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PROPAGATING FRONT SOLUTIONS IN A TIME-FRACTIONAL FISHER-KPP EQUATION

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ABSTRACT. In this paper, we treat the Fisher-KPP equation with a Caputo-type time fractional derivative and discuss the propagation speed of the solution. The equation is a mathematical model that describes the processes of sub-diffusion, proliferation, and saturation. We first consider a traveling wave solution in studying the propagation of the solution, but we cannot define it in the usual sense, owing to the time fractional derivative in the equation. We therefore assume that the solution asymptotically approaches a traveling wave solution, and the asymptotic traveling wave solution is formally introduced as a potential asymptotic form of the solution. The existence and properties of the asymptotic traveling wave solution are discussed using a monotone iteration method. Finally, the behavior of the solution is analyzed by conducting numerical simulations based on the results for asymptotic traveling wave solutions.

1. INTRODUCTION

In this paper, we consider the propagation of solutions to a time-fractional evolution equation

$$(1.1) \quad \partial_t^\alpha u = u_{xx} + f(u) \quad (t > 0, x \in \mathbb{R}),$$

where $u = u(t, x) \in \mathbb{R}$ ($t > 0, x \in \mathbb{R}$), and $f : \mathbb{R} \rightarrow \mathbb{R}$ is a monostable nonlinearity

$$f \in C^1(\mathbb{R}), \quad f(0) = f(1) = 0, \quad f(u) > 0 \quad (0 < u < 1),$$

satisfying the KPP-type condition (named after Kolmogorov, Petrovsky, and Piskunov [27])

$$f(u) \leq f'(0)u \quad (0 < u < 1).$$

One example is $f(u) = u(1 - u)$. Further, ∂_t^α is a Caputo derivative defined as

$$\partial_t^\alpha u(t, x) := \frac{1}{\Gamma(1 - \alpha)} \int_0^t \frac{\partial u}{\partial s}(s, x) \frac{ds}{(t - s)^\alpha},$$

where $\alpha \in (0, 1)$ and $\Gamma(1 - \alpha)$ is the gamma function with order $1 - \alpha$.

The equation with $\alpha = 1$ corresponds to the Fisher-KPP equation

$$(1.2) \quad u_t = u_{xx} + f(u) \quad (t > 0, x \in \mathbb{R}),$$

which is a well-known mathematical model used in the study of population dynamics [14, 15, 27]. A typical structure of propagation phenomena is the traveling wave solution.

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In the general case, a solution u is called a traveling wave solution if there exist a speed $c > 0$ and a wave profile $\phi \in C(\mathbb{R})$ such that $u(t, x) = \phi(x + ct)$. In the case of the Fisher-KPP equation, by substituting this expression into the equation and introducing the moving coordinate $\xi := x + ct$, $(c, \phi(\xi))$ must satisfy

$$c\phi'(\xi) = \phi''(\xi) + f(\phi(\xi)) \quad (\xi \in \mathbb{R}).$$

As a pioneering work, Fisher [15] considered the case that $f(u) = u(1 - u)$ and found that there is $c_1^* > 0$ such that (1.2) has a monotone traveling wave solution connecting from 0 to 1 when $c \geq c_1^*$, and no such solution when $c \in (0, c_1^*)$. In addition, Kolmogorov, Petrovsky, and Piskunov [27] revealed the existence of a traveling wave solution and its long-time behavior under more general conditions for $f(u)$. Building on this work, studies have expanded into various areas, such as traveling wave solutions (e.g. [31]), long-time behavior (e.g. [3]), and entire solutions (e.g. [26]). Here, we note that $c_1^* = 2\sqrt{f'(0)}$ holds in the case of KPP-type nonlinearity, and it is known that the propagation of the solution proceeds with the asymptotic highest term c_1^*t under suitable conditions [3, 27].

Meanwhile, research on different diffusion processes has expanded into various contexts. One extension is a spatially nonlocal model of diffusion

$$(1.3) \quad u_t = D[u] + f(u) \quad (t > 0, x \in \mathbb{R}).$$

Examples of the diffusion term $D[u]$ are the fractional Laplacian [6]:

$$D[u] = -(-\Delta)^s u \quad (0 < s < 1)$$

and nonlocal diffusion [2, 24]:

$$D[u] = \int_{\mathbb{R}} K(x - y)(u(t, y) - u(t, x))dy \quad (K \in L^1(\mathbb{R})).$$

In the case of nonlocal diffusion with a rapidly decaying kernel $K(x)$, it has been established that results of traveling wave solutions similar to (1.2) can be obtained [9, 11, 12, 34, 38]. In the case of the fractional Laplacian and nonlocal diffusion with slowly decaying kernels, there has been widespread analysis of traveling wave solutions and solution propagation [4, 5, 22, 35, 36]. In such cases with a monostable nonlinearity such as $f(u) = u(1 - u)$, it is known that (1.3) has no traveling wave solutions [7, 38]. Moreover, the long-time behavior of the solution has been studied, revealing that the solution propagates more rapidly than $O(t)$ [7, 21]. As it is not trivial that solution propagation occurs at $O(t)$ as the cited studies show, it is sometimes questionable whether it is appropriate to consider traveling wave solutions.

Returning to (1.1), the diffusion equation is modified to include a time-fractional derivative,

$$\partial_t^\alpha u = u_{xx} \quad (t > 0, x \in \mathbb{R}),$$

which is called the diffusion-wave equation or the time-fractional diffusion equation. The fundamental solution $G_\alpha(t, x)$ is defined by

$$G_\alpha(t, x) := \frac{1}{\sqrt{\pi|x|}} H_{1,2}^{2,0} \left[\frac{x^2}{4t^\alpha} \middle| \begin{matrix} (1, \alpha) \\ (1/2, 1) \end{matrix} \right. \left. \begin{matrix} (1, 1) \end{matrix} \right],$$

where $H_{p,q}^{\mu,\nu}(z)$ is Fox's H-function [13]. According to the properties of Fox's H function, the spatial decay rate is given by

$$G_\alpha \left(1, \frac{2x}{\alpha} \right) \simeq C_\alpha x^{(\alpha-1)/(2-\alpha)} \exp \left(-\frac{2-\alpha}{\alpha} x^{2/(2-\alpha)} \right)$$

(see [13, 32]). It is known that the diffusion is slower than local diffusion [13, 28]. In fact, as G_α satisfies $G_\alpha(t, x) = t^{-\alpha/2} G_\alpha(1, t^{-\alpha/2} x)$, the variance of the fundamental solution is

$$\int_{\mathbb{R}} x^2 G_\alpha(t, x) dx = \left(\int_{\mathbb{R}} x^2 G_\alpha(1, t^{-\alpha/2} x) dx \right) t^{-\alpha/2} = \left(\int_{\mathbb{R}} x^2 G_\alpha(1, x) dx \right) t^\alpha.$$

Therefore, the variance grows at a rate slower than $O(t)$ for the usual diffusion equation for sufficiently large times $t > 0$. In addition, it becomes slower as α decreases. To analyze phenomena with such sub-diffusion and nonlinearity, Equation (1.1) and similar models have been proposed (e.g., [16, 19, 20, 23, 37]). In particular, model (1.1) was proposed by [23] as a model of anomalous diffusion with linear reaction dynamics. In our context with $f(u) = u(1 - u)$, model (1.1) is a natural extension of their model with a newly added saturation effect.

In analyzing the solution to equation (1.1), we should keep in mind that the dynamical properties change owing to the time-fractional derivative. As an example, let us consider the time-fractional Malthus model as the simplest model:

$$\partial_t^\alpha v = \zeta v \quad (t > 0),$$

where $v = v(t) \in \mathbb{R}$, and ζ is a non-zero constant. The solution is given by

$$v(t) = v(0) E_{\alpha,1}(\zeta t^\alpha),$$

where $E_{\alpha,\beta}(z) := \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + \beta)}$ is the Mittag-Leffler function [25, 28]. Moreover, the asymptotic behavior of $E_{\alpha,1}$ is

$$(1.4) \quad E_{\alpha,1}(z) \simeq \frac{1}{\alpha} \exp(z^{1/\alpha}) - \frac{1}{z \Gamma(1 - \alpha)} + O(z^{-2}) \quad (|z| \rightarrow +\infty)$$

(see Section 6 in [25]). Hence, $v(t)$ converges to 0 with polynomial order when $\zeta < 0$, and grows exponentially when $\zeta > 0$. These properties often appear in time-fractional nonlinear equations when analyzing the decay and shape of characteristic solutions (e.g. [28, 37]).

In this paper, we consider propagating solutions of the front type, such as the traveling wave solutions we have been discussing so far. Long-time behavior has been considered in various settings for propagation and saturation processes under slow diffusion [16, 17, 19, 20]. However, to the best of our knowledge, no characterization of propagation solutions based on the discussion of traveling wave solutions has been done for equation (1.1) with a monostable nonlinearity. Therefore, this paper considers the relationship between the solution behavior and α by introducing the notion of a suitable solution like a traveling wave solution.

The paper is organized as follows. In Section 2, we introduce the concept of the traveling wave structure, termed the asymptotic traveling wave solution, and present the main result regarding the existence. The proof is given in Section 3. We then present numerical simulations under appropriate conditions in Section 4. Finally, we highlight unresolved issues in Section 5.

2. MAIN RESULTS

In analyzing the propagation of the solution, the concept of a traveling wave solution is considered. However, the usual method of obtaining traveling wave solutions cannot be immediately applied in the case of time-fractional equations. For $c > 0$ and $\phi \in C^2(\mathbb{R})$, the Caputo derivative of a traveling wave solution $u(t, x) = \phi(x + ct)$ is

$$\begin{aligned}\partial_t^\alpha u(t, x) &= \frac{c}{\Gamma(1-\alpha)} \int_0^t \frac{\phi'(x+cs)}{(t-s)^\alpha} ds \\ &= \frac{c}{\Gamma(1-\alpha)} \int_0^t \frac{\phi'(x+c(t-s))}{s^\alpha} ds \\ &= \frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{ct} \frac{\phi'(x+ct-s)}{s^\alpha} ds.\end{aligned}$$

We recall the moving coordinate $\xi = x + ct$ to obtain

$$(2.1) \quad \partial_t^\alpha u(t, x) = \frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{ct} \frac{\phi'(\xi-s)}{s^\alpha} ds.$$

This means that the term coming out of the time derivative depends on t . Hence, it is impossible to define a traveling wave solution in the classical sense.

On the other hand, in [37], they considered traveling wave solutions whose solutions might be approached asymptotically to investigate the final state of a front-type solution that connects different stable states. Based on this idea, we define an asymptotic traveling wave solution to analyze the final state. That is to say, the asymptotic traveling wave solution is given in the following form by setting $t \rightarrow +\infty$ in (2.1).

