



Title	Experimental Evidence of Cubic Rashba Effect in an Inversion-Symmetric Oxide
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Supplemental Material for “Experimental Evidence of Cubic Rashba Effect in an Inversion-Symmetric Oxide”

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In this Supplemental Material, we present details of the localization/antilocalization analysis using the Iordanskii, Lyanda-Geller, Pikus (ILP) theory. The general formula used for a theoretical fit is as follows [S1,S2].

$$\begin{aligned} \Delta\sigma(B) = & -\frac{e^2}{4\pi^2\hbar} \left[\frac{1}{a_0} + \frac{2a_0 + 1 + \frac{B_{\text{SO1}} + B_{\text{SO3}}}{B}}{a_1 \left(a_0 + \frac{B_{\text{SO1}} + B_{\text{SO3}}}{B} \right)} - 2 \frac{B_{\text{SO1}}}{B} \right. \\ & - \sum_{n=0}^{\infty} \left\{ \frac{3}{n} - \frac{3a_n^2 + 2a_n \frac{B_{\text{SO1}} + B_{\text{SO3}}}{B} - 1 - 2(2n+1) \frac{B_{\text{SO1}}}{B}}{\left(a_n + \frac{B_{\text{SO1}} + B_{\text{SO3}}}{B} \right) a_{n-1} a_{n+1} - 2 \frac{B_{\text{SO1}}}{B} [(2n+1)a_n - 1]} \right\} \\ & \left. + \Psi \left(\frac{1}{2} + \frac{B_\phi}{B} \right) \right], \end{aligned} \quad (\text{S1})$$

where

$$\begin{aligned} a_n = n + \frac{1}{2} + \frac{B_\phi}{B} + \frac{B_{\text{SO1}} + B_{\text{SO3}}}{B}, \quad B_\phi = \frac{\hbar}{4eD\tau_\phi}, \\ B_{\text{SO1}} = \frac{\hbar}{4eD} \cdot |\mathbf{\Omega}_1|^2 \tau_1, \quad B_{\text{SO3}} = \frac{\hbar}{4eD} \cdot |\mathbf{\Omega}_3|^2 \tau_3, \quad \frac{1}{\tau_m} = \int (1 - \cos m\theta) W(\theta) d\theta, \end{aligned} \quad (\text{S2})$$

Ψ is the digamma-function, D is the diffusion constant, τ_ϕ is the phase relaxation time, the index $m = 1$ or 3 , and $W(\theta)$ is the probability of scattering by an angle θ . Note that the terms in the original ILP theory which merely shift the overall curve in the y -direction is omitted. To calculate Eq. S1, the infinite sum was added up to $n = 2 \times 10^5$.

When $B_{\text{SO1}} = 0$ ($|\mathbf{\Omega}_1| = 0$), Eq. S1 has the following closed form [S2]:

$$\begin{aligned} \Delta\sigma(B) = & \frac{e^2}{2\pi^2\hbar} \left\{ \Psi \left(\frac{1}{2} + \frac{B_\phi}{B} + \frac{B_{\text{SO3}}}{B} \right) + \frac{1}{2} \Psi \left(\frac{1}{2} + \frac{B_\phi}{B} + 2 \frac{B_{\text{SO3}}}{B} \right) \right. \\ & \left. - \frac{1}{2} \Psi \left(\frac{1}{2} + \frac{B_\phi}{B} \right) - \ln \frac{B_\phi + B_{\text{SO3}}}{B} - \frac{1}{2} \ln \frac{B_\phi + 2B_{\text{SO3}}}{B} + \frac{1}{2} \ln \frac{B_\phi}{B} \right\}. \end{aligned} \quad (\text{S3})$$

In the main text, fitting was performed for the condition $|\mathbf{\Omega}_1| = 0$ or $|\mathbf{\Omega}_3| = 0$, leaving two free parameters B_ϕ and B_{SO} (B_{SO1} or B_{SO3}). Here, we analyze our data retaining all three fitting parameters (B_ϕ , B_{SO1} , and B_{SO3}). For this purpose, we first define the ratio $r = B_{\text{SO1}}/B_{\text{SO3}}$. Second, for several values of r which is set between 0 (*i.e.*, $|\mathbf{\Omega}_1| = 0$) and infinity ($|\mathbf{\Omega}_3| = 0$), parameters are adjusted to fit the data around $B = 0$. The result as well as the parameters used for the fitting are shown in Fig. S1 and Table SI, respectively.

