{2}\mathcal{D}_i \chi + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \chi_k$
N_i	$-\frac{i}{2}[\phi, \psi_i] + \frac{i}{2}[C_i, v] + \frac{i}{2}[C_4, v_i] + \frac{1}{2}[A_4, \lambda_i]$ $-\frac{i}{2}\epsilon_{ijk}[C_j, v_k] - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \lambda_k - \frac{i}{2}\mathcal{D}_i \lambda$	$-\frac{i}{2}[\bar{\phi}, \lambda_i] + \frac{i}{2}[\chi_i, v] - \frac{i}{2}[\chi, v_i] - \frac{1}{2}[A_4, \psi_i]$ $+\frac{i}{2}\mathcal{D}_i \psi_4 - \frac{i}{2}\epsilon_{ijk}[\chi_j, v_k] - \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j \psi_k$
L	$-\frac{i}{2}[\phi, \psi_4] + \frac{i}{2}[C_4, v] - \frac{1}{2}[A_4, \lambda] - \frac{i}{2}[C_j, v_j]$ $-\frac{i}{2}\mathcal{D}_j \lambda_j$	$\frac{i}{2}[\bar{\phi}, \lambda] - \frac{i}{2}[\chi, v] - \frac{1}{2}[A_4, \psi_4] - \frac{i}{2}[\chi_j, v_j]$ $-\frac{i}{2}\mathcal{D}_j \psi_j$

Supertransformation of $(\bar{s}, \bar{s}_4)$

	gh#	$\bar{s}_4$ (gh# = 1)	$\bar{s}$ (gh# = -1)
v	0	$-\frac{1}{2}C_4$	$\frac{1}{2}\chi$
v_i	0	$\frac{1}{2}C_i$	$\frac{1}{2}\chi_i$
λ	1	$-[A_4, \phi]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
λ_i	1	$-i\mathcal{D}_i\phi$	$-\frac{i}{2}\epsilon_{ijk}[v_j, v_k] + i[v_i, v]$ $-\mathcal{D}_iA_4 + \frac{1}{2}\epsilon_{ijk}F_{jk} + M_i$
ψ_i	-1	$\frac{i}{2}\epsilon_{ijk}[v_j, v_k] - i[v_i, v]$ $-\mathcal{D}_iA_4 - \frac{1}{2}\epsilon_{ijk}F_{jk} - M_i$	$i\mathcal{D}_i\bar{\phi}$
ψ_4	-1	$\frac{i}{2}[\phi, \bar{\phi}]$	$[A_4, \bar{\phi}]$
ϕ	2	0	λ
$\bar{\phi}$	-2	$-\psi_4$	0
χ	-1	$-[A_4, v] + i\mathcal{D}_kv_k - L$	$i[\bar{\phi}, v]$
χ_i	-1	$-i\mathcal{D}_iv - [A_4, v_i] - i\epsilon_{ijk}\mathcal{D}_jv_k + N_i$	$i[\bar{\phi}, v_i]$
C_i	1	$-i[\phi, v_i]$	$i\mathcal{D}_iv + i\epsilon_{ijk}\mathcal{D}_jv_k - [A_4, v_i] - N_i$
C_4	1	$i[\phi, v]$	$[A_4, v] + i\mathcal{D}_jv_j - L$
A_i	0	$-\frac{i}{2}\lambda_i$	$-\frac{i}{2}\psi_i$
A_4	0	$-\frac{i}{2}\lambda$	$-\frac{i}{2}\psi_4$