Definition 2.1. $(c, \phi) \in (0, +\infty) \times C^2(\mathbb{R})$ is an **asymptotic traveling wave solution of (1.1)** if it satisfies

$$(2.2) \quad \frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{\phi'(\xi-s)}{s^\alpha} ds = \phi''(\xi) + f(\phi(\xi))$$

for all $\xi \in \mathbb{R}$.

In particular, we focus on a front-type asymptotic traveling wave solution connecting two constant states:

$$(2.3) \quad \phi(-\infty) = 0, \quad \phi(+\infty) = 1.$$

Studies relating to equation (2.2) have been reported. In [33], the bistable nonlinear term $f(u)$, which is piecewise linear, was considered, and an attempt was made to analyze (c, ϕ) by describing its solution connecting two stable states explicitly. Furthermore, in [1, 10], traveling wave solutions in a nonlocal Korteweg–de Vries–Burgers equation, which

are analogous to equation (2.2), were considered, and their existence was explored by investigating the properties of the linear equation and constructing exponentially decaying solutions. In the case of the monostable nonlinear term $f(u)$ that we deal with, there are no results on the existence of solutions to equation (2.2) to the best of our knowledge.

Thus, first consider the existence of the asymptotic traveling wave solutions. Before stating the main theorem, we give the condition for $f(u)$ as

$$(2.4) \quad \begin{cases} f \in C^1(\mathbb{R}), f(0) = f(1) = 0, \\ \exists M > 0, \exists a > 0 \text{ s.t. } -Mu^{1+a} \leq f(u) - f'(0)u \leq 0 \ (0 \leq u \leq 1). \end{cases}$$

The condition includes the case that $f(u) = u(1-u)$. We then obtain the following theorem.

Theorem 2.2. *Let $\alpha \in (0, 1)$ and $c_\alpha^* := \frac{2^{1/\alpha}}{\sqrt{\alpha}} \left(\frac{f'(0)}{2-\alpha} \right)^{(2-\alpha)/2\alpha}$. Then, for any $c \geq c_\alpha^*$ there exists an asymptotic traveling wave solution of (1.1) with an increasing function ϕ satisfying (2.3). Furthermore, for $c > c_\alpha^*$, it can be normalized by*

$$\lim_{\xi \rightarrow -\infty} e^{-\lambda_1 \xi} \phi(\xi) = 1,$$

where λ_1 is the minimal positive root of $\lambda^2 - (c\lambda)^\alpha + f'(0) = 0$.

We have established the monotonicity of asymptotic traveling wave solutions and the exponential decay as $\xi \rightarrow -\infty$. Whether c_α^* is the minimum speed remains unresolved and is a problem to be considered in future work. The proof is given in Section 3.

Owing to the monotonicity, we obtain the correspondence of asymptotic traveling wave solutions to the time evolution equation (1.1).

Corollary 2.3. *For $\alpha \in (0, 1)$ and $c \geq c_\alpha^*$, we set $v(t, x) = \phi(x + ct)$, where ϕ is constructed in Theorem 2.2. Then, $v(t, x)$ is a sub-solution of (1.1), that is, $v(t, x)$ satisfies*

$$\partial_t^\alpha v < v_{xx} + f(v) \quad (t > 0, x \in \mathbb{R}).$$

Proof. Fix $\alpha \in (0, 1)$ and $c \geq c_\alpha^*$. Then, from (2.1) and the definition of the asymptotic traveling wave solution, we have

$$\partial_t^\alpha v - (v_{xx} + f(v)) = -\frac{c^\alpha}{\Gamma(1-\alpha)} \int_{ct}^{+\infty} \frac{\phi'(\xi-s)}{s^\alpha} ds < 0$$

for all $t > 0$ and $x \in \mathbb{R}$. \square

According to the comparison principle, solutions to Equation (1.1) are above the asymptotic traveling wave solution if the initial states are properly given. An example of an initial state is an upper solution constructed in Lemma 3.11. See [18, 28, 30] for results on comparison principles in fractional differential equations.

Remark 2.4. *By taking the limits, we obtain*

$$\lim_{\alpha \rightarrow 1} c_\alpha^* = 2\sqrt{f'(0)} = c_1^*, \quad \lim_{\alpha \rightarrow +0} c_\alpha^* = \begin{cases} +\infty & f'(0) \geq 1, \\ 0 & f'(0) < 1. \end{cases}$$

c_α^* connects to the minimum speed when $\alpha = 1$. In addition, the behavior depends on the value of $f'(0)$ as α approaches zero.

3. EXISTENCE OF ASYMPTOTIC TRAVELING WAVE SOLUTIONS

In this section, we give a proof of Theorem 2.2. Equation (2.2) is nonlocal, and it is not easy to perform an analysis that considers local properties, such as a phase-plane analysis, to investigate the existence and properties of asymptotic traveling wave solutions. Therefore, we apply a monotone iteration method based on the maximum principle.

Throughout this section, we suppose that $\alpha \in (0, 1)$ and $c > 0$. For convenience, we define the fractional derivative

$$\partial_\xi^\alpha \psi(\xi) := \frac{1}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{\psi'(\xi-s)}{s^\alpha} ds$$

and the set

$$C_b^k(\mathbb{R}) := \left\{ \psi \in C^k(\mathbb{R}) \mid \sup_{\xi \in \mathbb{R}} |\psi^{(j)}(\xi)| < +\infty \quad (j = 0, 1, \dots, k) \right\}.$$

Also, we denote $C_b(\mathbb{R}) := C_b^0(\mathbb{R})$.

3.1. Preliminaries. First, we analyze the existence of a solution to the linear problem

$$(3.1) \quad L\psi(\xi) + g(\xi) = 0$$

for $g \in C_b(\mathbb{R})$, where L is a linear operator

$$L\psi(\xi) := \psi''(\xi) - c^\alpha \partial_\xi^\alpha \psi(\xi) - \kappa^2 \psi(\xi)$$

where κ is any positive constant satisfying $\kappa^2 > \max_{u \in [0,1]} |f'(u)|$. We will fix it later. This problem corresponds to equation (2.2) if $g(\xi) = \kappa^2 \psi(\xi) + f(\psi(\xi))$.

To solve problem (3.1), we derive the integral equation and apply the Neumann series argument for the existence of a solution. Let us define the functions

$$K_0(\xi) := \frac{1}{2\kappa} e^{-\kappa|\xi|}, \quad K_\alpha(\xi) := -c^\alpha \partial_\xi^\alpha K_0(\xi) = -\frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{K_0'(\xi-s)}{s^\alpha} ds.$$

We note that for all $g \in C_b(\mathbb{R})$, $(K_0 * g)$ satisfies

$$(K_0 * g)''(\xi) - \kappa^2 (K_0 * g)(\xi) + g(\xi) = 0.$$

This implies $K_0 * g$ belongs to $C_b^2(\mathbb{R})$ for any $g \in C_b(\mathbb{R})$.

We first check the properties of K_α .

Lemma 3.1. K_α has the following properties:

- (i) $K_\alpha \in C(\mathbb{R})$;
- (ii) $K_\alpha(\xi) = -(c\kappa)^\alpha K_0(\xi)$ holds for $\xi \leq 0$;
- (iii) $\lim_{\xi \rightarrow +\infty} \xi^{1+\alpha} K_\alpha(\xi) = \frac{\alpha c^\alpha}{\kappa^2 \Gamma(1-\alpha)}$;

Proof. (i) $K_\alpha(\xi)$ is represented by

$$K_\alpha(\xi) = -\frac{c^\alpha}{2\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_{\max\{\xi,0\}}^{+\infty} \frac{e^{-\kappa s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^{\max\{\xi,0\}} \frac{e^{\kappa s}}{s^\alpha} ds \right\}.$$

It is easy to see that both integrals are well-defined for all $\xi \in \mathbb{R}$ and continuous with respect to ξ .

(ii) For $\xi \leq 0$, we have

$$K_\alpha(\xi) = -\frac{c^\alpha e^{\kappa\xi}}{2\Gamma(1-\alpha)} \int_0^{+\infty} \frac{e^{-\kappa s}}{s^\alpha} ds = -\frac{c^\alpha \kappa^{\alpha-1}}{2} e^{\kappa\xi} = -(c\kappa)^\alpha K_0(\xi).$$

(iii) When $\kappa\xi > 1$, we obtain

$$\begin{aligned} K_\alpha(\xi) &= -\frac{c^\alpha}{2\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_\xi^{+\infty} \frac{e^{-\kappa s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^\xi \frac{e^{\kappa s}}{s^\alpha} ds \right\} \\ &= -\frac{(c\kappa)^\alpha}{2\kappa\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^{\kappa\xi} \frac{e^s}{s^\alpha} ds \right\}. \end{aligned}$$

We note that

$$\begin{aligned} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^\alpha} ds &= \left[-\frac{e^{-s}}{s^\alpha} \right]_{s=\kappa\xi}^{s=+\infty} - \alpha \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds \\ &= \frac{e^{-\kappa\xi}}{(\kappa\xi)^\alpha} - \alpha \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds \end{aligned}$$

and for $\delta \in (0, 1)$,

$$\begin{aligned} \int_0^{\kappa\xi} \frac{e^s}{s^\alpha} ds &= \int_0^\delta \frac{e^s}{s^\alpha} ds + \int_\delta^{\kappa\xi} \frac{e^s}{s^\alpha} ds \\ &= \int_0^\delta \frac{e^s}{s^\alpha} ds + \frac{e^{\kappa\xi}}{(\kappa\xi)^\alpha} - \frac{e^\delta}{\delta^\alpha} + \alpha \int_\delta^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds. \end{aligned}$$

Then, we have

$$\begin{aligned} &e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^{\kappa\xi} \frac{e^s}{s^\alpha} ds \\ &= -\alpha e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds - \alpha e^{-\kappa\xi} \int_\delta^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds + \left(\int_0^\delta \frac{e^s}{s^\alpha} ds - \frac{e^\delta}{\delta^\alpha} \right) e^{-\kappa\xi}. \end{aligned}$$

Hence, we obtain

$$\lim_{\xi \rightarrow +\infty} \xi^{1+\alpha} K_\alpha(\xi) = \frac{\alpha(c\kappa)^\alpha}{\kappa\Gamma(1-\alpha)} \kappa^{-(1+\alpha)} = \frac{\alpha c^\alpha}{\kappa^2 \Gamma(1-\alpha)}.$$

from l'Hôpital's rule. \square

Next, let us derive the integral equation. For the derivation, we use the fact that K_0 is a Green's function of $(\kappa^2 - \partial_x^2)$. Then, we obtain the following integral equation.

Lemma 3.2. *Let $\psi \in C_b^2(\mathbb{R})$ be a solution of the integral equation*

$$(3.2) \quad \psi(\xi) = (K_\alpha * \psi)(\xi) + K_0 * g(\xi).$$

Then, ψ is a solution to (3.1).

