



HOKKAIDO UNIVERSITY

Title	Characterization of GaAs-Based Three-Branch Nanowire Junction Devices by Light-Induced Local Conductance Modulation Method
Author(s)	Sato, Masaki; Kasai, Seiya
Citation	Japanese Journal of Applied Physics(JJAP), 52(6S), 06GE08 https://doi.org/10.7567/JJAP.52.06GE08
Issue Date	2013-06
Doc URL	https://hdl.handle.net/2115/66634
Rights	© 2013 The Japan Society of Applied Physics
Type	journal article
File Information	JJAP2013_52_MN12077.pdf



Characterization of GaAs-Based Three-Branch Nanowire Junction Devices by Light-Induced Local Conductance Modulation Method

Masaki Sato and Seiya Kasai*

Graduate School of Information Science and Technology and Research Center for Integrated Quantum Electronics, Hokkaido University,
Sapporo 060-0814, Japan

Nonlinear voltage transfer characteristics in GaAs-based three-branch nanowire junction (TBJ) devices were investigated by a light-induced local conductance modulation method. In this measurement system, the conductance in the device was locally increased by focused laser light irradiation. The nonlinear transfer curve was greatly changed when the laser light was irradiated on the positively biased branch. The conductance domain was found to exist at the end of the positively biased branch of the TBJ by scanning the light position. When a SiN_x thin layer was deposited on the nanowire surface, the surface potential was increased and the nonlinearity in the transfer curve was reinforced simultaneously. The obtained results suggest that the asymmetric channel depletion model is appropriate for the observed nonlinearity mechanism in the GaAs TBJ at room temperature.

*E-mail: kasai@rciqe.hokudai.ac.jp

1. Introduction

The recent development of nanotechnology allows us to create various functional nanodevices. A semiconductor three-branch nano junction (TBJ) device is a typical and important example of such devices. The TBJ exhibits a unique nonlinear voltage transfer characteristic even with a simple structure¹⁾. Various analog and digital circuits integrating TBJs have been demonstrated thus far, including NOR logic gates²⁾, NAND logic gates³⁻⁶⁾, frequency mixers^{7, 8)}, half adders⁹⁾, rectifiers¹⁰⁾, and flip-flops^{11, 12)}. The TBJ also has the capability of ultrahigh-speed operation up to THz frequency¹³⁾. For the design and control of the device and circuits, understanding of the mechanism for the nonlinear characteristic is an important issue. A pioneering work by Xu indicates that the nonlinear characteristic appears when the TBJ is operated in the ballistic transport regime¹⁾. However, experimentally, it is also clearly observed at room temperature (RT)¹⁴⁻¹⁸⁾ where the carrier transport should be in the nonballistic transport regime. The mechanism in such a case has not been clarified yet, although several hypotheses have been introduced, including those regarding the effective mean free path extension¹⁶⁾ and the asymmetric channel depletion due to the surface potential¹⁹⁻²⁶⁾.

In this study, to identify the model for the nonlinear mechanism in the TBJ at room temperature, we characterized the GaAs-based TBJ device by local conductance modulation by focused laser light irradiation. This method revealed the conductance domain in the device, which is an important factor for understanding the nonlinear characteristic. We also fabricated and characterized a SiN_x-deposited TBJ to clarify the effect of the surface potential on the nonlinear characteristic in terms of channel depletion.

2. Device Operation and Possible Mechanisms

The TBJ with a typical measurement circuit is shown in Fig. 1(a). We consider the device with an n-type semiconductor. An equivalent circuit deduced from the structure is shown in Fig. 1(b). When voltage signals are applied to the left and right branches in push-pull fashion ($V_{\text{INL}} = -V_{\text{INR}}$), the output voltage V_{OUT} in the center branch shows a bell-shaped nonlinear voltage transfer curve, as shown in Fig. 1(c). To explain such a nonlinear behavior in terms of the equivalent circuit in Fig. 1(b), we consider the formula of V_{OUT} as

$$V_{\text{OUT}} = \frac{G_+ - G_-}{G_+ + G_-} V_{\text{IN}}, \quad (1)$$

where G_+ and G_- are the input branch conductance values with positive and negative biases, respectively. From the device configuration, the conductance in the input branches is expected to be the same and V_{OUT} is expected to be zero. In turn, to explain the nonlinear curve, we must accept that G_+ is always smaller than G_- .