Transformation of auxiliary fields in $(\bar{s}_4, \bar{s})$

	$\bar{s}_4$	$\bar{s}$
M_i	$\frac{i}{2}[\phi, \psi_i] - \frac{i}{2}[C_i, v] + \frac{i}{2}[v_i, C_4] + \frac{1}{2}[\lambda_i, A_4]$ $+\frac{i}{2}\epsilon_{ijk}[C_j, v_k] + \frac{i}{2}\mathcal{D}_i\lambda + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j\lambda_k$	$\frac{i}{2}[\bar{\phi}, \lambda_i] - \frac{i}{2}[\chi_i, v] - \frac{i}{2}[v_i, \chi] - \frac{1}{2}[\psi_i, A_4]$ $+\frac{i}{2}\epsilon_{ijk}[\chi_i, v_k] - \frac{i}{2}\mathcal{D}_i\psi_4 + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j\psi_k$
N_i	$-\frac{i}{2}[\phi, \chi_i] - \frac{i}{2}[\lambda_i, v] - \frac{i}{2}[\lambda, v_i] + \frac{1}{2}[A_4, C_i]$ $-\frac{i}{2}\mathcal{D}_iC_4 - \frac{i}{2}\epsilon_{ijk}[\lambda_j, v_k] + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_jC_k$	$-\frac{i}{2}[\bar{\phi}, C_i] - \frac{i}{2}[\psi_i, v] + \frac{i}{2}[\psi_4, v_i] - \frac{1}{2}[A_4, \chi_i]$ $-\frac{i}{2}\epsilon_{ijk}[\psi_j, v_k] + \frac{i}{2}\epsilon_{ijk}\mathcal{D}_j\chi_k + \frac{i}{2}\mathcal{D}_i\chi$
L	$\frac{i}{2}[\phi, \chi] + \frac{i}{2}[\lambda, v] + \frac{1}{2}[A_4, C_4] - \frac{i}{2}[\lambda_j, v_j]$ $+\frac{i}{2}\mathcal{D}_jC_j$	$-\frac{i}{2}[\bar{\phi}, C_4] - \frac{i}{2}[\psi_4, v] + \frac{1}{2}[A_4, \chi] - \frac{i}{2}[\psi_j, v_j]$ $+\frac{i}{2}\mathcal{D}_j\chi_j$

B.3 Supertransformation in $N = 4, D = 3$

A-model

	gh#	s	$\bar{s}$
ϕ	2	0	χ
N	0	$-\frac{i}{2}\chi$	$-\frac{i}{2}\psi$
$\bar{\phi}$	-2	ψ	0
ψ_μ	1	$-i\mathcal{D}_\mu\phi$	$H_\mu - \mathcal{D}_\mu N - \frac{1}{2}\epsilon_{\mu\rho\sigma}F_{\rho\sigma}$
χ_μ	-1	$-H_\mu - \mathcal{D}_\mu N + \frac{1}{2}\epsilon_{\mu\rho\sigma}F_{\rho\sigma}$	$-i\mathcal{D}_\mu\bar{\phi}$
χ	1	$-[N, \phi]$	$-\frac{i}{2}[\bar{\phi}, \phi]$
ψ	-1	$-\frac{i}{2}[\phi, \bar{\phi}]$	$-[N, \bar{\phi}]$
A_μ	0	$-\frac{i}{2}\psi_\mu$	$-\frac{i}{2}\chi_\mu$
H_μ	0	$\frac{i}{2}\mathcal{D}_\mu\chi - \frac{i}{2}\epsilon_{\mu\rho\sigma}\mathcal{D}_\rho\psi_\sigma$ $+\frac{i}{2}[\phi, \chi_\mu] + \frac{1}{2}[\psi_\mu, N]$	$-\frac{i}{2}\mathcal{D}_\mu\psi - \frac{i}{2}\epsilon_{\mu\rho\sigma}\mathcal{D}_\rho\chi_\sigma$ $-\frac{i}{2}[\bar{\phi}, \psi_\mu] - \frac{1}{2}[\chi_\mu, N]$

B-model

	gh#	s	$\bar{s}$
V_μ	0	λ_μ	$\tilde{\lambda}_\mu$
A_μ	0	$-i\lambda_\mu$	$-i\tilde{\lambda}_\mu$
λ	1	$-G$	$-\frac{i}{2}\mathcal{D}_\mu V_\mu + K$
$\tilde{\lambda}$	-1	$-\frac{i}{2}\mathcal{D}_\mu V_\mu - K$	$-\tilde{G}$
λ_μ	1	0	$-\frac{1}{4}\epsilon_{\mu\nu\rho}(F_{\nu\rho} + i[V_\nu, V_\rho] + 2i\mathcal{D}_\nu V_\rho)$
$\tilde{\lambda}_\mu$	-1	$\frac{1}{4}\epsilon_{\mu\nu\rho}(F_{\nu\rho} + i[V_\nu, V_\rho] + 2i\mathcal{D}_\nu V_\rho)$	0
G	2	0	$-i\mathcal{D}_\mu \lambda_\mu - i[V_\mu, \lambda_\mu]$
$\tilde{G}$	-2	$-i\mathcal{D}_\mu \tilde{\lambda}_\mu - i[V_\mu, \tilde{\lambda}_\mu]$	0
K	0	$-\frac{i}{2}\mathcal{D}_\mu \lambda_\mu - \frac{i}{2}[V_\mu, \lambda_\mu]$	$\frac{i}{2}\mathcal{D}_\mu \tilde{\lambda}_\mu + \frac{i}{2}[V_\mu, \tilde{\lambda}_\mu]$