Proof. We first show that

$$c^\alpha(K_0 * \partial_\xi^\alpha \psi) = -(K_\alpha * \psi).$$

Changing variables $z = y - s$ and integrating by parts, we obtain

$$\begin{aligned} c^\alpha(K_0 * \partial_\xi^\alpha \psi)(\xi) &= \frac{c^\alpha}{\Gamma(1-\alpha)} \int_{\mathbb{R}} \int_0^{+\infty} K_0(\xi - y) \frac{\psi'(y-s)}{s^\alpha} ds dy \\ &= \frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \int_{\mathbb{R}} \frac{K_0(\xi - z - s)}{s^\alpha} \psi'(z) dz ds \\ &= \frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \int_{\mathbb{R}} \frac{K'_0(\xi - z - s)}{s^\alpha} \psi(z) dz ds \\ &= -(K_\alpha * \psi)(\xi) \end{aligned}$$

for all $\xi \in \mathbb{R}$.

Next, let us prove that ψ is a solution to (2.2). From the properties of $K_0(\xi)$, we deduce

$$\psi''(\xi) - \kappa^2 \psi + g(\xi) = (K_\alpha * \psi)''(\xi) - \kappa^2 (K_\alpha * \psi)(\xi).$$

Then, we have

$$\begin{aligned} &(K_\alpha * \psi)''(\xi) - \kappa^2 (K_\alpha * \psi)(\xi) \\ &= -c^\alpha \left\{ (K_0 * \partial_\xi^\alpha \psi)''(\xi) - \kappa^2 (K_0 * \partial_\xi^\alpha \psi)(\xi) \right\} \\ &= c^\alpha \partial_\xi^\alpha \psi(\xi). \end{aligned}$$

Thus, ψ satisfies $L\psi(\xi) + g(\xi) = 0$. □

Let us show the existence of the solution to (3.2) in $L^\infty(\mathbb{R})$. To apply the Neumann series argument, we check the property of $\|K_\alpha\|_{L^1}$.

Lemma 3.3. *There is $\Theta > 0$ which is independent of κ such that $\|K_\alpha\|_{L^1} < \Theta \kappa^{\alpha-2}$ holds for all $\kappa > 1$.*

Proof. To obtain the desired assertion, we evaluate $\|K_\alpha\|_{L^1}$ as

$$\int_{\mathbb{R}} |K_\alpha(\xi)| d\xi = \left[\int_{-\infty}^0 + \int_0^{1/\kappa} + \int_{1/\kappa}^{+\infty} \right] |K_\alpha(\xi)| d\xi.$$

Since we have (ii) in Lemma 3.1, the first integral part is computed by

$$(3.3) \quad \int_{-\infty}^0 |K_\alpha(\xi)| d\xi = \frac{(c\kappa)^\alpha}{2\kappa^2}.$$

Next, when $\kappa\xi \in (0, 1)$, we obtain

$$\begin{aligned} |K_\alpha(\xi)| &\leq \frac{c^\alpha}{2\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_0^{+\infty} \frac{e^{-\kappa s}}{s^\alpha} ds + e^{-\kappa\xi} \int_0^{1/\kappa} \frac{e^{\kappa s}}{s^\alpha} ds \right\} \\ &\leq \frac{(c\kappa)^\alpha}{2\kappa\Gamma(1-\alpha)} \left\{ e\Gamma(1-\alpha) + \int_0^1 \frac{e^s}{s^\alpha} ds \right\}. \end{aligned}$$

Thus, the second integral part is evaluated as

$$(3.4) \quad \int_0^{1/\kappa} |K_\alpha(\xi)| d\xi \leq \frac{(c\kappa)^\alpha}{2\kappa^2\Gamma(1-\alpha)} \left\{ e\Gamma(1-\alpha) + \int_0^1 \frac{e^s}{s^\alpha} ds \right\}.$$

In the case that $\kappa\xi > 1$, we deduce

$$\begin{aligned} K_\alpha(\xi) &= -\frac{c^\alpha}{2\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_\xi^{+\infty} \frac{e^{-\kappa s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^\xi \frac{e^{\kappa s}}{s^\alpha} ds \right\} \\ &= -\frac{(c\kappa)^\alpha}{2\kappa\Gamma(1-\alpha)} \left\{ e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^\alpha} ds - e^{-\kappa\xi} \int_0^{\kappa\xi} \frac{e^s}{s^\alpha} ds \right\}. \end{aligned}$$

We know that

$$\begin{aligned} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^\alpha} ds &= \left[-\frac{e^{-s}}{s^\alpha} \right]_{s=-\kappa\xi}^{s=+\infty} - \alpha \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds \\ &= \frac{e^{-\kappa\xi}}{(\kappa\xi)^\alpha} - \alpha \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds \end{aligned}$$

and

$$\begin{aligned} \int_0^{\kappa\xi} \frac{e^s}{s^\alpha} ds &= \int_0^1 \frac{e^s}{s^\alpha} ds + \int_1^{\kappa\xi} \frac{e^s}{s^\alpha} ds \\ &= \int_0^1 \frac{e^s}{s^\alpha} ds + \frac{e^{\kappa\xi}}{(\kappa\xi)^\alpha} - e + \alpha \int_1^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds. \end{aligned}$$

Hence, we have

$$\begin{aligned} K_\alpha(\xi) &= \frac{(c\kappa)^\alpha}{2\kappa\Gamma(1-\alpha)} \left\{ \alpha e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds + \alpha e^{-\kappa\xi} \int_1^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds \right. \\ &\quad \left. - \left(\int_0^1 \frac{e^s}{s^\alpha} ds - e \right) e^{-\kappa\xi} \right\}. \end{aligned}$$

Since we know that

$$\begin{aligned} \int_{1/\kappa}^{+\infty} e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds d\xi &= -\frac{e}{\kappa} \int_1^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds + \int_{1/\kappa}^{+\infty} \frac{1}{(\kappa\xi)^{1+\alpha}} d\xi \\ &= -\frac{e}{\kappa} \int_1^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds + \frac{1}{\alpha\kappa} \\ &\leq \frac{1}{\alpha\kappa} \end{aligned}$$

and

$$\int_{1/\kappa}^{+\infty} e^{-\kappa\xi} \int_1^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds d\xi = \int_{1/\kappa}^{+\infty} \frac{d\xi}{(\kappa\xi)^{1+\alpha}} = \frac{1}{\alpha\kappa}$$

by using integration by parts, we evaluate

$$\begin{aligned} & \int_{1/\kappa}^{+\infty} |K_\alpha(\xi)| d\xi \\ & \leq \frac{(c\kappa)^\alpha}{2\kappa\Gamma(1-\alpha)} \left\{ \alpha \int_{1/\kappa}^{+\infty} e^{\kappa\xi} \int_{\kappa\xi}^{+\infty} \frac{e^{-s}}{s^{1+\alpha}} ds d\xi + \alpha \int_{1/\kappa}^{+\infty} e^{-\kappa\xi} \int_1^{\kappa\xi} \frac{e^s}{s^{1+\alpha}} ds d\xi \right. \\ & \quad \left. + \frac{e^{-1}}{\kappa} \left| \int_0^1 \frac{e^s}{s^\alpha} ds - e \right| \right\} \\ (3.5) \quad & \leq \frac{(c\kappa)^\alpha}{2\kappa^2\Gamma(1-\alpha)} \left\{ 2 + e^{-1} \left| \int_0^1 \frac{e^s}{s^\alpha} ds - e \right| \right\}. \end{aligned}$$

From inequalities (3.3), (3.4), and (3.5), we can choose a sufficiently large $\Theta > 0$ satisfying

$$\|K_\alpha\|_{L^1} \leq \Theta\kappa^{\alpha-2}.$$

□

Fix $\kappa > 0$ sufficiently large satisfying $\|K_\alpha\|_{L^1} < 1$. Define $\mathcal{K}_\alpha : L^\infty(\mathbb{R}) \rightarrow L^\infty(\mathbb{R})$ as

$$\mathcal{K}_\alpha\psi := (K_\alpha * \psi).$$

Also, we put I as the identity operator in $L^\infty(\mathbb{R})$. Then, \mathcal{K}_α is a bounded operator on $L^\infty(\mathbb{R}^n)$, and the operator norm of \mathcal{K}_α is equal to $\|K_\alpha\|_{L^1}$. Therefore, $(I - \mathcal{K}_\alpha)$ is invertible on $L^\infty(\mathbb{R})$, and the inverse is represented by

$$(I - \mathcal{K}_\alpha)^{-1} = \sum_{j=0}^{\infty} \mathcal{K}_\alpha^j$$

from the Neumann series theory.

Here, we apply the argument for (3.2) to obtain

$$\psi = \left(\sum_{j=0}^{\infty} (K_0 * K_\alpha^{[j]}) \right) * g,$$

where $K_\alpha^{[0]}(\xi) := \delta(\xi)$ and $K_\alpha^{[j]}(\xi) := (K_\alpha * K_\alpha^{[j-1]})(\xi)$ for $j \geq 1$.

For simplicity, we denote

$$(3.6) \quad G(\xi) := \sum_{j=0}^{\infty} (K_0 * K_\alpha^{[j]})(\xi) = \left[K_0 * \left(\sum_{j=0}^{\infty} K_\alpha^{[j]} \right) \right] (\xi).$$

Then, $G(\xi)$ has the following property.

Lemma 3.4. $G(\xi)$ belongs to $C(\mathbb{R}) \cap L^1(\mathbb{R})$ and satisfies $\int_{\mathbb{R}} G(\xi) d\xi = \frac{1}{\kappa^2}$.

Proof. We note that $K_0, K_\alpha \in C(\mathbb{R})$. Since the limit $G(\xi) = \lim_{J \rightarrow \infty} \sum_{j=0}^J (K_0 * K_\alpha^{[j]})$ holds in the sense of uniform convergence, G is continuous. Moreover, we obtain

$$\|G\|_{L^1} \leq \|K_0\|_{L^1} \left(\sum_{j=0}^{\infty} \|K_\alpha\|_{L^1}^j \right).$$

By using

$$\int_{\mathbb{R}} K_\alpha(\xi) d\xi = 0,$$

$$\text{we have } \int_{\mathbb{R}} G(\xi) d\xi = \int_{\mathbb{R}} K_0(\xi) d\xi = \frac{1}{\kappa^2}. \quad \square$$

From the construction of G and Lemma 3.2, we obtain the conclusion for the existence.

Proposition 3.5. *For $g \in C_b(\mathbb{R})$, $\psi(\xi) = (G * g)(\xi)$ belongs to $C_b^2(\mathbb{R})$ and is a solution to (3.1).*

In the next step, we establish the maximum principle for the operator L . Let us consider the sign of the fractional derivative ∂_ξ^α at the minimal point.