The nonlinearity in the TBJ is understood by the asymmetric conductance that depends on the polarity of the input voltage. The original model explains it by the ballistic transport of electrons¹⁾. In this model, the electrons emitted from the negatively biased branch reach another branch without scattering, as shown in Fig. 2(a). As a result, the resistance of the negatively biased branch is zero. Then, the output voltage follows that of the negatively biased branch, namely, $V_{\text{OUT}} = -|V_{\text{IN}}|$. This model is applicable to the low-temperature operation where the mean free path is longer than the device size. However, the nonlinear characteristic is observed at RT, which cannot be explained using this model. At present, there are two models for such a case. The mean free path extension model indicates that the effective mean free path is extended at a high electric field because the carrier velocity increases while the carrier relaxation time remains

constant¹⁶⁾. On the other hand, the asymmetric channel depletion model indicates that the potential difference between the surface and the biased branch causes asymmetric depletion in the channel, which is similar to the channel pinch off in a field-effect transistor (FET) with large drain bias. In the case of the n-type semiconductor, the channel depletion occurs in the positively biased branch, as shown in Fig. 2(b)¹⁹⁻²⁶⁾. This means that G_+ is small. The III-V semiconductor nanowire surface is known to behave similarly to a metal gate because high-density surface states fix the surface Fermi level at a certain energy level. From the size dependence of the nonlinear curve²⁶⁾ and its asymmetric change in the gate-controlled TBJ devices^{17, 25)}, we suppose that this model is applicable to the nonlinear characteristic at RT. To verify the applicability of this model, the identification of the existence of the conductance domain and its portion is necessary. The effect of the surface should also be clarified.

3. Experimental Procedure

In this study, GaAs-based TBJ devices were fabricated and characterized. The nanowire channel was formed on an AlGaAs/GaAs heterostructure on a (001) semi-insulating GaAs substrate by EB lithography and wet chemical etching. A Ni/Ge/Au/Ni/Au ohmic contact was formed on each branch. Figure 3(a) shows a scanning electron microscopy (SEM) image of a fabricated device. The two-dimensional electron gas (2DEG) density was $7.8 \times 10^{11} \text{ cm}^{-2}$. The 2DEG mobility (μ_e) values were $7,100 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ at 300 K and $11,000 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ at 77 K. The corresponding mean free paths were 100 nm at 300 K and 1600 nm at 77 K. To identify the conductance domain position distinctly, we designed a relatively large device. The input nanowire length was 20 μm and the output nanowire length was 2.5 μm . The

nanowire directions were $\langle -110 \rangle$ for input branches and $\langle 110 \rangle$ for output branches. In the case of the GaAs channel, the electron transport properties including the mobility, mean free path, and density of states do not depend on the nanowire direction, since the Fermi surface of GaAs has a spherical symmetry and the effective mass of the electron is isotropic. The width and cross section of the nanowire mainly affect the conductivity. The nanowire widths were 300 nm for the input branch and 500 nm for the output branch. The nanowire was larger than the mean free path and the device operated in the nonballistic transport regime at RT. To clarify the effect of the surface potential, we also prepared devices with and without a SiN_x layer on their surface. The SiN_x layer is known to change the surface potential in GaAs-based materials^{27, 28}). The SiN_x layer was deposited by electron cyclotron resonance chemical vapor deposition (ECR-CVD) at a substrate temperature of 260 °C. The gas flow conditions were a SiH_4 gas flow rate of 30 sccm and a N_2 gas flow rate of 1 sccm. The measured film thickness was 30 nm.

To find the conductance domain and its effect on the nonlinear characteristic, we developed the light-induced local conductance modulation measurement system shown in Fig. 3(b). A focused laser light was irradiated on the device. The diameter of the laser spot was 5 μm and the wavelength was 658 nm. The laser light excited carriers and modulated the conductance locally. The sample position could be controlled with a 10 nm resolution in this system. All measurements were performed at RT.