TABLE SI: Parameters used for the fitting shown in Fig. S1

$r = B_{\text{SO1}}/B_{\text{SO3}}$	B_ϕ (mT)	B_{SO1} (mT)	B_{SO3} (mT)
0	13	0	34
0.2	13	6.6	33
0.5	13	15.5	31
1	13	27.5	27.5
2	13	43	21.5
∞	13	78	0

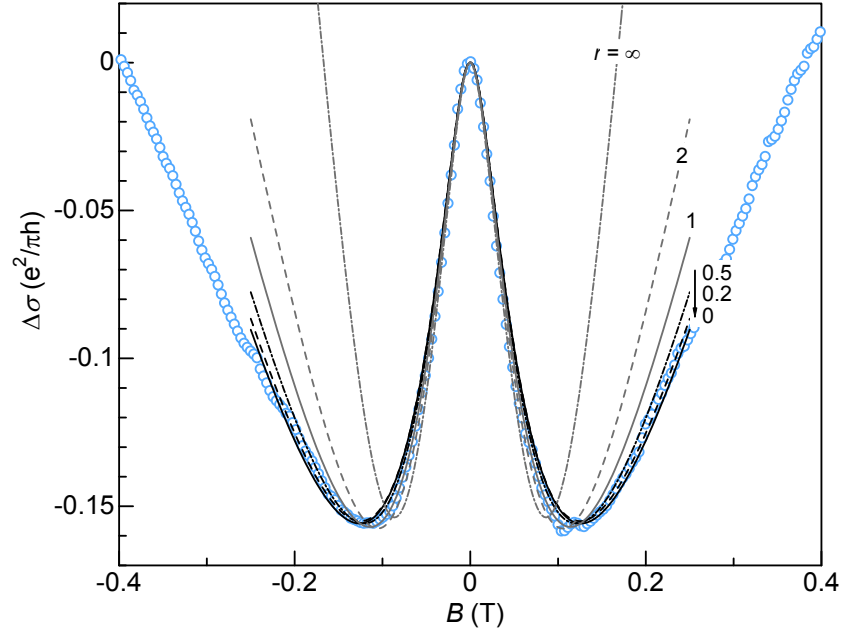


FIG. S1: Theoretical fits (curves) to experimental data $\Delta\sigma$ (open circles) of STO-B at $E_G = 4.38$ MV/cm given by the ILP theory. Several values of $r = \frac{B_{\text{SO1}}}{B_{\text{SO3}}}$ are considered. The theoretical curves approach the experimental data as $r \rightarrow 0$, *i.e.*, as $|\Omega_1| \rightarrow 0$.

Next, we show the results of changing B_ϕ for each value of r with fixed B_{SO1} and B_{SO3} (Figs.S2-S6). In these fittings, parameters are adjusted so that the magnetic field at the minimum of theoretical $\Delta\sigma$ curves coincides with that of experimental data.

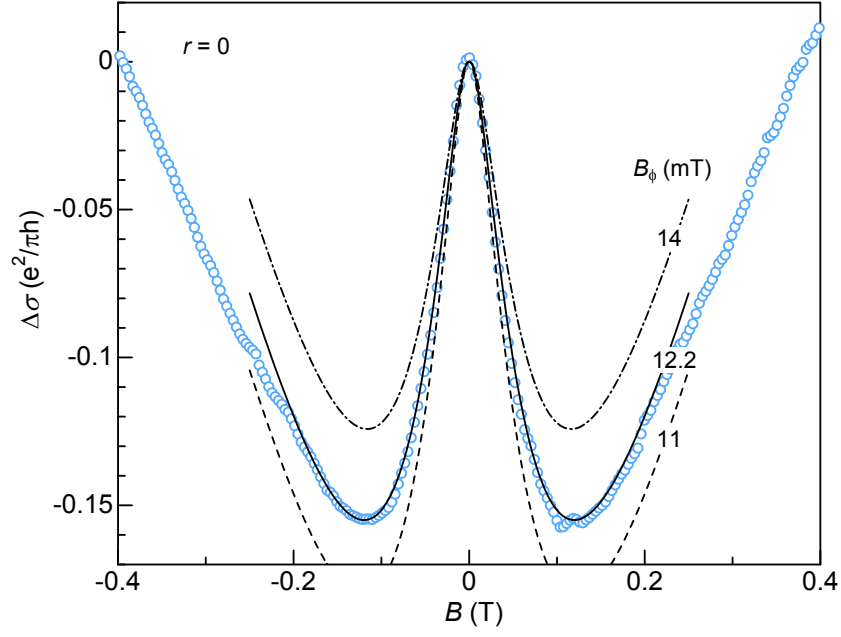


FIG. S2: The result of changing B_ϕ for $r = 0$ with the following fixed parameters: $B_{\text{SO1}}=0$ mT, $B_{\text{SO3}}=32.1$ mT.