Lemma 3.6. *Let $\psi \in C_b^1(\mathbb{R})$. Assume that $\psi(\xi)$ attains the minimum at $\xi^* \in \mathbb{R}$: $\psi(\xi^*) = \min_{\xi \in \mathbb{R}} \psi(\xi)$. Then, we have*

$$\partial_\xi^\alpha \psi(\xi^*) = -\frac{\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{(\psi(\xi^* - s) - \psi(\xi^*))}{s^{1+\alpha}} ds \leq 0.$$

Moreover, $\partial_\xi^\alpha \psi(\xi^*)$ attains 0 at $\xi^* \in \mathbb{R}$ if and only if $\psi(\xi) = \psi(\xi^*)$ for any $\xi \leq \xi^*$.

Proof. From the fact that

$$\frac{d}{ds} \left(\frac{\psi(\xi^* - s) - \psi(\xi^*)}{s^\alpha} \right) = -\frac{\psi'(\xi^* - s)}{s^\alpha} - \frac{\alpha(\psi(\xi^* - s) - \psi(\xi^*))}{s^{1+\alpha}},$$

we obtain

$$\begin{aligned}
\partial_\xi^\alpha \psi(\xi^*) &= -\frac{1}{\Gamma(1-\alpha)} \left[\frac{\psi(\xi^* - s) - \psi(\xi^*)}{s^\alpha} \right]_{s \rightarrow +0}^{s \rightarrow +\infty} \\
&\quad - \frac{\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{(\psi(\xi^* - s) - \psi(\xi^*))}{s^{1+\alpha}} ds \\
&= -\frac{\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{(\psi(\xi^* - s) - \psi(\xi^*))}{s^{1+\alpha}} ds \leq 0.
\end{aligned}$$

Furthermore, it is easy to see that if $\psi(\xi) > \psi(\xi^*)$ holds at the point $\xi \leq \xi^*$, then $\partial_\xi^\alpha \psi(\xi^*)$ must be negative from the continuity of $\psi(\xi)$. Thus, $\psi(\xi)$ must be constant on $(-\infty, \xi^*]$ if $\partial_\xi^\alpha \psi(\xi^*) = 0$. \square

Let us introduce the maximum principle.

Proposition 3.7. *Suppose that $\psi \in C_b^1(\mathbb{R}) \cap C^2(\mathbb{R})$ satisfies*

$$L\psi(\xi) \leq 0$$

and

$$\liminf_{\xi \rightarrow -\infty} \psi(\xi) \geq 0, \quad \liminf_{\xi \rightarrow +\infty} \psi(\xi) \geq 0.$$

Then, $\psi(\xi)$ is non-negative. Furthermore, if $\psi(\xi)$ attains 0 at the position $\xi^* \in \mathbb{R}$, then $\psi(\xi) \equiv 0$ for all $\xi \leq \xi^*$.

Proof. Suppose that $\psi(\xi)$ has a negative part. Then, from the assumption, there is $\xi^* \in \mathbb{R}$ such that

$$\psi(\xi^*) = \min_{\xi \in \mathbb{R}} \psi(\xi) < 0.$$

This implies

$$L\psi(\xi^*) = \psi''(\xi^*) - c^\alpha \partial_\xi^\alpha \psi(\xi^*) - \kappa^2 \psi(\xi^*) > 0$$

from Lemma 3.6. This is a contradiction of the assumption $L\psi(\xi) \leq 0$.

Next, we assume that $\psi(\xi)$ is non-negative and attains 0 at $\xi^* \in \mathbb{R}$. Then, we have the following.

$$0 \geq L\psi(\xi^*) = \psi''(\xi^*) - c^\alpha \partial_\xi^\alpha \psi(\xi^*) \geq -c^\alpha \partial_\xi^\alpha \psi(\xi^*) \geq 0.$$

Hence, ψ must be zero on $(-\infty, \xi^*]$ because $\partial_\xi^\alpha \psi(\xi^*) = 0$. \square

Finally, we construct an upper solution and a lower solution to (2.2). We introduce a notion of upper and lower solutions.

Definition 3.8. *For $c > 0$, $\bar{\phi} \in C^2(\mathbb{R})$ is an upper solution to (2.2) if $\bar{\phi}$ satisfies*

$$c^\alpha \partial_\xi^\alpha \bar{\phi}(\xi) \geq \bar{\phi}''(\xi) + f(\bar{\phi}(\xi))$$

for all $\xi \in \mathbb{R}$. Conversely, $\underline{\phi} \in C^2(\mathbb{R})$ is a lower solution to (2.2) if $\underline{\phi}$ satisfies

$$c^\alpha \partial_\xi^\alpha \underline{\phi}(\xi) \leq \underline{\phi}''(\xi) + f(\underline{\phi}(\xi))$$

for all $\xi \in \mathbb{R}$.

We focus on the property $u = 0$ to construct the upper and lower solutions. The linearized equation around $u = 0$ is given as

$$\frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{\phi'(\xi-s)}{s^\alpha} ds = \phi''(\xi) + f'(0)\phi(\xi).$$

Since we have

$$\frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{e^{-\lambda s}}{s^\alpha} ds = \frac{(c\lambda)^\alpha}{\lambda},$$

the characteristic equation around $u = 0$ is defined as follows:

$$V(\lambda) := \lambda^2 - (c\lambda)^\alpha + f'(0) = 0.$$

Then, we have

Lemma 3.9. *The following statements hold:*

- (i) *For $c > c_\alpha^*$, there are two positive roots $\lambda_1 < \lambda_2$ of $V(\lambda) = 0$. Moreover, $V(\lambda) < 0$ on (λ_1, λ_2) .*
- (ii) *When $c = c_\alpha^*$, there is a unique positive root λ_1 .*
- (iii) *For $c \in (0, c_\alpha^*)$, there are no positive roots.*

Proof. It is obvious that $V(\lambda)$ is smooth on $(0, +\infty)$ and satisfies $V(0) = f'(0) > 0$. Since we obtain

$$V'(\lambda) = 2\lambda - \frac{\alpha c^\alpha}{\lambda^{1-\alpha}}, \quad V''(\lambda) = 2 + \frac{\alpha(1-\alpha)c^\alpha}{\lambda^{2-\alpha}}$$

for $\lambda > 0$, $V(\lambda)$ is convex for $\lambda > 0$ and attains the minimum at $\lambda^* = (\alpha/2)^{1/(2-\alpha)} c^{\alpha/(2-\alpha)}$. Moreover, we have

$$V(\lambda^*) = 2^{-2/(2-\alpha)}(\alpha-2)(\alpha c^2)^{\alpha/(2-\alpha)} + f'(0).$$

Thus, $V(\lambda) = 0$ has two positive roots if $c > c_\alpha^*$, and no positive root when $c < c_\alpha^*$. Moreover, λ^* is a unique positive root of $V(\lambda) = 0$ if $c = c_\alpha^*$. \square

Remark 3.10. *When we consider the linearized equation for Equation (2.1) around $u = 0$,*

$$\partial_t^\alpha w = w_{xx} + f'(0)w \quad (t > 0, x \in \mathbb{R}),$$

it has a solution $w(t, x) = E_{\alpha,1}((c\lambda t)^\alpha) e^{\lambda x}$ for all $c \geq c_\alpha^$, where λ is a positive root of $V(\lambda) = 0$. From (1.4), the asymptotic of $w(t, x)$ is approximated by*

$$w(t, x) \simeq \frac{1}{\alpha} e^{\lambda(x+ct)} = \frac{1}{\alpha} e^{\lambda\xi} \quad (t \rightarrow +\infty).$$

This consideration suggests that the asymptotic traveling wave solution has an exponential tail at $u = 0$.

Next, let us construct an upper solution. For convenience, we set

$$\begin{aligned} N_1(\xi) &:= \bar{\phi}''(\xi) - c^\alpha \partial_\xi^\alpha \bar{\phi}(\xi) + f(\bar{\phi}(\xi)), \\ N_2(\xi) &:= \underline{\phi}''(\xi) - c^\alpha \partial_\xi^\alpha \underline{\phi}(\xi) + f(\underline{\phi}(\xi)). \end{aligned}$$

Let ρ be a positive mollifier, that is, $\rho(x)$ is an even function satisfying

$$\rho \in C_0^\infty(\mathbb{R}), \quad \rho(x) \begin{cases} > 0 & |x| < 1, \\ = 0 & |x| \geq 1, \end{cases} \quad \int_{\mathbb{R}} \rho(x) dx = 1.$$

We set $\rho_\epsilon(x) := \frac{1}{\epsilon} \rho\left(\frac{x}{\epsilon}\right)$. For $c > c_\alpha^*$, we define

$$\bar{\psi}(\xi) = \begin{cases} e^{\lambda_1 \xi} & \xi < 0, \\ 1 & \xi > 0, \end{cases}$$

where λ_1 is a minimal positive root of $V(\lambda)$. Then, we set $\bar{\phi}(\xi) = \bar{\phi}(\xi; \epsilon) := (\rho_\epsilon * \bar{\psi})(\xi)$. It is easy to see that we have

$$\bar{\phi}(\xi) = \begin{cases} R(\epsilon \lambda_1) e^{\lambda_1 \xi} & \xi \leq -\epsilon \\ \int_{-1}^{\xi/\epsilon} \rho(y) dy + e^{\lambda_1 \xi} \int_{\xi/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy & |\xi| < \epsilon \\ 1 & \xi \geq \epsilon, \end{cases}$$

where $R(\lambda) := \int_{-1}^1 \rho(y) e^{-\lambda y} dy$.

Lemma 3.11. *For $c > c_\alpha^*$ and $\epsilon > 0$, $\bar{\phi}$ defined as above satisfies the following properties:*

- (i) $\bar{\phi} \in C^\infty(\mathbb{R})$;
- (ii) $\bar{\phi}$ is non-decreasing;
- (iii) $\bar{\phi}(\xi) \in [0, 1]$ for all $\xi \in \mathbb{R}$.

Moreover, $\bar{\phi}$ is an upper solution to (2.2) if ϵ is sufficiently small.

Proof. (i) is derived from the general theory of the mollifier. Since $\bar{\psi}(\xi)$ is a non-decreasing function connecting 0 to 1, we have (ii) and (iii).

In the rest of the proof, we show that $\bar{\phi}$ is an upper solution to (2.2). For $\xi \leq -\epsilon$, we have

$$\begin{aligned} N_1(\xi) &= R(\epsilon \lambda_1) \left(\lambda_1^2 e^{\lambda_1 \xi} - \frac{c^\alpha \lambda_1}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{e^{\lambda_1(\xi-s)}}{s^\alpha} ds \right) + f(\bar{\phi}(\xi)) \\ &= R(\epsilon \lambda_1) e^{\lambda_1 \xi} \{ \lambda_1^2 - (c \lambda_1)^\alpha \} + f(\bar{\phi}(\xi)) \\ &= f(\bar{\phi}(\xi)) - f'(0) \bar{\phi}(\xi) \leq 0 \end{aligned}$$

from the fact that $V(\lambda_1) = 0$ and assumption (2.4).