4. Results and Discussion

4.1. Effect of focused light irradiation on branch conductance

To confirm the local conductance modulation in the developed system, we measured the I - V characteristics of the device in the beginning. Figure 4(a) shows the

measured I - V characteristics under dark and light irradiation conditions. The laser light was irradiated at the end of the left branch. The significant current increase was observed only when the laser light was irradiated on the positively biased branch. The results confirmed the local conductance modulation in the present system and also indicated the existence of the conductance domain in this portion. A small increase in current was also observed even in the branch without light irradiation, as shown in Fig. 4(b). This was because the photogenerated carriers in the laser-irradiated branch diffused to the current channel portion. The increase in conductance was 22 μ S. The estimated increase in carrier density, Δn , was $1.9 \times 10^{10} \text{ cm}^{-2}$. Theoretically, in the case of 658 nm light, the number of generated photons was $8.6 \times 10^{16} \text{ cm}^{-2} \text{ s}^{-1}$ considering the reflectance of 0.34 for GaAs and the laser power of 0.102 mW. The increase in carrier density is estimated using

$$\Delta n = \eta \tau f, \quad (2)$$

where τ is the carrier lifetime, η is the quantum efficiency, and f is the photon flux density. τ was estimated to be 105 ns by taking into account the surface recombination on the bulk structure²⁹⁾. Assuming $\eta = 30\%$, we obtained $\Delta n = 1.3 \times 10^{12} \text{ cm}^{-2}$ and the increase in conductance was estimated to be 1.5 mS. This value is two orders of magnitude larger than the experimental one. The difference is attributed to the small effective quantum efficiency due to the large surface recombination of the GaAs nanowire³⁰⁾.

4.2. Nonlinear voltage transfer characteristics

Figure 5 shows the measured V_{OUT} - V_{INL} characteristics with and without laser light irradiation. The output voltage increased when the laser light was irradiated on the

positively biased branch. In contrast, the output decreased when the laser light was irradiated on the negatively biased branch. The observed behaviors were explained by the increase in conductance in the laser-irradiated branch in terms of Eq. (1). The polarity of the output voltage was changed from negative to positive by the light on the positively biased branch. This was because the magnitude relationship between the conductance in the left and right branches was switched and the output branch potential followed the positively biased branch. The increase in conductance from the result was estimated to be 11 μS , which was in agreement with that from the I - V characteristics. This result indicates that the asymmetry of the input nanowire conductance caused the nonlinearity in the voltage transfer curve.

4.3 Conductance domain

To determine the position of the conductance domain and its relationship with the nonlinear curve, the laser light position dependences of the output voltage and the input-branch current I_{IN} were characterized. The measurement system is schematically shown in Fig. 6(a). In these measurements, constant positive and negative voltages were applied to the left and right branches, respectively ($V_{\text{INL}} = +1 \text{ V}$, $V_{\text{INR}} = -1 \text{ V}$). Figures 6(b) and 6(c) show the laser position dependences of V_{OUT} and I_{IN} , respectively. The output voltage increased as the laser position was moved from right ($x = 10 \mu\text{m}$) to left ($x = -10 \mu\text{m}$). The current also increased similarly to the output voltage. The highest voltage and current were observed at the end of the positively biased nanowire at $x = -10 \mu\text{m}$. These results indicate that the lowest conductance portion was the end of the positively biased nanowire. V_{OUT} showed a minimum value at $x = 2.5 \mu\text{m}$. This was because the conductance in the negatively biased branch, G_- , increased and it enhanced

the nonlinearity, as indicated by Eq. (1). V_{OUT} increased when the laser position was greater than 2.5 μm , since the light spot gradually left the nanowire and transferred to the wide channel portion, and the situation approached the dark condition.

We estimated the change in conductance in the two input branches as a function of the laser-irradiated position. Input branch conductance is expressed by

$$G_+ = \frac{I_{IN}}{V_{INL} - V_{OUT}}, \quad G_- = \frac{I_{IN}}{V_{OUT} - V_{INR}}. \quad (3)$$

The calculated result is shown in Fig. 7. Note that the largest conductance change occurs when the laser is irradiated on the conductance domain, since the lowest conductance dominates the current. Therefore, we could determine the position of the conductance domain from the plot in Fig. 7. The largest change in G_+ was obtained when the laser was irradiated at the end of the positively biased nanowire, $x = -10 \mu\text{m}$. Thus, the conductance domain existed at the end of the positively biased nanowire.