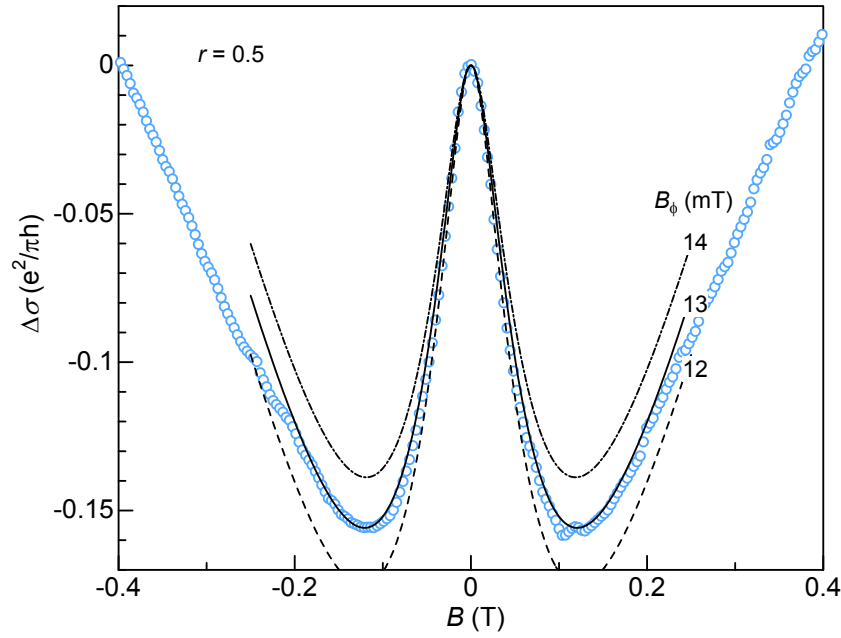


FIG. S3: The result of changing B_ϕ for $r = 0.5$ with the following fixed parameters: $B_{\text{SO1}}=15.5$ mT, $B_{\text{SO3}}=31$ mT.

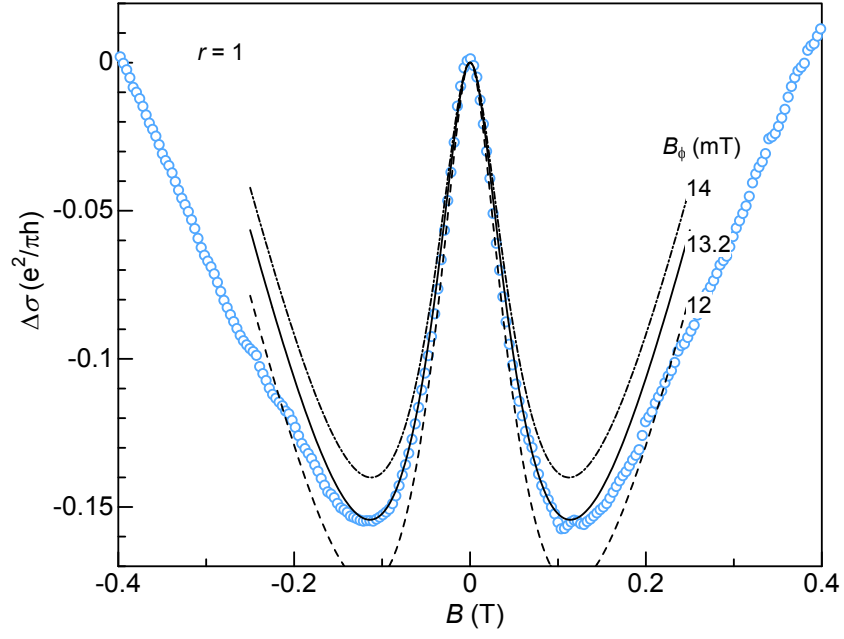


FIG. S4: The result of changing B_ϕ for $r = 1$ with the following fixed parameters: $B_{\text{SO}1}=27.5$ mT, $B_{\text{SO}3}=27.5$ mT.