When $\xi \geq \epsilon$, we obtain

$$N_1(\xi) = -\frac{c^\alpha}{\Gamma(1-\alpha)} \int_0^{+\infty} \frac{\bar{\phi}'(\xi-s)}{s^\alpha} ds < 0$$

because $\bar{\phi}$ is non-decreasing. Thus, for any $\epsilon > 0$, we have $N_1(\xi) \leq 0$ on $|\xi| \geq \epsilon$.

In the case that $|\xi| < \epsilon$, we put $\eta = \xi/\epsilon$. Then, the first and second derivatives satisfy

$$\begin{aligned}\bar{\phi}'(\xi) &= \lambda_1 e^{\lambda_1 \xi} \int_{\xi/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy \\ &= \lambda_1 e^{\epsilon \lambda_1 \eta} \int_{\eta}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy = \lambda_1 \int_{\eta}^1 \rho(y) dy + O(\epsilon), \\ \bar{\phi}''(\xi) &= \lambda_1^2 e^{\lambda_1 \xi} \int_{\xi/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - \frac{\lambda_1}{\epsilon} \rho\left(\frac{\xi}{\epsilon}\right) \\ &= \lambda_1^2 e^{\epsilon \lambda_1 \eta} \int_{\eta}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - \frac{\lambda_1}{\epsilon} \rho(\eta) = -\frac{\lambda_1}{\epsilon} \rho(\eta) + \lambda_1^2 \int_{\eta}^1 \rho(y) dy + O(\epsilon)\end{aligned}$$

for $\eta \in [-1, 1]$ if ϵ is sufficiently small. Furthermore, for any $\eta \in [-1, 1]$ and sufficiently small $\epsilon > 0$, we approximate

$$f(\bar{\phi}(\xi)) = f(1) + O(\epsilon) = O(\epsilon)$$

and

$$\begin{aligned}\partial_{\xi}^{\alpha} \bar{\phi}(\xi) &= \frac{1}{\Gamma(1-\alpha)} \left\{ \int_{\epsilon(\eta+1)}^{+\infty} \frac{\bar{\phi}'(\epsilon\eta-s)}{s^{\alpha}} ds + \int_0^{\epsilon(\eta+1)} \frac{\bar{\phi}'(\epsilon\eta-s)}{s^{\alpha}} ds \right\} \\ &= \frac{\lambda_1 R(\epsilon\lambda_1)}{\Gamma(1-\alpha)} \int_{\epsilon(\eta+1)}^{+\infty} \frac{e^{\epsilon\eta-s}}{s^{\alpha}} ds + O(\epsilon) = \lambda_1^{\alpha} + O(\epsilon).\end{aligned}$$

Thus, we deduce

$$\begin{aligned}N_1(\xi) &= \bar{\phi}''(\xi) - c^{\alpha} \partial_{\xi}^{\alpha} \bar{\phi}(\xi) + f(\bar{\phi}(\xi)) \\ &= -\frac{\lambda_1}{\epsilon} \rho(\eta) + \lambda_1^2 \int_{\eta}^1 \rho(y) dy - (c\lambda_1)^{\alpha} + O(\epsilon)\end{aligned}$$

for all $\eta \in (-1, 1)$ and sufficiently small $\epsilon > 0$. Since we know $\rho(\eta) \geq 0$ and have

$$\lambda_1^2 \int_{\eta}^1 \rho(y) dy - (c\lambda_1)^{\alpha} \leq \lambda_1^2 - (c\lambda_1)^{\alpha} = -f'(0) < 0,$$

$N_1(\xi)$ is negative on $(-\epsilon, \epsilon)$ as long as ϵ is sufficiently small. \square

Next, we construct a lower solution. For $c > c_{\alpha}^*$, let $\psi(\xi) = e^{\lambda_1 \xi} - h e^{\nu \lambda_1 \xi}$ with $\nu \in (1, \min\{1+a, \lambda_2/\lambda_1\})$, where $h > 1$ is a sufficiently large constant satisfying

$$(3.7) \quad \frac{1}{h^2} \leq -\frac{-h R(\epsilon \nu \lambda_1) V(\nu \lambda_1)}{M \{R(\epsilon \lambda_1)\}^{1+a}} e^{(\nu-1)/(1+a-\nu)}.$$

We remark that the right-hand side is strictly positive for any $\epsilon \in [0, 1)$. We fix h , satisfying (3.7) for all $\epsilon \in [0, 1)$. When we set $\xi_0 := -\frac{\log h}{\lambda_1(\nu-1)} < 0$, we have

$$\psi(\xi) \begin{cases} > 0 & \xi < \xi_0 \\ = 0 & \xi = \xi_0 \\ < 0 & \xi > \xi_0. \end{cases}$$

Now, we set

$$\underline{\psi}(\xi) = \begin{cases} e^{\lambda_1 \xi} - h e^{\nu \lambda_1 \xi} & \xi < \xi_0, \\ 0 & \xi > \xi_0 \end{cases}$$

and define $\underline{\phi}(\xi) = \underline{\phi}(\xi; \epsilon) := (\rho_\epsilon * \underline{\psi})(\xi)$. From direct computation, we obtain

$$\underline{\phi}(\xi) = \begin{cases} R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi} & \xi \leq \xi_0 - \epsilon, \\ e^{\lambda_1 \xi} \int_{(\xi - \xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - h e^{\nu \lambda_1 \xi} \int_{(\xi - \xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \nu \lambda_1 y} dy & |\xi - \xi_0| < \epsilon \\ 0 & \xi \geq \xi_0 + \epsilon. \end{cases}$$

Lemma 3.12. *For $c > c_\alpha^*$ and $\epsilon > 0$, $\underline{\phi}$ defined as above satisfies the following properties:*

- (i) $\underline{\phi} \in C^\infty(\mathbb{R})$;
- (ii) $\underline{\phi}$ is non-negative on \mathbb{R} and positive on $(-\infty, \xi_0 + \epsilon)$;
- (iii) $\underline{\phi}(\xi) \leq \bar{\phi}(\xi)$ for all $\xi \in \mathbb{R}$.

Moreover, $\underline{\phi}$ is a lower solution to (2.2) if ϵ is sufficiently small.

Proof. (i) holds from the general theory of the mollifier. $\underline{\phi}$ is non-negative because $\underline{\psi}$ is non-negative. Also, $\underline{\psi}$ is positive on $(-\infty, \xi_0]$, and thus $\underline{\phi}$ is positive on $(-\infty, \xi_0 + \epsilon)$. Hence, we have (ii). (iii) is derived from the fact that $\underline{\psi} < \bar{\psi}$ on \mathbb{R} .

Next, when $\xi < \xi_0 - \epsilon$, we deduce

$$\begin{aligned} N_2(\xi) &= \lambda_1^2 R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h(\nu \lambda_1)^2 R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi} + f(R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi}) \\ &\quad - \frac{c^\alpha}{\Gamma(1-\alpha)} \left\{ \lambda_1 R(\epsilon \lambda_1) \int_0^{+\infty} \frac{e^{\lambda_1(\xi-s)}}{s^\alpha} ds - h \nu \lambda_1 R(\epsilon \nu \lambda_1) \int_0^{+\infty} \frac{e^{\nu \lambda_1(\xi-s)}}{s^\alpha} ds \right\} \\ &= \lambda_1^2 e^{\lambda_1 \xi} - h(\nu \lambda_1)^2 e^{\nu \lambda_1 \xi} + f(R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi}) \\ &\quad - (c \lambda_1)^\alpha e^{\lambda_1 \xi} - h(c \nu \lambda_1)^\alpha e^{\nu \lambda_1 \xi} \\ &= f(R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi}) - f'(0)(R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi}) \\ &\quad - R(\epsilon \nu \lambda_1) h V(\nu \lambda_1) e^{\nu \lambda_1 \xi} \\ &\geq -M(R(\epsilon \lambda_1) e^{\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) e^{\nu \lambda_1 \xi})^{1+a} - R(\epsilon \nu \lambda_1) h V(\nu \lambda_1) e^{\nu \lambda_1 \xi} \\ &\geq -M\{R(\epsilon \lambda_1)\}^{1+a} e^{(1+a)\lambda_1 \xi} - h R(\epsilon \nu \lambda_1) V(\nu \lambda_1) e^{\nu \lambda_1 \xi}. \end{aligned}$$

We note that $V(\nu \lambda_1) < 0$. Since it is easy to show that the last function is positive if

$$\xi \leq \frac{1}{(1+a-\nu)\lambda_1} \log \left(\frac{-h R(\epsilon \nu \lambda_1) V(\nu \lambda_1)}{M\{R(\epsilon \lambda_1)\}^{1+a}} \right) =: \xi^*,$$

we obtain $N_2(\xi) \geq 0$ for all $\xi < \xi_0$ when $\xi_0 \leq \xi^*$. It is easy to show that $\xi_0 \leq \xi^*$ holds if and only if h satisfies (3.7). Then, we obtain $N_2(\xi) \geq 0$ for $\xi < \xi_0$.

Next, we consider the case that $\xi \geq \xi_0 + \epsilon$. Since $\underline{\phi}(\xi) \equiv 0$ on $[\xi_0 + \epsilon, +\infty)$ and $\min_{\xi \in \mathbb{R}} \underline{\phi}(\xi) = 0$, we have

$$N_2(\xi) = -c^\alpha \partial_\xi^\alpha \underline{\phi}(\xi) > 0$$

by using Lemma 3.6.