4.4 Mechanism for conductance domain formation

To clarify the mechanism of the formation of the conductance domain, we investigated the effect of the nanowire surface on the electric characteristics in the TBJ. Figure 8(a) shows the currents in the nanowire before and after the SiN_x deposition. The current decreased after the SiN_x deposition. This showed the increase in surface potential that promoted the channel depletion. Figure 8(b) shows the voltage transfer curves before and after the SiN_x deposition. The reinforcement of the nonlinearity was observed in the device after the SiN_x deposition. Therefore, the surface condition affected the nonlinearity in the TBJ. The study on the electric field concentration in MESFETs³¹⁾ showed that the large surface potential of the channel semiconductor shifted the position of the high field to the drain electrode. In the case of the TBJ, the

increased surface potential enhanced the channel depletion in the positively biased branch similar to the drain in a FET and the conduction was decreased and the nonlinearity was reinforced. Considering the position of the conductance domain, the results suggest that the asymmetry channel depletion model is feasible for the nonlinear characteristic of the TBJ in the nonballistic transport regime.

5. Conclusion

Nonlinear voltage transfer characteristics in GaAs-based three-branch nanowire junction (TBJ) devices were investigated by a light-induced local conductance modulation method. The nonlinear transfer curve greatly changed only when the laser was irradiated on the positively biased branch. The conductance domain was found to exist at the end of the positively biased branch of the TBJ. A SiN_x layer deposited on the TBJ increased the surface potential and reinforced the nonlinearity in the transfer curve. The obtained results indicate that the asymmetric channel depletion model is appropriate for the observed nonlinearity mechanism in the GaAs TBJ at room temperature.

Acknowledgements

We thank Dr. T. Maemoto of Osaka Institute of Technology for various useful discussions. We also thank Professor T. Fukui and Professor T. Hashizume of Hokkaido University for their continuous support. This work was partly supported by a Grant-in-Aid for Scientific Research (No. 22310084), the Japan Society for the Promotion of Science.

References

- 1) H. Q. Xu: Appl. Phys. Lett. **78** (2001) 2064.
- 2) C. R. Müller, L. Worschech, P. Höpfner, S. Höfling, and A. Forchel: IEEE Electron Device Lett. **28** (2007) 859.
- 3) H. Q. Xu, I. Shorubalko, D. Wallin, I. Maximov, P. Omling, L. Samuelson, and W. Seifert: IEEE Electron Device Lett. **25** (2004) 164.
- 4) H. Q. Xu: Nat. Mater. **4** (2005) 649.
- 5) S. Reitzenstein, L. Worschech, C. R. Müller, and A. Forchel: IEEE Electron Device Lett. **26** (2005) 142.
- 6) S. F. Abd Rahman, D. Nakata, Y. Shiratori, and S. Kasai: Jpn. J. Appl. Phys. **48** (2009) 06FD01.
- 7) J. Sun, D. Wallin, P. Brushiem, I. Maximov, Z. G. Wang, and H. Q. Xu: Nanotechnology **18** (2007) 195205.
- 8) I. Shorubalko, H. Q. Xu, I. Maximov, D. Nilsson, P. Omling, L. Samuelson, and W. Seifert: IEEE Electron Device Lett. **23** (2002) 377.
- 9) S. Reitzenstein, L. Worschech, and A. Forchel: IEEE Electron Device Lett. **25** (2004) 462.
- 10) L. Bednarz, Rashmi, B. Hackens, G. Farhi, V. Bayot, and I. Huynen: IEEE Trans. Electron Devices **4** (2005) 576.
- 11) J. Sun, D. Wallin, I. Maximov, and H. Q. Xu: IEEE Electron Device Lett. **29** (2008) 540.
- 12) H. Shibata, Y. Shiratori, and S. Kasai: Jpn. J. Appl. Phys. **50** (2011) 06GF03.
- 13) S. Bollaert, A. Cappy, Y. Roelens, J. S. Gallo, C. Gardes, Z. Teukam, X. Wallart, J. Mateos, T. González, B. G. Vasallo, B. Hackens, L. Berdnarz, and I. Huynen: Thin

Solid Films **515** (2007) 4321.