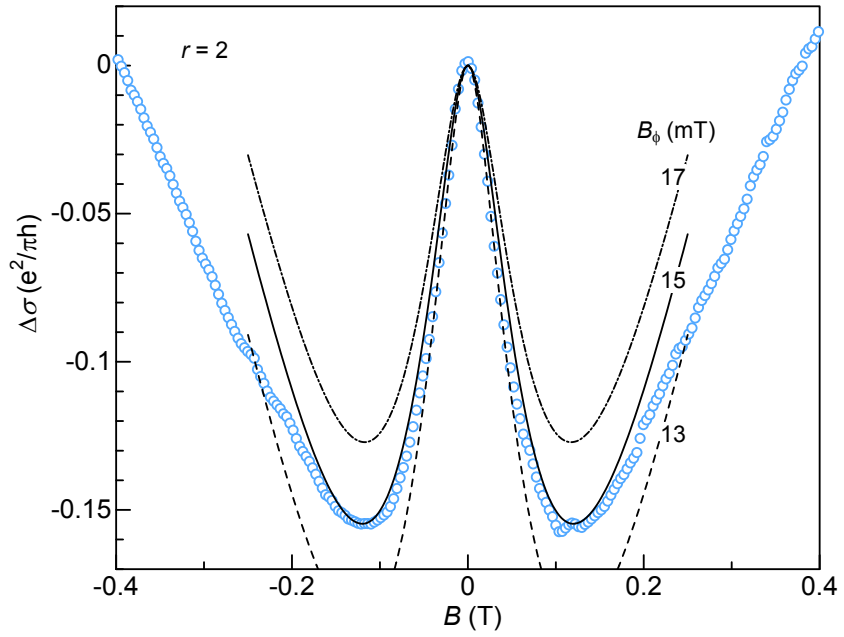


FIG. S5: The result of changing B_ϕ for $r = 2$ with the following fixed parameters: $B_{\text{SO}1}=49$ mT, $B_{\text{SO}3}=24.5$ mT.

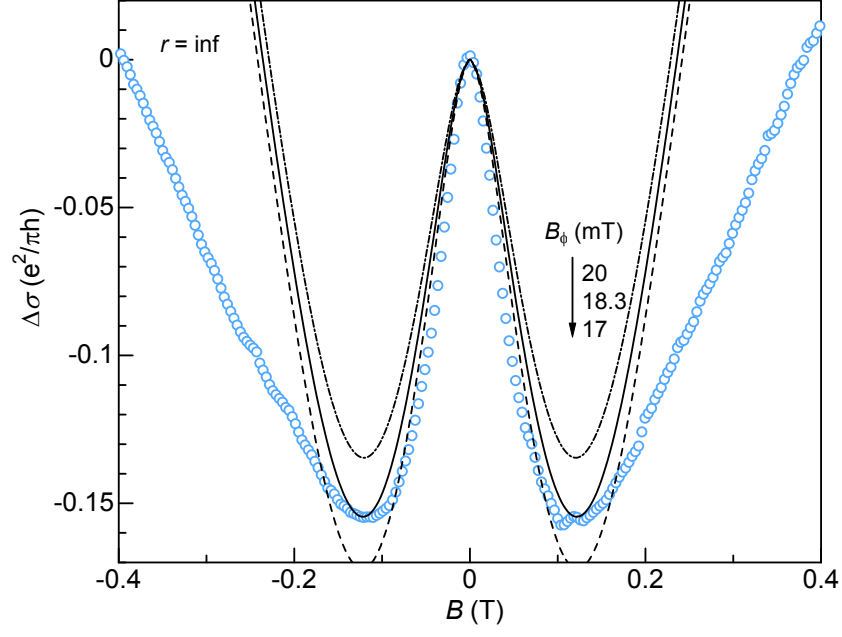


FIG. S6: The result of changing B_ϕ for $r = \infty$ with the following fixed parameters: $B_{\text{SO}1}=110$ mT, $B_{\text{SO}3}=0$ mT.

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- [S1] S. V. Iordanskii, Y. B. Lyanda-Geller, and G. E. Pikus, JETP Lett. **60**, 206 (1994).
[S2] W. Knap, *et al.*, Phys. Rev. B **53**, 3912 (1996).