Finally, in the case that $|\xi - \xi_0| < \epsilon$, we apply the same argument of the proof of Lemma 3.11 to show. Let $\eta = (\xi - \xi_0)/\epsilon$. Then, for sufficiently small $\epsilon > 0$, the first derivative is approximated by

$$\begin{aligned}\underline{\phi}'(\xi) &= \lambda_1 e^{\lambda_1 \xi} \int_{(\xi-\xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - h \nu \lambda_1 e^{\nu \lambda_1 \xi} \int_{(\xi-\xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \nu \lambda_1 y} dy \\ &= \lambda_1 e^{\lambda_1 (\xi_0 + \epsilon \eta)} \int_{\eta}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - h \nu \lambda_1 e^{\nu \lambda_1 (\xi_0 + \epsilon \eta)} \int_{\eta}^1 \rho(y) e^{-\epsilon \nu \lambda_1 y} dy \\ &= \lambda_1 e^{\lambda_1 \xi_0} (1 - \nu) \int_{\eta}^1 \rho(y) dy + O(\epsilon).\end{aligned}$$

Here, we use the fact that $e^{\lambda_1 \xi_0} = h e^{\nu \lambda_1 \xi_0}$. Also, the approximation of the second derivative is computed as

$$\begin{aligned}\underline{\phi}''(\xi) &= \lambda_1^2 e^{\lambda_1 \xi} \int_{(\xi-\xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy - h(\nu \lambda_1)^2 e^{\nu \lambda_1 \xi} \int_{(\xi-\xi_0)/\epsilon}^1 \rho(y) e^{-\epsilon \nu \lambda_1 y} dy \\ &\quad + \frac{\lambda_1 e^{\lambda_1 \xi_0} (1 - \nu)}{\epsilon} \rho\left(\frac{\xi - \xi_0}{\epsilon}\right) \\ &= e^{\lambda_1 \xi_0} \left\{ \lambda_1^2 e^{\epsilon \lambda_1 \eta} - (\nu \lambda_1)^2 e^{\epsilon \nu \lambda_1 \eta} \right\} \int_{\eta}^1 \rho(y) e^{-\epsilon \lambda_1 y} dy + \frac{\lambda_1 e^{\lambda_1 \xi_0} (\nu - 1)}{\epsilon} \rho(\eta) \\ &= \frac{\lambda_1 e^{\lambda_1 \xi_0} (\nu - 1)}{\epsilon} \rho(\eta) + e^{\lambda_1 \xi_0} \lambda_1^2 (1 - \nu^2) \int_{\eta}^1 \rho(y) dy + O(\epsilon).\end{aligned}$$

Moreover, we obtain $f(\underline{\phi}(\xi)) = O(\epsilon)$ and

$$\begin{aligned}\partial_{\xi}^{\alpha} \underline{\phi}(\xi) &= \frac{1}{\Gamma(1 - \alpha)} \left[\int_{\xi - \xi_0 + \epsilon}^{+\infty} + \int_0^{\xi - \xi_0 + \epsilon} \right] \frac{\underline{\phi}'(\xi - s)}{s^{\alpha}} ds \\ &= \frac{1}{\Gamma(1 - \alpha)} \left[\int_{\epsilon(\eta+1)}^{+\infty} + \int_0^{\epsilon(\eta+1)} \right] \frac{\underline{\phi}'(\epsilon \eta + \xi_0 - s)}{s^{\alpha}} ds \\ &= \frac{1}{\Gamma(1 - \alpha)} \int_{\epsilon(\eta+1)}^{+\infty} \frac{\underline{\phi}'(\epsilon \eta + \xi_0 - s)}{s^{\alpha}} ds + O(\epsilon) \\ &= \frac{\lambda_1 e^{\lambda_1 \xi_0}}{\Gamma(1 - \alpha)} \left\{ R(\epsilon \lambda_1) \int_{\epsilon(\eta+1)}^{+\infty} \frac{e^{\lambda_1(\epsilon \eta - s)}}{s^{\alpha}} ds - \nu R(\epsilon \nu \lambda_1) \int_{\epsilon(\eta+1)}^{+\infty} \frac{e^{\nu \lambda_1(\epsilon \eta - s)}}{s^{\alpha}} ds \right\} + O(\epsilon) \\ &= \lambda_1^{\alpha} e^{\lambda_1 \xi_0} (1 - \nu^{\alpha}) + O(\epsilon)\end{aligned}$$

for all $\eta \in (-1, 1)$ and sufficiently small $\epsilon > 0$. Thus, we have

$$N_2(\xi) \leq \frac{\lambda_1 e^{\lambda_1 \xi_0} (\nu - 1)}{\epsilon} \rho(\eta) - e^{\lambda_1 \xi_0} (\nu^2 - 1) \left\{ \lambda_1^2 \int_{\eta}^1 \rho(y) dy - (c \lambda_1)^{\alpha} \right\} + O(\epsilon).$$

Since the $O(1)$ term satisfies

$$\begin{aligned} e^{\lambda_1 \xi_0} (\nu^2 - 1) \left\{ \lambda_1^2 \int_{\eta}^1 \rho(y) dy - (c\lambda_1)^\alpha \right\} &\leq e^{\lambda_1 \xi_0} (\nu^2 - 1) \{ \lambda_1^2 - (c\lambda_1)^\alpha \} \\ &\leq -f'(0) e^{\lambda_1 \xi_0} (\nu^2 - 1) < 0, \end{aligned}$$

$N_2(\xi)$ is positive on $|\xi - \xi_0| < \epsilon$ as long as ϵ is sufficiently small. \square

3.2. Proof of Theorem 2.2 in the case $c > c_\alpha^*$. Fix $c > c_\alpha^*$. We give a sufficiently small $\epsilon > 0$ such that $\bar{\phi}$ and $\underline{\phi}$ become upper and lower solutions to (2.2), respectively. Let us define the sequence of functions

$$\begin{cases} \psi_0 = \bar{\phi}, \\ \psi_j = G * (\kappa^2 \psi_{j-1} + f(\psi_{j-1})) \quad (j = 1, 2, \dots), \end{cases}$$

where G is defined by (3.6).

Lemma 3.13. *For $j \in \mathbb{N} \cup \{0\}$, the following hold:*

- (i) $\psi_j \in C_b^2(\mathbb{R})$;
- (ii) $\lim_{\xi \rightarrow -\infty} \psi_j(\xi) = 0$ and $\lim_{\xi \rightarrow +\infty} \psi_j(\xi) = 1$ hold;
- (iii) $\underline{\phi}(\xi) \leq \psi_j(\xi) \leq \bar{\phi}(\xi)$ ($\xi \in \mathbb{R}$);
- (iv) ψ_j is non-decreasing function;

Furthermore, ψ_j is an increasing function for $j \in \mathbb{N}$.

Proof. Since it is obvious that (i)-(iv) hold in the case $j = 0$, we prove them by induction. Suppose that the statements (i)-(iv) hold for $j = k \geq 0$.

- (i) From the definition of ψ_{k+1} and Proposition 3.5, $\psi_{k+1} \in C_b^2(\mathbb{R})$ holds.
- (ii) Since we know that $\psi_k \in C_b$ and $\lim_{\xi \rightarrow +\infty} (\kappa^2 \psi_k(\xi) + f(\psi_k(\xi))) = \kappa^2$, we deduce

$$\lim_{\xi \rightarrow +\infty} \psi_{k+1}(\xi) = \lim_{\xi \rightarrow +\infty} G * (\kappa^2 \psi_k + f(\psi_k))(\xi) = \kappa^2 \int_{\mathbb{R}} G(\xi) d\xi = 1$$

from the dominated convergence theorem. The same argument also provides the result of the case $\xi \rightarrow -\infty$.

- (iii) Since ψ_{k+1} is a solution to

$$L\psi_{k+1} + \kappa^2 \psi_k + f(\psi_k) = 0,$$

we have

$$L\underline{\phi} \geq L\psi_{k+1} \geq L\bar{\phi}$$

from the monotonicity of $\kappa^2 u + f(u)$. Applying Proposition 3.7, statement (iii) holds.

- (iv) Let $\psi_{k+1,p}(\xi) := \psi_{k+1}(\xi - p)$ for $p > 0$. Then, we deduce

$$\begin{aligned} L\psi_{k+1,p}(\xi) &= -\kappa^2 \psi_k(\xi - p) - f(\psi_k(\xi - p)) \\ &\geq -\kappa^2 \psi_k(\xi) - f(\psi_k(\xi)) \\ &= L\psi_{k+1}(\xi) \end{aligned}$$

for $\xi \in \mathbb{R}$. From Proposition 3.7, we obtain $\psi_{k+1,p}(\xi) \leq \psi_{k+1}(\xi)$ for all $\xi \in \mathbb{R}$. This implies that ψ_{k+1} is non-decreasing.

Furthermore, if there is $\xi^* \in \mathbb{R}$ such that $\psi_{k+1,p}(\xi^*) = \psi_{k+1}(\xi^*)$ holds, then by applying Proposition 3.7, we deduce

$$\psi_{k+1}(\xi) = \psi_{k+1}(\xi - p)$$

for all $\xi \leq \xi^*$. Since we have Lemma 3.13 (ii) and $\lim_{\xi \rightarrow -\infty} \bar{\phi}(\xi) = \lim_{\xi \rightarrow -\infty} \underline{\phi}(\xi) = 0$, ψ_{k+1} must be zero on $(-\infty, \xi^*]$. However, this is a contradiction of Lemma 3.13 (ii) because ϕ is non-zero on $(-\infty, \xi_0 + \epsilon]$. Thus, $\psi_{k+1}(\xi) > \psi_{k+1}(\xi - p)$ holds for all $p > 0$ and $\xi \in \mathbb{R}$. This implies that ψ_{k+1} is an increasing function. \square

Lemma 3.14. *For $j \in \mathbb{N}$, we have $\psi_j(\xi) \leq \psi_{j-1}(\xi)$ for all $\xi \in \mathbb{R}$.*

Proof. The statement holds for the case $j = 1$ from Lemma 3.13 (iii). Suppose that the statement holds for $j = k \geq 1$. From the definition of ψ_j ($j \in \mathbb{N}$) and Proposition 3.5, we deduce

$$L\psi_{k+1} = -\kappa^2\psi_k - f(\psi_k) \geq -\kappa^2\psi_{k-1} - f(\psi_{k-1}) = L\psi_k.$$

By applying Proposition 3.7, we obtain the statement. \square

Since $\{\psi_j\}_{j \in \mathbb{N}}$ is a monotone sequence from Lemmas 3.13 and 3.14, it converges pointwise to a non-decreasing function ϕ satisfying

$$(3.8) \quad \underline{\phi}(\xi) \leq \phi(\xi) \leq \bar{\phi}(\xi)$$

from the monotone convergence theorem and Helly's selection theorem. Furthermore, we have

$$\phi(\xi) = \lim_{j \rightarrow \infty} \psi_j(\xi) = \lim_{j \rightarrow \infty} G * (\kappa^2\psi_{j-1} + f(\psi_{j-1}))(\xi) = G * (\kappa^2\phi + f(\phi))(\xi)$$

for all $\xi \in \mathbb{R}$ by using the dominated convergence theorem. Since G is continuous and ϕ is bounded, $\phi \in C_b(\mathbb{R})$ holds. Hence, ϕ belongs to $C_b^2(\mathbb{R})$ and is a solution to (2.2) from Proposition 3.5.

It is easy to see that ϕ is an increasing function from the same argument of the proof of Lemma 3.13. From (3.8) and Lemmas 3.11 and 3.12, we obtain $\lim_{\xi \rightarrow -\infty} \phi(\xi) = 0$. Moreover, ϕ is bounded and increasing, thus we know that there exists a limit $\phi^+ := \lim_{\xi \rightarrow +\infty} \phi(\xi)$ satisfying $0 \leq \phi^+ \leq 1$. Then, the condition of ϕ^+ is derived as

$$\phi^+ = \lim_{\xi \rightarrow +\infty} \phi(\xi) = \lim_{\xi \rightarrow +\infty} G * (\kappa^2\phi + f(\phi))(\xi) = \frac{1}{\kappa^2}(\kappa^2\phi^+ + f(\phi^+)).$$

This implies $f(\phi^+) = 0$, and so $\phi^+ = 1$ holds. Therefore, ϕ is a solution to (2.2) satisfying (2.3).