- 14) I. Shorubalko, H. Q. Xu, I. Maximov, P. Omling, L. Samuelson, and W. Seifert: Appl. Phys. Lett. **79** (2001) 1384.
- 15) I. Maximov, P. Carlberg, I. Shorubalko, D. Wallin, E-L. Sarwe, M. Beck, M. Graczyk, W. Seifert, H.Q. Xu, L. Montelius, and L. Samuelson: Microelectron. Eng. **67** (2003) 196.
- 16) D. Wallin, I. Shorubalko, H. Q. Xu, and A. Cappy: Appl. Phys. Lett. **89** (2006) 092124.
- 17) T. Nakamura, S. Kasai, Y. Shiratori, and T. Hashizume: Appl. Phys. Lett. **90** (2007) 102104.
- 18) M. Frimmer, J. Sun, I. Maximov, and H. Q. Xu: Appl. Phys. Lett. **93** (2008) 133110.
- 19) J. Mateos, B. G. Vsasallo, D. Pardo, T. González, J.-S. Galloo, S. Bollaert, Y. Roelens, and A. Cappy: Nanotechnology **14** (2003) 117.
- 20) J. Mateos, B. G. Vsasallo, D. Pardo, T. González, E. Pichonat, J.-S. Galloo, S. Bollaert, Y. Roelens, and A. Cappy: IEEE Electron Device Lett. **25** (2004) 235.
- 21) J. Mateos, B. G. Vsasallo, D. Pardo, T. González, J.-S. Galloo, S. Bollaert, Y. Roelens, and A. Cappy: IEEE Trans. Electron Devices **50** (2003) 1897.
- 22) I. Iñiguez-de-la-Torre, J. Mateos, T. González, D. Pardo, J.-S. Galloo, S. Bollaert, Y. Roelens, and A. Cappy: Semicond. Sci. Technol. **22** (2007) 663.
- 23) T. Sadi, F. Dessenne, and J. Thobel: J. Appl. Phys. **105** (2009) 053707.
- 24) I. Iñiguez-de-la-Torre, J. Mateos, Y. Roelens, C. Gardés, S. Bollaert, and T. González: Nanotechnology **22** (2011) 445203.
- 25) S. Kasai, T. Nakamura, S. F. Abd Rahman, and Y. Shiratori: Jpn. J. Appl. Phys. **49**

(2008) 4958.

- 26) D. Nakata, H. Shibata, Y. Shiratori¹, and S. Kasai: Jpn. J. Appl. Phys. **49** (2010) 06GG03.
- 27) E. Y. Chang, G. T. Cibuzar, and K. P. Pande: IEEE Trans. Electron Devices **35** (1988) 1412.
- 28) S. Pal and D. N. Bose: Appl. Surf. Sci. **181** (2001) 179.
- 29) J. M. Borrego and S. K. Ghandhi: Solid-State Electron. **33** (1990) 733.
- 30) C. Chang, C. Chi, M. Yao, N. Huang, C. Chen, J. Theiss, A. W. Bushmaker, S. LaLumondiere, T. Yeh, M. L. Povinelli, C. Zhou, P. D. Dapkus, and S. B. Cronin: Nano Lett. **12** (2012) 4484.
- 31) H. Mizuta, K. Yamaguchi, and S. Takahashi: IEEE Trans. Electron Devices **34** (1987) 2027.

Figure captions

Fig. 1. Three-branch nanowire junction (TBJ) device: (a) measurement circuit, (b) equivalent circuit of the device, and (c) typical voltage transfer characteristic.

Fig. 2. (a) Ballistic transport model and (b) asymmetric channel depletion model.

Fig. 3. (a) SEM image of a fabricated GaAs TBJ device and (b) a light-induced local conductance modulation measurement system.

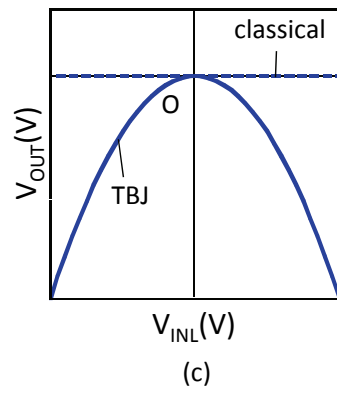
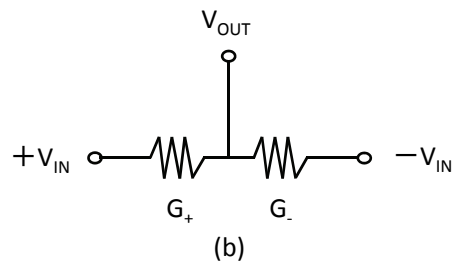
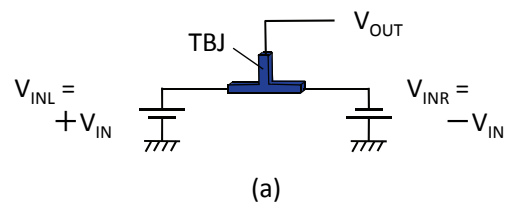
Fig. 4. I - V characteristics with laser light irradiation: (a) irradiation on the current path side and (b) irradiation on the outside of the path.