B.4 Supertransformation in $N = 4, D = 2$ A-model SYM

A-model

	gh#	$s (+1)$	$\bar{s} (-1)$
ϕ	0	0	χ
N	0	$-\frac{i}{2}\chi$	$-\frac{i}{2}\psi$
$\bar{\phi}$	0	ψ	0
A_3	0	$-\frac{i}{2}\psi_3$	$-\frac{i}{2}\chi_3$
ψ_μ	1	$-i\mathcal{D}_\mu \phi$	$H_\mu - \mathcal{D}_\mu N - \epsilon_{\mu\nu}\mathcal{D}_\nu A_3$
χ_μ	-1	$-H_\mu - \mathcal{D}_\mu N + \epsilon_{\mu\nu}\mathcal{D}_\nu A_3$	$-i\mathcal{D}_\mu \bar{\phi}$
ψ_3	1	$-[A_3, \phi]$	$H_3 + i[A_3, N] - \frac{1}{2}\epsilon_{\mu\nu}F_{\mu\nu}$
χ_3	-1	$-H_3 + i[A_3, N] + \frac{1}{2}\epsilon_{\mu\nu}F_{\mu\nu}$	$-[A_3, \bar{\phi}]$
χ	1	$-[N, \phi]$	$-\frac{i}{2}[\bar{\phi}, \phi]$
ψ	-1	$-\frac{i}{2}[\phi, \bar{\phi}]$	$-[N, \bar{\phi}]$
A_μ	0	$-\frac{i}{2}\psi_\mu$	$-\frac{i}{2}\chi_\mu$
H_μ	0	$\frac{i}{2}\mathcal{D}_\mu \chi - \frac{i}{2}\epsilon_{\mu\nu}\mathcal{D}_\nu \psi_3$ $+\frac{1}{2}\epsilon_{\mu\nu}[A_3, \psi_\nu] + \frac{i}{2}[\phi, \chi_\mu] + \frac{1}{2}[\psi_\mu, N]$	$-\frac{i}{2}\mathcal{D}_\mu \psi - \frac{i}{2}\epsilon_{\mu\nu}\mathcal{D}_\nu \chi_3$ $+\frac{1}{2}\epsilon_{\mu\nu}[A_3, \chi_\nu] - \frac{i}{2}[\bar{\phi}, \psi_\mu] - \frac{1}{2}[\chi_\mu, N]$
H_3	0	$\frac{1}{2}[A_3, \chi] - \frac{i}{2}\epsilon_{\mu\nu}\mathcal{D}_\mu \psi_\nu$ $+\frac{i}{2}[\phi, \chi_3] + \frac{1}{2}[\psi_3, N]$	$-\frac{1}{2}[A_3, \psi] - \frac{i}{2}\epsilon_{\mu\nu}\mathcal{D}_\mu \chi_\nu$ $-\frac{i}{2}[\bar{\phi}, \psi_3] - \frac{1}{2}[\chi_3, N]$