Finally, from Lemmas 3.11 and 3.12, we have

$$\lim_{\xi \rightarrow -\infty} e^{-\lambda_1 \xi} \phi(\xi) = R(\epsilon \lambda_1).$$

By translating ϕ appropriately, the limit can be determined to be 1. Thus, Theorem 2.2 for $c > c_\alpha^*$ is proved.

3.3. Proof of Theorem 2.2 in the case $c = c_\alpha^*$. Let $\{c_j\}_{j \in \mathbb{N}}$ be a minimizing sequence for c_α^* and $\{\phi_j\}_{j \in \mathbb{N}}$ be the solution to (2.2) corresponding to $c_j > c_\alpha^*$ constructed by Subsection 3.2. For convenience, we take $\phi_j(0) = \frac{1}{2}$ for all $j \in \mathbb{N}$ by shifting the spatial variable. Since $\{\phi_j\}_{j \in \mathbb{N}}$ is a bounded sequence of increasing functions in $C^{2,\beta}(\mathbb{R})$ from the similar argument for the elliptic regularity, there exists a subsequence $\{\phi_{j_k}\}_{k \in \mathbb{N}}$ such that ϕ_{j_k} converges to a non-decreasing function ϕ^* in $C_{loc}^{2,\beta}(\mathbb{R})$ which satisfies (2.2) with $c = c_\alpha^*$. Furthermore, ϕ^* have the limits $\phi^{*\pm} = \lim_{\xi \rightarrow \pm\infty} \phi^*(\xi)$, and so $\phi^{*\pm}$ satisfy $f(\phi^{*\pm}) = 0$ and $0 \leq \phi^{*-} < \phi^{*+} \leq 1$. This implies $\phi^{*-} = 0$ and $\phi^{*+} = 1$. Moreover, ϕ^* is an increasing function from the same argument of the proof of Lemma 3.13. Thus, ϕ^* is a desired solution.

4. NUMERICAL EXPERIMENTS

Numerical experiments are performed to investigate the spatio-temporal dynamics of the solution to equation (1.1).

4.1. Numerical scheme and setting. We first introduce the numerical scheme for the simulation.

Let $\delta_t > 0$ and $u_j(x) := u(j\delta_t, x)$ for $j \geq 0$. We apply the time difference proposed by [29]:

$$\begin{aligned} \partial_t^\alpha u((j+1)\delta_t, x) &= \frac{1}{\Gamma(1-\alpha)} \sum_{m=0}^j \int_{m\delta_t}^{(m+1)\delta_t} \frac{\partial u}{\partial s}(s, x) \frac{ds}{\{(j+1)\delta_t - s\}^\alpha} \\ &\simeq \frac{1}{\Gamma(1-\alpha)} \sum_{m=0}^j \frac{u_{m+1}(x) - u_m(x)}{\delta_t} \int_{m\delta_t}^{(m+1)\delta_t} \frac{ds}{\{(j+1)\delta_t - s\}^\alpha}. \end{aligned}$$

Following the straightforward computation of [29], the discretization is represented by

$$\partial_t^\alpha u((j+1)\delta_t, x) \simeq \frac{1}{\Gamma(2-\alpha)} \sum_{m=0}^j \frac{u_{j+1-m}(x) - u_{j-m}(x)}{\delta_t} [(m+1)^{1-\alpha} - m^{1-\alpha}].$$

Thus, when $u_m(x)$ ($m = 0, 1, \dots, j$) are given, we obtain $u_{j+1}(x)$ by using the time difference equation

$$\frac{1}{\Gamma(2-\alpha)} \sum_{m=0}^j \frac{u_{j+1-m}(x) - u_{j-m}(x)}{\delta_t} [(m+1)^{1-\alpha} - m^{1-\alpha}] = (u_{j+1})_{xx}(x) + f(u_j(x)).$$

Next, let $\delta_x > 0$, $k_x \in \mathbb{N}$, and $u_{j,k} := u(j\delta_t, k\delta_x)$ for $0 \leq k \leq k_x$. For $0 \leq k \leq k_x$, we approximate $(u_{j+1})_{xx}(k\delta_x)$ as

$$(u_{j+1})_{xx}(k\delta_x) \simeq \frac{u_{j+1,k+1} - 2u_{j+1,k} + u_{j+1,k-1}}{\delta_x^2}.$$

Thus, spatial discretization yields

$$\begin{aligned} \frac{1}{\Gamma(2-\alpha)} \sum_{m=0}^j \frac{u_{j+1-m,k} - u_{j-m,k}(x)}{\delta_t} [(m+1)^{1-\alpha} - m^{1-\alpha}] \\ = \frac{u_{j+1,k+1} - 2u_{j+1,k} + u_{j+1,k-1}}{\delta_x^2} + f(u_{j,k}). \end{aligned}$$

When considering the Neumann boundary condition

$$\frac{\partial u}{\partial x}(0) = \frac{\partial u}{\partial x}(k_x \delta_x) = 0,$$

we set $u_{j,-1} := u_{j,1}$ and $u_{j,k_x+1} := u_{j,k_x-1}$ for $j \geq 0$. Numerical experiments are performed with this scheme.

For the simulation, we solve the problem

$$\begin{cases} \partial_t^\alpha u = u_{xx} + u(1-u) & (t > 0, 0 < x < l), \\ \frac{\partial u}{\partial x}(t, 0) = \frac{\partial u}{\partial x}(t, l) = 0 & (t > 0), \\ u(0, x) = u_0(x) & (0 < x < l), \end{cases}$$

where $l > 0$ and

$$u_0(x) = \begin{cases} 0 & 0 < x < l_0, \\ \omega & l_0 \leq x < l \end{cases}$$

for some $l_0 \in (0, l)$. We compute the profile of the solution when $T > 0$ in the following step.

- (i) Compute the point $x^*(t)$ such that $u(t, x^*(t)) = 0.1$ for all $t > 0$;
- (ii) Find the minimal $T > 0$ such that $x^*(T) < x_0$ for a given constant $x_0 \in (0, l)$.

Moreover, the propagation speed of a numerical solution for α is computed by

$$c_\alpha^{\text{num}} := \frac{x^*(T) - x^*(T - \delta_t)}{\delta_t}.$$

4.2. Numerical results. The profile of the numerical solution is shown in Fig. 1. Initially, the front of the solution appears and moves from right to left in the spatio-temporal plots. In addition, the position of the front is linear with respect to time, indicating that the propagation advances at a rate $O(t)$. This suggests that the solutions exhibit behavior similar to that of a traveling wave solution as time progress. Moreover, we see that the front profile decays to 0 rapidly and approaches 1 slowly. This expected behavior reflects kinetics, where the system leaves the unstable state exponentially and approaches the stable state with an algebraic decay.

Next, we discuss the propagation speed of the numerical solution. In the case of the Fisher-KPP equation (i.e. $\alpha = 1$), it is known that when the initial state has a compact support or is a Heaviside function, the propagation speed of the solution is characterized by a minimal speed at which a traveling wave solution with a monotone positive profile exists [3, 27]. In KPP-type nonlinearity, the minimal speed for traveling waves is equal to

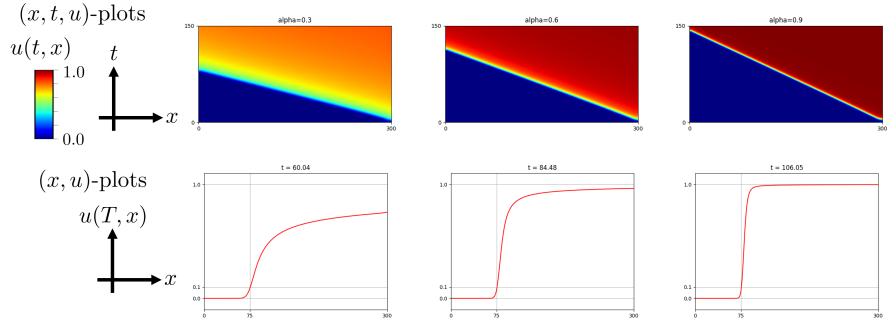


FIGURE 1. Spatio-temporal plots (top panels) and profiles of the numerical solutions (bottom panels). The parameters used are $l = 300$, $l_0 = 295$, $\omega = 0.01$, and $x_0 = 75$. The numerical simulations correspond to $\alpha = 0.3$ (left), $\alpha = 0.6$ (middle), and $\alpha = 0.9$ (right). The solution profiles are shown at $T = 60.04$ (left), $T = 84.48$ (middle), and $T = 106.05$ (right). See Subsection 4.1 for how to determine T .

the minimal linear speed, $c_1^* = 2\sqrt{f'(0)}$. Here, the minimal linear speed is the minimal speed at which a positive solution $\phi(\xi) = e^{\lambda\xi}$ ($\lambda > 0$) exists in the linearized equation

$$c\phi'(\xi) = \phi''(\xi) + f'(0)\phi(\xi) \quad (\xi \in \mathbb{R})$$

around the unstable state $u = 0$. Then, $\lambda > 0$ must satisfy $\lambda^2 - c\lambda + f'(0) = 0$, and it thus has a positive root λ if and only if $c \geq c_1^*$. In general, the property that the minimal speed of a traveling wave coincides with the minimal linear speed is called linear determinacy, and it has been confirmed in the propagation phenomena of various reaction-diffusion systems [3, 8, 27].

Let us assume that a similar hypothesis holds for the behavior of the solution of equation (1.1) by replacing the traveling wave solution with an asymptotic traveling wave solution. In other words, the propagation speed of the solution is regarded as being characterized by the minimal speed for the asymptotic traveling wave solution. We have not yet obtained results on the minimal speed, but we know from Theorem 2.2 that there exists at least an asymptotic traveling wave solution with $c \geq c_\alpha^*$. Moreover, from Lemma 3.9, we have the linear minimal speed c_α^* . Therefore, if linear determinacy holds, the propagation speed is expected to be characterized by c_α^* .

A graph of c_α^* and the propagation speed of the numerical solution c_α^{num} is shown in Fig. 2 (a). It is seen that c_α^{num} is close to c_α^* . The relative errors are shown in Fig. 2 (b). The error is less than 0.1 and comparatively small in the range where α is close to 1. This result suggests that c_α^* corresponds to the propagation speed of the solution in some sense.