Fig. 5. V_{OUT} - V_{INL} characteristics with and without the laser light irradiation.

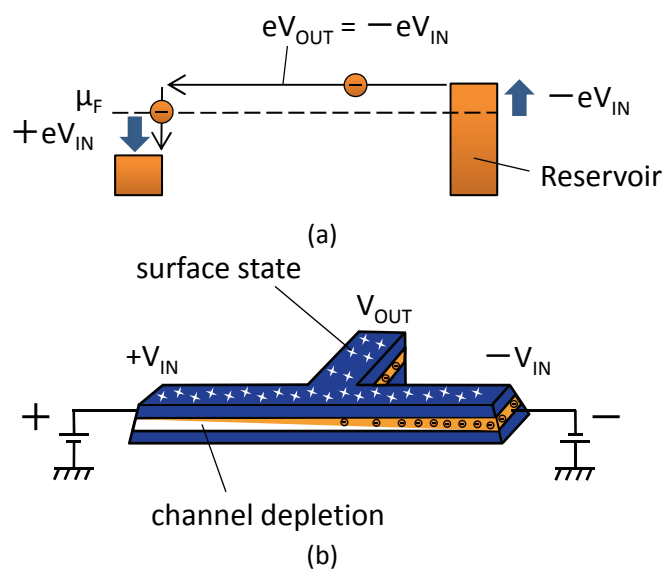
Fig. 6. (a) Measurement system for conductance domain identification and (b) laser light positional dependences of (b) V_{OUT} and (c) I_{IN} .

Fig. 7. Evaluated conductance values in the left and right branches.

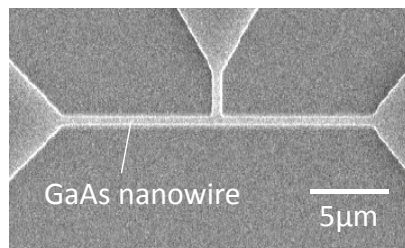
Fig. 8. (a) I - V and (b) V_{OUT} - V_{INL} characteristics before and after SiN_x deposition.



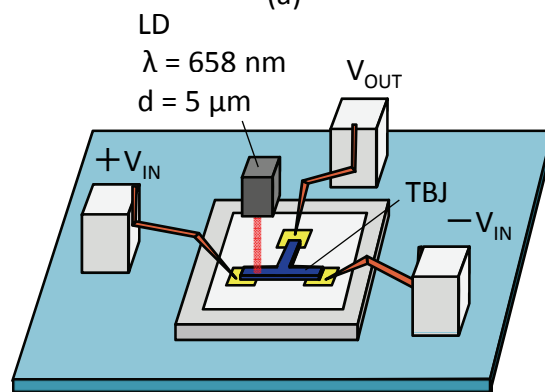
Sato et al., Figure. 1.



Sato et al., Figure. 2.

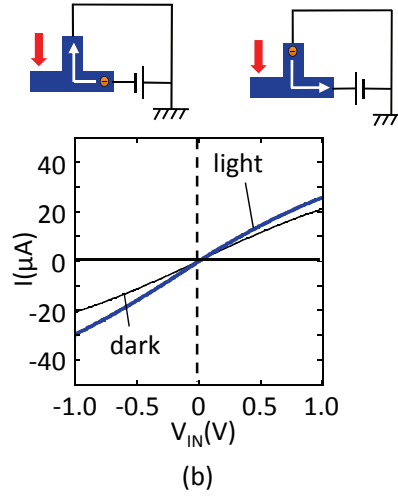
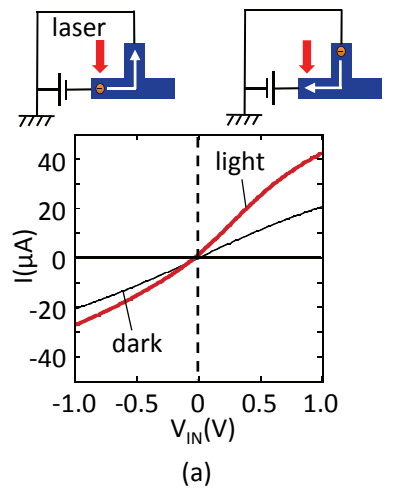


(a)