	gh#	$s_3 (-1)$	$\bar{s}_3 (+1)$
ϕ	2	ψ_3	0
N	0	$\frac{i}{2}\chi_3$	$\frac{i}{2}\psi_3$
$\bar{\phi}$	-2	0	χ_3
A_3	0	$-\frac{i}{2}\psi$	$-\frac{i}{2}\chi$
ψ_μ	1	$\epsilon_{\mu\rho}H_\rho - \mathcal{D}_\mu A_3 + \epsilon_{\mu\rho}\mathcal{D}_\rho N$	$-i\epsilon_{\mu\rho}\mathcal{D}_\rho\phi$
χ_μ	-1	$i\epsilon_{\mu\rho}\mathcal{D}_\rho\bar{\phi}$	$\epsilon_{\mu\rho}H_\rho - \mathcal{D}_\mu A_3 - \epsilon_{\mu\rho}\mathcal{D}_\rho N$
ψ_3	1	$-\frac{i}{2}[\phi, \phi]$	$[N, \phi]$
χ_3	-1	$[N, \bar{\phi}]$	$-\frac{i}{2}[\phi, \bar{\phi}]$
χ	1	$-H_3 - i[A_3, N] + \frac{1}{2}\epsilon_{\rho\sigma}F_{\rho\sigma}$	$-[A_3, \phi]$
ψ	-1	$-[A_3, \bar{\phi}]$	$H_3 - i[A_3, N] - \frac{1}{2}\epsilon_{\rho\sigma}F_{\rho\sigma}$
A_μ	0	$-\frac{i}{2}\epsilon_{\mu\rho}\chi_\rho$	$\frac{i}{2}\epsilon_{\mu\rho}\psi_\rho$
H_μ	0	$\frac{i}{2}(\epsilon_{\mu\rho}\mathcal{D}_\rho\psi - i[A_3, \chi_\mu] - \mathcal{D}_\mu\chi_3)$ $-\frac{1}{2}\epsilon_{\mu\rho}[N, \chi_\rho] + \frac{i}{2}\epsilon_{\mu\rho}[\bar{\phi}, \psi_\rho]$	$\frac{i}{2}(\epsilon_{\mu\rho}\mathcal{D}_\rho\chi + i[A_3, \psi_\mu] + \mathcal{D}_\mu\psi_3)$ $-\frac{1}{2}\epsilon_{\mu\rho}[N, \psi_\rho] + \frac{i}{2}\epsilon_{\mu\rho}[\phi, \chi_\rho]$
H_3	0	$\frac{i}{2}\mathcal{D}_\rho\chi_\rho + \frac{1}{2}[A_3, \chi_3] + \frac{i}{2}[\bar{\phi}, \chi] - \frac{1}{2}[\psi, N]$	$-\frac{i}{2}\mathcal{D}_\rho\psi_\rho - \frac{1}{2}[A_3, \psi_3] - \frac{i}{2}[\phi, \psi] + \frac{1}{2}[\chi, N]$

B.5 Supertransformation in $N = 4, D = 2$ B-model SYM

	gh#	s	$\bar{s}$
V_μ	0	λ_μ	$\tilde{\lambda}_\mu$
V_3	0	λ_3	$\tilde{\lambda}_3$
A_μ	0	$-i\lambda_\mu$	$-i\tilde{\lambda}_\mu$
A_3	0	$-i\lambda_3$	$-i\tilde{\lambda}_3$
λ	1	$-G$	$-\frac{i}{2}\mathcal{D}_\mu V_\mu - \frac{1}{2}[A_3, V_3] + K$
$\tilde{\lambda}$	-1	$-\frac{i}{2}\mathcal{D}_\mu V_\mu - \frac{1}{2}[A_3, V_3] - K$	$-G$
λ_μ	1	0	$-\epsilon_{\mu\nu}(\frac{1}{2}\mathcal{D}_\nu A_3 + \frac{i}{2}[V_\nu, V_3]$ $+\frac{i}{2}\mathcal{D}_\nu V_3 - \frac{1}{2}[A_3, V_\nu])$
$\tilde{\lambda}_\mu$	-1	$\epsilon_{\mu\nu}(\frac{1}{2}\mathcal{D}_\nu A_3 + \frac{i}{2}[V_\nu, V_3]$ $+\frac{i}{2}\mathcal{D}_\nu V_3 - \frac{1}{2}[A_3, V_\nu])$	0
λ_3	1	0	$-\frac{1}{4}\epsilon_{\mu\nu}(F_{\mu\nu} + i[V_\mu, V_\nu] + 2i\mathcal{D}_\mu V_\nu)$
$\tilde{\lambda}_3$	-1	$\frac{1}{4}\epsilon_{\mu\nu}(F_{\mu\nu} + i[V_\mu, V_\nu] + 2i\mathcal{D}_\mu V_\nu)$	0
G	2	0	$-i\mathcal{D}_\mu\lambda_\mu - [A_3, \lambda_3]$ $-i[V_\mu, \lambda_\mu] - i[V_3, \lambda_3]$
$\tilde{G}$	-2	$-i\mathcal{D}_\mu\tilde{\lambda}_\mu - [A_3, \tilde{\lambda}_3]$ $-i[V_\mu, \tilde{\lambda}_\mu] - i[V_3, \tilde{\lambda}_3]$	0
K	0	$-\frac{i}{2}\mathcal{D}_\mu\lambda_\mu - \frac{1}{2}[A_3, \lambda_3]$ $-\frac{i}{2}[V_\mu, \lambda_\mu] - \frac{i}{2}[V_3, \lambda_3]$	$\frac{i}{2}\mathcal{D}_\mu\tilde{\lambda}_\mu + \frac{1}{2}[A_3, \tilde{\lambda}_3]$ $+\frac{i}{2}[V_\mu, \tilde{\lambda}_\mu] + \frac{i}{2}[V_3, \tilde{\lambda}_3]$