5. DISCUSSION

In this paper, we discussed the effect of the time-fractional derivative on front solution propagation. By introducing a monostable nonlinear term to the time-fractional diffusion

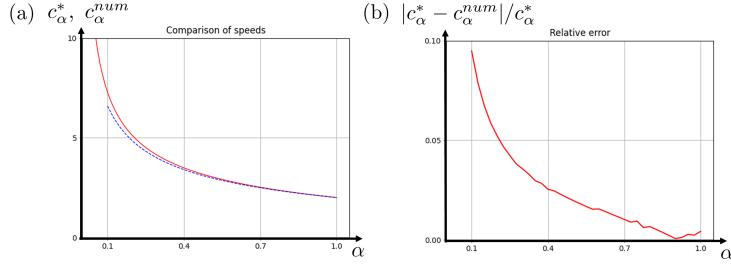


FIGURE 2. (a) Comparison of c_α^* (red line) and the propagation speed of the numerical solution (dashed line). The parameters are $l = 300$, $l_0 = 295$, $\omega = 0.01$, and $x_0 = 75$. The blue dashed line is constructed by connecting straight line segments computed for $\alpha = 0.1 + 0.025 \times k$ ($0 \leq k \leq 36$). (b) Relative errors.

equation, we demonstrated that the sub-diffusion property is not conserved with respect to the time order of propagation. Specifically, we introduced the concept of an asymptotic traveling wave solution to characterize solution propagation. Through the analysis of the asymptotic traveling wave solution, we gained insights into the dependence of α on the propagation speed and the shape of the solution, and thus gained an understanding the effect of the time-fractional derivative.

Although the asymptotic traveling wave solution appears essential object for analyzing solution propagation, we have not yet constructed a solution that converges to this solution. One challenge in proving convergence lies in determining the appropriate initial state. As the asymptotic traveling wave solution is not a solution to equation (1.1), estimating the initial conditions for convergence backward is not straightforward. A useful piece of information is that the asymptotic traveling wave solution serves as a sub-solution. The existence of a converging solution might be established if a well-constructed super-solution is introduced.

Several issues remain unresolved regarding the asymptotic traveling wave solution. One such issue is the solution's asymptotic behavior as the solution approaches a stable equilibrium. As mentioned in Section 1, convergence to a stable equilibrium at an algebraic decay rate is anticipated in the case of time-fractional derivatives. In a related work [33], although for a different problem setting, a bistable discontinuous nonlinear term was used to provide a concrete representation of the traveling wave solution and to construct a solution that decays algebraically to a constant state. It might be feasible to achieve asymptotic behavior that decays algebraically by adopting a constructive method, even in this problem setting, but this remains an open question for future investigation.

Finally, research related to sub-diffusion and proliferation/saturation effects is discussed. This includes studies on proliferation models with random walks, which differ from normal diffusion [16, 19, 20]. The time order of the propagation speed of the solution has been investigated, and was found to be $O(t^\alpha)$ for the model in [20] and $O(t)$ for the model in [16, 19]. These differences arise from the nonlocality in the time direction that describes the effect of sub-diffusion. As kinetics have the effect of propagating the solution in $O(t)$, it is necessary to model the time nonlocality more strongly to propagate the solution in $O(t^\alpha)$. The finding that no sub-diffusion properties remain for the time

order of the propagation in our setup is important information from a modeling point of view. A mathematically rigorous approach to incorporating time non-locality into the proliferation model for sub-diffusion behavior would provide insights for modeling.

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REFERENCES

- [1] F. ACHLEITNER, C. M. CUESTA, S. HITTMEIR, *Travelling waves for a non-local Korteweg–de Vries–Burgers equation*, Journal of Differential Equations **257** (2014), 720–758.
- [2] F. ANDREU-VAILLO, J. MAZÓN, J. D. ROSSI, J. J. TOLEDO-MELERO, Nonlocal diffusion problems, Math. Surveys Monogr. **165**, AMS, Providence, RI, 2010.
- [3] D.G. ARONSON, H.F. WEINBERGER, MULTIDIMENSIONAL NONLINEAR DIFFUSIONS ARISING IN POPULATION GENETICS, *Adv. Math.*, **30** (1978), 33–76.
- [4] P. W. BATES, X. CHEN, A. CHMAJ, *Heteroclinic solutions of a van der Waals model with indefinite nonlocal interactions*, Calc. Var., **24** (2005), 261–281.
- [5] P. W. BATES, P. C. FIFE, X. REN, X. WANG, *Traveling waves in a convolution model for phase transitions*, Arch. Ration. Mech. Anal., **138** (1997), 105–136.
- [6] C. BUCUR, E. VALDINOCI, Nonlocal diffusion and applications, Springer, 2016.
- [7] X. CABRÉ, J.-M. ROQUEJOFFRE, *The influence of fractional diffusion in Fisher-KPP equations*, Commun. Math. Phys., **320** (2013), 679–722.
- [8] C. CASTILLO-CHAVEZ, B. LI, H. WANG, *Some recent developments on linear determinacy*, Mathematical Biosciences and Engineering, **10** (2013), 1419–1436.
- [9] X. CHEN, J.-S. GUO *Existence and Asymptotic Stability of Traveling Waves of Discrete Quasilinear Monostable Equations*, Journal of Differential Equations, **184** (2002), 549–569
- [10] C. M. CUESTA, F. ACHLEITNER, *Addendum to “Travelling waves for a non-local Korteweg–de Vries–Burgers equation”* [J. Differential Equations 257 (3) (2014) 720–758], Journal of Differential Equations **262** (2017), 1155–1160.
- [11] J. COVILLE, L. DUPAIGNE, *On a non-local equation arising in population dynamics*, Proc. R. Soc. Edinb. **137A** (2007), 727–755.
- [12] S.-I. EI, J.-S. GUO, H. ISHII, C.-C. WU, *Existence of traveling waves solutions to a nonlocal scalar equation with sign-changing kernel*, Journal of Mathematical Analysis and Applications, **487(2)**, (2020), 124007.
- [13] S. EIDELMAN, A. KOCHUBEI, *Cauchy problem for fractional diffusion equations*, J. Differential Equations, **199** (2004), 211–255.
- [14] R. A. FISHER, *The genetical theory of natural selection: a complete variorum edition* (ed. J. H. Bennett), (Oxford University Press), (1999).
- [15] R. A. FISHER, The wave of advance of advantageous genes, Annals of eugenics **7** (1937), 355–369.
- [16] S. FEDOTOV, *Non-markovian random walks and nonlinear reactions: subdiffusion and propagating fronts*, Physical Review E, **81** (2010), 011117.
- [17] S. FEDOTOV, V. MÉNDEZ, *Continuous-time random walks and traveling fronts*, Physical Review E, **66** (2002), 030102(R).
- [18] G. FLORIDIA, Y. LIU AND M. YAMAMOTO, *Blowup in $L^1(\omega)$ -norm and global existence for time-fractional diffusion equations with polynomial semilinear terms*, Adv. Nonlinear Anal. **12** (2023), 20230121, 15 pp.

- [19] D. FROEMBERG, H. SCHMIDT-MARTENS, I. M. SOKOLOV, F. SAGUÉS, *Front propagation in $A + B \rightarrow 2A$ reaction under subdiffusion*, Physical Review E, **81** (2008), 011128.
- [20] D. FROEMBERG, H. SCHMIDT-MARTENS, I. M. SOKOLOV, F. SAGUÉS, *Asymptotic front behavior in an $A + B \rightarrow 2A$ reaction under subdiffusion*, Physical Review E, **83** (2011), 031101.
- [21] J. GARNIER, *Accelerating solutions in integro-differential equations*, SIAM J. Math. Anal. **43** (2011), 1955–1974.
- [22] C. Gui, M. Zhao, *Traveling wave solutions of Allen-Cahn equation with a fractional Laplacian*, Ann. I. H. Poincaré, **32** (2015), 785-812.
- [23] B. I. HENRY, T. A. LANGLANDS, S. L. WEARNE *Anomalous diffusion with linear reaction dynamics: From continuous time random walks to fractional reaction-diffusion equations*, Physical Review E, **74** (2006), 031116.
- [24] V. HUTSON, S. MARTINEZ, K. MISCHAIKOW, G.T. VICKERS, *The evolution of dispersal*, J. Math. Biol. **47** (2003), 483-517.
- [25] H. J. HAUBOLD, A. M. MATHAI, R. K. SAXENA, *Mittag-Leffler functions and their applications*, J. Appl. Math. (2011) 298628.
- [26] F. HAMEL AND N. NADIRASHVILI, *Entire solutions of the KPP equation*, Comm. Pure Appl. Math. **52** (1999), 1255–1276.
- [27] A. N. KOLMOGOROV, I. G. PETROVSKY AND N. S. PISKUNOV, *Etude de l'équation de la diffusion avec croissance de la quantité de matière et son application à un problème biologique*, Bjul. Moskow. Gos. Univ. **A1** (1937), 1-26.
- [28] A. KUBICA, K. RYSZEWSKA, M. YAMAMOTO, *Time-Fractional Differential Equations: A theoretical introduction*, Springer 2020.
- [29] Y. LIN, C. XU, *Finite difference/spectral approximations for the time-fractional diffusion equation*, J. Computational Physics, **225** (2007), 1533-1552.
- [30] Y. LUCHKO AND M. YAMAMOTO, *Comparison principles for the time-fractional diffusion equations with the Robin boundary conditions. Part I: Linear equations*, Fract. Calc. Appl. Anal., **26** (2023), 1504-1544.
- [31] S. MA, *Traveling wavefronts for delayed reaction-diffusion systems via a fixed point theorem*, J. Differential Equations, **171** (2001), 294-314.
- [32] F. MAINARDI, *The fundamental solutions for the fractional diffusion-wave equation*, Appl. Math. Lett. **9** (1996), 23–28.
- [33] Y. NEC, V. A. VOLPERT, A. A. NEPOMNYASHCHY, *Front propagation problems with sub-diffusion*, Discrete and Continuous Dynamical Systems, **27** (2010), 827-846.
- [34] Y.-J. SUN, W.-T. LI, Z.-C. WANG, *Traveling waves for a nonlocal anisotropic dispersal equation with monostable nonlinearity*, Nonlinear Analysis, **74** (2011), 814-826.
- [35] V.A. VOLPERT, A.A. NEPOMNYASHCHY, A. A. GOLOVIN, *Front-type solutions of fractional Allen-Cahn equation*, Physica D: Nonlinear Phenomena, **237** (2008), 3237-3251.
- [36] V.A. VOLPERT, Y. NEC, A.A. NEPOMNYASHCHY, *Exact solutions in front propagation problems with superdiffusion*, Physica D: Nonlinear Phenomena, **239** (2010), 134-144.
- [37] V. A. VOLPERT, Y. NEC, A. A. NEPOMNYASHCHY, *Fronts in anomalous diffusion-reaction systems*, Phil Trans R Soc A, **371** (2013), 20120179.
- [38] H. YAGISITA, *Existence and nonexistence of traveling waves for a nonlocal monostable equation*, Publ. Res. Inst. Math. Sci., **45** (2009), 925-953.

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