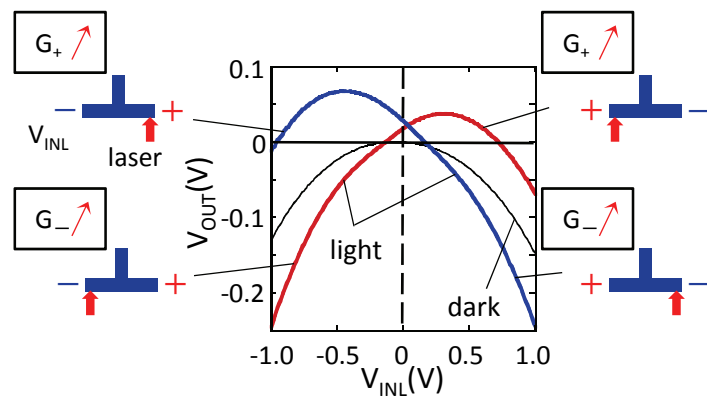


(b)

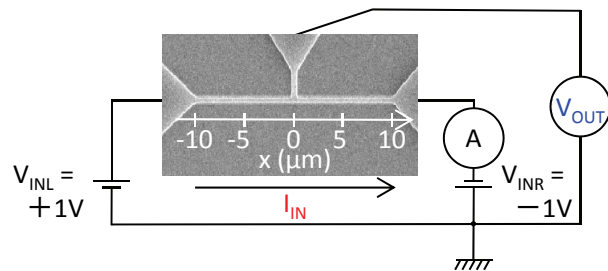
Sato et al., Figure. 3.



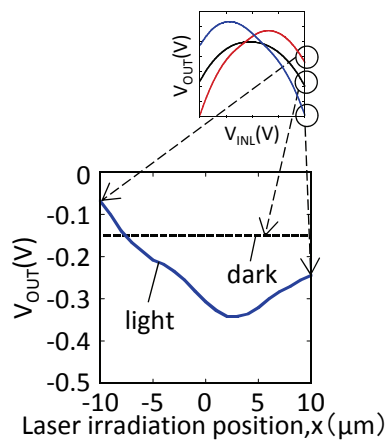
Sato et al., Figure. 4.



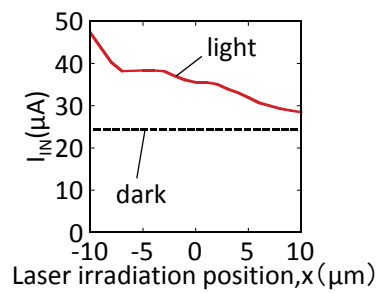
Sato et al., Figure. 5.



(a)

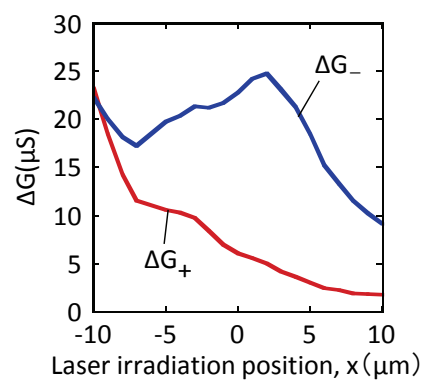


(b)

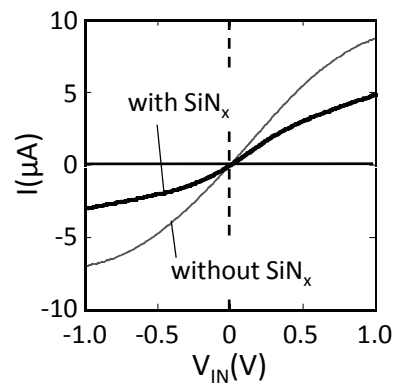


(c)

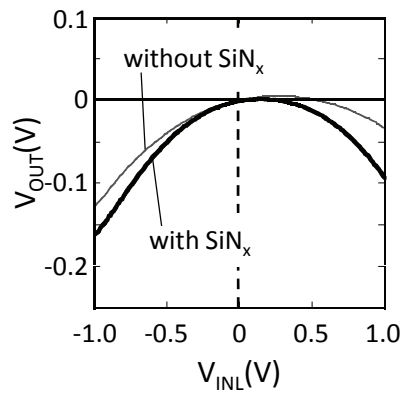
Sato et al., Figure. 6.



Sato et al., Figure. 7.



(a)



(b)

Sato et al., Figure. 8.