	gh#	s_3	$\bar{s}_3$
V_μ	0	$-\epsilon_{\mu\rho}\tilde{\lambda}_\rho$	$\epsilon_{\mu\rho}\lambda_\rho$
V_3	0	$\tilde{\lambda}$	λ
A_μ	0	$-i\epsilon_{\mu\rho}\tilde{\lambda}_\rho$	$i\epsilon_{\mu\rho}\lambda_\rho$
A_3	0	$-i\tilde{\lambda}$	$-i\lambda$
λ	1	$\frac{1}{4}\epsilon_{\rho\sigma}(F_{\rho\sigma} + i[V_\rho, V_\sigma] - 2i\mathcal{D}_\rho V_\sigma)$	0
$\tilde{\lambda}$	-1	0	$-\frac{1}{4}\epsilon_{\rho\sigma}(F_{\rho\sigma} + i[V_\rho, V_\sigma] - 2i\mathcal{D}_\rho V_\sigma)$
λ_μ	1	$-\frac{1}{2}\mathcal{D}_\mu A_3 - \frac{i}{2}[V_3, V_\mu] - \frac{1}{2}[A_3, V_\mu] - \frac{i}{2}\mathcal{D}_\mu V_3$	0
$\tilde{\lambda}_\mu$	-1	0	$-\frac{1}{2}\mathcal{D}_\mu A_3 - \frac{i}{2}[V_3, V_\mu] - \frac{1}{2}[A_3, V_\mu] - \frac{i}{2}\mathcal{D}_\mu V_3$
λ_3	1	$-\frac{1}{2}[A_3, V_3] + \frac{i}{2}\mathcal{D}_\rho V_\rho + K$	G
$\tilde{\lambda}_3$	-1	$\tilde{G}$	$-\frac{1}{2}[A_3, V_3] + \frac{i}{2}\mathcal{D}_\rho V_\rho - K$
G	2	$[A_3, \lambda] - i\epsilon_{\rho\sigma}\mathcal{D}_\rho\lambda_\sigma + i[V_3, \lambda] + i\epsilon_{\rho\sigma}[V_\rho, \lambda_\sigma]$	0
$\tilde{G}$	-2	0	$[A_3, \tilde{\lambda}] + i\epsilon_{\rho\sigma}\mathcal{D}_\rho\tilde{\lambda}_\sigma + i[V_3, \tilde{\lambda}] - i\epsilon_{\rho\sigma}[V_\rho, \tilde{\lambda}_\sigma]$
K	0	$\frac{1}{2}[A_3, \tilde{\lambda}] + \frac{i}{2}\epsilon_{\rho\sigma}\mathcal{D}_\rho\tilde{\lambda}_\sigma + \frac{i}{2}[V_3, \tilde{\lambda}] - \frac{i}{2}\epsilon_{\rho\sigma}[V_\rho, \tilde{\lambda}_\sigma]$	$-\frac{1}{2}[A_3, \lambda] + \frac{i}{2}\epsilon_{\rho\sigma}\mathcal{D}_\rho\lambda_\sigma - \frac{i}{2}[V_3, \lambda] - \frac{i}{2}\epsilon_{\rho\sigma}[V_\rho, \lambda_\sigma]$

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